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## Comments on superconductivity in $AE_2CuO_{4\pm\delta}$ (AE = Sr, Ba)

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Li et al. (1) report on superconductivity ( $T_c \sim 73$  K) that emerged in Ba<sub>2</sub>CuO<sub>4-y</sub> (Ba214) synthesized at extremely high pressures (~18 GPa). They claim that one of the unique features of this compound is an exceptionally compressed local octahedron. I read this article with great interest. However, I am afraid they stretch things a bit. While they claim the crystal structure of the compound is of the  $K_2NiF_4$  (coined as T) type, as is the case with (La, Ba)<sub>2</sub>CuO<sub>4</sub>, their X-ray powder diffraction data resemble more the Nd<sub>2</sub>CuO<sub>4</sub> (coined as T')-type structure (2): In the latter, no local octahedron exists. When lattice parameters are fixed to the reported values (a = 4.003 Å and c = 12.942 Å), a distinction between T- and T'-Ba214 may not be simple using X-ray diffraction. Nevertheless, the large a/c ratio (~0.309) suggests fluorite-structure Ba<sub>2</sub>O<sub>2</sub> layers, equivalently T'-Ba214: At the least, Rietveld refinements assuming the T' structure may be worth being considered.

Superconducting T-Sr<sub>2-x</sub>Ba<sub>x</sub>CuO<sub>3+δ</sub> (3) (or Sr<sub>2-x</sub>Ba<sub>x</sub>CuO<sub>4-x</sub>) (4), where a substantial number of O vacancies exist in the CuO<sub>2</sub> planes, is challenging the notion that the CuO<sub>2</sub> planes are the playground of high- $T_c$  superconductivity. Out-of-plane lattice constants of these materials are smaller than that for T-La<sub>2</sub>CuO<sub>4</sub> (5) despite the larger ionic radii of Sr<sup>2+</sup> and Ba<sup>2+</sup> compared to La<sup>3+</sup>. This reduced lattice length stems from the reduced Jahn–Teller distortions due to the O vacancies in the CuO<sub>2</sub> planes. The material (Ba214) reported by Li et al. (1) may be related to those largely O-deficit T-Sr<sub>2-x</sub>Ba<sub>x</sub>CuO<sub>4-v</sub> cuprates as their Rietveld refinement results (table S1 in supporting information of ref. 1) indicate (~40% of the oxygen vacancies in the CuO<sub>2</sub> planes). However, the *c* value (12.942 Å) of the Ba<sub>2</sub>CuO<sub>4-v</sub> presented in ref. 1 is far off the extrapolation to x = 2.0 in T-Sr<sub>2-x</sub>Ba<sub>x</sub>CuO<sub>4-v</sub> with various *x* (3): *c* would be ~13.4 Å. This may indicate that the Ba<sub>2</sub>O<sub>2</sub> layer does not crystallize in the rock-salt structure as should be in the T structure. Altogether, a formation of a local octahedron in their Ba<sub>2</sub>CuO<sub>4-v</sub> remains elusive and the concept of an exceptionally compressed local octahedron may lose its vindicability.

Finally it is worth noting that we have synthesized Sr<sub>2</sub>CuO<sub>4±δ</sub> (c = 13.55 Å) (6) and Ba<sub>2</sub>CuO<sub>4±δ</sub> (c = 14.6 Å) (7) whose c-axis lattice constants agree well with those expected from the variation of ionic radius [ $r_i$ (La<sup>3+</sup>) <  $r_i$ (Sr<sup>2+</sup>) <  $r_i$ (Ba<sup>2+</sup>)] and also from the trends derived from the  $n \ge 2$  members of 0<sup>(Sr)</sup>2(n-1)n and 0<sup>(Ba)</sup>2(n-1)n homologous series materials in the Sr-Ca-Cu-O (8) and Ba-Ca-Cu-O (9) systems. Also,  $T_c$  of ~73 K reported by Li et al. (1) roughly coincides with the onset  $T_c$  in the resistivity-temperature curves for our T-Ba<sub>2</sub>CuO<sub>4±δ</sub> specimens (7).

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