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A Simulation Study of the Influence of the Traveling Wave Patterns on Ion Mobility Separations in Structures for Lossless Ion Manipulations

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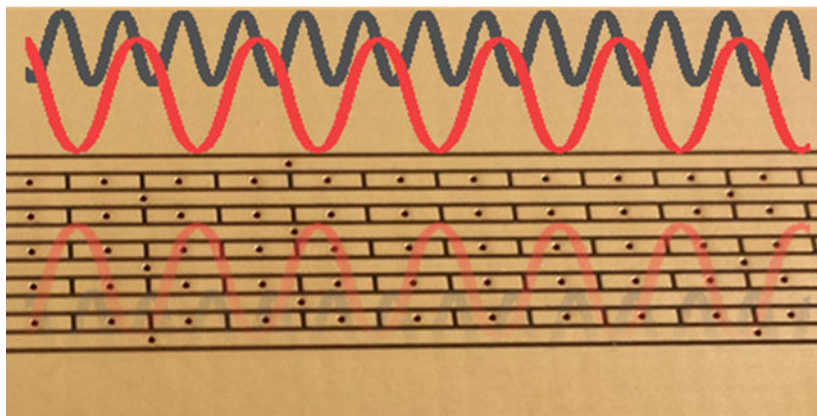
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Abstract

Probing molecular properties in the gas phase requires the integration of complementary ion manipulation approaches such as ion mobility spectrometry. Structures for lossless ion manipulations (SLIM) has recently been developed to perform ultra-high resolution ion mobility separations using traveling waves as well as provide other advanced capabilities. Despite its success, the design aspects of SLIM have not been fully explored and remained largely unchanged. Here, we report on a computational study using SIMION simulations of a number of traveling wave (TW) patterns that can be used in SLIM. The TW pattern used in the current SLIM device is a set of 8 electrodes where, at any time, 4 electrodes are held at high voltage (i.e. 1111) while the other 4 electrodes are held at low voltage (i.e. 0000), forming one micro-trapping region of 11110000 pattern. Ion trajectory simulations demonstrated the feasibility to simplify the 8-electrode set to a shorter pattern (e.g. 6-electrode or 4-electrode set) while maintaining or improving the performance. The RF and TW amplitude, guard voltage, and TW speed were optimized subsequently on the symmetric patterns of the 4-, 6-, and 8-electrode sets to further improve the performance. The resolution, peak broadening, peak capacity, and peak generation rate of each pattern were evaluated, showing that the 111000 pattern of the 6-electrode set has comparable performance as the current 11110000 pattern, and is always better than the 1100 pattern. This work provides insight and feasibility for simplification and modification of TW configuration in SLIM and other traveling wave devices.

Graphical Abstract

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Keywords

Traveling wave patterns; Ion trajectory simulations; ion mobility; resolution; peak broadening; peak capacity; peak generation rate

Introduction

Ion mobility (IM) coupled to mass spectrometry (MS) is a powerful technique for analytical separations, detection, and structural characterization by measuring collision cross sections (CCS) and mass-to-charge (m/z) ratio^{1–6}. IM-MS has been widely used for the analysis of a broad range of analytes including chemical warfare agents⁶, volatile organic compounds⁷, drugs⁸, and biomolecules⁹. Many types of IM have been developed in the past decades, including drift tube ion mobility spectrometry (DTIMS)¹⁰, high field asymmetric ion mobility spectrometry (FAIMS)¹¹, differential ion mobility spectrometry (DMS)¹², trapped ion mobility spectrometry (TIMS)¹³ and traveling-wave ion mobility spectrometry (TWIMS)¹⁴. A DTIMS is the traditional type of IM where ions are separated according to their collision cross sections with a buffer gas under the influence of a homogeneous electric field. The DTIMS separation power is directly correlated to the length of the drift tube and the electrical field applied. Owing to diffusional broadening, the ability to resolve two peaks using DTIMS can be obtained from the diffusion-limited resolving power equation

$$R = \sqrt{\frac{qEL}{16 k_B T \ln(2)}},$$

where q is the ion charge, E is the applied constant electric field, L is the length of the drift tube, k_B is the Boltzmann constant, and T is the temperature. The performance of DTIMS is limited by the maximum electric field that can be applied (typically few kV) and geometry constraints imposed by the linear drift tube length (more than 2 m length becomes cumbersome and impractical). This effectively limits the maximal achievable resolving power from DTIMS instruments. Unlike DTIMS, which requires a fixed high voltage, TWIMS uses a low amplitude traveling wave (TW) to create a dynamic field to drive and separate ions. Ions with higher mobility can keep up with TW than lower mobility ions that tend to fall behind by the passing wave. This form of ion separation enables the use of a fixed voltage without the debilitating limitation of high voltage drops associated with DTIMS. When the TW approach is coupled to a flexible ion path design (e.g. serpentine path), enabled by Structures for Lossless Ion Manipulations (SLIM), there is

potential for ultra-long path length separations. SLIM¹⁵ use two planar surfaces patterned with arrays of electrodes that are supplied with RF, DC, and TW potentials to confine and drive ions in gases at a pressure of few Torr. The electrode arrays can be patterned on printed circuit boards (PCBs) at low cost and are easily fabricated. The ability to pattern such electrodes on surfaces enables complex ion manipulations such as turning around a corner¹⁶, switching ions to orthogonal path¹⁷, and moving ions between levels using elevators and escalators¹⁸. These manipulations enable extremely long path length and ultra-high resolution ion mobility separations through multi-pass operations^{15, 19}. The serpentine path design and multi-pass operation have been demonstrated for ultra-high resolution with more than 500 m ion path length, indicating the SLIM ability to achieve high resolution IM separation¹⁵.

SLIM instruments built thus far have involved high resolution instrumentation targeted for biological analysis. The nominal path length was 13 m with the ability for multi-pass. The TW patterns applied used 8 electrodes with wave profile that were 45° phase-shifted to create a TW based IM separation. However, the performance of TW created with different number of electrode sets to create a TW was not investigated. One of the most effective ways to influence the performance of SLIM is to optimize the TW configuration. TW configuration is defined as the number of TW electrodes in a set while the TW pattern is the number of electrodes held at high and low voltage. The current SLIM design has 8 electrodes forming a single TW configuration which is repeated throughout the entire path length. In the 8-electrode set, 4 consecutive electrodes are held at high potential while the remaining 4 electrodes are held at low potential (i.e. a pattern of 11110000, where 1 represents an electrode with high potential and 0 an electrode with a low potential). The TW is created when the pattern is stepped through the 8 electrodes one electrode at a time. However, this configuration requires a power supply with 8 separate channels. If the TW configuration can be reduced to less than 8 electrodes, power supply requirements can be reduced, SLIM can be simplified, and tighter corners can be designed. These features enable the construction of a longer path length SLIM in a smaller footprint and provide the basis for a small and portable SLIM device. Previously published work from our group compared the mobility resolution for different patterns of the 8-electrode set and indicated that the current pattern of 11110000 is optimum²⁰. Other work has also demonstrated that a 4 repeat pattern of 8 TW configuration has 20% increased field strength in Synapt G2 IM-MS instrument compared to the 6 repeat pattern of 8 TW configuration in the Synapt G1 IM-MS^{3, 21}.

The investigation of different traveling wave configurations and patterns would enable designing simpler waveforms with similar or better peak characteristics. For instance, the nature of the serpentine path separation in SLIM requires ions to turn around corners. While straight motion of ions entails some diffusional broadening, turning around corners has to be accomplished with minimization of “race track” effects which could lead to a reduced separation performance. This is typically overcome in SLIM by making sharp 90° turns. This is accomplished by intersecting ions traversing a path with an orthogonally moving wave. Precise alignment of the TW electrodes is required to maintain wave continuity at the turn, which is simpler if there is flexibility in the TW configuration.

The efforts in SLIM development so-far have not examined a more diverse range of TW electrodes, simpler patterns, or smaller number of electrodes. Here, we perform an ion simulation study to explore the feasibility of simplifying the TW configuration and its effect on the peak properties. We demonstrate that other TW configurations can be used without loss of resolution via optimizing the operating parameters. The resolution trends on different TW patterns are calculated to investigate the impact of the TW patterns. Other operating parameters including RF and TW amplitude, guard voltage, and TW speed of 6- and 4-electrode sets were optimized. The resolution of different TW patterns has been demonstrated to be enhanced and comparable to the currently used 11110000 pattern. This work provides insight into the feasibility of simplifying and modifying the TW configuration in SLIM and other traveling wave devices while maintaining or improving the performance.

Methods

Ion motion and optimizations of the SLIM array design were simulated using SIMION 8.1 (Scientific Instrument Services Inc., Ringoes, NJ, USA) incorporated with statistical diffusion simulation (SDS)²² collision model for ions drifting in 4 Torr nitrogen gas. The electrode geometries were generated based on the current design of SLIM device. User programs were imported to operate the simulation via controlling the ion-neutral collision, voltage and frequency variation for ion confinement and separation in the created electric field. All ions were created at an initial condition of 0.1 eV kinetic energy with a Gaussian 3-dimensional distribution at the center of the TW microtrap. The SLIM design for the 8-electrode set is shown in Figure 1. In this design, an independent voltage is applied to each of the 8 electrodes. Two mirror-image planar surfaces are arranged in parallel at a spacing of 3 mm. The model consists of 6 RF strips interleaved with 5 arrays of TW electrodes with a 0.1 mm electrode gap between the electrodes. The length, width, and thickness of each TW electrode are 1 mm, 0.5 mm and 0.5 mm, respectively. The width and thickness of each RF electrode are 0.5 and 0.5 mm, respectively. There are two “guard” strip electrodes with a width of 1 mm adjacent to the two outer RF electrodes, supplied with a DC voltage for ion confinement in the lateral direction. Each RF electrode is supplied with the same amplitude but 180° out-of-phase sinusoidal waveform between each adjacent RF electrode to generate a pseudopotential for ion confinement between the two SLIM surfaces. A DC potential applied on each TW electrode varies from 0 to the assigned amplitude. The TW is stepped one electrode at a time with a constant speed.

Results and Discussion

1. TW patterns comparison

The effect of different TW patterns on the IM resolution of different TW sets of 4-, 6-, 8-, and 12-electrode sets were evaluated. An example of the simulated patterns of 4- and 6-electrode sets are shown in Figure 2. Singly charged phosphazine ions of m/z 622 and 922 from Agilent tuning mixture were simulated for a path length of 1320 mm (except when noted). Optimum guard electrodes voltage (15 V), RF (1 MHz, 350 V_{p-p}), and TW (92 m/s, 25 V_{0-p}) conditions were chosen for 100% ion transmission.

The mean arrival time and its standard deviation for the ions that transverse the full path was recorded. The resolution of m/z 622 and 922 ions was calculated using equation 1:

$$R_{622-922} = \frac{2(t_{922} - t_{622})}{\Delta t_{922} + \Delta t_{622}} \quad (1)$$

where t_{622} and t_{922} are the arrival times of m/z 622 and 922 ions measured at the apex of the peak and t_{622} and t_{922} are the full widths at half-maximum (FWHM) of m/z 622 and 922 peaks, respectively. As shown in Figure 3, the resolution of each TW set is the highest when TW patterns are symmetric. This is because asymmetric patterns generate a low potential wall of TW to move ions forward (i.e. ions are easily rolling over). Another consequence of asymmetric waves is a longer field-free region. For example, 10000000 pattern results in a large portion of the wave form at the same voltage leading to a long field-free region where ions execute diffusional broadening for much longer time. Thus the two extremes of asymmetric wave patterns (for example 10000000 and 11111110) will cause longer separation times and more ion diffusion with relatively low resolution and more peak broadening effect as compared to symmetric or near symmetric wave forms.

As shown in Figure 3, the ion mobility resolution of the symmetric of 6-, 8-, and 12-electrode sets are close, which is promising for using a lower number of TW electrodes. The resolution of 1100 pattern is ~20% lower than the other symmetric patterns which can be further optimized and improved. However, the 11100000 and 11111000 patterns show similar (within error bars) resolution compared to that of the 11110000 pattern. The error bars were calculated from the results of three replicates. The variations in the simulation arise from the computation of collisional statistics by the SDS model, wherein the variance in the arrival time distribution decreases with increased ion populations at the cost of increased computational time.

2. Ion mobility resolution optimization

Many parameters affect the IM performance in SLIM, especially the RF and TW amplitude which can impact the ion motion and TW separations. In this work, the amplitude of RF and TW, guard voltage, and TW speed were optimized for optimum resolution. Peak capacity and peak broadening effects were also compared before and after the optimization.

The RF potentials are utilized in SLIM to confine ions between the two surfaces. When the RF amplitude is not sufficiently high, ions are not effectively confined, especially ions with high m/z . TW is used to separate and move ions along the axial direction. When the TW amplitude is too high, microtraps are formed where ions are pushed forward by the TW voltage “wall”, causing ions to move at the same speed as the TW, termed “surfing”. However, if the TW amplitude is too low, the ion drift velocity is low and the moving wave frequently overtakes the ions, resulting in more rollover events. This causes the effective flight time of ions within the SLIM device to increase, leading to increased peak broadening.

The RF and TW amplitude (for TW patterns of 1100, 111000 and 1111000) were optimized by calculating the resolution for a range of RF and TW amplitudes. The RF amplitude was ramped from 250 V_{p-p} (in 50 V steps) to reach a practical voltage limit of 400 V_{p-p} (onset of

electrical breakdown in our current setup), while the TW amplitudes of 10, 20, 25, 30, 35, 40, and 50 V were used. A total of 100 ions of each of m/z 622 and 922 was utilized in a path length of 152 mm. A longer path length of 1320 mm (as used in other simulations here) and a larger number of ions was unrealistic in this case owing to extremely long computation time and the large number of simulations conducted. Therefore, while the trends of computed resolutions should be similar, the absolute values are lower compared to results in Figure 3. As shown in Figure 4a–c, the optimal RF amplitude is $\sim 400 V_{p-p}$ for all the TW patterns, and the optimal TW amplitude for 1100, 111000, and 11110000 are 40 V, 30 V, and 25 V, respectively. The resolution increased slightly for all the TW patterns with the increase in RF amplitude. A higher RF amplitude results in a better confinement and reduced electric field variation causing a slightly higher resolution. This was demonstrated by comparing the electric field of ion plume at two different RF amplitudes $250 V_{p-p}$ and $400 V_{p-p}$ with the same TW amplitude (30 V) using 111000 TW pattern. Figure 4d shows a cross section view of SLIM where a tighter ion plume is observed at higher RF amplitude. Figure 4e shows the TW electric field distribution an ion experiences moving in axial- x dimension (i.e. the direction of ion motion). The results indicate that at higher RF amplitudes, ions experience a tighter TW electric field distribution, which provides less deviation in ion speed and less peak broadening, allowing higher resolution to be achieved than with lower RF amplitude.

After the critical parameters have been optimized, guard voltage was optimized to investigate its impact on resolution. Guard voltage is a low DC voltage that is used to constrain the ions along the two sides of the ion path in the lateral direction (i.e. perpendicular to ion motion). If the guard voltage is too low, ions are not efficiently constrained, leading to ion leaking from the two sides. If the guard voltage is too high, there may be a field penetration from the guard electrode sides to the electric field at the ion trapping center, causing a field distortion at the trapping region.

An RF amplitude of $400 V_{p-p}$ was used for all of the TW patterns, and the optimized TW amplitude 40 V, 30 V, and 25 V were applied for 1100, 111000 and 11110000 TW patterns, respectively. Figure 5a shows the resolution trends of m/z 622 and 922 ions at the different guard voltages. The optimal guard voltage is 15 V, 17.5 V and 20 V for the 1100, 111000 and 11110000 TW patterns, respectively. The resolution was not calculated at lower voltages since the guard voltage was too low to prevent ion losses. The resolution trends indicate a slight impact of guard voltage on the resolution change, which is somewhat negligible in the resolution optimization.

To better understand the small effect of guard voltage, the ion plume distribution was plotted at 4 different guard voltages (12.5 V, 17.5 V, 22.5 V, and 27.5 V) using the same RF amplitude ($400 V_{p-p}$), TW amplitude (30 V), and 111000 pattern. Figure 5b shows the ion plume distribution in the yz plane, presenting a more laterally-constrained ion plume at higher guard voltages. Figure 5c shows similar TW electric field profiles in the axial- x dimension at different guard voltages, which confirms the negligible effect of guard voltage on the resolution. Similar results were obtained at an RF amplitude of $250 V_{p-p}$ (Supplementary information Figure 1) but with a wider ion plume distribution due to the weaker RF field. Overall, these results indicate that the guard voltage does not significantly

affect the RF or TW electric field but slightly affects the ion confinement in the lateral direction.

TW speed was optimized as another parameter to improve IM resolution for each TW pattern at their optimal RF and TW amplitudes. As shown in Figure 6, the observed trend of the TW speed on the resolution is similar to previous experimental findings and theoretical predictions^{15, 20, 23}. When the TW speed is too low (or ions' mobility is sufficiently high), ions will surf and no separation will occur. When the TW speed is increased, m/z 622 ions with higher mobility move at the same velocity as the traveling wave without rolling over, while m/z 922 ions with lower mobility roll over and begin to separate. The optimal resolution for a specific m/z is obtained at the TW speed where ion motion transitions from surfing to separation mode. When the TW speed is increased further, the resolution shows a decreasing trend to a plateau. At this high TW speed range, rolling over commonly occurs, which does not significantly change the resolution as shown in Figure 6. The optimal TW speed of 1100, 111000 and 11110000 TW patterns is 82 m/s, 88 m/s and 82 m/s, respectively.

3. Performance evaluation

The resolution between m/z 622 and 922 ions was compared before and after the optimization. In the unoptimized setting, the resolution of 111000 pattern is comparable to 11110000 pattern, and higher than 1100 pattern (Figure 7). Optimization of the RF, TW, and guard potentials resulted in resolution improvements of 11%, 35% and 33% for the 1100, 111000 and 11110000 patterns, respectively. Additionally, the resolution of 111000 shows similar improvement like that of the 11110000 pattern, and is always better than the 1100 pattern.

The improvement of resolution can be also inferred from understanding the spatial distribution of ions across the TW bins. Ions experiencing a large number of rolling over events tend to spread across a large number of TW bins and vice versa. To only understand the impact of rolling over events, we fixed the number of TW bins to 14 for 1100, 111000, and 11110000 patterns. 1000 ions were allowed to fly the entire 14 bins for each configuration. Shown in Figure 8 a–f is the spatial distribution of ions after traversing the 14 TW bins before and after the optimization. It is clear the spatial distribution of ions for all three patterns studied is much narrower after optimization.

Device performance is commonly evaluated by resolution, which gives information about how well two charged species can be separated according to their mobility. Other metrics such as peak capacity are also useful in the performance evaluation for TWIMS. Peak capacity is a dimensionless measurement representing the maximum number of peaks that can be resolved in a single separation. Peak capacity can be calculated using equation 2^{24, 25}:

$$peak\ capacity = \frac{t_{1822} - t_{622}}{(\Delta t)_{average}} \quad (2)$$

where t_{1822} and t_{622} represent the arrival times of m/z 1822 and 622 ions (range of ions studied here), and $(\Delta t)_{\text{average}}$ is the average full width at half maximum of m/z 622, 922, 1222, 1522 and 1822 phosphazine ions. Table 1 shows the peak capacity of 1100, 111000, and 11110000 patterns measured before and after the optimization. Although no improvement in peak capacity was observed for the 1100 pattern, the 111000 pattern dropped by ~6% while the 11110000 pattern improved by ~20%. Regardless, the 111000 pattern resulted in a similar performance to the 11110000 pattern. This result indicates that the 111000 pattern can generate comparable peak capacity to the 11110000 pattern.

Peak generation rate is another figure of merit calculated as the ratio of the peak capacity to the range of the separation time, shown by equation 3^{20, 26}:

$$\text{peak generation rate} = \frac{\text{peak capacity}}{t_{1822} - t_{622}} \quad (3)$$

We estimated the peak generation rate of each TW pattern before and after the optimization (Table 2). Before the optimization, the measured peak generation rates of 1100, 111000, and 11110000 TW pattern were 258 s^{-1} , 402 s^{-1} , and 527 s^{-1} , respectively. The 11110000 pattern shows a higher peak generation rate than the other patterns, while after optimizing the parameters, the peak generation rate of 1100 and 111000 has been significantly improved by 120–200 %, showing a much higher improvement than the 11% of the 11110000 pattern. The highest peak generation rate among the three TW patterns is 910 s^{-1} for the 111000 pattern, which is 55 % higher than that of the 11110000 pattern, and approximately 3 orders of magnitude greater than those in condensed phases (e.g., The average peak generation rate of liquid chromatography is $\sim 0.1\text{--}1.2 \text{ s}^{-1}$) due to the fast speeds of gas-phase ions^{25, 27}. However, ion mobility separation is post-ionization technique while liquid chromatography is pre-ionization technique and their integration is beneficial for maximum separation power. This result indicates a high peak generation rate can be achieved from the 111000 pattern than the 1100 and 11110000 patterns after the optimization. This increase in peak generation rate will maximize the separation power of SLIM when integrated with other separation approaches such liquid chromatography.

Conclusion

SLIM utilizes the traveling wave to perform ultra-high resolution ion mobility separations. RF and DC guard voltages are used in SLIM for efficient ion confinement with lossless ion transmission. In this work, the effect of TW patterns on the resolution was studied for different TW configurations. In addition to the currently utilized 11110000 pattern shorter patterns (1100 and 111000) were also investigated. It was observed that the symmetric patterns, in general, provided the highest resolution. Operational parameters such as RF and TW amplitude, guard voltage, and TW speed were then optimized subsequently for the 1100, 111000 and 11110000 TW patterns to further enhance the resolution. The resolution was compared before and after the optimization, showing an improvement of 11%, 35% and 33% for the 1100, 111000 and 11110000 TW patterns, respectively. The resolution of 111000 is comparable to 11110000 before and after the optimization, and higher than the

resolution of 1100 pattern. The spatial ion distribution for each TW pattern was simulated to understand the effect of ions rollover which can cause peak broadening. Upon optimizing the resolution, the spatial distribution of ions narrowed than before optimization which translates to a tighter ion packet. Peak capacities and peak generation rates were estimated before and after the optimization as well. The peak capacity of 111000 pattern is comparable to 11110000 pattern and higher than 1100 pattern before and after the optimization. The peak generation rate of 111000 is significantly increased after the optimization by 17% and 55% higher than 1100 and 11110000 pattern, respectively. The peak capacity and peak generation rate comparison indicates that the improvement trend is not following the same trend of resolution change, but 111000 always shows comparable or better results than 11110000. The performance evaluation of resolution, ion distribution, and peak capacities and peak generation rates demonstrated that the 111000 pattern has the ability to provide comparable performance as the current 11110000 pattern. Additionally, the 111000 pattern always shows higher performance than 1100 pattern, implying the 111000 pattern would be a better candidate for SLIM simplification than 1100 pattern. On the other hand, the 1100 pattern requires fewer channels (and smaller power supply to operate), has higher peak generation rate than the 11110000 pattern, and the resolution can be comparable to the other patterns by doubling the path length. This would be a promising step and insight for the simplification of SLIM and other traveling based devices.

Supplementary Material

Refer to Web version on PubMed Central for supplementary material.

Acknowledgment

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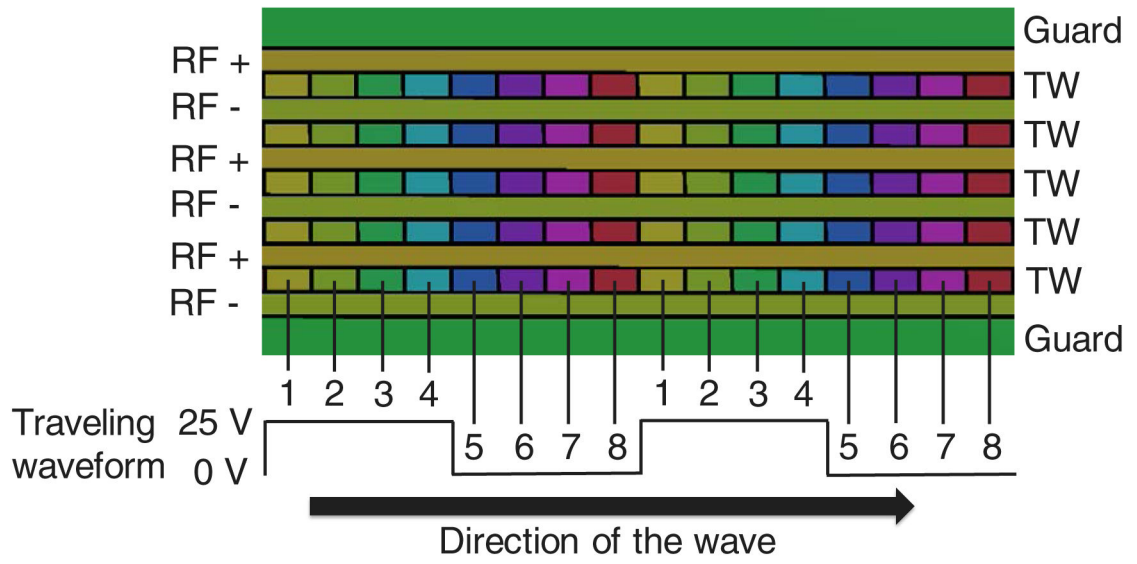


Figure 1.
The current SLIM design of an 8-electrode set using 11110000 pattern

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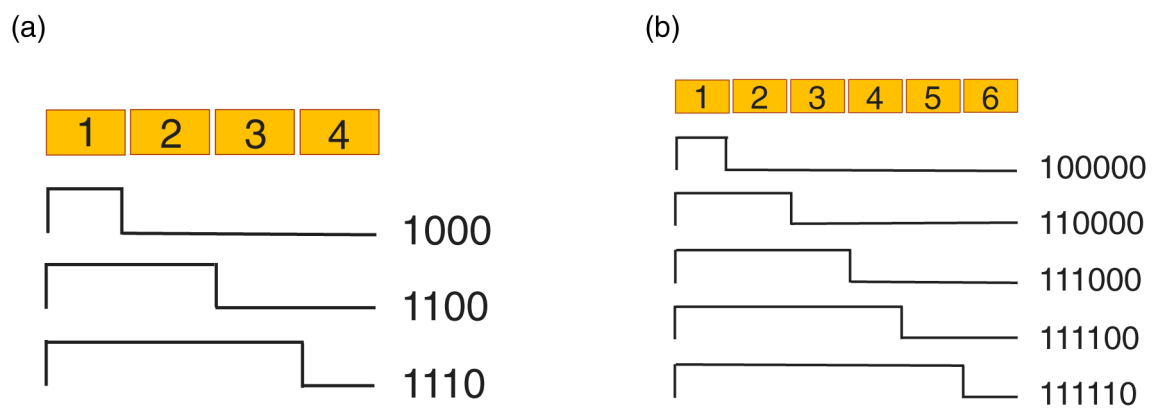


Figure 2.
Potential TW patterns for 4- and 6-electrode sets.

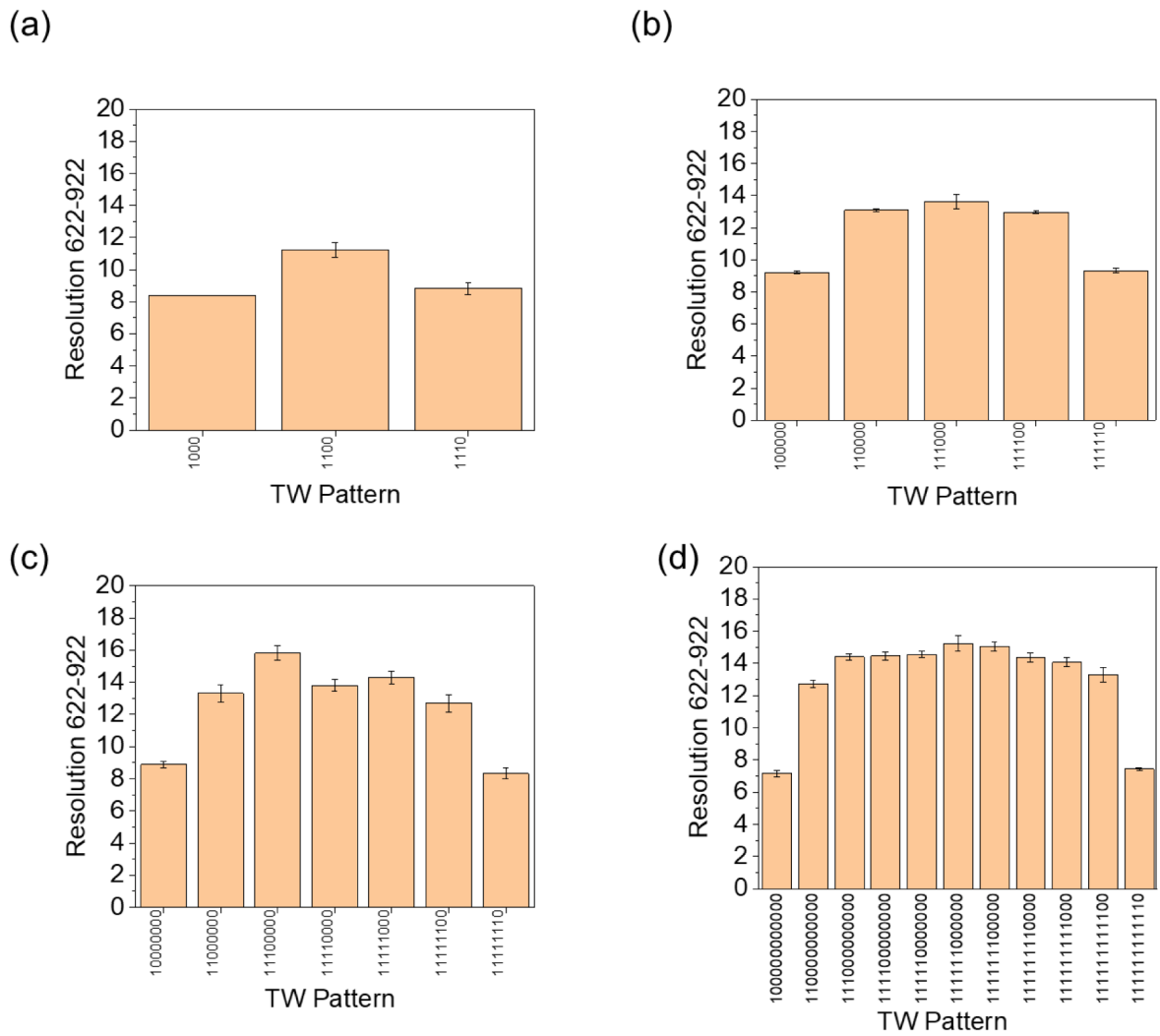


Figure 3. Resolution of each pattern of different TW electrode sets (3 replicate runs for each). (a) 4-electrode set; (b) 6-electrode set; (c) 8-electrode set; (d) 12-electrode set.

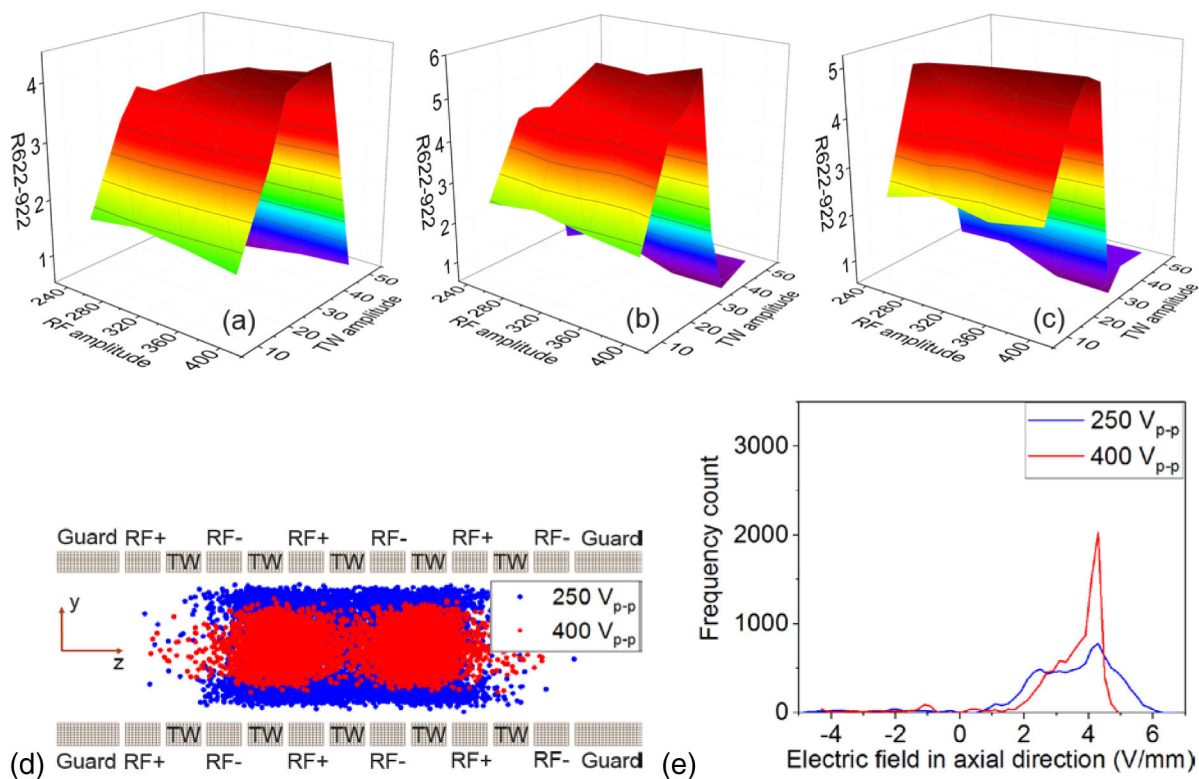


Figure 4. Optimization of RF amplitude and TW voltage. (a) 1100 TW pattern; (b) 111000 TW pattern; (c) 1111000 TW pattern. Electric field comparison using two different RF amplitude (250 V_{p-p} and 400 V_{p-p}) (d) cross section view of SLIM and ion plume confinement distribution of 10000 ions; (e) distribution of TW electric field in axial-x dimension as experienced by an ion.

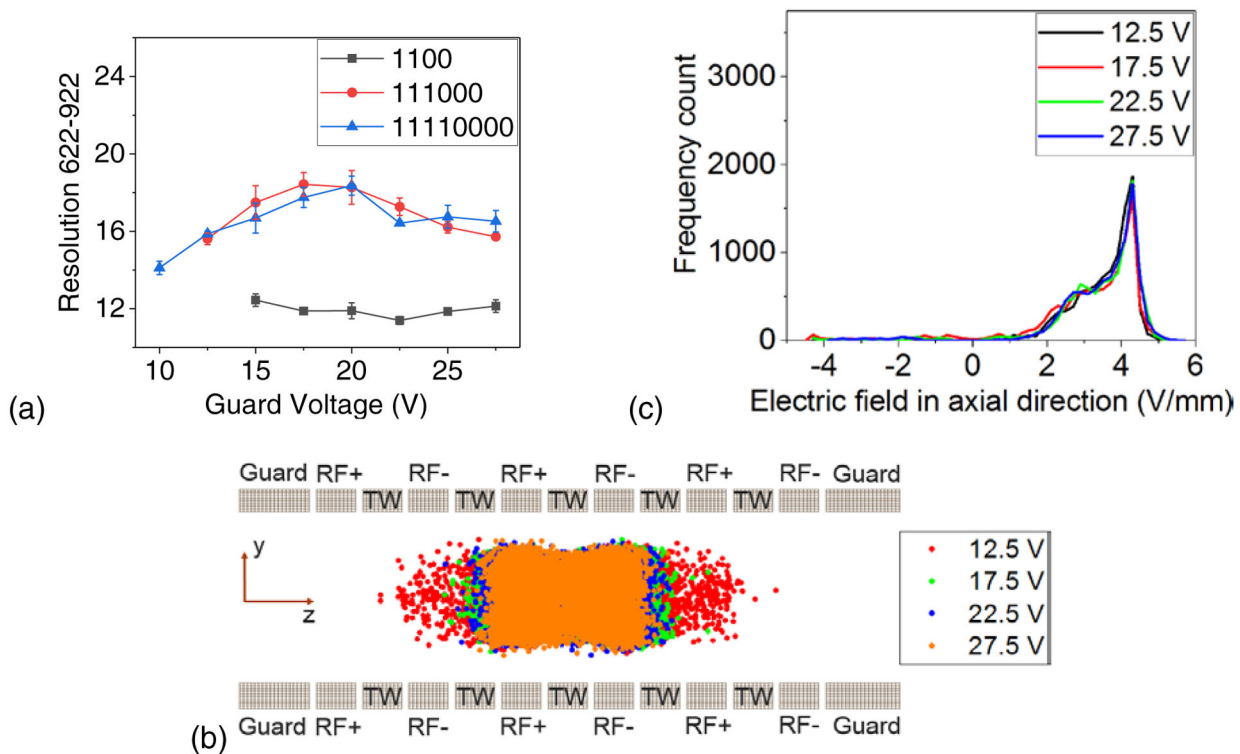


Figure 5. (a) Guard voltage optimization of 1100, 111000 and 11110000 TW patterns. (b) Plume confinement distribution of 10000 ions; (c) Variation of TW electric field in axial-x dimension for 111000 pattern.

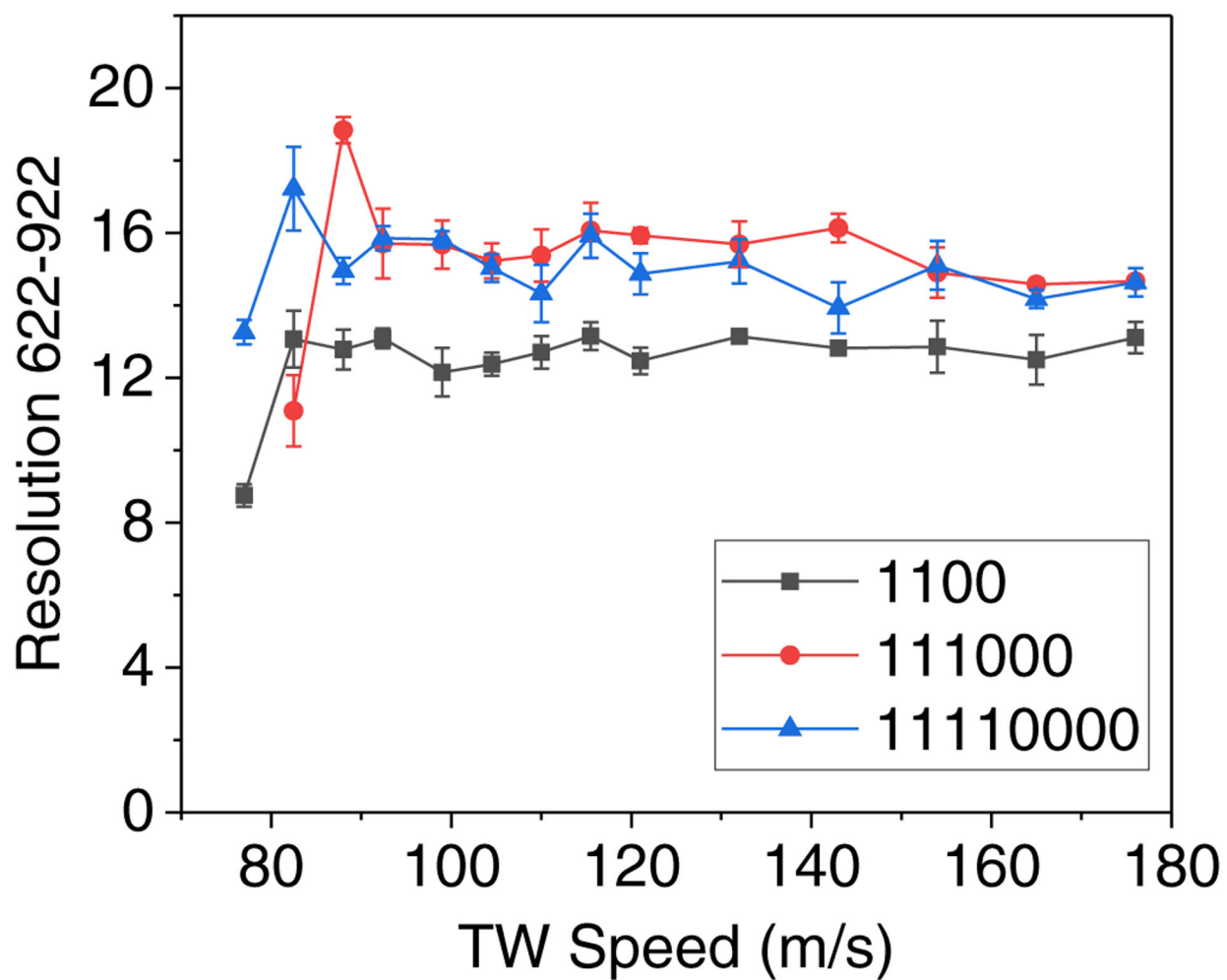


Figure 6.
TW speed optimization of 1100, 111000 and 11110000 TW patterns.

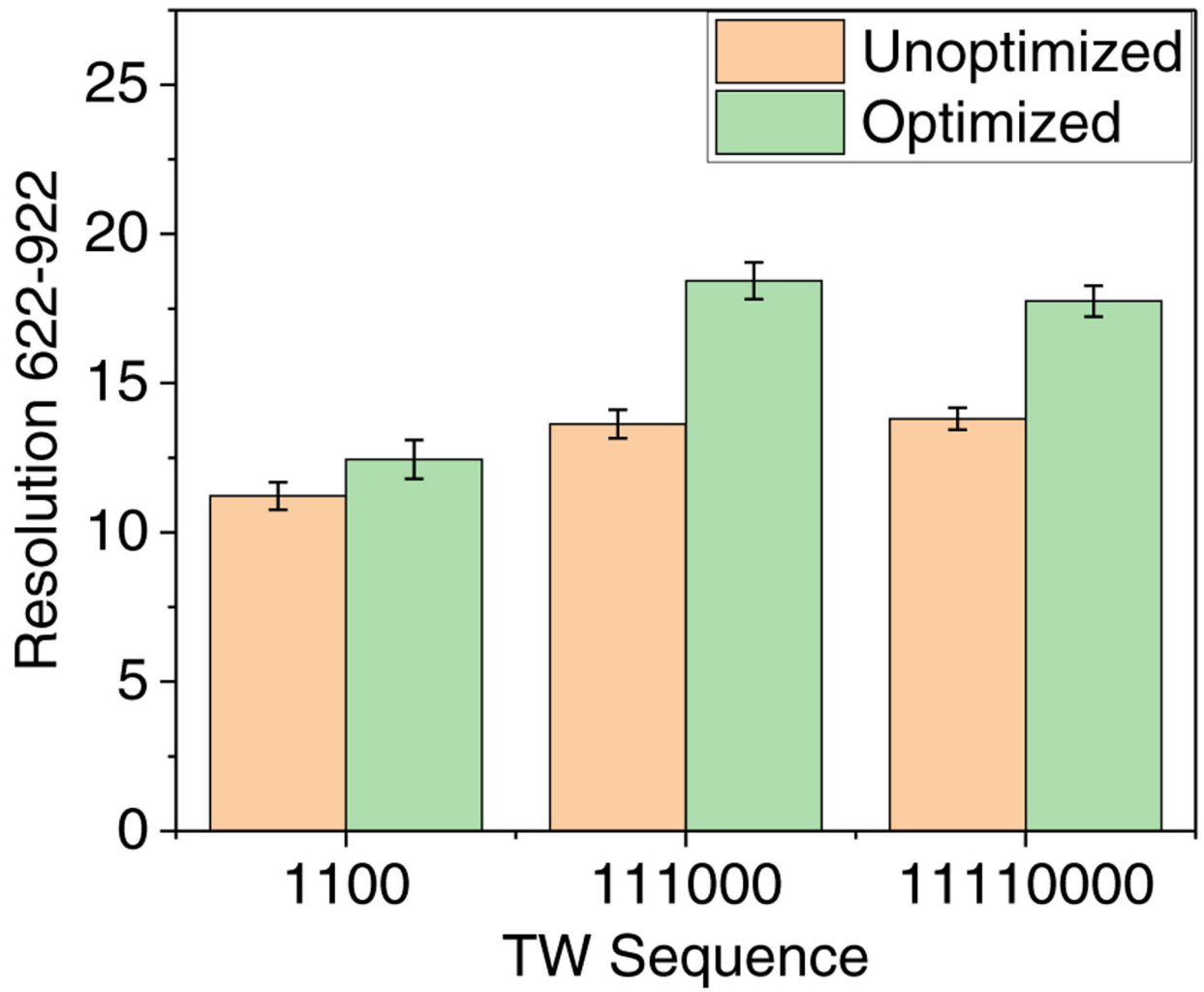


Figure 7. Resolution comparison of 1100, 111000 and 11110000 TW patterns before and after the optimization.

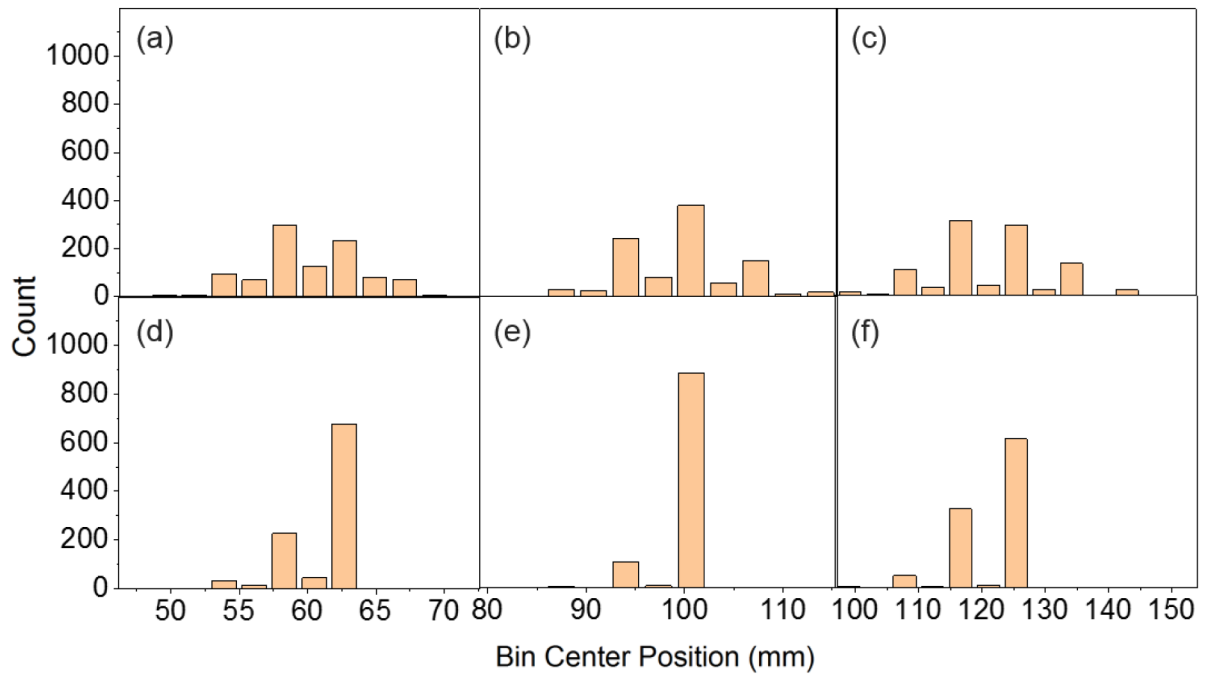


Figure 8. Ion distributions histograms recorded for unoptimized (top) and optimized (bottom) wave patterns of (a) 1100 (b) 111000 (c) 11110000.

Table 1.

The peak capacity of 1100, 111000, and 11110000 TW pattern measured before and after the optimization

TW pattern	Peak capacity (unoptimized)	Peak capacity (optimized)	Improvement %
1100	30	30	0
111000	46	43	-6
11110000	37	45	22%

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Table 2.

The peak generation rate of 1100, 111000, and 11110000 TW patterns measured before and after the optimization

TW pattern	Peak generation rate (before)	Peak generation rate (after)	Improvement %
1100	258 s ⁻¹	775 s ⁻¹	200 %
111000	402 s ⁻¹	910 s ⁻¹	126 %
11110000	527 s ⁻¹	587 s ⁻¹	11 %