Cellular/Molecular

Kv7/KCNQ Channels Control Action Potential Phasing of Pyramidal Neurons during Hippocampal Gamma Oscillations *In Vitro*

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While the synaptic mechanisms involved in the generation of *in vitro* network oscillations have been widely studied, little is known about the importance of voltage-gated currents during such activity. Here we study the role of the M-current (I_M) in the modulation of network oscillations in the gamma-frequency range (20 – 80 Hz). Kv7/KCNQ subunits, the molecular correlates of I_M , are abundantly expressed in CA1 and CA3 pyramidal neurons, and I_M is an important modulator of pyramidal neuron firing. Using hippocampal slices, we recorded field activity and pyramidal neuron action potential timing during kainate-induced gamma oscillations. Application of the specific I_M blocker XE991 causes a significant reduction of gamma oscillation amplitude with no significant change in oscillation frequency. Concomitant CA3 pyramidal neuron recordings show a significant increase in action potential frequency during ongoing gamma oscillations after the application of XE991. This increase is associated with a significant loss of periodicity of pyramidal neuron action potentials relative to the phase of the gamma oscillations. Using dynamic clamp, we show that I_M acts to improve the periodicity of action potential timing and to decrease action potential frequency. We further validate these results in a compartmental model of a pyramidal neuron. Our work suggests that I_M modulates gamma oscillations by regulating the phasing of action potential firing in pyramidal neurons.

Introduction

Acute brain slice preparations can provide reduced but highly accessible models of neural networks of higher order. The ability to induce gamma oscillations in brain slices (Fisahn et al., 1998) has allowed a dissection of the neural mechanisms involved in the generation of oscillatory patterns (Fisahn, 2005). Cholinergically induced in vitro gamma oscillations show several features that are observed in in vivo intrahippocampal gamma oscillations (e.g., pyramidal neurons fire at low rates but phase locked to the oscillation, rhythm is generated in area CA3 and propagates to CA1, etc.) (Mann and Paulsen, 2005). Gamma oscillations can also be induced by activation of kainate receptors, and the resulting rhythmic activity seems to involve synaptic mechanisms similar to cholinergically induced oscillations (Traub et al., 2002; Fisahn et al., 2004). While several previous studies have focused on the role of synaptic activity and neural connectivity in the generation of gamma rhythms in vitro (Fisahn et al., 2004; Mann et al., 2005; Oren et al., 2006), fewer studies addressed the importance of subcellular mechanisms, in particular K⁺ currents, for the modulation of gamma oscillations (Fisahn et al., 2002).

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Together with the hyperpolarization-activated current (I_h) , the M-current (I_M) is a prime candidate for a subcellular mechanism capable of modulating network synchrony. The M-current is a slowly activating and non-inactivating voltage-dependent K⁺ current abundantly expressed in hippocampal pyramidal neurons (Halliwell and Adams, 1982; Pedarzani and Storm, 1995). In the hippocampus, perisomatic I_M is known to control the excitability of pyramidal neurons (Gu et al., 2005; Yue and Yaari, 2006) and blocking $I_{\rm M}$ in CA3 pyramidal neurons acts to increase action potential (AP) afterdepolarization and promotes burst firing (Vervaeke et al., 2006). Recently an increasing number of studies has focused on $I_{\rm M}$, as a new generation of antiepileptic and cognitive-enhancing drugs targeting this current became available. However, $I_{\rm M}$ is present in numerous different brain nuclei and its role in network function is complex and poorly understood. A global attenuation or augmentation of $I_{\rm M}$ to treat conditions that affect particular brain areas may cause unwanted effects in nonaffected regions and attempts to improve cognition using I_M blockers largely failed, underscoring our lack of understanding of I_M function. For example, linopirdine, an I_M blocker and a neurotransmitter release enhancer, appeared to be a promising therapeutic agent for memory deficits a decade ago (Lamas et al., 1997) but failed to deliver meaningful results in clinical trials (Rockwood et al., 1997; Börjesson et al., 1999). Another more potent I_M blocker and neurotransmitter release enhancer, XE991 [10,10-bis(4-pyridinylmethyl)-9(10H)-anthracenone], is currently undergoing clinical trials, with no significant results obtained so far (Romero et al., 2004). To understand the discrepancy between the known effects of $I_{\rm M}$ on the cellular level and the apparent ineffectiveness of therapies targeting $I_{\rm M}$ it is fundamen-

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tal to address its role in cognition-relevant activity in local networks. Here we demonstrate that kainate-induced gamma oscillation in hippocampal slices are reduced by $I_{\rm M}$ antagonists as a result of decreased phase coupling of CA3 pyramidal neuron APs when Kv7/KCNQ channels are blocked. Furthermore, we show that selective restoration of $I_{\rm M}$ in single pyramidal neurons by dynamic clamp (DC) increases AP phase coupling back to control levels.

Materials and Methods

Electrophysiology. Electrophysiological experiments were performed in hippocampal slices of Sprague Dawley rats (postnatal day 21-24). All procedures followed Swedish and Karolinska Institute guidelines for the care and use of laboratory animals. Rats were decapitated and the brains rapidly removed and placed in ice-cold artificial CSF (ACSF) containing the following (in mM): 2.49 KCl, 1.43 NaH₂PO₄, 26 NaHCO₃, 10 glucose, 252 sucrose, 1 CaCl₂, 4 MgCl₂. The brain was then hemisected and trimmed and horizontal 400-µm-thick hippocampal slices were obtained using a vibratome. Slices were kept in a submerged holding chamber containing ACSF (in mM: 124 NaCl, 3.5 KCl, 1.25 NaH₂PO₄, 1.5 MgCl₂, 1.5 CaCl₂, 30 NaHCO₃, 10 glucose), constantly bubbled with 95% O2 and 5% CO2, incubated at 35°C for 1 h, and then allowed to cool to room temperature. For recordings, slices were transferred to a submerged recording chamber (28°C) under an upright microscope equipped with infrared and differential interference contrast optics and perfused with oxygenated ACSF at a rate of 1-1.25 ml/min. We opt for the usage of a submerged chamber rather than an interface chamber to be able to perform patch-clamp recordings from cells under visual guidance. While oscillations have a better signal-to-noise ratio in interface chambers, there are no differences in oscillation frequency and in pyramidal neuron firing (Fisahn et al., 2004). The differences in power observed between submerged versus interface recording chambers appear to reflect different oxygenation levels but synaptic transmission is not significantly different between the two recording configurations (Hájos et al., 2009). Extracellular field and intracellular whole-cell patch-clamp recordings were performed using either a Multiclamp 200B or an Axopatch 200B and an Axoclamp 2A amplifier (Molecular Devices). Glass pipettes with tip resistances between 3 and 5M Ω were used. Data were digitized using WinWCP (Dr. John Dempster, http://spider.science. strath.ac.uk/sipbs/software_ses.htm) or a custom acquisition software developed in LabView (National Instruments) and analyzed in Matlab (Mathworks). Extracellular field recordings were made in stratum radiatum of the hippocampal area CA3 and whole-cell patch-clamp recordings from somata of pyramidal neurons of area CA3 as described previously (Fisahn et al., 2002). For field and cell-attached recordings pipettes were filled with ACSF and, for whole-cell patch-clamp recordings, with a K-gluconate-based internal solution (in mM: 17.5 KCl, 122.5 K-gluconate, 9 NaCl, 1 MgCl₂, 3 Mg-ATP, 0.3 GTP-Tris, 1 HEPES, 0.2 EGTA; pH was adjusted to 7.2 using KOH). Linopirdine, ZD7288, XE991, and tetrodotoxin were purchased from Tocris Bioscience.

Analysis. Fast Fourier transformations for power spectra were computed from 120-s-long data traces and gamma oscillation power was measured as the integrated power spectra between 20 and 80 Hz. The low cutoff frequency of 20 Hz was chosen because of the low recording temperature (28°C) used in our experiments (at this temperature M-current traces could be fit with a single exponential; see Results). It has been demonstrated that the frequency of gamma oscillations decreases by 3.3 Hz for each 1°C temperature decrease (Dickinson et al., 2003; Fisahn et al., 2009). The phase of APs with regard to field oscillation cycle was analyzed using a template-based algorithm for field cycles (Clements and Bekkers, 1997) and a derivative threshold-based algorithm for detecting action potentials. AP-phase histograms were fitted with third-order polynomial functions as follows:

$$y = ax^3 + bx^2 + cx + d,$$
 (1)

and the third-order coefficient *a* used as a measure of the histogram "peakedness" (AP synchronicity or randomness). Data values are re-

ported as means \pm SEM. Student's paired *t* test was used for statistics unless otherwise stated.

Dynamic clamp. To assess the functional effects of the M-current in CA3 pyramidal neurons during gamma oscillations, we have simulated $I_{\rm M}$ in pyramidal neurons using the dynamic clamp technique. This method was implemented on a second computer running a Linux kernel modified by the Real Time Application Interface for Linux from the Politecnico di Milano Institute-Dipartimento di Ingegneria Aerospaziale (http://www.rtai.org) and custom-made software that reads membrane voltage and generates current commands at a 40 kHz refresh rate using a Data Acquisition data card (National Instruments) and drivers from the Linux Control And Measurement Device Interface (COMEDI, www. comedi.org) (Leao et al., 2005). To block the native $I_{\rm M}$ of single cells we supplemented our internal solution with 0.5 mM BaCl₂ (Lamas et al., 1997). We used BaCl₂ instead of XE991 in our internal solution since XE991 did not block $I_{\rm M}$ significantly (even at high concentrations of 40 $\mu_{\rm M}$; data not shown). Our non-inactivating $I_{\rm M}$ model followed a Hodgkin and Huxley notation, i.e., $I_{\rm M} = \bar{g}_{\rm M} u (V - V_{\rm K})$, where $\bar{g}_{\rm M}$ is the maximal $I_{\rm M}$ conductance, V is the membrane voltage, and $V_{\rm K}$ is the K⁺ reversal potential. The evolution variable u was numerically solved using the following equation:

$$\frac{du}{dt} = \frac{u_{\infty}(V) - u}{\tau_{u}(V)}.$$
(2)

The kinetic functions of our model $[\tau_u(V) \text{ and } u_\infty(V)]$ were constructed from the time constants-versus-voltage and current-versus-voltage relationships of XE991-sensitive current traces. Time constants for each voltage step were obtained by fitting the XE991-sensitive currents with single exponentials and time constants-versus-voltage relationships were fitted with a function of the following form:

$$\tau_{u}(V) = \frac{SF}{C_{a}e^{V_{a}^{(V+V_{a})}} + C_{b}e^{V_{b}^{(V+V_{a})}}} + M \tau_{u}(V)$$
$$= \frac{SF}{C_{a}e^{V_{a}^{(V+V_{a})}} + C_{b}e^{V_{b}^{(V+V_{a})}}} + M. \quad (3)$$

The normalized conductance g_M/\bar{g}_M versus voltage relationship, where \bar{g}_M is the maximum conductance and $g_M = I_M (V - V_K)$, the 20 μ M XE991-sensitive M-current, was fitted with the following function:

$$g/\bar{g} = (1 + e^{-(V+V_{1/2})/k})^{-1}.$$
 (4)

The results of the least-square fit of these functions are shown in Results.

Modeling. We also tested the importance of $I_{\rm M}$ in a pyramidal neuron model based on the modified Pinsky–Rinzel model (Pinsky and Rinzel, 1994) as implemented by Tiesinga et al. (2001). A pyramidal neuron consisted of a soma with voltage equal $V_{\rm S}$ and a single dendritic compartment with voltage equal to $V_{\rm D}$. All voltage-gated currents except $I_{\rm M}$ were implemented according to Tiesinga et al. (2001) as follows:

$$C_{\rm M} \frac{dV_{\rm S}}{dt} = -g_{\rm L}(V_{\rm S} - E_{\rm L}) - I_{\rm Na} - I_{\rm KDR} - I_{\rm M} - \frac{1}{2}g_{\rm c}(V_{\rm S} - V_{\rm D})$$
(5)

and

$$C_{\rm M} \frac{dV_{\rm D}}{dt} = -g_{\rm L}(V_{\rm S} - E_{\rm L}) - I_{\rm Ca} - I_{\rm KAHP} - I_{\rm KC} - I_{\rm NaO} - I_{\rm KO} - I_{\rm syn} - \frac{1}{2}g_{\rm c}(V_{\rm S} - V_{\rm D}), \quad (6)$$

where $g_{\rm L}$ is the leak conductance, $E_{\rm L}$ is the leak reversal potential, $I_{\rm KDR}$ is a delayed-rectifier potassium current, a fast sodium $I_{\rm Na}$, and the "coupling current" between the dendrite and soma compartments. The dendrite voltage was determined by calcium current $I_{\rm Ca}$, afterhyperpolarization (AHP) potassium current $I_{\rm KAHP}$, calcium-dependent potas-



Figure 1. XE991 decreases gamma oscillation amplitude. *A*, Representative examples of field recordings from hippocampal slices in the presence of 100 nm kainate. Left, Control recordings of kainate-induced gamma oscillations; middle, field activity of the same slices in the presence of 10 or 20 μ m XE991 or 20 μ m XE991 when the bath temperature was 35°C; right, respective power spectra (solid lines: control; dotted lines: XE991). *B*, Summary of integrated gamma oscillation power in control conditions and after the application of 10 or 20 μ m XE991 or 20 μ m XE991 in experiments performed at 35°C. *p < 0.03, **p = 0.04.

sium current $I_{\rm KC}$, persistent sodium $I_{\rm NaO}$ and potassium $I_{\rm KO}$ currents, and a leak current. For the implementation of the M-current we used our own experimental results (see above) and placed that current in the somatic compartment. Details for the model implementation and parameter values can be found in the study by Tiesinga et al. (2001). The dendritic compartment also received 10 excitatory inputs and the soma was innervated by 5 inhibitory inputs modeled by α functions ($x_{syn} =$ $t_{\text{syn}}e^{(1-t_{\text{syn}})/t_p} t_p^{-1} t_{\text{syn}} = t_e, t_i$) with a time constant t_p ($t_e = 0.15$ ms for excitatory and $t_i = 0.5$ ms for inhibitory synapses) and t_{syn} as the time of the last synaptic event (t_e and t_i , excitatory and inhibitory times, respectively). Synaptic currents I_{syn} (I_e , excitatory; I_i , inhibitory) were modeled by the following equation: $I_{syn} = \bar{g}_{syn} x_{syn} (V - V_{syn})$, $I_{syn} = I_e$, I_i , $V_{syn} = V_e$, V_i , $\bar{g}_{syn} = \bar{g}_e$, \bar{g}_i , where V_{syn} is the reversal potential for the excitatory ($V_e = 0 \text{ mV}$) and inhibitory ($V_i = -75 \text{ mV}$) synapses and \bar{g}_{syn} is the total synaptic conductance ($g_e = 0.05$ ns and $g_i = 0.075$ ns for excitatory and inhibitory synapses, respectively). Rhythmic inputs were modeled by a sinusoid function of 30 Hz frequency that determined the probability $[P(t) = 0.05 \sin(60\pi t)]$ of release of all presynaptic terminals j (j = 1, 2, ..., 15), at the most negative value of the sinusoid (valley), p =

0 and p = 0.0025 at the sinusoid peak. Each terminal had a absolute refractory time of 1 ms and a relative refractory time $[P_i(t) =$ $R_i(t)P(t)$] that obeyed a exponential function of the following form: $R_i(t) = e^{-0.1(t-t_{syn})}$. Synaptic parameters were chosen arbitrarily to produce a firing rate similar to the rate observed in kainate-induced gamma oscillations. In the study by Tiesinga et al. (2001) as well as in this work, the pyramidal cell model did not include the axon initial segment, the zone where APs are likely to be initiated (Kole et al., 2008). Data on the currents present on the initial segment are too sparse for the construction of a realistic axon initial segment compartment, and the scope of this paper was to assess the role of perisomatic $I_{\rm M}$ on the pyramidal cell function during gamma oscillations.

Results

M-current block decreases the power of kainate-induced gamma oscillations

Gamma oscillations in 400 µm hippocampal slices were induced by bath application of 100 nM kainate in a submerged recording chamber at 28°C. Slices that produced peak-to-peak oscillations smaller than 200 μ V (in our experiments, ~30%) were discarded. It is possible that further optimized recording chambers would augment our success rate (Hájos et al., 2009) but oscillation amplitudes likely would not differ considerably from the amplitudes reported here (Hájos et al., 2009). Once induced and stabilized (~15 min after kainate perfusion), gamma oscillations lasted >1 h without any noticeable change in amplitude if perfusion rate and oxygenation were maintained (longer durations were not tested). To assess the effect of I_M block on gamma oscillations we bath applied 10 μ M XE991, which caused a reduction of 32% in oscillation power (from 95 \pm 32 to 56 \pm 15 μ V², n = 20, p = 0.03) without affecting the peak frequency (Fig. 1). In a second set of experiments, 20 µM bath-applied XE991 caused

a reduction of 42% in oscillation power (from 86.7 \pm 4.5 to $50.6 \pm 5.4 \ \mu V^2$, n = 10, p < 0.01), with no change in peak frequency or increase of power in low-frequency bands (Fig. 1). The higher concentration of 20 μ M XE991 also reduced the power of gamma oscillations when slices were recorded at 35°C (39%, from 171.9 \pm 41.95 to 104.4 \pm 28.84 μ V², n = 6, p = 0.04) (Fig. 1). In agreement with a previous study (Fisahn et al., 2002), the addition of 10 μ M linopirdine did not affect either the power or the frequency of gamma oscillations (106.2 \pm 13.2 to 108.6 \pm 11.3 μV^2 , n = 4) (supplemental Fig. 1, available at www. ineurosci.org as supplemental material). However, addition of either 20 or 40 µM linopirdine caused a significant reduction (29 or 30%, respectively) in gamma oscillation power (to 75.7 \pm 7 or 74.3 \pm 4.6 μ V², respectively, n = 4, p < 0.01) (supplemental Fig. 1, available at www.jneurosci.org as supplemental material). We also observed a 23% decrease in gamma oscillation power when

10 μ M linopirdine application was supplemented with 5 μ M XE991 (from 102.76 ± 16.41 to 77.38 ± 7.51 μ V², n = 4, p = 0.03, data not shown).

M-current block changes the firing pattern of CA3 pyramidal neurons during gamma oscillations

Next, we investigated the effect of XE991 on the firing pattern of CA3 pyramidal neurons during kainate-induced gamma oscillations (Fig. 2A). Once oscillations were established and showed constant amplitude for 5 min, we recorded 2 min of concomitant field and current-clamp data (control). After 15 min of application of 20 µM XE991 another 2 min of concomitant field and current-clamp data was recorded. Single gamma cycles were detected off-line using a sliding template function (Clements and Bekkers, 1997) and AP times were calculated in a time window equivalent to each gamma cycle (Fig. 2B). The mean membrane potentials (in the absence of APs, during interspike intervals) in the presence of 100 nm kainate were -56 ± 1.3 and -51.6 ± 0.6 mV when 20 μ M XE991 was also present (n =8, *p* < 0.01) (Fig. 2*C*,*D*). XE991 increased the firing rate of CA3 pyramidal neurons during gamma oscillations (Fig. 2B-D). Neurons fired APs in 6.2 \pm 0.9% of the detected cycles before the addition of 20 μ M XE991 and in 7.1 ± 1.2% after (n = 8, p < 0.05) (Fig. 2*C*,*D*). We have fitted normalized AP-firing histograms with a third-order polynomial of the form ax^3 + $bx^2 + cx$ (Eq. 1) to quantify the phase response of CA3 pyramidal neurons during a gamma cycle (Fig. 2B). If neurons fire more randomly, AP histograms would be expected to flatten and the value of a would move closer to zero. Average a value was 0.0086 ± 0.0011 before and 0.0036 ± 0.001 after the addition of 20 μ M XE991 (n = 8, p < 0.01) (Fig. 2B,D). Average gamma cycles were aligned with AP-firing histograms to calculate the "preferred" firing phase angle of CA3 pyramidal neurons (angles >90° indicate that the peak firing occurred after the positive peak of the field oscillation) (Fig. 2B-D). The mean preferred firing phase angle



Figure 2. XE991 desynchronizes firing times of CA3 pyramidal neurons during gamma oscillations. *A*, Example of concomitant field and current-clamp recordings from pyramidal neurons showing the decrease of oscillation power associated with the increase of CA3 pyramidal neuron firing frequency after the application of 20 μ m XE991. *B*, Average gamma cycle (top) and AP time histograms before (middle) and after (bottom) the addition of 20 μ m XE991. *B*, Average gamma cycle (top) and AP time histograms before (middle) and after (bottom) the addition of 20 μ m XE991. AP time histograms were fitted with a third-order polynomial (gray traces) to quantify the "randomness" of AP times in relation to the gamma cycle (flatter histograms generate fits with smaller third-order coefficients). A 90° angle (vertical dashed line) was assigned to the peak of the oscillation cycle and the angle of the valley was equal to -90° . The arcsin (added by 90 or -90° if the AP-firing time occurred after the peak or the valley, respectively) of the normalized field oscillation was used to calculate the preferred firing angle. *C*, Raw voltage traces of a pyramidal neuron in the presence of kainate showing differential AP alignment to the gamma cycle shown in *B* before and after the addition of XE991 (the layout of the average gamma cycle is also shown in gray between the voltage traces). *D*, Summary of changes in basal membrane potential, percentage of gamma cycles in which APs were detected, third-order polynomial coefficient *a* (see Results), and position of AP time histogram peak in relation to the gamma cycle angle. Contr., Control. **p* < 0.03.

was 117.8 \pm 6.8° before and 26.8 \pm 24.9° after 20 μ M XE991 application (n = 8, p < 0.01) (Fig. 2*B*,*D*). In addition, application of 20 μ M XE991 resulted in a greater variance of the peak firing angle compared with kainate-only control. The SD of the peak firing angle in control conditions was 19.1 but increased to 70.5 when XE991 was present (p < 0.01, F test) (Fig. 2*D*). We obtained similar results using field recordings concomitantly with cell-attached recordings. In cell-attached configuration, pyramidal cells fired APs in 6.7 \pm 1.1% of the cycles (supplemental Fig. 2, available at www.jneurosci.org as supplemental material)

in control and 12.13 \pm 2.9% of the gamma cycles in the presence of 20 μ M XE991 (n = 4, p < 0.03, paired t test) (supplemental Fig. 2, available at www.jneurosci.org as supplemental material). The preferred firing angle was equal to 133.4 \pm 11.9° in control and 2.48 \pm 45.14° in the presence of 20 μ M XE991 (n =4, p < 0.03, paired t test) (supplemental Fig. 2*B*,*D*, available at www.jneurosci.org as supplemental material) and *a* was equal to 0.18 \pm 0.03 and 0.06 \pm 0.01 for control and 20 μ M XE991 (n = 4, p < 0.05, paired t test) (supplemental Fig. 2*B*,*D*, available at www.jneurosci.org as supplemental material).



Figure 3. XE991 increases burst firing of CA3 pyramidal neurons. *A*, Example of gap-free recordings of CA3 pyramidal neurons in the absence of kainate and presence of different concentrations of XE991. Histograms on the left show the interspike interval in control conditions (white bars) and in the presence of 10 μ M XE991 (gray, top) and 20 μ M XE991 (gray, bottom). Insets show the percentage of spikes that occurred <20 ms apart from each other. *B*, Same as in *A* but during kainate-induced gamma oscillations. In these experiments only 20 μ M XE991 was used. *p < 0.02.

XE991 increases burst firing in pyramidal neurons

During gap-free current-clamp recordings in the absence of kainate, pyramidal neurons can fire AP bursts (quantified by the percentage of APs that occur <20 ms from each other). Application of 10 μ M XE991 increased the percentage of APs occurring <20 ms apart from 39 \pm 4% of the total AP count to 89 \pm 4% (n = 4, p < 0.01) (Fig. 3A). If 20 μ M XE991 was applied to the bath, the percentage of closely spaced APs increased from 39 \pm 3% to 97 \pm 3% (n = 4, p < 0.01) (Fig. 3A). To avoid rundown of $I_{\rm M}$ due to prolonged recordings we tested the effect of 10 or 20 μ M XE991 in different neurons instead of applying both concentrations in succession to the same neuron. In contrast to what is seen for theta oscillations (4-12 Hz) (Williams and Kauer, 1997), during gamma oscillations, pyramidal neurons do not fire in bursts (Fisahn et al., 1998). Our data confirm that after induction of gamma oscillations by kainate the percentage of APs fired in bursts (i.e., <20 ms apart) decreased from $39 \pm 4\%$ to $6 \pm 1\%$. Subsequent application of 20 μ M XE991 caused an increase in the percentage of APs fired ≤ 20 ms apart from 6 \pm 1% to 37 \pm 4% (n = 8, p < 0.01) (Fig. 3*B*).

XE991 and linopirdine differentially affect I_M

The effect of the $I_{\rm M}$ blocker XE991 and linopirdine on outward currents in voltage-clamp recordings of CA3 pyramidal neurons was further investigated. We first tested if we could record $I_{\rm M}$ in whole-cell configuration without rundown of the current by

measuring the ratio between the instantaneous and steady-state outward current in response to a 0 mV voltage step in 5 min intervals. Prolonged whole-cell recordings caused a significant rundown of $I_{\rm M}$. After 1 h of whole-cell recording, there was a decrease of 22.47% of the ratio between the steady-state and instantaneous outward current (see below) (0.27 \pm 0.06 to 0.21 ± 0.04 , p = 0.04, n = 4, paired t test) but no significant decrease was observed until 30 min from the establishment of whole-cell configuration (data not shown). We tried to minimize the effect of inward currents by perfusing the slices with 1 µM TTX, 50 µM CdCl₂, 20 µM ZD7288. Despite the usage of these blockers, we observed a substantial variability in the shape of outward currents recorded from CA3 pyramidal neurons (Fig. 4A). One-second-long voltage steps varying from -90 to 0 mV were applied after a -100 mV/1-s-long pre test step. XE991sensitive currents were obtained by digitally subtracting the outward current before and 20 min after the addition of XE991. We recorded current responses to voltages up to 0 mV as the preparation generally became unstable after the application of more depolarizing voltage steps in the presence of XE991. At 0 mV there was a decrease of the total steady-state current by $40.8 \pm 2.3\%$ (measured at 990 ms from the test step onset) after the application of 10 μ M XE991 (from 2.68 \pm 0.5 nA to 1.44 \pm 0.3 nA, p < 0.03, n = 6). A concentration of 20 µM XE991 caused a

reduction of the steady-state current by $46 \pm 5.3\%$ (from 2.31 ± 0.12 nA to 1.25 ± 0.16 nA, p < 0.02, n = 5) and 10μ M linopirdine reduced the current by $36.1 \pm 5.7\%$ (from 2.53 ± 0.48 nA to 1.84 ± 0.49 nA, p < 0.03, n = 5) (Fig. 4A,B). There were no significant differences between the percentages of block (at 0 mV) after the application of 10 or 20 μ M XE991 or 10 μ M linopirdine. However, at -40 mV, the current sensitive to 10μ M linopirdine (after digital subtraction of the traces before and after linopirdine addition) represented only $6.4 \pm 0.1\%$ of the total linopirdine-sensitive current recorded at 0 mV, whereas these percentages were equal to $14.6 \pm 3.3\%$ and $23.5 \pm 2.5\%$ of the 10 and 20 μ M XE991-sensitive currents, respectively (Fig. 4B).

Using our voltage-clamp recordings (20 μ M XE991-sensitive currents) (Fig. 4*C*), we constructed an M-current model. The average time constants-versus-voltage and conductance-versus-voltage relationships are shown in Figure 4*D*. Least-mean square fit (using Eqs. 2 and 3, respectively) of these two relationships yields the following values: SF = 4135.76, Ca = 164.64 ms⁻¹, Cb = 0.33 ms⁻¹, V_a = 0.118 mV, V_b = 0.097 mV, M = 35.66 ms, V_c = -0.0497 mV, $V_{1/2}$ = 39.1465 mV, and k = 10.7530 mV. Immunohistochemistry data suggest that the main source of I_M is located at the axon initial segment (Devaux et al., 2004). However, we were able to reasonably clamp the perisomatic I_M using somatic electrodes. Space-clamp errors in non-inactivating K⁺ current recordings appear as small inactivation of the current and a nonsaturating g/V relationship (Bar-Yehuda and Korngreen,

2008) but our recordings clearly demonstrate a non-inactivating XE991-sensitive current (Fig. 4*C*) and a strong trend to saturation in the g/V relationship for the XE991-sensitive current (Fig. 4*D*).

We also tested the direct effect of 100 nM kainate and 20 µM muscarine on the $I_{\rm M}$. Both these compounds can induce gamma-frequency field oscillations in hippocampal slices. The relative effect on $I_{\rm M}$ was estimated by the ratio between the time-dependent current (steady-state minus instantaneous current) and steadystate current (the minimum value of current after the onset of the current step) of current traces recorded in the presence of 1 µm TTX, 50 µm CdCl₂, 20 µm ZD7288 with or without 100 nM kainate or 20 µM muscarine. The time-dependent/steady-state current ratio was equal to 0.25 \pm 0.03 before and 0.24 \pm 0.05 after the application of 100 nM kainate (not significant, n = 4) (supplemental Fig. 3, available at www.jneurosci.org as supplemental material). In the presence of 20 μ M muscarine, we observed a significant decrease in the time-dependent/steady-state current ratio (from 0.26 \pm 0.06 to 0.15 \pm 0.04, p < 0.01, n = 4) (supplemental Fig. 3, available at www.jneurosci.org as supplemental material).

XE991 increases the excitability of pyramidal neurons

Current-clamp recordings in the absence of kainate showed that the addition of XE991 can cause a depolarization of membrane potential (V_{rest}), a decrease in the fast AHP and an increase in the num-

ber of APs fired in response to depolarizing current steps. In control conditions, the $V_{\rm rest}$ of CA3 pyramidal neurons was -62.4 ± 0.7 mV and the application of 10 μ M XE991 depolarized the $V_{\rm rest}$ to -57.0 ± 1.2 mV (n = 4, p < 0.01) (Fig. 5). In control, AHP amplitude (the difference to the potential at the onset of the AP) was 2.9 ± 0.5 mV and decreased to 2.1 ± 0.28 mV after the application of 10 μ M XE991 (n = 4, p = 0.02) (Fig. 5). Pyramidal neurons fired 1.5 ± 0.5 APs in control conditions in response to a 200-ms-long 100 pA current step and 1.75 ± 0.5 APs after the application of 20 μ M XE991 (n = 4, not significant) (Fig. 5). The application of 20 μ M XE991 caused a depolarization of $V_{\rm rest}$ from -62.5 ± 0.3 to -55 ± 0.7 mV, a decrease in AHP amplitude from 3.4 ± 0.3 to 1.8 ± 0.24 mV, and an increase of AP firing (in response to a 200-ms-long 100 pA current step) from 1.4 ± 0.22 to 3.1 ± 0.3 APs (n = 10, p < 0.02 for all cases) (Fig. 5).

$I_{\rm M}$ is important for the control of the transfer function of pyramidal neurons

Currents active around threshold are important determinants of the bandpass filtering and input/output gain control performed by pyramidal neurons (Cook et al., 2007) and, consequently, these currents can modulate subcellular oscillatory activity that, in turn, will modulate field oscillations. The role of the M-current as a modulator of resonant frequency (band-



Figure 4. Linopirdine (10 μ M) is insufficient to block M-currents at subthreshold potentials. *A*, Example of a current response of a CA3 pyramidal neuron to voltage steps ranging from -90 to 0 mV in control conditions (top; in the presence of 1 μ M TTX, 50 μ M CdCl₂, and 20 μ M ZD7288) and in the presence of 10 or 20 μ M XE991 or 10 μ M linopirdine (bottom). *B*, Current–voltage relationships of outward currents (measured 50 ms before the end of the voltage test step) in control conditions and after the addition of 10 or 20 μ M XE991 or 10 μ M linopirdine (bottom). *B*, Current–voltage relationships of outward current (measured 50 ms before the end of the voltage test step) in control conditions and after the addition of 10 or 20 μ M XE991 or 10 μ M linopirdine. *C*, Digitally subtracted current trace showing an example of the 20 μ M XE991-sensitive component of the outward current of a CA3 pyramidal neuron. *D*, Left, Summary of activation time constant-versus-voltage relationship (solid line shows the fit with Eq. 2). Note that the time constant value at 10 mV was added to improve the fit. Right, Normalized conductance of 20 μ M XE991-sensitive currents-versus-voltage relationship (solid line shows the fitting with a Boltzmann function: Equation 3).

pass filtering) in the theta range (4–12 Hz) in pyramidal neurons has been shown in other studies (Hu et al., 2002; Peters et al., 2005). However, pyramidal neurons receive excitatory and inhibitory currents at frequencies much higher than the theta range during gamma oscillations (Fisahn et al., 1998). Additionally, Piccinin and colleagues (2006) have shown that $I_{\rm M}$ is essential for nonsynaptic gamma oscillations in the hippocampus. Thus, the filtering properties of pyramidal neurons for inputs at frequencies around the gamma range may also be important in determining firing properties of pyramidal neurons. To analyze the filtering properties of the cell membrane and the role of $I_{\rm M}$ in determining these properties, we have applied a chirp signal (120 pA peak to peak) with a linear increase of frequency from 5 to 40 Hz over 10 s in current clamp (in the presence of 1 μ M TTX). The normalized transfer function (the relation between the input current and output voltage) in the frequency domain was 3.5 ± 0.3 (at 30 Hz) in control conditions and 3 \pm 0.5 after the addition of 10 $\mu{\rm M}$ XE991 (n = 4, not significant) (Fig. 6). However, if 20 μ M XE991 was applied to the bath, the normalized transfer function amplitude decreased from 3.3 \pm 0.2 to 2.1 \pm 0.3 (n = 4, p < 0.03) (Fig. 6). Hence, the addition of 20 μ M XE991 augments the filtering of subthreshold inputs at frequencies in the gamma range.



Figure 5. Firing properties of CA3 pyramidal neurons are altered by M-current block. *A*, Examples of voltage responses of CA3 pyramidal neurons to 100 pA current steps in control (solid traces) and in the presence of external 10 μ m XE991 (dashed trace, left) and 20 μ m XE991 (dashed trace, right). *B*, Summary of resting membrane potential (V_{rest}), AHP voltage, and number of APs in a 200-ms-long 100 pA current step in control and external 10 μ m XE991 (top) or 20 μ m XE991 (bottom). *p < 0.01.



Figure 6. Blocking $I_{\rm M}$ augments the low-pass filtering of injected currents in CA3 pyramidal neurons. Left, Examples of the voltage response of a pyramidal neuron to the injection of a 10-s-long chirp current waveform (frequency changing linearly from 5 to 40 Hz) in control conditions (black traces) and in the presence of 10 μ M XE991 (green trace, top) or 20 μ M XE991 (red trace, bottom). The membrane potentials of these two cells were equal to -61.6 mV (control, top black trace), -62.1 mV (control, bottom trace), and -58.4 mV (top green trace) after the addition of 10 μ M XE991 and -56.1 mV (bottom red trace) after the addition of 20 μ M XE991. Right, Examples of transfer function estimates (input/output relationship in the frequency domain) in control conditions (black traces) and in the presence of 10 μ M XE991 (green trace, top) or 20 μ M XE991 (red trace, bottom). Both current and voltage were normalized before estimation of the transfer function.

Addition of a simulated M-current improves phasic firing of pyramidal neurons

To further elucidate the role of pyramidal neuron I_M in shaping neuronal firing during gamma oscillations we have used dynamic clamp to simulate $I_{\rm M}$ in neurons in which the native current was suppressed. Using an internal recording solution containing 0.5 mM BaCl₂ we blocked the native $I_{\rm M}$ of single pyramidal neurons (Lamas et al., 1997). At this concentration the barium salt did not significantly affect the hyperpolarization-activated current (data not shown). BaCl₂ can also block a variety of other K⁺ channels (e.g., inward rectifier, leak channels, etc.) (Coetzee et al., 1999). In the presence of kainate, pyramidal neurons are depolarized to potentials considerably above the range in which inward rectifier channels are active. Regarding the BaCl₂ effect on leak channels, we did not observe a significant difference between input resistance between recordings in the presence of 20 μ M XE991 and experiments performed with BaCl₂ in the internal recording solution (data not shown). Outward currents recorded in the presence of TTX and ZD7288 were similar to outward currents observed when 20 µM XE991 was present in the bath (Figs. 4A, 7A). To demonstrate this similarity, we measured the ratio between the time-dependent current and steady-state current of current traces recorded in the presence of 20 µM XE991 in the bath and compared it to the same ratio when the internal recording solution was supplemented with BaCl₂ instead of XE911 application. At 0 mV, these current ratios were 0.25 ± 0.05 (control), 0.05 ± 0.02 (XE991 bath application), and 0.07 ± 0.02 (intracellular BaCl₂ application; n = 8) (Fig. 7B). While there was a significant difference between the control group (neither XE991 nor BaCl₂) and both the external 20 µM XE991 and internal BaCl₂ groups (p < 0.01 and p < 0.05, respectively), no significant differences were observed in the current ratios between the recordings performed with 20 µM XE991 added to the bath and with an internal recording solution supplemented with $BaCl_2$ (Fig. 7B).

By blocking the native $I_{\rm M}$ of single neurons, we could test whether postsynaptic $I_{\rm M}$ contributes to the firing pattern of pyramidal neurons during gamma oscillations. By subsequently reintroducing $I_{\rm M}$ using the dynamic clamp technique, we could test whether this current could rescue the native firing phenotype of CA3 pyramidal neurons. The ability of

our simulated $I_{\rm M}$ to emulate the native $I_{\rm M}$ was tested by applying current steps (-300 to 300 pA in 50 pA increments, 200 ms long) in normal ACSF using 0.5 mM BaCl₂ in the internal solution (for reasons see Materials and Methods page 5) with or without the simulated $I_{\rm M}$ (dynamic clamp, $\bar{g}_{\rm M} = 13$ nS). When our internal solution contained 0.5 mM BaCl₂, V_{rest} was -55.2 ± 2 mV without and $-62.1 \pm$ 1 mV with the $I_{\rm M}$ dynamic clamp (n = 7, p < 0.01) (Fig. 7*C*,*D*). In the same conditions, AHP amplitude was $2.29 \pm 0.2 \text{ mV}$ without and 1.29 \pm 0.2 mV with the $I_{\rm M}$ dynamic clamp (n = 7, p < 0.01) (Fig. 7C,D). In addition, when the internal recording solution contained 0.5 mM BaCl₂, pyramidal neurons responded with 2.3 \pm 0.2 APs but the same neurons fired 1.3 \pm 0.2 APs if the simulated $I_{\rm M}$ was applied to the cell (n = 7, p < 0.01) (Fig. 7*C*,*D*).

During kainate-induced gamma oscillations CA3 pyramidal neurons that had BaCl₂ dialyzed into their cytoplasm showed an atypical firing pattern in relation to the field oscillation cycle. Without the $I_{\rm M}$ dynamic clamp, cells fired APs in 8.8 ± 0.9% of the cycles, whereas with the dynamic clamp ($g_{\rm M} = 13$ nS), APs were fired only in 4.8 ± 0.6% of the cycles (n = 6, p < 0.01) (Fig. 8*A*,*C*). The potential was -52.4 ± 1.15 mV in the absence and -56.4 ± 1.23 mV in the presence of $I_{\rm M}$ dynamic clamp (n = 6, p < 0.01) (Fig. 8*A*,*C*). The third-order coefficient *a* (cal-

culated by fitting the normalized AP time histogram with Eq. 1) was 0.01 \pm 0.003 before and 0.03 \pm 0.008 after the addition of the simulated $I_{\rm M}$ (n = 6, p < 0.02) (Fig. 8*B*, *C*). Mean preferred firing angle was 39.3 \pm 47° before and 98.2 \pm 2.8° after the $I_{\rm M}$ dynamic clamp (n = 6, p < 0.01) (Fig. 8*B*, *C*). SDs were also significantly larger for preferred firing phase angle before (282.2) compared with after (17) the addition of the $I_{\rm M}$ dynamic clamp (n = 6, p < 0.01) (Fig. 8*C*).

The role of $I_{\rm M}$ in determining pyramidal neuron firing in response to noisy (or sparsely synchronized) inputs was further evaluated using a simplified CA3 pyramidal neuron model. IPSP and EPSP times were determined by a 30 Hz sinusoid. Because our model used a low synaptic probability, the neuron fired APs with no evident rhythm (Fig. 9). When $\bar{g}_{\rm M}$ equaled 13 nS, the modeled neuron fired APs in 4% of the cycles, whereas when $\bar{g}_{\rm M}$ equaled 0 nS, APs were fired in 11% of the cycles (Fig. 9). The peaks of the AP time histogram coincided with an oscillation phase of -12 and 102° if $\bar{g}_{\rm M}$ equaled 13 nS and 36 and 102° when $\bar{g}_{\rm M}$ equaled 0 (Fig. 9*B*). We simulated the effect of a voltagedependent block of $I_{\rm M}$ by linopirdine by multiplying Equation 4 with a voltage-dependent scaling function $\gamma(V)$ as follows:

$$\gamma(V) = 1 - 0.76/(1 + e^{-(V+51.4061)0.2149}).$$
(7)

The function $\gamma(V)$ is based in the work of Romero et al. (2004). Multiplying Equation 4 with $\gamma(V)$ results in a reduction of only 10% of $I_{\rm M}$ at -60 mV and 76% at 0 mV. When $I_{\rm M}$ was affected by a voltage-dependent blocker, the simulated neuron



Figure 7. Altered firing properties due to M-current block in the absence of kainate can be rescued by an $I_{\rm M}$ dynamic clamp. **A**, Example of a voltage-clamp recording (same protocol as in Fig. 2) showing the effect of 0.5 mM BaCl₂ in the internal recording solution on the M-current. Arrows indicate the regions where the instantaneous (lnst) and steady-state (SS) currents were measured. **B**, Summary of (SS – Inst)/SS ratios of outward currents (in response to a 0 mV voltage step) in control conditions, after the addition of 20 μ M XE991, and when 0.5 mM BaCl₂ was present in the internal solution. **C**, Example of a current-clamp recording (200 ms, 100 pA current step) of a pyramidal neuron using an internal solution supplemented with 0.5 mM BaCl₂ with (solid trace) and without (dashed trace) the simulated $I_{\rm M}$ (dynamic clamp). **D**, Summary of $V_{\rm rest}$, AHP voltage, and numbers of APs in response to a 100 pA current step recorded with internal BaCl₂ before and after the DC. *p < 0.03.

fired in 5% of the cycles and the peak of AP time histogram coincided with an oscillation angle of 30° (Fig. 9B). We also tested the importance of the compartmentalization of $I_{\rm M}$ by moving $I_{\rm M}$ to the dendritic compartment. When I_M was present only in the dendrite the neuron fired APs in 9% of the oscillation cycles and the peak of the AP time histogram coincided with an oscillation angle of 54° (Fig. 9B). Third-order coefficients of the polynomial fit of AP time histogram were a = 0.18 for the model with a somatic $\bar{g}_{\rm M} = 13$ nS, a = 0.05 for the model with a somatic $\bar{g}_{\rm M} =$ 0 nS, and a = 0.14 for the model with a somatic $\bar{g}_{M} = 13$ nS but with the voltage-dependent block implemented. The third-order coefficient of the polynomial fit was a = 0.06 when $I_{\rm M}$ was placed in the dendrite instead of the soma. Another aspect explored in our simulations was the effect of eliminating $I_{\rm M}$ when the cell received tonic excitation and rhythmical inhibition (Traub et al., 2000; Hájos et al., 2004). Tonic excitation was provided by the injection of 100 pA current at the soma of the modeled cell, whereas inhibitory synaptic inputs were as described in Materials and Methods. When $I_{\rm M}$ was intact ($\bar{g}_{\rm M} = 13$ nS), the modeled neuron fired APs in 11% of the cycles, the peak of the AP time histogram coincided with an oscillation angle of -114° , and the third-order coefficient of the polynomial fit was a = 0.11. For $\bar{g}_{M} = 0$ nS, those values were equal to 28%, -90 and 138° (two equal AP time histogram peaks), and a = 0.04, respectively (Fig. 9C). The model data suggest that a perisomatic I_{M} is more efficient in controlling the firing phase of the neuron in response to noisy but rhythmical inputs than a dendritic $I_{\rm M}$ and that, to unmask the effect of $I_{\rm M}$, it is necessary to fully block KCNQ channels at voltages around the resting potential.



Figure 8. Addition of a modeled M-current to CA3 pyramidal neurons improves phase preference firing during gamma oscillations. *A*, Example of a field recording and a simultaneous CA3 whole-cell current-clamp recording of a pyramidal neuron (0.5 mm BaCl₂ was added to the internal solution to block the native M-current). A simulated M-current was added to the neuron using the average M-current conductance g_M (DC) and $2 \times g_M$ ($2 \times DC$). *B*, Average gamma cycle (top) and AP time histograms with their respective third-order polynomial fits (gray traces) before (middle) and after (bottom) the addition of the modeled M-current. The vertical dashed line shows the position of the positive oscillation peak (90°). *C*, Summary of changes in basal membrane potential, percentage of gamma cycles with APs, third-order polynomial coefficient, and position of AP time histogram peak in relation to the gamma cycle angle before and after the addition of the artificial M-current. *p < 0.02.

Discussion

In this study we have shown that $I_{\rm M}$ is an important modulator of the firing phase of CA3 pyramidal neurons during *in vitro* gamma oscillations in the hippocampus. We first demonstrated that application of XE991, an $I_{\rm M}$ blocker, decreases field oscillation amplitude. We next described a significant loss of synchrony between field activity and pyramidal cell firing in the presence of XE991. In addition, the decrease in synchrony observed when $I_{\rm M}$ was decremented in a single cell was reversed if a simulated $I_{\rm M}$ was applied to the cell via dynamic clamp. Whereas the results of a previous study suggested that $I_{\rm M}$ was not important for the main-

tenance of in vitro gamma oscillations (but rather I_h and I_{cat}) (Fisahn et al., 2002), this discrepancy can be explained by different induction protocols used (carbachol- vs kainate-induced gamma oscillations) and the higher potency of the M-current blocker used in this study. Carbachol/muscarine induction protocols involve the activation of metabotropic receptors that have well described effects on $I_{\rm M}$ (Brown and Adams, 1982), while the application of kainate activates mostly ionotropic kainate receptors, and we did not observe any direct effects on I_M. However, it is interesting to note how gamma oscillations induced by the activation of different receptor families appear to be similar macroscopically (Traub et al., 2002; Fisahn et al., 2004). Also, in the study by Fisahn et al. (2002), the authors applied 10 μ M linopirdine to block $I_{\rm M}$, and at this concentration linopirdine has a very small effect on M-currents at resting membrane potential as a later study by Romero et al. (2004) showed. The larger difference in the percentage of I_M block by these two compounds at -40 mV when compared with the percentage of block at 0 mV can be explained by the voltagedependent effect of both XE991 and linopirdine (Romero et al., 2004). Due to the higher potency of XE991 we could observe a significant decrease of $I_{\rm M}$ at voltages close to $V_{\rm rest}$. Hence, the 10 μ M linopirdine used in the study by Fisahn et al. (2002) to examine the effect of $I_{\rm M}$ block on gamma oscillations was insufficient to affect I_M significantly at V_{rest}. Our pyramidal neuron model also shows that a voltage-dependent block of $I_{\rm M}$ does not change the pyramidal neuron transfer function as dramatically as a total block of that current.

Here we have used XE991, a compound $10 \times$ more potent than linopirdine (Romero et al., 2004). XE991 affected V_{rest} and altered AHP size and bursting activity, whereas $10 \ \mu\text{M}$ linopirdine showed almost no effect on the resting potential of pyramidal neurons (data not shown). Ten micromoles XE991 caused a substantial reduction in gamma oscillations (by 32%)

and double that concentration decreased the oscillation power even further (by 42%). Ten micromoles XE991 is >10× the IC₅₀ for this compound when measured at a membrane potential of -30 mV (Romero et al., 2004) and one would expect this concentration to block the majority (if not all) KCNQ channels in hippocampal slices. In fact, both concentrations caused an increase in burst firing in the absence of current injection and a depolarization of $V_{\rm rest}$ (~5 and 7 mV, respectively). However, these effects were generally greater when 20 μ M XE991 was applied to the bath, indicating the presence of an unblocked fraction of $I_{\rm M}$ at resting potentials if 10 μ M XE991 was used. This fraction A Control

No IM

Control

40 No I_M

Count

AP

0**⊾**

В

might be responsible for the greater decrease in gamma oscillation power observed in the presence of 20 μ M XE991. Small variations in $V_{\rm rest}$ can alter substantially the firing properties of pyramidal neurons during gamma oscillations (Fisahn et al., 2004); hence, the further block of KCNQ channels near resting potential caused by 20 µM XE991 could explain the larger effect we observed on pyramidal neuron firing.

Our voltage-clamp data suggest that at 20 µM, XE991 did not affect currents other than $I_{\rm M}$. However, XE991 is not as specific as previously suggested. At greater concentrations (40 µM) XE991 causes a reduction in the fast inactivating K⁺ current (I_A) (R. N. Leão and A. Fisahn, unpublished data). It is also known that concentrations of $\sim 100 \ \mu\text{M}$ XE991 can block ERG channels (Elmedyb et al., 2007). Furthermore, it is important to note that there was a large variability in the kinetics of outward currents recorded from CA3 neurons (Fig. 4A). However, the kinetics of XE991-sensitive currents appear to be more homogeneous among different CA3 pyramidal neurons than other outward currents (Fig. 4C). To obtain a "clean" $I_{\rm M}$ from digital subtraction, we had to apply TTX and CdCl₂ to block fast inactivating Na+ and Ca2+ (and Ca2+-dependent) currents. Vanner and colleagues (1993) have shown that 200 μ M $CdCl_2$ can partially block (~60%) M-like currents. We used $4 \times \text{less CdCl}_2$ in this study and to our knowledge this concentration has little effect on $I_{\rm M}$. Also, the ki-



AP Count

0°

00

100

AP Count

35

VD Block

Figure 9. In a compartmental pyramidal neuron model /_M decreases firing rate and randomness of action potentials. *A*, Model voltage response to rhythmic synaptic activity (see Materials and Methods) in control conditions ($g_M = 13$ nS; top left), when no $I_{\rm M}$ was implemented in the neuron ($g_{\rm M} = 0$ nS; bottom left), when a voltage-dependent block was applied to M-channels (top right), and when M-channels were implemented in the dendrite rather than the soma (bottom right). B, AP time histograms in relation to a gamma oscillation cycle (dotted lines) for the conditions shown in A. C, AP time histogram of a model cell that received rhythmical inhibition and tonic excitation when $g_{\rm M}=$ 13 nS (top) and 0 nS (bottom).

Time (ms)

VD Block

Dendritic I_M

40

Count

AP

35

Time (ms)

0<mark></mark>⊾

netics of our XE991-sensitive currents are highly suggestive of a "pure" $I_{\rm M}$.

In hippocampal neurons, I_M is responsible for spike frequency adaptation (Storm, 1990), for the maintenance of the resting potential (Brown et al., 1990), for preventing burst firing (Vervaeke et al., 2006) and, together with $I_{\rm b}$, for the generation of subcellular theta oscillations (Gu et al., 2005). Therefore, $I_{\rm M}$ can alter the input/output relationship of pyramidal neurons contributing to the sparse single firing displayed by these cells during in vitro gamma oscillations (Fisahn et al., 2004). Our experiments have shown that without $I_{\rm M}$, low-pass filtering for gamma oscillationfrequencies is increased, indicating that $I_{\rm M}$ can modulate the way pyramidal neurons process subthreshold inputs. Also, there is an increase in burst firing after the addition of XE991 that could disrupt oscillations by altering the AP timing of pyramidal neurons.

Obviously, interneurons will be affected by XE991 also, and this effect conceivably could contribute to the decrease in gamma power we describe here (Fuchs et al., 2007). Lawrence and colleagues (2006) have shown that $I_{\rm M}$ is a major determinant of interspike interval of hippocampal interneurons and future experiments should address the effect of I_M block in different interneuron populations during gamma oscillations. The scope of this study was to demonstrate that I_M modulates pyramidal neuron firing and that this modulation can explain the decrease in

gamma oscillations caused by XE991. To separate the subcellular and the network effects of I_M , we created an I_M model based on our voltage-clamp results and applied the modeled current to pyramidal neurons that had their native I_M blocked by intracellularly applied BaCl₂. Current-clamp recordings with Ba²⁺ in the intracellular solution have shown that the preferred firing angle of pyramidal neurons during gamma oscillations is quite different from cells recorded with intact I_M despite intact ongoing oscillations and synaptic input. A firing behavior similar to pyramidal neurons in control conditions was observed when we added artificial $I_{\rm M}$ to cells using dynamic clamp. Blocking $I_{\rm M}$ in a single cell did not produce the dramatic change in the AP time histogram observed when XE991 was applied globally. However, in the latter case, there is a global decrease in the field oscillation power while in the former experiments, the synaptic drive is intact. Using a simplified pyramidal neuron model, we have shown that $I_{\rm M}$ can alter the preferred firing phase angle of a pyramidal neuron and that a model without $I_{\rm M}$ displays a flatter AP time histogram in response to rhythmical inputs.

It is important to note that the subcellular localization of $I_{\rm M}$ dramatically affects the function of this current (Hu et al., 2007). Immunohistochemistry experiments have shown that KCNQ3 channels are mostly localized in the axons of CA3 pyramidal neurons (Devaux et al., 2004). Our dynamic clamp experiments were performed in somatic recordings, simulating only periso-

Control

No I_{M}

Time (ms)

35

35

matic I_{M.} However BaCl₂ could have diffused to distal compartments, producing a distal block of I_M that could not be compensated for by the dynamic clamp technique. I_M channels responsible for controlling somatic excitability are located proximal to the soma region (Devaux et al., 2004; Hu et al., 2007). Also, by rescuing the normal response of CA3 neurons in field oscillations, our dynamic clamp experiments have demonstrated that $I_{\rm M}$ is essential for pyramidal neurons to maintain coherent firing during gamma oscillations. These experiments have also shown that I_M can modulate pyramidal neuron firing only by its voltage dependence. Additionally, using a two-compartment model we have demonstrated that perisomatic $I_{\rm M}$ is more efficient than dendritic $I_{\rm M}$ in decreasing AP frequency and phase locking the pyramidal neuron firing to synchronous inputs. It will be interesting in future experiments to reassess the cholinergic modulation of $I_{\rm M}$ during gamma oscillations.

In this work, we demonstrate that $I_{\rm M}$, a current active around threshold, is an important modulator of neuronal excitability during oscillatory activity. While this study is devoted to the role of I_M in the computation of dendrosomatic inputs, the axonal location of KCNQ suggests that this current might also modulate axonal excitability (Devaux et al., 2004; Pan et al., 2006). KCNQ channels have motifs similar to Na⁺ channel proteins that cause the retention of these channels in the axon (Pan et al., 2006) and the strategic placement of $I_{\rm M}$ in the initial segment could serve as a blocker of backpropagating action potentials and a decoupling between soma and axon. The axonal placement of $I_{\rm M}$ is particularly important for the analysis of in vitro oscillations as some studies have suggested that intrinsic axonal activity in pyramidal cells is essential to the generation of rhythmic activity (Traub et al., 2002, 2003). However, the role of axons and backpropagating action potentials in gamma oscillations still needs to be addressed experimentally before studying the importance of $I_{\rm M}$ in the axonal compartment. Nevertheless, we have shown that perisomatic $I_{\rm M}$ can modulate the transfer function of the somatodendritic compartment. The results shown here demonstrate that intrinsic neuronal properties are important determinants of network oscillations in vitro and that I_M has a significant impact on gamma oscillations, which, in turn, play a significant role in cognitive function.

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