

# Sn-Doped Rutile TiO<sub>2</sub> Hollow Nanocrystals with Enhanced Lithium-Ion Batteries Performance

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**S** [Supporting Information](#page-6-0)

ABSTRACT: Hollow structures and doping of rutile  $TiO<sub>2</sub>$  are generally believed to be effective ways to enhance the performance of lithium-ion batteries. Herein, uniformly distributed Sn-doped rutile  $TiO<sub>2</sub>$  hollow nanocrystals have been synthesized by a simple template-free hydrothermal method. A topotactic transformation mechanism of solid TiOF, precursor is proposed to illustrate the formation of rutile  $TiO<sub>2</sub>$  hollow nanocrystals. Then, the Sn-doped rutile  $TiO<sub>2</sub>$  hollow nanocrystals are calcined and tested as anode in the lithium-ion battery. They deliver a highly reversible specific capacity of 251.3 mA h  $g^{-1}$  at 0.1 A  $g^{-1}$  and retain ∼110 mA h g $^{-1}$  after 500 cycles at a high current rate 5 A g $^{-1}$  (30 C), which is much higher than most of the reported work.

# 1. INTRODUCTION

Development of high-performance lithium-ion batteries (LIBs) with wider work-temperature range, increased power density, and improved safety is highly demanded for future portable electronics and electric and hybrid vehicles. $1,2$  $1,2$  $1,2$  Among various anode materials,  $TiO<sub>2</sub>$  has been considered as a potential anode candidate because of excellent Li-ion insertion/extraction reversibility with a neglectable volume expansion (3−4%) and a high operating voltage (>1 V vs Li<sup>+</sup>/Li),<sup>[3](#page-6-0)–[7](#page-6-0)</sup> which reduces the excessive formation of solid electrolyte interphase (SEI) on anodes. In addition,  $TiO<sub>2</sub>$  is suitable for large-scale application because of its natural abundance, low cost, and nontoxicity.<sup>[8](#page-6-0)</sup> Therefore, various polymorphs of  $TiO<sub>2</sub>$  (rutile, anatase, brookite, and bronze) have recently been studied as active materials for LIBs.<sup>[9](#page-6-0)</sup> Among them, rutile  $TiO<sub>2</sub>$  is the most thermodynamically stable form.<sup>[10](#page-6-0)</sup> Unfortunately, the lithium storage performance of rutile  $TiO<sub>2</sub>$  is still unsatisfactory, which is mainly derived from its highly anisotropic diffusion of Li<sup>+</sup> in bulk rutile  $(10^{-6} \text{ cm}^2 \text{ s}^{-1}$  along c-axis;  $10^{-15} \text{ cm}^2 \text{ s}^{-1}$  along abplane) $11,12$  and intrinsic limitation of storage capacity.<sup>[13,14](#page-6-0)</sup> To enhance the rate capacity, enormous researches have been devoted to overcome its intrinsic disadvantages. One of those effective approaches focuses on preparing nanosized hollow rutile  $TiO<sub>2</sub>$  to shorten the transport path for both  $Li<sup>+</sup>$  ions and electrons in TiO<sub>2</sub> nanostructures.<sup>[15](#page-6-0)−[17](#page-7-0)</sup> The hollow interior also provides an extra free space that can alleviate the structural strain and accommodate the volume variation. It is associated with the repeated Li<sup>+</sup> ions insertion/extraction process, leading to the improved cycling stability.<sup>[18](#page-7-0),[19](#page-7-0)</sup> Another effective strategy



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is to dope alien atoms (e.g., C, N, B, Fe, Nb, Mn, Sn, etc.) in  $TiO<sub>2</sub>$ .<sup>[8](#page-6-0),[20](#page-7-0)−[27](#page-7-0)</sup> The doping method is attracting more attention because it not only benefits the electron-transfer process but also offers more channels for Li<sup>+</sup> diffusion due to the slightly modified lattice of the doped  $TiO_2$ .<sup>[24](#page-7-0)</sup> Sn<sup>4+</sup> is taken as an example because the theoretical capacity of  $SnO<sub>2</sub>$  ( $\approx$ 790 mA h  $g^{-1}$ ) is much higher than that of rutile TiO<sub>2</sub>.<sup>[28,2928,29](#page-7-0)</sup> Sn(IV)ion doping into rutile  $TiO<sub>2</sub>$  at the atomic level can anticipate improved electrochemical performance with increased capacity and higher retention.

Although many kinds of rutile  $TiO<sub>2</sub>$  nanostructures with different morphologies have been intensively investigated, synthesis of hollow nanocrystals is still a great challenge by a template-free method.<sup>[30](#page-7-0)</sup> Herein, we report the preparation of Sn-doped rutile  $TiO<sub>2</sub>$  hollow nanocrystals by a facile templatefree hydrothermal method. When measured as anode in lithium-ion batteries, the Sn-doped rutile  $TiO<sub>2</sub>$  hollow nanocrystals exhibit a greatly improved electrochemical performance, including a high specific capacity, an excellent rate property, and a long-life cycling stability. This material showed a much enhanced capacity of 251.3 mA h  $g^{-1}$  after 150 cycles under 0.1 A g<sup>-1</sup>. Furthermore, a reversible capacity of ~110 mA h g<sup>−1</sup> after 500 cycles at 5 A g<sup>−1</sup> (30 C) can still be retained and the Coulombic efficiency (CE) maintained almost 100% during the 500 cycles, confirming the excellent rate property of the Sn-

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Figure 1. (a) XRD pattern of the Sn-doped rutile TiO<sub>2</sub> sample and the pure rutile TiO<sub>2</sub>. (b) TEM image of the Sn-doped rutile TiO<sub>2</sub> hollow nanocrystals. (c) Magnified TEM image of single Sn-doped rutile TiO<sub>2</sub> hollow nanocrystal. (d) Top view of the Sn-doped rutile TiO<sub>2</sub> hollow nanocrystal. (e) HRTEM image of the Sn-doped rutile TiO<sub>2</sub> hollow nanocrystal. Inset is the corresponding SAED pattern of the nanocrystal.

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doped rutile  $TiO<sub>2</sub>$  hollow nanocrystals. The excellent lithiumion storage performance benefits from hollow nanostructures, high crystallinity, and uniform Sn-ion doping.

# 2. EXPERIMENTAL SECTIONS

2.1. Preparation of the Sn-Doped Rutile TiO<sub>2</sub> Hollow Nanocrystals. The Sn-doped rutile  $TiO<sub>2</sub>$  hollow nanocrystals were obtained via a facile template-free hydrothermal method. In a typical process, 4 mmol  $SnCl<sub>4</sub>·5H<sub>2</sub>O$  was dissolved in 20 mL tetrabutyl titanate (TBOT) and vigorously stirred for 10 min. 4.5 mL HF solution (30%) was then added drop by drop in it. The solution was then transferred into a 40 mL Teflon tube and sealed into the autoclave, which was heated to 200 °C for 24 h. The obtained products were washed several times with ethanol and water after cooling to the room temperature. Finally, they were annealed at 500 °C for 2 h to remove the impurities. To investigate its growth process, time-dependent experiments were carried out for 10, 18, and 24 h, respectively. In addition, different amounts (0, 1, 2, 4, 6, and 8 mmol) of  $SnCl<sub>4</sub>·5H<sub>2</sub>O$  were also introduced to illustrate the formation of pure Sn-doped rutile  $TiO<sub>2</sub>$  hollow nanocrystals.

2.2. Electrochemical Characterization. Electrochemical measurements were carried out by using LIR2032 coin cells.<sup>[31](#page-7-0)</sup> The working electrode consisted of active material, acetylene black (super-P) and poly(vinylidene fluoride) with the weight ratio of 8:1:1. The assembly of the coin cells was performed in a glove box full of Ar. The Celgard 2400 film was used as separator. The electrolyte was composed of ethylene carbonate,

Accelerated Surface Area and Porosimeter 2000 analyzer was

was dried at 80 °C for 24 h before examination.

dimethyl carbonate (DMC), and ethylene methyl carbonate, with the volume ratio of 1:1:1 as solvent and  $LIPF<sub>6</sub>$  as solute with the concentration of 1 mol  $L^{-1}$ . The counter electrode of the coin cells was mainly composed of pure lithium foil. The galvanostatic charge/discharge curve of the coin cells was obtained on a LAND Cell Test system in a setting voltage range (2001A, Wuhan, China). A CHI760D electrochemical working station was used to carry out the cyclic voltammetry (CV) tests in the same setting voltage range with the charge/ discharge test rate of 0.2 mV s<sup>-1</sup>. Electrochemical impedance spectroscopy (EIS) patterns were studied by using the same working station in the CV tests. The frequency range was set from 100 kHz to 0.01 Hz, with an amplitude of 5 mV. After the cycle performance test, the coin cells were disassembled and washed with dimethyl carbonate (DMC) to remove the electrolyte from the electrode surface. After that, the sample

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Figure 2. TEM images of TiO<sub>2</sub> samples synthesized with (a) 0 mmol, (b) 2 mmol, (c) 4 mmol, (d) 6 mmol, and (e) 8 mmol SnCl<sub>4</sub>·5H<sub>2</sub>O; (f) HRTEM image of  $SnO<sub>2</sub>$  quantum dots from (e).



Figure 3. (a) STEM image of single Sn-doped TiO<sub>2</sub> nanocrystal. (b-d) EDS mappings of Ti, O, and Sn elements. The scale bar is 100 nm.

used to measure the Brunauer−Emmett−Teller specific surface area of the sample.

#### 3. RESULTS AND DISCUSSION

The phase composition of Sn-doped  $TiO<sub>2</sub>$  sample was confirmed by XRD ([Figure 1](#page-1-0)a). All of the diffraction peaks can be indexed to rutile  $TiO<sub>2</sub>$  (JCPDS No. 21-1276), and no other peaks corresponding to  $SnO<sub>2</sub>$  were observed in the XRD pattern, preliminary indicating effective doping of tin in rutile  $TiO<sub>2</sub>$  lattice. It is noted that the peaks slightly shifted to smaller angles, as observed by Rogach et al. $^{25,32}$  $^{25,32}$  $^{25,32}$  The lattice expansion is attributed to a substitution of the Ti site in the rutile by Sn due to a larger ionic radius of  $Sn^{4+}$  (55 pm) than that of  $Ti^{4+}$  (42 pm).[33](#page-7-0) [Figure 1](#page-1-0)b−d shows the typical TEM images of Sndoped rutile TiO2 with a size of 100−150 nm. The nanocrystals have clear hollow interior voids, with some holes on their surfaces ([Figure S1\)](http://pubs.acs.org/doi/suppl/10.1021/acsomega.7b01340/suppl_file/ao7b01340_si_001.pdf), which is possibly derived from the etching of HF solution. The selected area electron diffraction (SAED) pattern corresponding to the single hollow nanocrystal demonstrates its single-crystalline nature (inset in [Figure 1](#page-1-0)e). The HRTEM image exhibits clear lattice fringes, indicating high

crystallinity of the nanocrystal ([Figure 1e](#page-1-0)). The lattice spacing is ~0.33 nm, corresponding to the (110) plane of rutile  $\text{TiO}_2$ .<sup>[13](#page-6-0)</sup>

Actually, when no Sn precursor was added to the starting solution, only anatase  $TiO<sub>2</sub>$  nanosheets were prepared (Figures 2a and [S2\)](http://pubs.acs.org/doi/suppl/10.1021/acsomega.7b01340/suppl_file/ao7b01340_si_001.pdf). Incorporation of Sn element into  $TiO<sub>2</sub>$  lattice was accompanied with a phase transformation from anatase to rutile (Figure 2b,c). When 2 mmol  $SnCl<sub>4</sub>·5H<sub>2</sub>O$  was dissolved in TBOT, it is noticed that some hollow nanocrystals (labeled by arrow in Figure 2b) appeared in the sample of anatase  $TiO<sub>2</sub>$ nanosheets. The analysis of phase composition reveals the coexistence of anatase and rutile  $TiO<sub>2</sub>$  ([Figure S2](http://pubs.acs.org/doi/suppl/10.1021/acsomega.7b01340/suppl_file/ao7b01340_si_001.pdf)). Furthermore, the anatase nanosheets were completely converted into Sn-doped rutile hollow nanocrystals with the injection of 4 mmol  $SnCl<sub>4</sub>·5H<sub>2</sub>O$ , which was also supported by the XRD pattern in [Figure S2](http://pubs.acs.org/doi/suppl/10.1021/acsomega.7b01340/suppl_file/ao7b01340_si_001.pdf). Only rutile  $TiO<sub>2</sub>$  sample was synthesized in the condition. However, excessive Sn precursor easily resulted in phase separation of  $TiO<sub>2</sub>$  and  $SnO<sub>2</sub>$ . As shown in Figure 2d–f, a large amount of  $SnO<sub>2</sub>$  nanocrystals uniformly dispersed on the surfaces of  $TiO<sub>2</sub>$  nanosheets. The emergence of  $TiO<sub>2</sub>$  nanosheets instead of hollow nanocrystals means that the rutile  $TiO<sub>2</sub>$  returned back to anatase phase. The

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Figure 4. Patterns of (a) Sn 3d peaks, (b) O 1s peaks, and (c) Ti 2p peaks of the Sn-doped rutile TiO<sub>2</sub>.



Figure 5. The TEM images and XRD patterns of the as-obtained products at different time duration: (a, b) 10 h, (c, d) 18 h, and (e, f) 24 h. The scale bar is 200 nm.

corresponding XRD patterns in [Figure S2](http://pubs.acs.org/doi/suppl/10.1021/acsomega.7b01340/suppl_file/ao7b01340_si_001.pdf) supported the reverse transformation and the appearance of  $SnO<sub>2</sub>$ . The quantum dots, uniformly deposited on anatase  $TiO<sub>2</sub>$  nanosheets, were measured by HRTEM ([Figure 2](#page-2-0)f). The lattice spacing was ∼0.339 nm, corresponding to the (110) plane of  $SnO<sub>2</sub>$ . Therefore, a Sn precursor, however, with very narrow range of  $Sn^{4+}$  concentration, plays a crucial role in the formation of pure Sn-doped rutile  $TiO<sub>2</sub>$  hollow nanocrystals. In addition, for eliminating the effect of Cl<sup>−</sup> for the formation of rutile  $TiO<sub>2</sub>$  hollow nanocrystals, several kinds of metal chlorides were introduced instead of  $SnCl<sub>4</sub>$  in this reaction. As a result, only anatase  $TiO<sub>2</sub>$  products could be synthesized [\(Figure S3\)](http://pubs.acs.org/doi/suppl/10.1021/acsomega.7b01340/suppl_file/ao7b01340_si_001.pdf), which further illustrates the effect of doping of Sn.

To further validate the Sn doping in rutile  $TiO<sub>2</sub>$  hollow nanocrystals, a STEM image of single nanocrystal was recorded [\(Figure 3](#page-2-0)a). The corresponding EDS elemental mappings were conducted to reveal the elemental composition and distribution status in the hollow nanocrystal. [Figure 3b](#page-2-0)−d confirms that Sn element was homogeneously dispersed in the  $TiO<sub>2</sub>$  hollow nanocrystal. As mentioned above, the as-obtained rutile  $TiO<sub>2</sub>$ hollow nanocrystals have smooth surfaces and clear lattice fringes [\(Figure S4](http://pubs.acs.org/doi/suppl/10.1021/acsomega.7b01340/suppl_file/ao7b01340_si_001.pdf)). Therefore, it can be concluded that the Sn element was well doped into the rutile  $TiO<sub>2</sub>$  lattice.

The molar ratio of Sn element in TiO<sub>2</sub> sample was 8.7%, as calculated by fitting of XPS spectra. The peaks centered at 487.0 and 494.9 eV can be indexed to Sn 3d peaks (Figure 4a), $34,35$  $34,35$  $34,35$  corresponding to  $Sn^{4+}$  bound to oxygen in the Sndoped rutile  $TiO<sub>2</sub>$ . Because of the structural similarity between the rutile  $SnO<sub>2</sub>$  and  $TiO<sub>2</sub>$  ([Figure S5 and Table S1](http://pubs.acs.org/doi/suppl/10.1021/acsomega.7b01340/suppl_file/ao7b01340_si_001.pdf)),  $^{10,14,36}$  $^{10,14,36}$  $^{10,14,36}$  $^{10,14,36}$  $^{10,14,36}$  $^{10,14,36}$  the Sn elements would occupy a part of the Ti sites in the doping process. In Figure 4b, the slightly asymmetric spectrum of O 1s can be resolved into three peaks. The O 1s peaks centered at 529.5 and 530.1 eV can be typically assigned to Ti−O bond in the rutile TiO<sub>2</sub> structure<sup>[37,38](#page-7-0)</sup> and the Sn−O bond,<sup>[31](#page-7-0)</sup> respectively. The peak at 532.05 may result from the band of the adsorbed hydroxyl (−OH) groups on the nanomaterial surface.<sup>[31](#page-7-0)</sup> Ti 2p peaks (Figure 4c) of the Sn-doped rutile  $TiO<sub>2</sub>$ show two peaks, corresponding to  $Ti^{4+}$  (Ti  $2p_{1/2}$  at 464.6 eV and Ti 2p<sub>3/2</sub> at 458.7 eV).<sup>[37](#page-7-0)–[39](#page-7-0)</sup> The UV–vis spectrum of Sndoped rutile TiO<sub>2</sub> presented a narrowed band gap of ~0.1 eV compared to that of pure rutile  $TiO<sub>2</sub>$  ([Figure S6\)](http://pubs.acs.org/doi/suppl/10.1021/acsomega.7b01340/suppl_file/ao7b01340_si_001.pdf), which was consistent with the theoretical simulation [\(Figures S7 and S8\)](http://pubs.acs.org/doi/suppl/10.1021/acsomega.7b01340/suppl_file/ao7b01340_si_001.pdf).

To shed light on the formation process of the rutile  $TiO<sub>2</sub>$ hollow nanocrystals, experiments were carried out at 200 °C for different reaction times to monitor the morphology evolution of nanoparticles. As shown in Figure 5a, solid nanocrystals were



Figure 6. Schematic illustration of forming the rutile Sn-doped TiO<sub>2</sub> hollow nanocrystals via the template-engaged topotactic transformation process from Sn-doped TiOF<sub>2</sub> precursor.

obtained for 10 h, which were indexed to be  $TiOF<sub>2</sub>$  (JCPDS No. 08-0060) by the corresponding XRD pattern ([Figure 5b](#page-3-0)). When the reaction time was extended to 18 h [\(Figure 5c](#page-3-0),d), some new diffraction peaks, corresponding to the rutile  $TiO<sub>2</sub>$ , also appeared besides the peaks of  $TiOF<sub>2</sub>$ , indicating the conversion from  $TiOF<sub>2</sub>$  precursor to rutile  $TiO<sub>2</sub>$ . This is possibly due to the hydrolysis of  $TiOF_2$ .<sup>[40](#page-7-0)-[43](#page-7-0)</sup> More importantly, a part of solid  $TiOF<sub>2</sub>$  nanocrystals transformed into hollow structures, which still kept the similar external morphologies. The conversion process reflected that the formation of hollow nanocrystals through in situ transformation of TiOF<sub>2</sub> precursor.<sup>[43](#page-7-0)</sup> The phase transformation from TiOF<sub>2</sub> to rutile TiO<sub>2</sub> was completed within 24 h. Then, uniform TiO<sub>2</sub> hollow nanocrystals were obtained ([Figure 5e](#page-3-0),f).

Furthermore, a sacrificial self-template mechanism was proposed to illustrate the formation process of rutile  $TiO<sub>2</sub>$ hollow nanocrystals.<sup>[44,45](#page-7-0)</sup> First, the Sn-doped TiOF<sub>2</sub> crystallites (Figures 6 and [S9\)](http://pubs.acs.org/doi/suppl/10.1021/acsomega.7b01340/suppl_file/ao7b01340_si_001.pdf) were gradually hydrolyzed to form Sn-doped TiO<sub>2</sub> nanocrystals in a hydrothermal reaction.<sup>[42](#page-7-0)</sup> These primary nanocrystals uniformly deposited on the surfaces of Sn-doped TiOF<sub>2</sub> crystallites ([Figure S10\)](http://pubs.acs.org/doi/suppl/10.1021/acsomega.7b01340/suppl_file/ao7b01340_si_001.pdf) and then formed TiO<sub>2</sub> shells. Because of the restriction from the template, the asobtained rutile  $TiO<sub>2</sub>$  showed the same shape with the  $TiOF<sub>2</sub>$ precursor. The  $TiO<sub>2</sub>$  shells became thicker with the increasing reaction time, whereas the solid  $TiOF<sub>2</sub>$  templates were continually consumed inward. Finally, the  $TiOF<sub>2</sub>$  precursor completely disappeared. Therefore, the template-engaged topotactic transformation mechanism is responsible for the

formation of rutile  $TiO<sub>2</sub>$  hollow nanocrystals.<sup>[46,47](#page-7-0)</sup> When no SnCl<sub>4</sub> was added in the reaction system, only anatase  $TiO<sub>2</sub>$ nanosheets can be obtained. Excessive  $SnCl<sub>4</sub>$  easily led to the formation of  $SnO<sub>2</sub>$  quantum dots instead of doping in the rutile  $TiO<sub>2</sub>$  hollow nanocrystals.

The lithium storage properties of the Sn-doped rutile  $TiO<sub>2</sub>$ hollow nanocrystals were then evaluated. [Figure 7a](#page-5-0) shows the representative discharge−charge curves of the crystals during the 1st, 2nd, 5th, and 150th cycle at a voltage range from 0.01 to 3.0 V with a current rate of 0.1 A  $g^{-1}$ . Two major voltage plateaus at ∼1.4 and ∼1.0 V vs Li/Li<sup>+</sup> are presented in the first discharge curve, consistent with other rutile  $TiO<sub>2</sub>$  electro- $des.<sup>22,48</sup>$  $des.<sup>22,48</sup>$  $des.<sup>22,48</sup>$  The results were also supported by the CV results [\(Figure S11\)](http://pubs.acs.org/doi/suppl/10.1021/acsomega.7b01340/suppl_file/ao7b01340_si_001.pdf), which are related to the Li<sup>+</sup> insertion into the octahedral sites of the TiO<sub>6</sub> octahedrons.<sup>[49](#page-7-0)</sup> Despite the two major peaks in the CV patterns, some unknown peaks have also been observed, which are possibly derived from some side reactions in the electrode. In addition, the absence of peaks at ∼0.6 and ∼1.3 V, which are characteristic of the dealloying process of  $Li<sub>x</sub>Sn$  and the partially reversible reaction of  $SnO<sub>2</sub>$ with  $Li^{+,50-53}$  $Li^{+,50-53}$  $Li^{+,50-53}$  $Li^{+,50-53}$  $Li^{+,50-53}$  thus, indicates the uniform doping of tin in the , rutile TiO<sub>2</sub> lattice. A high initial capacity of 596.9 mA h  $g^{-1}$  is exhibited in the initial discharge process, which is much higher than the theoretical capacity of the pure rutile  $TiO<sub>2</sub>$  (TiO<sub>2</sub> +  $x$ Li +  $xe^-$  → Li<sub>x</sub>TiO<sub>2</sub> ( $x \le 1$ )).<sup>[8](#page-6-0)</sup> The hollow structures with a high specific surface area and a large pore volume [\(Figure S12\)](http://pubs.acs.org/doi/suppl/10.1021/acsomega.7b01340/suppl_file/ao7b01340_si_001.pdf) and doping of tin offer more active sites for  $Li^{+}$  storage.<sup>[18](#page-7-0),[54,55](#page-8-0)</sup> In the following process, a capacity of 337.9 mA h  $g^{-1}$  was

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Figure 7. Electrochemical properties of the Sn-doped rutile TiO<sub>2</sub> sample: (a) discharge−charge curves of them during the 1st, 2nd, 5th, and 150th cycle; (b) cycle performances at 100 mA  $g^{-1}$  in the 0.01–3 V range of the annealed and unannealed Sn-doped rutile TiO<sub>2</sub>; (c) electrochemical impedance spectroscopies of annealed and unannealed Sn-doped rutile TiO<sub>2</sub>; (d) rate capabilities of the annealed and unannealed rutile TiO<sub>2</sub> at a current density ranging from 100 to 5000 mA h g<sup>−1</sup>; (e) the cycling performance of annealed Sn-doped rutile TiO<sub>2</sub> at the rate of 5000 mA h g<sup>−1</sup> (30 C, 1 C = 173 mA  $g^{-1}$ ).

obtained. The irreversible capacity loss is about 43.4%. This can be derived from the deep trapping of inserted Li<sup>+</sup> after large volume strain in the crystal structure, as well as the formation of SEI film.<sup>[48](#page-7-0)</sup> After the first several cycles, the specific capacity of the Sn-doped rutile  $TiO<sub>2</sub>$  hollow nanocrystals electrode is stabilized (Figure 7b). In the second charge/discharge cycle, a discharge capacity of 337.9 mA h  $g^{-1}$  is delivered, followed by a charge capacity of 333.9 mA h  $g^{-1}$ , leading to a high Coulombic efficiency (CE) of about 98.8%. The discharge capacity sustains at about 251.3 mA h g<sup>−</sup><sup>1</sup> after 150 cycles at a current rate of 0.1 A  $g^{-1}$ , which is superior to most other nanostructured rutile  $TiO<sub>2</sub>$ -based anode materials [\(Table S2\)](http://pubs.acs.org/doi/suppl/10.1021/acsomega.7b01340/suppl_file/ao7b01340_si_001.pdf). Because SnO<sub>2</sub> has a higher theoretical capacity than  $TiO<sub>2</sub>$ ,  $Sn<sup>4+</sup>$  doping in the rutile  $TiO<sub>2</sub>$  lattice naturally leads to an enhanced reversible capacity  $(Li^+ + SnO_2 \rightarrow Sn + Li_2O; Sn + xLi \rightarrow Li_sSn (x < 4.4)).$ <sup>[51](#page-7-0),[53](#page-7-0)</sup> In addition, annealing treatment also benefited the increase in  $Li^{+}$ storage property in the prepared Sn-doped  $TiO<sub>2</sub>$  hollow nanocrystals (Figure 7b). This is due to the improved crystallinity and absence of side reactions, related with the oxygen vacancies, after annealing. In comparison, the undoped  $TiO<sub>2</sub>$  nanosheets were also synthesized via a similar hydrothermal method without  $SnCl<sub>4</sub>$  in the reaction system, which exhibited inferior electrochemical activity [\(Figure S13](http://pubs.acs.org/doi/suppl/10.1021/acsomega.7b01340/suppl_file/ao7b01340_si_001.pdf)). Electrochemical impedance spectroscopy (EIS) measurement is carried out to study the charge-transfer impedance of Sn-

doped rutile  $TiO<sub>2</sub>$  hollow nanocrystals electrode (Figure 7c). The Nyquist plots of the annealed and unannealed hollow nanocrystals are obtained with the amplitude of 5 mV and a frequency ranging from 0.01 Hz to 100 kHz. It is clear that the semicircle of the annealed Sn-doped rutile  $TiO<sub>2</sub>$  hollow nanocrystals is smaller than that of the unannealed ones, demonstrating a smaller solid-state interface resistance of the sample after annealing. Inset in Figure 7c shows an equivalent circuit. The small intercept  $(R_s)$  at the  $Z'$  axis represents the total resistance of the electrolyte, separator, and electrical contacts.  $53$  The semicircle corresponds to  $R_1$ , representing charge-transfer impedance on the electrode−electrolyte inter-face.<sup>[4](#page-6-0)</sup> The sloping line, corresponding to  $W$  (Warburg impedance), indicates the diffusion of  $Li<sup>+</sup>$  into the active materials.<sup>[2](#page-6-0)</sup> The CPE<sub>1</sub> is the constant phase-angle element, involving double-layer capacitance, which is relevant with the surface property of the electrode.<sup>[52](#page-7-0)</sup>  $C_1$  represents the capacitance caused by Li<sup>+</sup> transport process in the electrode material.<sup>[16](#page-6-0)</sup> The  $R_1$  of annealed Sn-doped TiO<sub>2</sub> hollow nanocrystals was 64.11  $\Omega$  cm<sup>-2</sup>, which is much smaller than that of unannealed electrode (160.9  $\Omega$  cm<sup>-2</sup>) and means well charge transport property on the electrode−electrolyte surface. It is due to the doping of  $Sn^{4+}$  and high crystallinity. The improved conductivity benefits the electrons and Li<sup>+</sup> transport in the electrode.<sup>[8](#page-6-0)</sup> Meanwhile, the hollow nanostructures can <span id="page-6-0"></span>offer more sites for ions insertion. Therefore, a high specific capacity can be expected.

The excellent rate behavior of the sample was further investigated at different current rates (0.1−5  $\rm \AA{g^{-1}};$  [Figure 7d](#page-5-0)). Ten charge/discharge cycles were tested at each step from 0.1 to 5 A  $g^{-1}$ . Obviously, the Sn-doped rutile TiO<sub>2</sub> hollow nanocrystals exhibited much more improved rate performances than unannealed sample at each current rate. The specific capacities for Sn-doped rutile  $TiO<sub>2</sub>$  hollow nanocrystals are 313.4, 291.4, 256.5, 235.5, and 201.4 mA h g<sup>-1</sup> at the rate of 0.1, 0.2, 0.5, 1, and 2 A  $g^{-1}$ , respectively. Moreover, even at the highest current of 5 A  $\rm g^{-1}$ , a high capacity of 155.6 mA h  $\rm g^{-1}$  is still obtained. Furthermore, the capacity can be returned to the initial value of 295.6 mA h  $g^{-1}$  after measuring at high current rate of 30 C, indicating their good structural stability. To evaluate the cyclability at a high current rate, the Sn-doped rutile TiO<sub>2</sub> hollow nanocrystals electrode was discharged− charged at 30 C over 500 cycles. It is clear that this sample retained a reversible capacity of ~110 mA h g<sup>-1</sup> after 500 cycles at about 30 C [\(Figure 7e](#page-5-0)). The CE remained almost 100% during the 500 cycles, indicating excellent stability of the electrode. More importantly, the hollow structure of Sn-doped rutile  $TiO<sub>2</sub>$  nanocrystals was well maintained after 100 cycles under 1 A  $g^{-1}$ , which further illuminates the stability of the electrode<sup>[56](#page-8-0)</sup> [\(Figure S14](http://pubs.acs.org/doi/suppl/10.1021/acsomega.7b01340/suppl_file/ao7b01340_si_001.pdf)). In comparison with the reported rutile  $TiO<sub>2</sub>$ , the Sn-doped rutile  $TiO<sub>2</sub>$  hollow nanocrystals exhibited significant rate capabilities and cycling stability ([Table](http://pubs.acs.org/doi/suppl/10.1021/acsomega.7b01340/suppl_file/ao7b01340_si_001.pdf) [S2\)](http://pubs.acs.org/doi/suppl/10.1021/acsomega.7b01340/suppl_file/ao7b01340_si_001.pdf).

### 4. CONCLUSIONS

In summary, Sn-doped rutile  $TiO<sub>2</sub>$  hollow nanocrystals were successfully synthesized by a facile hydrothermal method. The hollow nanostructures were transformed from solid  $TiOF<sub>2</sub>$ nanocrystals by hydrolysis. Owing to tin doping, uniform hollow structure, and high crystallinity, the  $TiO<sub>2</sub>$  products demonstrated remarkable rate capabilities and excellent longlife cycling stability at a high current rate. They delivered a high reversible specific capacity of 251.3 mA h  $g^{-1}$  at the current rate of 0.1 A g<sup>−1</sup>. Moreover, they still exhibited as high as ~110 mA h g<sup>-1</sup> over 500 cycles at the current rate of 5 A g<sup>-1</sup> (30 C). These electrochemical performances are superior to those of most rutile  $TiO<sub>2</sub>$  reported so far. Therefore, it is expected that the Sn-doped rutile  $TiO<sub>2</sub>$  hollow nanocrystals will be a potential candidate as the anode material.

# ■ ASSOCIATED CONTENT

#### **S** Supporting Information

The Supporting Information is available free of charge on the [ACS Publications website](http://pubs.acs.org) at DOI: [10.1021/acsomega.7b01340](http://pubs.acs.org/doi/abs/10.1021/acsomega.7b01340).

Figures S1−S14; Tables S1 and S2; references ([PDF\)](http://pubs.acs.org/doi/suppl/10.1021/acsomega.7b01340/suppl_file/ao7b01340_si_001.pdf)

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#### Author Contributions

The manuscript was written through contributions of all of the authors. All of the authors have given approval to the final version of the manuscript.

# Notes

The authors declare no competing financial interest.

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