



# NMR and computational methods for molecular resolution of allosteric pathways in enzyme complexes

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Received: 16 September 2019 / Accepted: 5 December 2019 / Published online: 14 December 2019

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## Abstract

Allostery is a ubiquitous biological mechanism in which a distant binding site is coupled to and drastically alters the function of a catalytic site in a protein. Allostery provides a high level of spatial and temporal control of the integrity and activity of biomolecular assemblies composed of proteins, nucleic acids, or small molecules. Understanding the physical forces that drive allosteric coupling is critical to harnessing this process for use in bioengineering, de novo protein design, and drug discovery. Current microscopic models of allostery highlight the importance of energetics, structural rearrangements, and conformational fluctuations, and in this review, we discuss the synergistic use of solution NMR spectroscopy and computational methods to probe these phenomena in allosteric systems, particularly protein-nucleic acid complexes. This combination of experimental and theoretical techniques facilitates an unparalleled detection of subtle changes to structural and dynamic equilibria in biomolecules with atomic resolution, and we provide a detailed discussion of specialized NMR experiments as well as the complementary methods that provide valuable insight into allosteric pathways *in silico*. Lastly, we highlight two case studies to demonstrate the adaptability of this approach to enzymes of varying size and mechanistic complexity.

**Keywords** Allostery · NMR · Molecular dynamics · Protein dynamics · Community network analysis

## Introduction

Allostery is a fundamental biomolecular regulatory mechanism characterized by communication between spatially distinct sites within a protein. The binding of an allosteric effector (i.e., peptide, small molecule) modulates substrate binding affinity ( $K_d$ , K-type) and/or enzymatic activity ( $V_{max}$ , V-type) by altering the structure and/or dynamics of the protein matrix (Fenton 2008). The idea that subtle conformational motions affect the energetic landscape of a protein to transmit chemical information has evolved with experimental technology, as

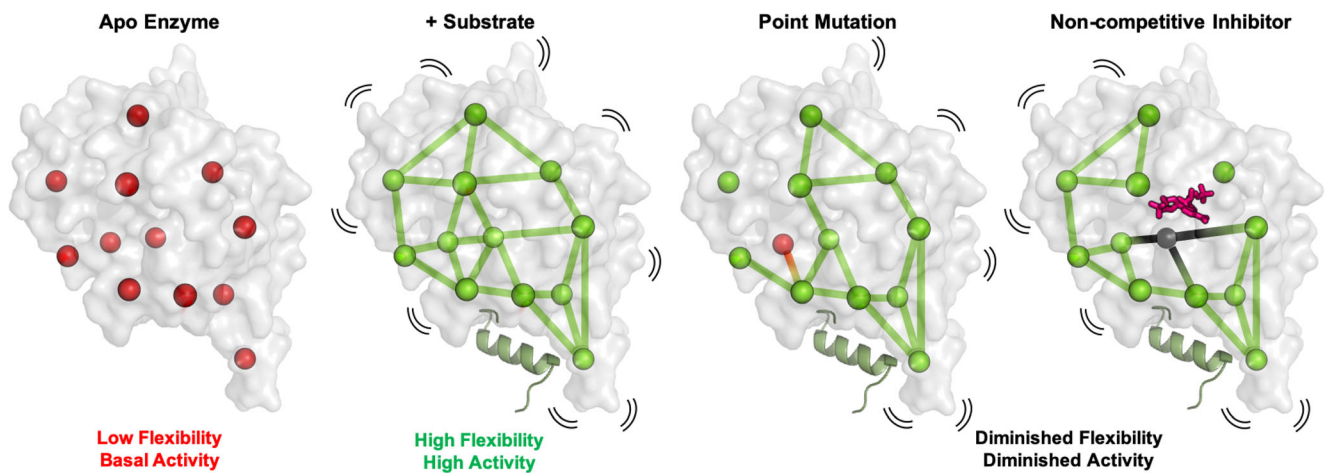
novel allosteric systems do not always conform to classical paradigms from phenomenological models (Koshland Jr. et al. 1966; Monod et al. 1965). Coupled to this observation is the notion that amino acid “networks” intrinsic to the protein are activated by endogenous or exogenous stimuli (Fig. 1). These *allosteric pathways* present an opportunity for fine-tuning or controlling biological responses; thus, ensemble models of allostery, where proteins sample microstates along a free energy continuum (Motlagh et al. 2014), have replaced a purely structural view of discrete conformational changes. However, a unifying model for all allosteric systems remains elusive. Ensemble models describe differing proteins with the same thermodynamic parameters, but such models generally exclude communicative pathways between active and regulatory sites, even though such a connection is necessary from an experimental point-of-view. Coupled communication organizes the active and allosteric sites of enzymes for substrate binding and mediates proper functionality. Despite advancements in biochemical and biophysical probes, the complexity of these mechanisms is such that allosteric pathways remain largely uncharacterized, especially in high molecular weight proteins.

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**Fig. 1** Allosteric pathways are composed of amino acid nodes that rely on the binding of a substrate or activator molecule to engage the network, often by stimulating local or global flexibility of the protein structure. Alteration of the allosteric pathway, shown here as a point mutation or

the introduction of a non-competitive inhibitor, can abolish connections made by critical nodes, resulting in attenuated structural flexibility and catalytic activity. Hijacking these routes of chemical information transfer for distal control of protein function is a promising therapeutic approach

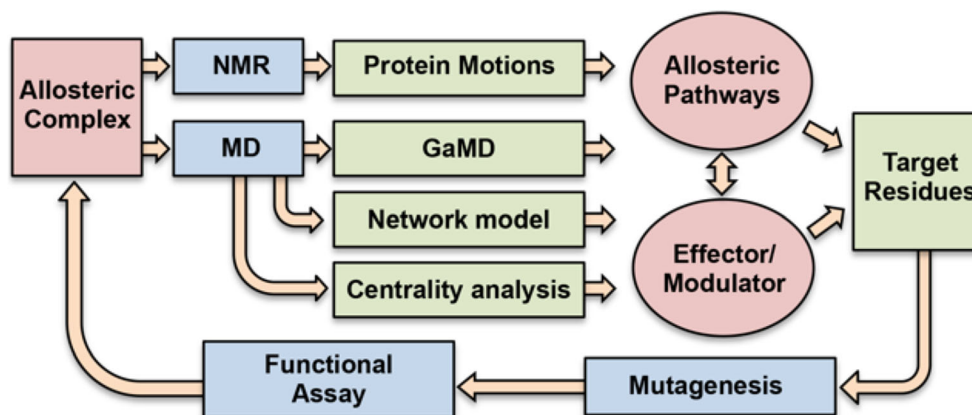
Identifying critical nodes along these pathways is desirable in drug discovery and tailored therapeutic design, and it is critical to engage a multitude of techniques, both complementary and orthogonal, to fully investigate allosteric mechanisms. Here, we highlight synergistic solution nuclear magnetic resonance (NMR) and computational studies used to elucidate structural and dynamic changes resulting from allosteric signaling. NMR is highly sensitive to subtle changes in protein structure and is extremely powerful for quantifying dynamic equilibria on a wide range of timescales (ps–sec). NMR is also the preferred method to validate computational predictions in ligand screening/docking and molecular dynamics (MD) simulations. Advanced computational techniques such as community network analysis and eigenvector centrality (EC) have become essential for the prediction and validation of allosteric pathways (Negre et al. 2018b; Rivalta et al. 2012), particularly since sophistication of modern computational tools expands the range of dynamic timescales that can be reliably probed, allowing access to slower dynamics utilized by large enzyme complexes for long-range communication.

Although other structural techniques such as free-electron laser crystallography can probe dynamic processes on timescales similar to those of NMR (Mizohata et al. 2018; Nango et al. 2016), its connection to MD simulations is not as well-established and crystallography still requires multiple static snapshots to infer solution-like behavior. Cryo-electron microscopy (EM), by contrast, is adept at probing dynamics in very large complexes (Kujirai et al. 2018), but lacks the atomistic resolution of NMR, the ability to quantitate motional timescales, and is not well-suited to studies of biomolecules < 40 kDa. NMR is able to accurately quantitate both the ensemble structure and dynamics across many timescales, and its coupling to MD simulations to drastically improve the

identification and characterization of allostery in protein complexes > 50 kDa is well-established. These studies, aided by modern experimental practices such as perdeuteration (Venters et al. 1996), transverse relaxation-optimized spectroscopy (TROSY) (Pervushin et al. 1997), sparse isotopic labeling (Tugarinov et al. 2006; Tugarinov and Kay 2003),  $^{15}\text{N}$ -detection (Takeuchi et al. 2016), and non-uniform sampling (NUS) (Barna et al. 1987; Delaglio et al. 2017), have facilitated NMR studies of much larger systems by preserving signal-to-noise and deconvoluting crowded spectra (Grishaev et al. 2005).

### Scope of the review

In this review, we discuss advances in solution NMR and computation that facilitate characterization of allosteric pathways and lead to functional insights. A large number of elegant reports utilizing these methods have proven the efficiency of a combined NMR/computational approach (Fizil et al. 2018; Jensen et al. 2014; Turupcu et al. 2019). These studies have paved the way for the investigation of systems of increasing size and complexity, such as protein-nucleic acid complexes where allosteric effects are at the basis of mechanistic function. Here, we focus on two important protein-nucleic acid complexes—DNA polymerase  $\beta$  and the Clustered Regularly Interspaced Palindromic Repeats (CRISPR)-Cas9—as case studies to detail how multi-domain structures transmit allosteric signals. In these systems, the complexity of the allosteric response is being tackled through a highly integrated approach, harnessing solution NMR and advanced MD methodologies, in combination with network models derived from graph theory (Fig. 2).



**Fig. 2** Integration of experimental and computational methods for studies of protein allostery. NMR spin relaxation experiments are conducted in parallel with MD simulations to identify clusters of amino acids exhibiting flexibility and inter-residue correlations (via network analysis and centrality). The data are pooled and compared to identify overlapping sites of flexibility that may be relevant to an allosteric pathway. Site-

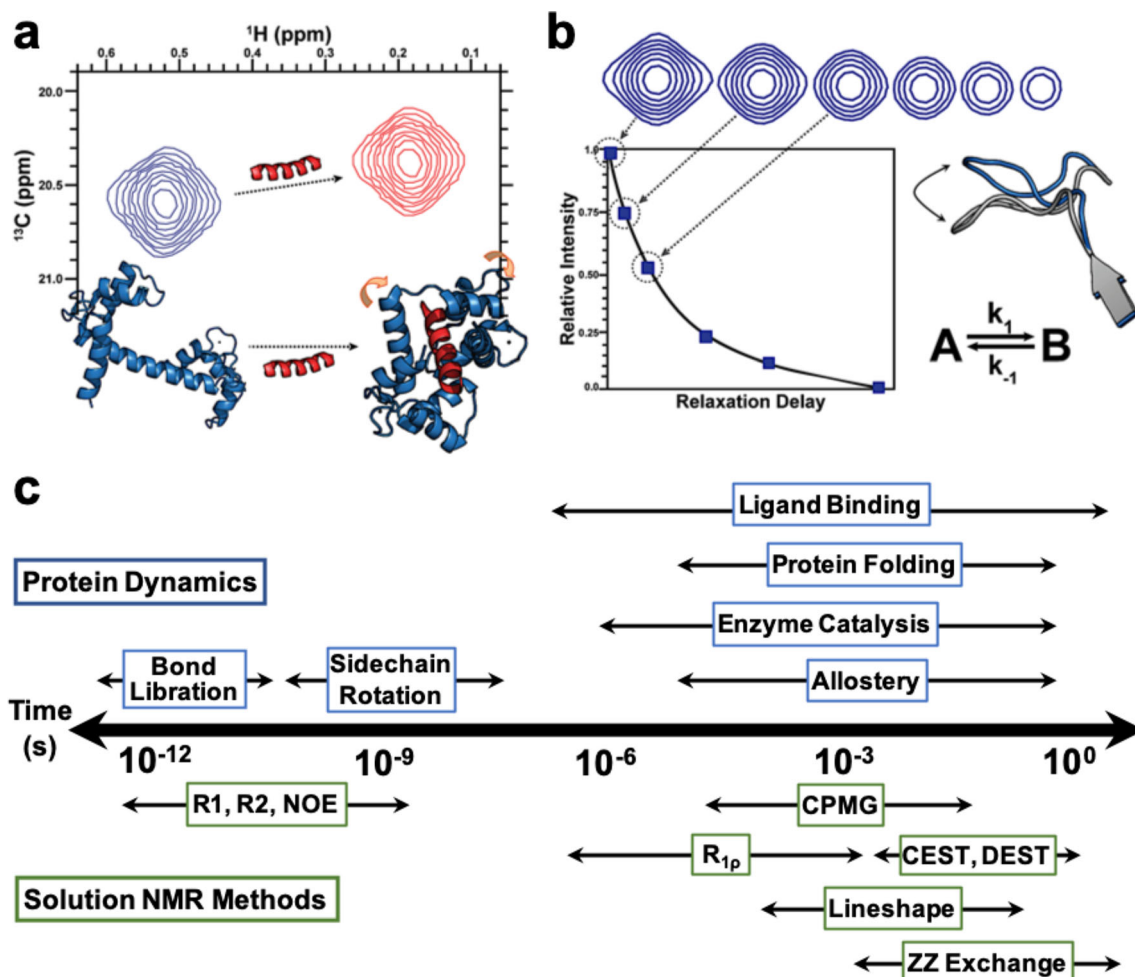
directed mutagenesis is conducted *in vitro* and *in silico*, and the effect on the structure, dynamic signal and correlations, and functional response are quantified. The process is iterated (i.e., carried out on mutants) to pinpoint changes in the pathway as well as critical allosteric nodes, after which allosteric modulators, such as drug-like small molecules, can be introduced into the system and tested biochemically

### Solution NMR studies of protein dynamics and allostery

Conformational rearrangements occur frequently during enzymatic mechanisms, often as a result of ligand binding where the associated conformational changes can be rate-limiting to catalysis (Benkovic and Hammes-Schiffer 2003; Lisi and Loria 2017; Watt et al. 2007; Whittier et al. 2013). Thus, characterization of relevant dynamics associated with these transitions is essential for understanding the biochemical mechanism. The equilibria between energetically similar apo (ligand-free) and liganded enzymatic states can be characterized by numerous biophysical techniques; however, in many cases, it is unclear whether the enzyme in question is intrinsically capable of sampling the relevant conformations required for allosteric activation. To address this question in molecular detail, experimental and computational approaches that can accurately quantitate multi-timescale protein motions are required (Yuan et al. 2015). It is now well-established that allosteric communication contains enthalpic and entropic thermodynamic components (Tsai et al. 2009), the latter of which is strongly influenced by conformational fluctuations (Henzler-Wildman and Kern 2007; Tsai et al. 2008; Volkman et al. 2001; Wand 2001). Well-vetted NMR methods for the study of protein dynamics have been reported (Dyson and Wright 2004; Loria et al. 2008; Palmer 3rd 2004; Palmer III 2015), and we will briefly describe how these methods can be used in concert with molecular simulations to probe biological processes that access weakly populated conformational states (Baldwin and Kay 2009; Loria et al. 2008; Mittermaier and Kay 2009; Palmer 3rd 2004; Palmer III 2015) and equilibrium dynamics that contribute to the configurational entropy of the system (Igumenova et al. 2006; Mittermaier and Kay 2006; Trbovic et al. 2009).

Solution NMR spectroscopy is performed under physiologically relevant conditions and is capable of determining atomic resolution protein structures, capturing site-specific conformational changes, and quantifying protein motions spanning the ps–sec timescales. The contributions of many laboratories to the application of novel NMR experiments and elucidation of mechanisms by which proteins propagate chemical signals have demonstrated the diversity of allosteric systems, though only a few examples can be discussed in detail. Novel examples of dynamically-driven allostery continue to surface where NMR experiments and/or molecular dynamics (MD) simulations have been used to establish the mechanism (Dhulesia et al. 2008; Hwang et al. 2004; Jacoby et al. 1996; Jarymowycz and Stone 2008; Kalodimos 2011; Liu et al. 2008; Wiesner et al. 2006; Zhang et al. 2012; Fuentes et al. 2004; Masterson et al. 2010; Petit et al. 2009; Coyne and Giedroc 2013; Lee 2015; Chakravorty et al. 2013; Shi and Kay 2014).

NMR probes local and global biomolecular structure and dynamics (Shi and Kay 2014) (Velyvis et al. 2007, 2009; Lipchock and Loria 2010) to illuminate pathways of allosteric cross-talk (Figs. 1 and 3). Dynamic processes specifically related to allostery occur primarily in the  $\mu\text{s}$ – $\text{ms}$  time regime, though contributions of faster (ps–ns) timescale motions have been noted, particularly as it relates to the optimization of the protein scaffold for effector binding and quantitation of configurational entropy in protein systems (Capdevila et al. 2017; Caro et al. 2017; Tzeng and Kalodimos 2012). The suite of NMR spin relaxation experiments designed to probe these time regimes have been used to interrogate internal protein motions (Akke and Palmer 1996; Loria et al. 1999a; Mulder et al. 2002) (Grey et al. 2003; Massi et al. 2005), protein folding (Hill et al. 2000; Korzhnev et al. 2004c; Tang et al.



**Fig. 3** Solution NMR is an ideal tool to probe biomolecular structure and dynamics. **a** Interaction of a protein with a binding partner (i.e., substrate or inhibitor) can cause a conformational transition that alters the local chemical environments of nearby amino acids or domains, perturbing NMR chemical shifts. **b** Conformational exchange motions can be readily measured with NMR by sampling peak intensities as a non-

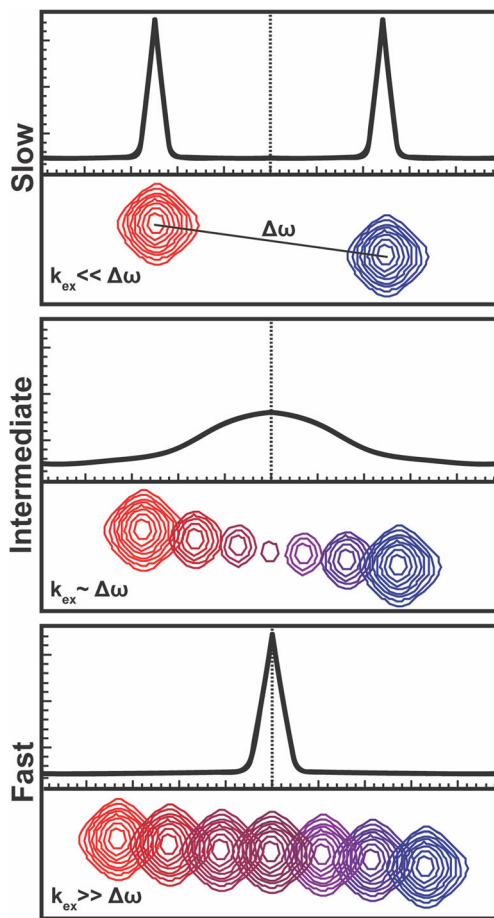
equilibrium state undergoes relaxation. Changes in peak intensities as a function of a variable relaxation delay provide signatures of flexible amino acids within a protein matrix. **c** Protein motions can be measured by NMR across the entire timescale of biology. Critically for this review, NMR experiments developed to probe slow timescale ( $\mu$ s–ms) dynamics are discussed in detail below

2006), ligand binding events (Mittag et al. 2003; Tolkathev et al. 2003), and enzymatic mechanisms (Cole and Loria 2002; Beach et al. 2005; Kovrigin and Loria 2006a, b; Kempf and Loria 2002; Berlow et al. 2007). The theoretical details of these solution NMR experiments are discussed in the next section, as are useful NMR observables that have powerful synergy with computation.

### Chemical shift perturbations

The resonant frequency ( $\omega_0$ ), or Larmor precession, of an NMR-active nucleus is dependent on the local magnetic field ( $B$ ), with contributions from the external field ( $B_0$ ) and local electronic environment ( $B_i$ ), where  $B = B_0 + B_i$ . In practice, resonant frequencies are referenced to a standard frequency in order to quantitate the chemical shift ( $\delta$ , in parts-per-million or Hertz) and minimize the contribution of  $B_0$ . Thus, the

chemical shift is extremely sensitive to changes in the local electronic environment and can be measured with high precision. Small changes in protein structure caused by ligand binding or a transition between inactive and active states ( $T \rightarrow R$  in allosteric nomenclature) are easily detectable with NMR chemical shifts. In a simple two-site exchange model describing the interaction of an enzyme with an allosteric effector, NMR-detectable nuclei sample two unique structural (and magnetic) environments as the allosteric effector binds/unbinds. The populations of each state,  $p_T$  (tense or inactive) and  $p_R$  (relaxed or active), have chemical shifts  $\delta_T$  and  $\delta_R$ , respectively. The chemical shift difference between these two states is given by  $\Delta\omega = |\delta_R - \delta_T|$  and the rate of chemical exchange between states  $T$  and  $R$  is given by  $k_{ex} = k_{T \rightarrow R} + k_{R \rightarrow T}$ . In the case of fast-to-intermediate exchange between conformers, where  $k_{ex} \geq \Delta\omega$ , the observed chemical shift ( $\delta_{obs} = p_R\delta_R + p_T\delta_T$ ) is a population weighted average of the



**Fig. 4** Exchange regimes dictate NMR chemical shift behavior based on the relationship between  $k_{ex}$  and  $\Delta\omega$ . For three distinguishable time regimes: slow (top), intermediate (middle), and fast (bottom), the 1D projection of a 1:1 mixture is shown for a 2D titration series. Analysis of the NMR lineshapes provide information about  $k_{ex}$  and  $\Delta\omega$

$T$  and  $R$  states, where  $p_{R/T}$  are the equilibrium populations of each conformational state and  $\delta_{R/T}$  are their respective NMR chemical shifts (Fig. 4). Under slow exchange conditions ( $k_{ex} \ll \Delta\omega$ ), two resonances are observed and the equilibrium populations are proportional to the peak volumes. Thus, changes in NMR chemical shifts for any allosteric system broadly represent a shift in the  $T$ -to- $R$  equilibrium.

Using this theoretical model of NMR chemical shifts, several methods have been developed to identify allosteric residues in proteins (Cui et al. 2017). One example, Chemical Shift Covariance Analysis (CHESCA) first reported by Melacini and coworkers (Selvaratnam et al. 2011), analyzes the chemical shifts of multiple allosteric states of a protein to identify pairwise correlations in covariant residues that compose an allosteric network. Alternatively, protocols to fit NMR lineshapes in a series of 2D spectra collected as allosteric effectors are titrated into protein samples have been described to extract both thermodynamic ( $K$  and  $\Delta G$ ) and kinetic information ( $k_{ex}$ ,  $k_{T \rightarrow R}$ , and  $k_{R \rightarrow T}$ ) (Kovrigin 2012; Shinya et al. 2017; Waudby et al. 2016). The sensitivity of these methods to communication

between active and allosteric sites is limited by the assumption that extrinsic factors such as buffer conditions, temperature, and pH have little effect on the allosteric ensemble, as well as by the fact that the chemical shift itself contains contributions from ligand binding, structural reshuffling at protein interfaces, and small  $pK_a$  shifts as hydrogen bond networks are altered. Coupling NMR chemical shift analysis to molecular simulations (vide infra) provides a solution to the inherent complexities of a purely chemical shift-based treatment of allosteric equilibria by utilizing MD trajectories to visualize structural states (i.e., conformational changes) along the dynamic ensemble.

### Fast timescale conformational motions (ps–ns)

Several thorough reviews of NMR spin relaxation have been given (Cavanagh et al. 2007; Lisi and Loria 2016a, b; Palmer 3rd 2004; Palmer 3rd and Massi 2006), so here we present a focused overview of useful experiments for biological investigations of allostery. Motions on the ps–ns timescale, which are faster than the rotational diffusion of a protein, are examined to quantitate equilibrium fluctuations in bond vectors of individual amino acids and changes in chemical shift anisotropy or dipolar interactions between nuclei (Clare et al. 1990; Peng and Wagner 1992; Ishima and Nagayama 1995; Mandel et al. 1995; Farrow et al. 1995; d’Auvergne and Gooley 2003; Cole and Loria 2003). Fast timescale dynamics associated with entropically driven allostery have been described for several systems (Lisi et al. 2016; Tzeng and Kalodimos 2012; Wand 2017) and are a critical component of the free energy landscape that regulates the population of allosterically activated (i.e.,  $R$ ) and inactive ( $T$ ) states. Pico-nanosecond motions are readily measured by NMR and classical MD, providing an avenue to synergistically probe the changing energy landscapes and three-dimensional structures associated with protein flexibility. For a heteronuclear spin-1/2 pair, i.e., an amide proton-nitrogen ( $^1H$ - $^{15}N$ ) of the protein backbone, longitudinal ( $S_z$ ,  $R_1$ ) and transverse ( $S_{xy}$ ,  $R_2$ ) relaxation of non-equilibrium magnetization to Boltzmann equilibrium are described by (Abragam 1961)

$$R_1 = d[3J(\omega_S) + J(\omega_I - \omega_S) + 6J(\omega_I + \omega_S)] + cJ(\omega_S) \quad (1)$$

$$R_2 = \frac{d}{2}[4J(0) + 3J(\omega_S) + J(\omega_I - \omega_S) + 6J(\omega_I) + 6J(\omega_I + \omega_S)] \quad (2)$$

$$+ \frac{C}{6}[4J(0) + 3J(\omega_S)] + R_{ex}$$

and the steady-state heteronuclear nuclear overhauser effect (NOE) is given by

$$^1H-[^{15}N] \text{ NOE} = \frac{\sigma_{NOE} \gamma_I}{R_2 \gamma_S} + 1 \quad (3)$$

where  $\omega_I$  and  $\omega_S$  are the Larmor frequencies of the  $I$  ( $^1H$ ) and  $S$  ( $^{15}N$ ) nuclei,  $c = (2/15)\Delta\sigma^2\omega_S^2$ ,  $\gamma_I$  and  $\gamma_S$  are the

gyromagnetic ratios of the  $I$  ( $^1\text{H}$ ) and  $S$  ( $^{15}\text{N}$ ) nuclei,  $\Delta\sigma$  is the chemical shift anisotropy of the  $S$  ( $^{15}\text{N}$ ) nucleus and  $\sigma_{\text{NOE}}$  is the cross-correlation relaxation term from  $I \rightarrow S$  ( $^1\text{H} \rightarrow ^{15}\text{N}$ ). The contribution of  $\mu\text{s}$ – $\text{ms}$  conformational motion to transverse relaxation is contained in  $R_{\text{ex}}$ , which is negligible in many cases. The dipolar coupling constant ( $d$ ) in the  $R_1$  and  $R_2$  cases is described by Eq. 4,

$$d = \frac{1}{10} \left( \frac{\mu_0}{4\pi} \right)^2 \hbar^2 \gamma_I^2 \gamma_S^2 \langle r_{IS}^{-6} \rangle \quad (4)$$

where  $\mu_0$  is the permeability of free space,  $\hbar$  is Planck's constant divided by  $2\pi$ ,  $\gamma_I$  and  $\gamma_S$  are again the gyromagnetic ratios of the  $I$  ( $^1\text{H}$ ) and  $S$  ( $^{15}\text{N}$ ) nuclei, and  $\langle r_{IS} \rangle$  is the average bond length between  $I$  and  $S$ . The spectral density function describing overall and internal bond vector fluctuations is written as

$$J(\omega) = \frac{2}{5} \left( \frac{S^2 \tau_c}{1 + (\tau_c \omega)^2} + \frac{(1-S^2)\tau}{1 + (\tau \omega)^2} \right) \quad (5)$$

where  $\tau_c$  is the rotational correlation time of the biomolecule,  $\tau = \tau_c^{-1} + \tau_e^{-1}$  and  $\tau_e$  is the effective correlation time for internal motions.  $S^2$  is a generalized order parameter ranging from zero-to-one that is commonly related to the configurational entropy of protein bond vectors, where lower values indicate heightened flexibility (Akke et al. 1993; Li et al. 1996; Yang and Kay 1996). Qualitative assessments of conformational entropy through the average order parameter,  $\langle S^2 \rangle$ , have been described by Bracken et al. (Kneller et al. 2002) using the product of experimentally measured  $R_1$  and  $R_2$  relaxation rates

$$\langle S^2 \rangle = \sqrt{\frac{\langle R_1 R_2 \rangle}{R_1 R_2^{\text{max}}}} \quad (6)$$

where  $\langle R_1 R_2 \rangle$  is the mean value and  $R_1 R_2^{\text{max}}$  is the calculated maximum value. Thus,  $R_1 R_2$  values below the mean of the data correspond to sites with lower order parameters (i.e., heightened flexibility).

Characterization of the amino acid side chain dynamics (i.e., in methyl groups of hydrophobic residues such as Ile- $\delta^{\text{CH}_3}$ , Leu- $\delta^{\text{CH}_3}$ , and Val- $\gamma^{\text{CH}_3}$ ) is an increasingly popular probe in large allosteric systems with dense backbone  $^1\text{H}$ - $^{15}\text{N}$  NMR spectra. Fast motions of methyl-containing side chains can be characterized by  $^1\text{H}$ - $^1\text{H}$  dipolar cross-correlated relaxation rate constant,  $\eta$ , expressed in Eq. 7 and described in studies of malate synthase G (Tugarinov et al. 2003), ribonuclease (Gill et al. 2019), and fatty acid binding protein (Liu et al. 2003).

$$\eta = \frac{R_{2,H}^F - R_{2,H}^S}{2} \approx 0.9 \left[ P_2 \left( \cos^2 \theta_{\text{axis,HH}}^2 - 1 \right) \right]^2 \frac{S_{\text{axis}}^2 \gamma_H^4 \hbar^2 \tau_c}{r_{\text{HH}}^6} \quad (7)$$

Here,  $r_{\text{HH}}$  is the distance between methyl protons,  $\gamma_H$  is the proton gyromagnetic ratio, and  $\tau_c$  is once again the rotational

correlation time.  $P_2(\cos \theta_{\text{axis,HH}}, \text{HH}^2)$  is the second Legendre polynomial in which  $\theta_{\text{axis,HH}}$  is the angle between the 3-fold symmetry axis of the methyl group and the vector between two methyl protons (i.e.,  $90^\circ$ ). The order parameter  $S^2$  describes the amplitude of the equilibrium fluctuations at the methyl symmetry axis.  $R_{2,H}^F$  and  $R_{2,H}^S$  are the relaxation rates for the single quantum  $^1\text{H}$  coherences for fast and slow relaxation, respectively.

### Slower timescale motions ( $\mu\text{s}$ – $\text{ms}$ )

Molecular motions occurring on the  $\mu\text{s}$ – $\text{ms}$  timescale are associated with conformational changes that participate in allosteric transitions and can be commensurate with rates of enzymatic catalysis (Whittier et al. 2013). Transitions between inactive and active enzymatic states (i.e.,  $T \rightleftharpoons R$  where  $k_{\text{ex}} = k_{T \rightarrow R} + k_{R \rightarrow T}$ ) cross significant energy barriers, modulate the isotropic chemical shift of the atoms involved, broaden the associated NMR resonances, and increase the transverse relaxation rate constant ( $R_2$ ) by an amount  $R_{\text{ex}}$  where  $R_2 = R_2^0 + R_{\text{ex}}$  and  $R_2^0$  is the “motion-free” transverse relaxation rate that describes motion at or near the Larmor frequency. In the fast exchange limit,  $R_{\text{ex}} = p_{\text{T}} p_{\text{R}} \Delta\omega^2 / k_{\text{ex}}$ , where  $\Delta\omega$  is the chemical shift difference between states  $T$  and  $R$ ,  $p_{\text{T/R}}$  are the equilibrium populations of the (assumed) two conformations, and  $k_{\text{ex}}$  is the exchange rate constant of the dynamic process. Relaxation experiments probing  $\mu\text{s}$ – $\text{ms}$  dynamics have been widely employed to investigate allosteric mechanisms in kinases (Masterson et al. 2010; Srivastava et al. 2014), molecular machines, and DNA binding proteins (Capdevila et al. 2017; Vise et al. 2005).

Micro-millisecond conformational exchange is characterized by relaxation-compensated Carr-Purcell-Meiboom-Gill (rcCPMG) (Loria et al. 1999a, b) and/or off-resonance rotating frame relaxation ( $R_{1\rho}$ ) experiments (Akke and Palmer 1996; Deverell et al. 1970), both of which quantify the rate of conformational exchange ( $k_{\text{ex}}$ ), the equilibrium populations of each conformer ( $p_a p_b$ ), and the chemical shift difference between conformers ( $\Delta\omega$ ) (Lisi and Loria 2016a). The measured value of  $R_{\text{ex}}$  in the rcCPMG experiment depends on  $\tau_{\text{cp}}$ , the repetition rate of a  $180^\circ$  pulse (i.e.,  $[\tau - 180 - \tau]_n$ ) (Carver and Richards 1972; Jen 1978; Davis et al. 1994) or the effective field strength  $\omega_e$  in the  $R_{1\rho}$  experiment (Deverell et al. 1970). The variation of the measured transverse relaxation rate ( $R_2(1/\tau_{\text{cp}})$ ) with  $180^\circ$  pulse spacing ( $\tau_{\text{cp}}$ ) in the single quantum rcCPMG experiment applied to backbone amide groups in biomolecules (Loria et al. 1999a) is expressed in Eq. 8 (Carver and Richards 1972; Jen 1978; Davis et al. 1994)

$$R_2(1/\tau_{\text{cp}}) = \frac{1}{2} \left( R_{2A}^0 + R_{2B}^0 + k_{\text{ex}} - \frac{1}{\tau_{\text{cp}}} \cosh^{-1} [D_+ \cosh(\eta_+) - D_- \cosh(\eta_-)] \right) \quad (8)$$

where  $R_{2A}^0$  and  $R_{2B}^0$  are the intrinsic (i.e., motion-free) transverse relaxation rates of the two sites and  $D_{\pm}$  and  $\eta_{\pm}$  are

$$D_{\pm} = \frac{1}{2} \left[ \pm 1 + \frac{\Psi + 2\Delta\omega^2}{(\Psi^2 + \zeta^2)^{1/2}} \right], \eta_{\pm} = \frac{\tau_{cp}}{\sqrt{2}} \left[ \pm \Psi + (\Psi^2 + \zeta^2)^{1/2} \right]^{1/2} \tag{9}$$

and  $\Psi = (R_{2A}^0 - R_{2B}^0 - p_A k_{ex} + p_B k_{ex})^2 - \Delta\omega^2 + 4p_A p_B k_{ex}^2$ ;  $\zeta = 2\Delta\omega (R_{2A}^0 - R_{2B}^0 - p_A k_{ex} + p_B k_{ex})$ .

Equation 8 is required when conformational motion is in the slow-to-intermediate exchange regime ( $k_{ex} \leq \Delta\omega$ ); however, under fast exchange conditions ( $k_{ex} > \Delta\omega$ ), the expression simplifies to (Luz and Meiboom 1963)

$$R_2(1/\tau_{cp}) = R_2^0 + \phi_{ex}/k_{ex} \left[ 1 - 2 \tanh\left(\frac{k_{ex}\tau_{cp}}{2}\right) / (k_{ex}\tau_{cp}) \right] \tag{10}$$

where  $\phi_{ex} = p_A p_B \Delta\omega^2$  (Ishima and Torchia 1999; Millet et al. 2000; Kovrigin et al. 2006; Lisi and Loria 2016a). Adaptation of the rcCPMG experiment for multiple quantum (MQ) relaxation dispersion studies of Ile- $\delta$ -CH<sub>3</sub>, Leu- $\delta$ -CH<sub>3</sub>, and Val- $\gamma$ -CH<sub>3</sub> methyl groups (Korzhnev et al. 2004a, b; Isaacson et al. 2007) has facilitated NMR investigations of large multimeric proteins (Fig. 5) and machines.

Molecular motions occurring faster than can be quantified by rcCPMG experiments ( $\sim 10^5$  s<sup>-1</sup>) are probed by off-resonance R<sub>1ρ</sub> (Akke and Palmer 1996; Trott and Palmer 3rd 2002)

$$R_{1\rho} = R_1 \cos^2\theta + R_2 \sin^2\theta + \frac{\sin^2\theta p_A p_B \Delta\omega^2 k_{ex}}{\frac{\omega_{Ae}^2 \omega_{Be}^2}{\omega_e^2} + k_{ex}^2} \tag{11}$$

where effective magnetic fields for conformers A and B are given by  $\omega_{Ae/Be}^2 = \omega_{A/B}^2 + \omega_1^2$  and  $\omega_e^2 = \omega_{iso}^2 + \omega_1^2$ , and the offsets

from the radio-frequency (RF) carrier for resonances A and B, and population-averaged resonances are represented as  $\omega_A$ ,  $\omega_B$ , and  $\omega_{iso}$ , respectively. The RF field strength  $\omega_1$  has a tilt angle  $\theta = \arctan(\omega_1/\omega_{iso})$ , and in the fast exchange limit, the R<sub>1ρ</sub> experiment is expressed in Eq. 12

$$R_{1\rho} = R_1 \cos^2\theta + R_2 \sin^2\theta + R_{ex} \tag{12}$$

where

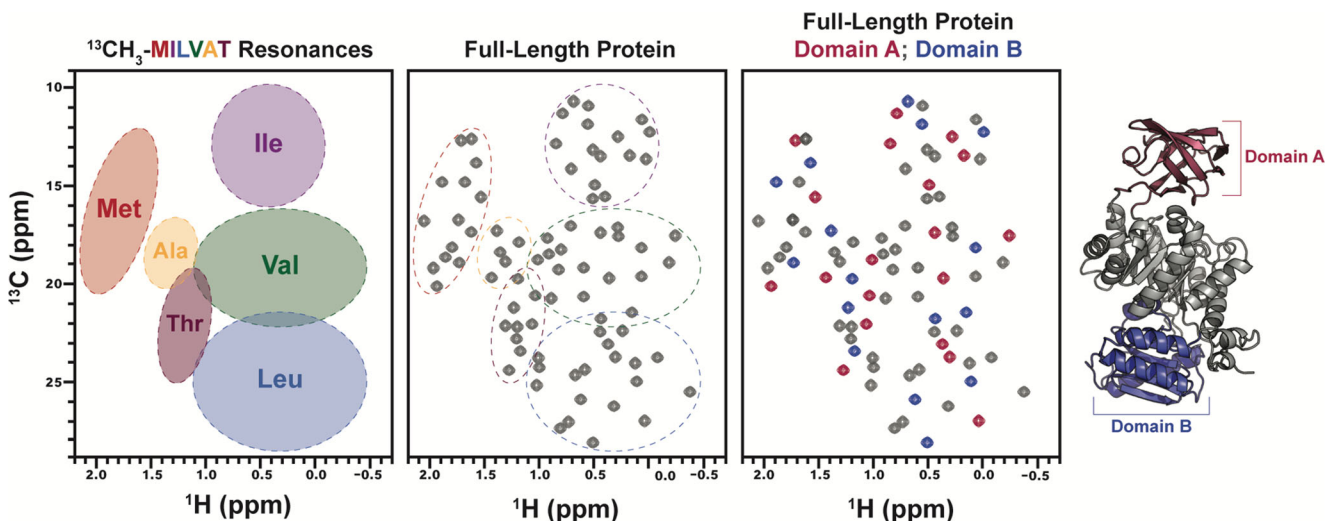
$$R_{ex} = \frac{\phi_{ex} k_{ex}}{k_{ex}^2 + \omega_e^2} \sin^2\theta^3 \tag{13}$$

The slower molecular motions probed by rcCPMG and R<sub>1ρ</sub> can be accessed, to some extent, by long MD simulations that require high computing power. Pushing computation further into the micro-millisecond regime is a new frontier and would allow for a more direct comparison with NMR data on the most critical timescale for allostery.

### Computational approaches for elucidation of allosteric pathways

#### Molecular dynamics simulations

To date, MD simulations have been powerful for investigations of conformational dynamics and subtle fluctuations associated with allostery (Dokholyan 2016; Guo and Zhou 2016; Wagner et al. 2016; Wodak et al. 2019). Notably, the group of Luthey-Schulten combined MD simulations with graph theory (Sethi et al. 2009), first describing intricate allosteric responses as a communication network. In large biomolecular complexes, i.e., a multi-domain protein with nucleic acid elements, the conformational landscape is characterized



**Fig. 5** <sup>13</sup>CH<sub>3</sub>-labeling of Met-Ile-Leu-Val-Ala-Thr residues can deconvolute NMR spectra of large proteins and effectively resolve amino acid types by chemical shift. NMR studies of small subdomains within larger proteins and complexes remain popular, and the ability to

use <sup>13</sup>C-methyl labeling to build up NMR spectra of large complexes from smaller fragments can aid in the assignment and interpretation of NMR data

by slow dynamical motions, thereby affecting the transmission of the allosteric response over longer timescales (Kern and Zuiderweg 2003). The synergistic use of MD and NMR to quantify contributions of fast timescale dynamics to protein function and allosteric regulation, most notably via order parameters ( $S^2$ ), is well-documented (Salvi et al. 2016; Wand 2001, 2017). However, a remarkable constraint to MD simulations in capturing allosteric responses of large biomolecular systems has been the difficulty of quantifying slow timescale ( $\mu\text{s}$ – $\text{ms}$ ) dynamics with computation. The recently introduced the application of Gaussian-accelerated MD (GaMD) (Miao et al. 2015) has now enabled unconstrained enhanced sampling that captures  $\mu\text{s}$ – $\text{ms}$  timescale events, facilitating a stronger degree of complementarity between simulation and NMR in studies of dynamic allostery.

GaMD adds a harmonic boost potential that smoothes the potential energy surface of the simulation system, accelerating the transitions between low-energy states. For a system of  $N$  atoms at positions  $\vec{r} = \{\vec{r}_1, \dots, \vec{r}_N\}$ , when the system potential  $V(\vec{r})$  is lower than a threshold energy  $E = V_{\max}$ , the energy surface is modified by adding a boost potential as

$$V^*(\vec{r}) = V(\vec{r}) + \Delta V(\vec{r}) \text{ for } V(\vec{r}) < E, \quad \Delta V(\vec{r}) = \frac{1}{2}k(E - V(\vec{r}))^2 \quad (14)$$

where  $k$  is the harmonic force constant. The two adjustable parameters,  $E$  and  $k$ , are determined from classical MD to collect the maximum ( $V_{\max}$ ), minimum ( $V_{\min}$ ), average ( $V_{\text{avg}}$ ), and standard deviation ( $\sigma_V$ ) of the system potential energies. When  $E$  is set to the lower bound  $E = V_{\max}$ ,  $k_0$  can be calculated as

$$k_0 = \min(1.0, k'_0) = \min\left(1.0, \frac{\sigma_0}{\sigma_V} \cdot \frac{V_{\max} - V_{\min}}{V_{\max} - V_{\text{avg}}}\right) \quad (15)$$

where  $\sigma_0$  is a user-specified upper limit (i.e., 10 kBT) for accurate energetic reweighting. Alternatively, when the threshold energy  $E$  is set to its upper bound,

$$E = V_{\min} + 1/k, \quad k_0 \text{ is } k_0 = k''_0 \equiv \left(1 - \frac{\sigma_0}{\sigma_V}\right) \cdot \frac{V_{\max} - V_{\min}}{V_{\text{avg}} - V_{\min}} \quad (16)$$

with  $k''_0$  calculated between 0 and 1. The boost potential follows a near-Gaussian distribution, allowing for robust sampling, while ensemble canonical average is reached by reweighting each point in the configuration space on the modified potential by the strength of the Boltzmann factor of the bias energy,  $\exp[\beta \Delta V(r_{(i)})]$  at that particular point. GaMD has been successfully employed to describe slow timescale biomolecular processes in remarkably large biological systems, such as the CRISPR-Cas9 complex (Palermo 2019; Palermo et al. 2017a; Ricci et al. 2019) and G protein-

coupled receptors (Miao and McCammon 2016), as well as the intracellular signaling pathway of a medically relevant G protein mimetic nanobody (Miao and McCammon 2018). GaMD is highly effective at preserving the information arising from slow timescale motion that is necessary for tracing allosteric signals in large biomolecular complexes, which has been typically carried out solely through network models.

## Network models

Network models are efficient analyses to characterize allosteric effects (Dokholyan 2016; Guo and Zhou 2016; Sethi et al. 2009; Wagner et al. 2016; Wodak et al. 2019) that rely on correlation matrices that inform whether spatially distant sites in a protein are coupled in motion over dynamic trajectories. Generalized correlation ( $\mathbf{GC}_{ij}$ ) methods (Lange and Grubmuller 2006) are often employed to capture non-collinear correlations between residue pairs ( $i, j$ ) based on Shannon's entropy and describe both linear and non-linear coupled motions. Based on  $\mathbf{GC}_{ij}$ , a dynamical network model is built describing the system as a graph of nodes (i.e., amino acid  $C\alpha$ ), connected by edges, with a distance weighted according to  $w_{ij} = -\log \mathbf{GC}_{ij}$ . The resulting weighted graph is used to structure the system's correlations through community network analysis (CNA) (Sethi et al. 2009), which defines groups of highly correlated residues and the strength of their interconnectivity to quantify chemical information flow. Residues clustered together by CNA have shown similar conformational exchange parameters in NMR experiments (Lisi et al. 2016; Rivalta et al. 2012). Further, the reshuffling of dynamic community networks upon allosteric ligand binding in MD simulations has informed NMR studies of allosteric pathways (Lisi et al. 2017).

To extract the optimal pathways for information transfer across protein networks, several approaches have been proposed. In a seminal paper (Sethi et al. 2009), Luthey-Schulten and coworkers analyzed the weighted graph by computing the "shortest pathways" connecting pairs of functionally important residues. In these weighted dynamic networks, the "shortest pathways" were calculated using the Floyd-Warshall algorithm (Floyd 1962), which sums the lengths ( $w_{ij}$ ) of all edges of various pathways connecting two distant residues and identifies the route displaying the shortest total length. Another strategy utilizes the Dijkstra algorithm (Dijkstra 1959), which is widely employed in graph theory, particularly in cartography, to determine shortest routes of travel. This algorithm uses the  $\mathbf{GC}_{ij}$  coefficients as a metric to identify routes composed by inter-node connections ( $w_{ij}$ ), that minimize the total distance (and maximize correlation) between amino acids. Thus, the Dijkstra algorithm optimizes motional momentum transport that facilitates efficient signaling between selected residues. Clusters of flexible residues identified in NMR experiments that fall along the



computationally derived route of information transfer are likely to be directly involved in the allosteric mechanism.

### Eigenvector centrality

One of the cornerstones of network theory is the concept of centrality (Doshi et al. 2016), i.e., the relative influence of a node or cluster of nodes to the network (Alvarez-Socorro et al. 2015; Borgatti 2005; Hanke and Foraita 2017; Newman 2010). A recently proposed method harnesses the so-called eigenvector centrality (EC) (Negre et al. 2018b), where the EC of a node,  $c_i$ , is defined as the sum of the centralities of all nodes that are connected to it by an edge,  $A_{ij} c_j = \frac{1}{\lambda} \sum_{j=1}^n A_{ij} c_j$ , where the edges  $A_{ij}$  are elements of the adjacency matrix  $A$  (based on  $GC_{ij}$ ), and  $\lambda$  is the eigenvalue associated with the eigenvector composed by  $c_i$  elements. This approach relies on assigning functional dynamics to the major collective mode of the system (i.e., the first eigenvector of  $A$ ), and hence,  $\lambda$  is the highest eigenvalue of  $A$ . The EC estimation quantifies the degree of connectivity of each amino acid or nucleobase in the system, measuring how well the nodes of the protein network are connected to other well-connected nodes. This enriches the information derived from betweenness centrality (BC) metrics, which measures instead how information flows between nodes (or edges) in a network (Sethi et al. 2009). Distinct from other centrality descriptors, EC also serves as a measure of the connectivity against a fixed scale when normalized and can reliably compare the connectivity of a system following allosteric effector binding. This concept has been especially critical to identifying targets for site-directed mutagenesis and inhibitor docking within allosteric networks. Residues with high EC may also be dynamic transducers in proteins, where mutations at these sites stimulate or attenuate protein motions (Lisi et al. 2017; Negre et al. 2018b). EC is also valuable to obtain the main mode of collective correlation responsible for the allosteric signal, beyond the capabilities of standard principal component methods.

### Reliability of simulated ensembles with respect to NMR experiments

One critical aspect of a combined NMR-computational approach lies in assessing the reliability of the simulated ensemble with respect to the NMR experiment. A useful way to gain insight into how well the structural and dynamical features captured by MD represent NMR data is to compute the NMR chemical shifts based on the conformational ensemble obtained via MD simulations. A straightforward way to predict the  $^1\text{H}$ ,  $^{13}\text{C}$ , and  $^{15}\text{N}$  chemical shifts from an ensemble of structures derived via MD simulations relies on the use of the SHIFTX2 code (Han et al. 2011), which employs ensemble machine learning combined with a mixed sequence, structure-

based method. The SHIFTX2 algorithm has been trained and tested with high-resolution X-ray structures ( $< 2.1 \text{ \AA}$ ) and verified chemical shifts assignments. SHIFTX2 combines two ensemble machine learning algorithms: bagging and boosting. Bagging algorithm trains “base learners” (i.e., the individual learning algorithms of the ensemble) from a random sample of the original dataset, averaged over predictions of all the individual base learners. In contrast, boosting algorithm trains subsequent base learners on mistakes of the previous base learner. SHIFTX2 further implements an algorithm to select the optimal set of features (i.e.,  $\chi^2$  and  $\chi^3$  angles, solvent accessibility, H-bond geometry, pH, temperature). By applying SHIFTX2 over MD simulation runs, one can generate average chemical shifts that are comparable to experimental values, thereby providing an assessment of how a structural ensemble derived from MD relates to that of NMR at the molecular level. It should be noted, however, that the comparison of computed and experimental chemical shifts does not probe allosteric effects directly. Indeed, chemical shifts are a local property, susceptible to changes in the electronic environment, as well as local events such as substrate binding and shifts in amino acid  $\text{pK}_a$ . Hence, allosteric signaling itself is described through the analysis of the *slow- and fast-timescale dynamics* captured through solution NMR and simulations. Nonetheless, SHIFTX2 remains an important tool to predict which amino acids may be most sensitive to changes in a dynamic ensemble.

### Case study 1: DNA polymerase $\beta$ (pol $\beta$ )

DNA polymerases are a well-established family of enzymes that have been extensively studied for their role in replicating and maintaining genomic DNA. The fidelity of DNA polymerases, or ability to select and insert correct nucleotides from an intracellular pool (i.e., A, T, G, C) during catalysis, is essential for multicellular organisms. Polymerases rely on various mechanisms for fidelity, and larger replicative polymerases contain accessory proofreading domains that confirm the inserted nucleotide bases are not mismatched. However, the maintenance of fidelity for smaller polymerases lacking these accessory domains is poorly understood. In this case study, we highlight the 39 kDa DNA polymerase (pol)  $\beta$  that lacks a proofreading domain and functions in the mammalian base excision repair (BER) pathway (Krokan and Bjoras 2013; Wallace et al. 2012). Pol  $\beta$  is composed of two domains with distinct enzymatic functions, a deoxyribose phosphate (dRP) lyase domain (8 kDa) and a nucleotide transferase domain (31 kDa) composed of a DNA binding subdomain (thumb), catalytic subdomain (palm), and nucleotide selection subdomain (fingers). Within short patch BER, pol  $\beta$  identifies and repairs gapped DNA by recruiting and inserting a deoxynucleoside triphosphate (dNTP). The ability of pol  $\beta$  to accurately discern matched (correct) versus mismatched

(incorrect) base pairs is integral to its function, highlighted by the fact that over 30% of human cancers have mutations in the gene encoding pol  $\beta$  (Marsden et al. 2017; Starcevic et al. 2004). Many of these cancer-linked mutations lead to a loss of function; however, several mutations maintain near wild-type (WT) catalytic efficiency with vastly altered fidelity and have been termed “mutator variants.” Expression of pol  $\beta$  mutants within healthy mammalian cells has been shown to cause cellular transformation and metastasis (Marsden et al. 2017; Murphy et al. 2012; Wallace et al. 2012); thus, there is great interest in elucidating the mechanistic underpinnings of the pol  $\beta$  selection process, particularly the allosteric component by which binding of the appropriate nucleotide is regulated. Several reviews (Barakat et al. 2012; Beard and Wilson 2006, 2014; Hakem 2008; Yamtich and Sweasy 2010) have rigorously described the known molecular mechanism and structure of pol  $\beta$  and related polymerases; therefore, we will only highlight important NMR and computational studies that have improved the understanding of allosteric regulation underlying its nucleotide selection.

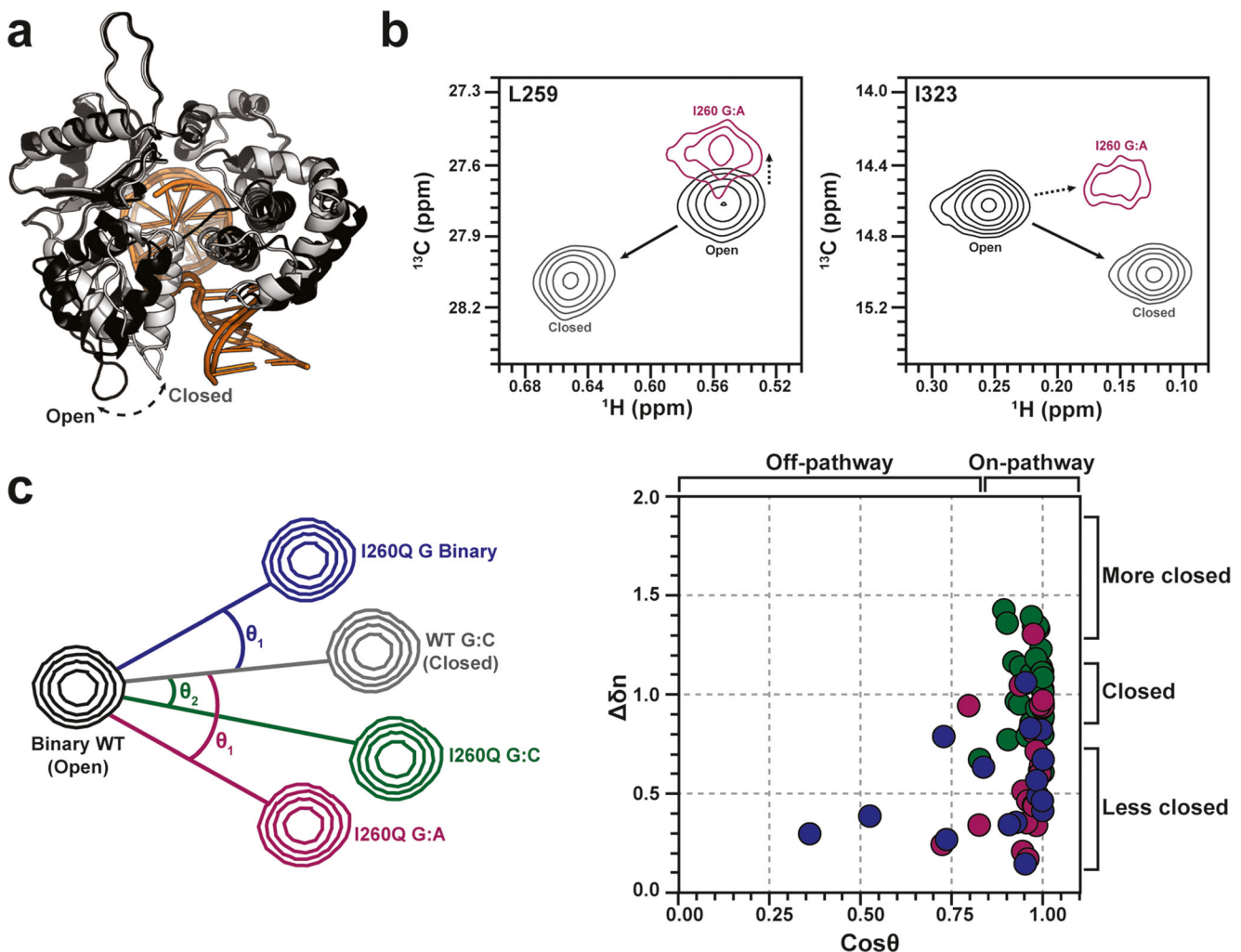
Foundational understanding of four precatalytic states of pol  $\beta$ ; apo enzyme, gapped DNA bound (binary), as well as two ternary states (gapped DNA + dNTP), in “open” and “closed” forms, has been established through X-ray crystallography and Förster Resonance Energy Transfer (FRET) (Towle-Weicksel et al. 2014). The ternary states are distinguished by 7–10 Å closure of the C-terminal nucleotide selection domain around the active site complex (Fig. 6a). However, time-averaged snapshots of these complexes in open and closed conformations do not inform the allosteric mechanism of nucleotide selection (Freudenthal et al. 2012; Sawaya et al. 1997). FRET studies explored global motions of the precatalytic transitions of pol  $\beta$  following dNTP binding to the binary (i.e., DNA-bound) enzyme (Towle-Weicksel et al. 2014) and suggested an induced fit mechanism for allosteric control of pol  $\beta$ , where the allosteric effector, a matched dNTP, causes a conformational change and subsequent non-covalent mechanistic step that regulates catalysis (Towle-Weicksel et al. 2014). This obligate non-covalent step has been suggested to be the binding of a magnesium ion or an additional conformational change. Recent NMR and computational investigations have provided further insight into intra- and intermolecular structural and dynamical processes underlying nucleotide selection and pol  $\beta$  activation.

Due to the size of pol  $\beta$  (~40 kDa), early NMR studies utilized sparse  $^1\text{H}$ ,  $^{13}\text{C}$ -methyl labeling of methionine residues to probe DNA binding to the apo enzyme and subsequent binding of a non-hydrolyzable dNTP to the binary complex. This work highlighted specific residues within pol  $\beta$  that participate in the formation of binary and ternary complexes as well as spectral signatures of the “open” and “closed” allosteric states sampled in the presence of mismatched and matched dNTPs, respectively (Bose-Basu et al. 2004). More recent

$^1\text{H}$ - $^{15}\text{N}$  experiments probing backbone amides highlighted activation of the pol  $\beta$  allosteric network upon matched dNTP binding due to changes in the flexibility of the protein backbone (Berlow et al. 2012; Loria et al. 2008), where NMR relaxation dispersion revealed specific amino acids in apo and binary pol  $\beta$  exhibiting millisecond timescale motions. Interestingly, several of these dynamic sites are coincident with mutations found in certain cancers (Starcevic et al. 2004), suggesting that flexibility of pol  $\beta$  precatalytic states, allosteric regulation of pol  $\beta$ , and its role in disease may be linked (Berlow et al. 2012; Moscato et al. 2016).

Additional support for an allosteric network in pol  $\beta$  comes from MD simulations of cancer-associated mutants with disrupted hydrogen bonding between N279 and a bound dCTP (C = cytosine) that showed strongly attenuated catalytic efficiencies ( $k_{cat}/K_M$ ) despite the fact that N279 does not make direct contact with the primer or  $\alpha$ -phosphate of dCTP (Martinek et al. 2007). Corresponding kinetic studies note the overall change in free energy ( $\Delta G$ ) more closely approximates  $\Delta G$  of dNTP binding than that of the catalytic reaction (Xiang et al. 2006), supporting a mechanism in which pol  $\beta$  is preorganized for dNTP binding and subsequent allosteric changes propagate to improve or impede pol  $\beta$  catalytic efficiency. In addition to the importance of N279 as an allosteric modulator of pol  $\beta$ , computational studies have established the importance of magnesium ions in both the active site and the allosteric network in general (Palermo et al. 2015). Computational modeling of pre- and post-chemistry steps by Schlick and coworkers showed that in the absence of magnesium, dCTP triphosphate oxygens interact with different amino acid side chains than in the presence of magnesium and suggested that closure of the nucleotide selection domain around the active site requires two magnesium ions (Yang et al. 2004). Opening of the enzyme post-chemistry is concomitant with the release of one magnesium ion from the active site in MD simulations.

Recent NMR and FRET experiments by Loria, Sweasy, and coworkers on an I260Q pol  $\beta$  mutant that retains catalytic activity but with greatly diminished dNTP discrimination, assessed the allosteric role of the hydrophobic hinge connecting the nucleotide selection and catalytic subdomains (Liptak et al. 2018; Starcevic et al. 2005). In order to understand how dNTP binding altered the structure (i.e., the degree of enzymatic closure) of I260Q relative to the WT enzyme, Liptak et al. used a vector analysis method to measure the “on” and “off” pathway trajectories of colinear NMR chemical shifts (Fig. 6b, c). In this approximation, trajectories of NMR chemical shifts that deviated from those of WT pol  $\beta$ , quantified by the angle between WT and I260Q vectors ( $\theta$ , Fig. 6c, left), were suggestive of an altered closed conformation of the enzyme, distinct from that WT pol  $\beta$ . Further, diminished magnitudes of linear I260Q chemical shifts suggested this mutation and altered structure hindered the ability



**Fig. 6** **a** X-ray crystal structures of open (black) and closed (gray) pol  $\beta$ . DNA molecules are shown in orange shades. **b** NMR chemical shift signatures of the open (black, WT without nucleotide), closed (gray, WT with matched nucleotide), and “off-pathway” (purple, I260Q mutant with mismatched nucleotide) structural states of pol  $\beta$ . **c** Schematic representation of chemical shift vector analysis reported by Liptak et al., accounting for various degrees of enzymatic closure around matched and mismatched nucleotides in WT and I260Q mutant pol  $\beta$ . WT pol  $\beta$  adopts its native, closed structure upon matched dNTP binding (black,

gray), while the cancer-associated I260Q variant populates altered structures with various dNTPs (red, blue, green) that deviate from the WT pathway. The normalized vector magnitude was calculated as  $\Delta\delta_n = \Delta\delta_{\text{exptl}}/\Delta\delta_{\text{ref}}$ , where  $\Delta\delta_{\text{ref}}$  is the WT trajectory. The degree of I260Q pol  $\beta$  closure and its overall structure relative to WT pol  $\beta$  is summarized on the correlation plot (right). I260Q data points deviating from the 1,1-cross-section display either altered structures (horizontal) or altered levels of enzymatic closure (vertical). Dots are colored according to the I260Q:DNA complexes indicated in the vector diagram at left

of pol  $\beta$  to sample a fully closed structure (Fig. 6c, right). Diminished closure in the nucleotide selection subdomain in the presence of dNTP mismatches is known to be important for fidelity of the WT enzyme, as addition of a mismatched nucleotide does not induce large chemical shift perturbations in its  $^1\text{H}$ - $^{13}\text{C}$  NMR spectrum. However, addition of a mismatched nucleotide to the I260Q mutant induces large chemical shift perturbations in the nucleotide selection domain, many of which fall along the pathway to closure seen for WT pol  $\beta$  in the presence of a matched nucleotide. The chemical shift vector methodology in Fig. 6c allowed for identification of specific residues that contribute to the pol  $\beta$  fidelity based on their ability (or lack thereof) to sample closed conformations in various dNTP-bound states.

The use of NMR and computation in the study of DNA pol  $\beta$  fidelity was crucial to refining the current view of allostery that governs the transferase activity of the enzyme. Initial X-ray crystallographic work highlighted the induced fit of a dNTP into the active site via closure of the nucleotide selection domain, while NMR and computation clarified the molecular details of the process, namely the regions of pol  $\beta$  that composed an allosteric network and displayed the flexibility necessary to alter the open/closed ensemble. In the current model of pol  $\beta$  nucleotide selection, allosteric activation is achieved by formation of a Watson-Crick pair between the incoming nucleotide and templating base, which is controlled by the affinity of pol  $\beta$  for matched/mismatched nucleotides. A chemical signal is then propagated through the hydrophobic

hinge to the active site, where the incorporation of magnesium ions modulates catalytic function. Since enzymatic catalysis ( $k_{\text{cat}}$ ,  $k_{\text{pol}}$ ) is affected by preceding steps in this information transfer, pol  $\beta$  appears to display mixed K-type and V-type allostery.

### Case study 2: CRISPR-Cas9

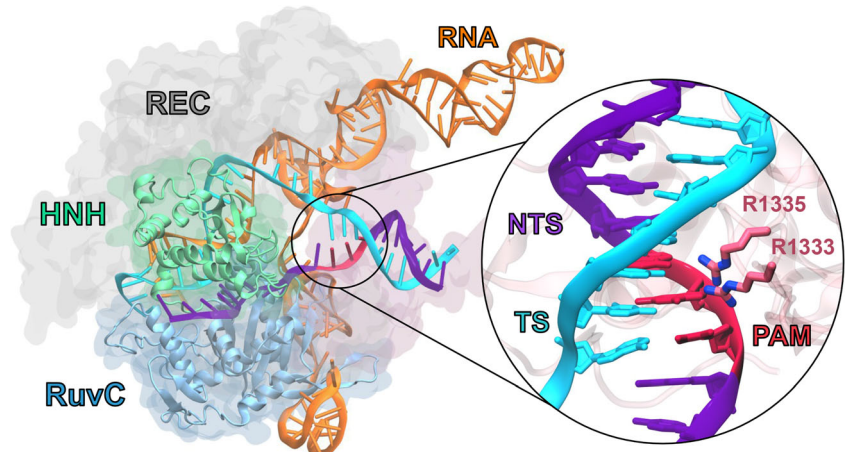
The CRISPR-Cas9 (clustered regularly interspaced short palindromic repeats–associated protein 9) system is crucial to prokaryotic organisms as an adaptive immune response against invading bacteriophages (Doudna and Charpentier 2014; Jinek et al. 2012) and is widely utilized as a genome editing tool in biotechnology and medicine (Adli 2018). Cas9 is a 160 kDa RNA-guided endonuclease that creates double-strand breaks in DNA upon site-specific recognition and binding of a short 2–5 nucleotide Protospacer Adjacent Motif (PAM) that precedes the cleavage site (Fig. 7). The multi-domain Cas9 enzyme is composed of a large recognition lobe (REC) that accommodates the formation of a RNA:DNA hybrid through three subdomains (REC1-3), and a nuclease lobe including two domains HNH and RuvC that cleave the DNA strand complementary to the guide RNA and the non-complementary strand, respectively.

The spatially separated, yet functionally connected recognition and nuclease sites of Cas9 implied an allosteric relay, which was initially probed by the insertion of exogenous PDZ domains into the native Cas9 structure by Doudna and co-workers, highlighting regions of the enzyme as “hotspots” (Oakes et al. 2016). Biochemical studies by Doudna and co-workers (Sternberg et al. 2014, 2015) have indicated that allostery in Cas9 synchronizes DNA binding, recognition, and concerted double-stranded cleavage. Early computational investigations revealed that allosteric cross-talk between the HNH and RuvC catalytic domains is essential for activation of concerted DNA cleavage (Palermo et al. 2017b) and more recently, Chen and colleagues revealed by MD that motions of the REC lobe govern the conformational changes of the HNH

domain toward cleavage (Chen et al. 2017). While X-ray structures have shown “open” (apo) or “closed” (RNA/DNA-bound) snapshots of Cas9 (Huai et al. 2017; Jiang and Doudna 2015; Jinek et al. 2014; Nishimasu et al. 2014), it is widely hypothesized that Cas9, like other allosteric proteins (Koshland Jr. et al. 1966; Lukin et al. 2003; Monod et al. 1965; Yuan et al. 2015) populates numerous microstates between these structures (Dagdas et al. 2017).

In this respect, the nature of allosteric communication within this protein-nucleic acid complex is poorly understood, though recent investigations have provided intricate detail about how nucleic acid binding information is transmitted across the multi-domain Cas9 structure to control the activation of the catalytic HNH nuclease. HNH displays a remarkable conformational mobility, observed by cryo-EM (CryoEM; EMD-3277) to be at lower resolution (8–10 Å) than the overall structure (6 Å) (Jiang et al. 2016). The first all-atom MD simulation of Cas9 (Palermo et al. 2016) also highlighted the “striking” flexibility of HNH with respect to the remaining protein subunits that enable HNH to rapidly sample its activated state primed for DNA cleavage. HNH activation is dependent on the conformational dynamics of the REC lobe of Cas9, where allosteric function relies on the opening of the REC3 region (amino acids 480–718) to integrate the incoming RNA:DNA hybrid (Chen et al. 2017). This interaction is thought to propagate a biological signal to REC2 (amino acids 180–308), which directly contacts HNH, enabling its approach to the DNA cleavage site on the complementary strand. In spite of extensive structural characterization of Cas9, high-resolution snapshots of its fully activated state have yet to be observed, perhaps due to the inherent flexibility of HNH when engaged in enzymatic function. However, early computational studies employing targeted approaches and accelerated MD have provided the first structural information on the activated complex, including the structure of an activated HNH domain (Palermo et al. 2017a; Zuo and Liu 2017). Extensive MD simulations have also revealed that the REC1-3 regions collectively move with HNH during its

**Fig. 7** Structure of the *Sp* CRISPR-Cas9 system, composed of the endonuclease Cas9 in complex with a guide RNA and a target DNA. The Cas9 enzyme is shown in molecular surface, highlighting the HNH (green), RuvC (blue), and REC (gray) domains with different colors. The RNA (orange) as well as the target DNA strand (TS, cyan) and the non-target DNA strand (NTS, violet) are shown as ribbons. A close-up view of the PAM-binding region is shown on the right

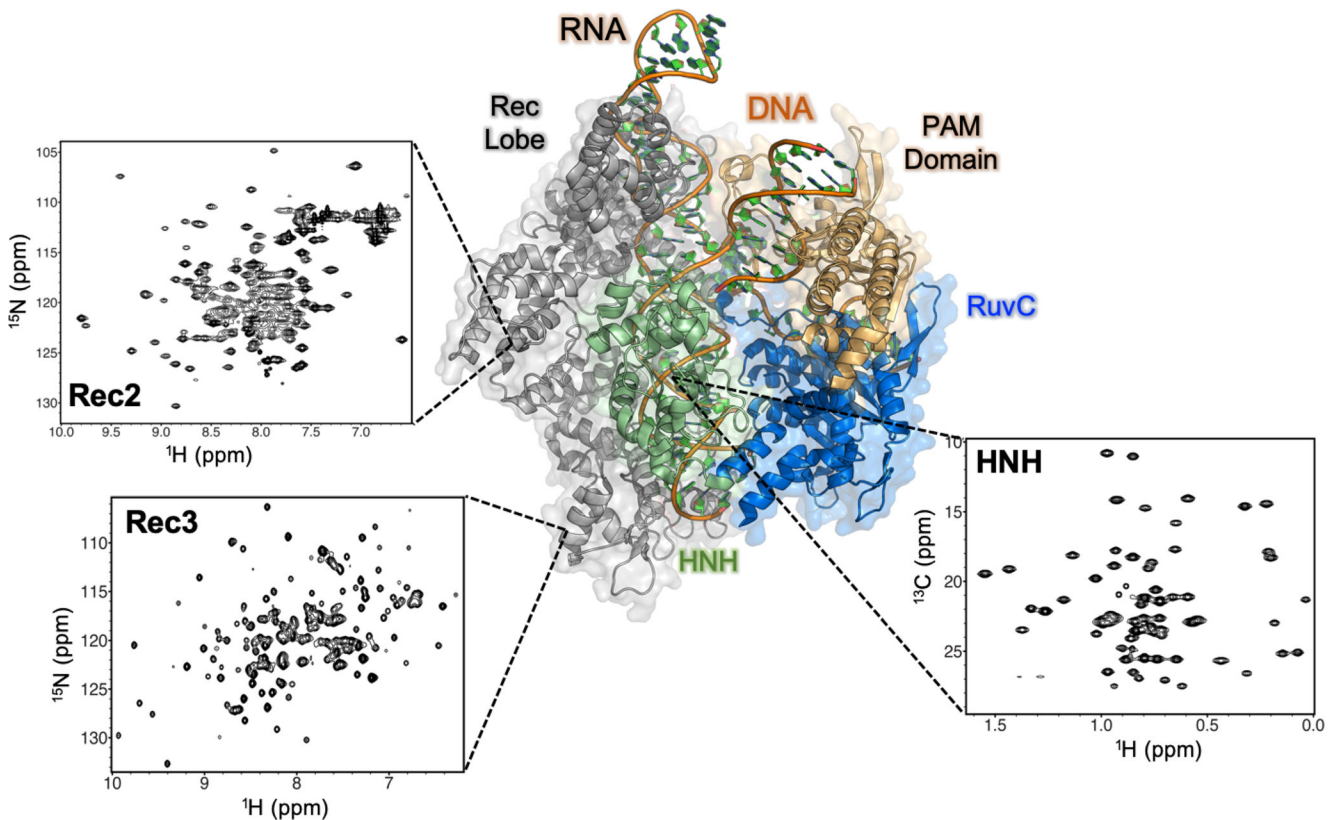


activation (Palermo et al. 2018), supporting the “allosteric hypothesis” of Chen and colleagues (Chen et al. 2017). However, the nature of the allosteric control that REC exerts on HNH and the mechanism of its signal transduction remains unclear. Each of these functionalities relies on molecular motion and is thus amenable to characterization by synergistic NMR and MD approaches. Thus, the most recent biophysical probes of Cas9 are integrating solution NMR (Lisi and Loria 2016a) and MD or novel GaMD (Miao et al. 2015) with network models (Sethi et al. 2009) and newly developed centrality analysis methods (Negre et al. 2018a) to track the molecular motions at the core of the allosteric communication between REC, HNH, and RuvC.

Since coupling of NMR data to all-atom MD simulations has significantly increased our understanding of dynamic allostery in smaller protein-nucleic acid complexes (Adhikarsan et al. 2017; Lisi and Loria 2016a; Wodak et al. 2019), these methodologies are being extended to the Cas9 machinery (Belato et al. 2019), building on outcomes of prior MD simulations (Palermo 2019; Palermo et al. 2016, 2017a, b, 2018; Ricci et al. 2019; Zuo and Liu 2017). Several domains critical to Cas9 allosteric control have been studied by  $^1\text{H}$ - $^{15}\text{N}$  and  $^1\text{H}$ - $^{13}\text{C}$  NMR (Belato 2019) (Fig. 8) in order to generate fingerprints of the key players in Cas9 allostery and connect

with previous computational work highlighting the dynamic interplay of HNH, REC2, and REC3. NMR spin relaxation experiments reveal multi-timescale dynamics within this axis of Cas9, which may be involved in information transfer for concerted cleavage of the two DNA strands. Regions of fast (ps–ns) dynamics are consistent with those proposed by early MD studies of Cas9 allostery, but further investigation of these dynamic properties is required to decipher the mechanism of communication between these domains.

The role of the PAM sequence is also intriguing, as its binding initiates DNA association and cleavage, triggering inter-dependent molecular motions of the Cas9 domains (Palermo et al. 2017b; Sternberg et al. 2014). Specifically, PAM has been identified as an “allosteric activator” of Cas9 function from MD simulations performed of Cas9 binding to a ‘5-TGG-3’ PAM sequence (i.e., Cas9–PAM) (Anders et al. 2014), and on its analogue, crystallized without PAM (Nishimasu et al. 2014). A total of  $\sim 13 \mu\text{s}$  of aggregate MD runs probing the impact of PAM binding on the dynamics of Cas9 revealed that PAM induces an “open-to-close” conformational transition, as indicated by principal component analysis. This conformational change agrees with the transition hypothesized for nucleic acid binding (Jiang et al. 2016; Palermo et al. 2017a) and reveals that PAM is essential in

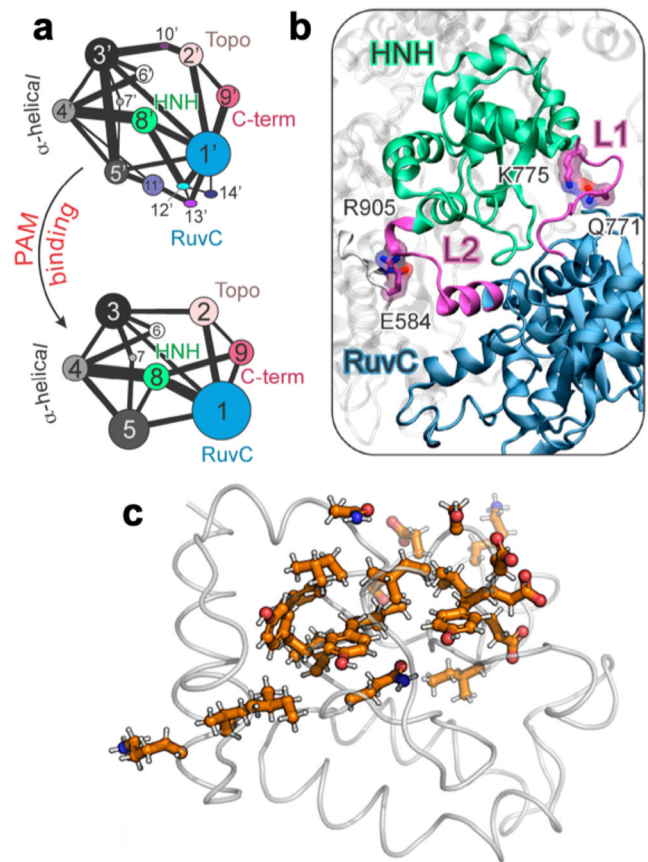


**Fig. 8** NMR spectra of critical domains of Cas9 implicated in its allosteric mechanism.  $^1\text{H}$ - $^{15}\text{N}$  correlation spectra of backbone amide residues are shown for the REC2 and REC3 regions, while a  $^1\text{H}$ - $^{13}\text{C}$  methyl spectrum reporting on Ile, Leu, Val, Ala, Met, and Thr side chains

is shown for HNH. NMR studies of the structure and dynamics of these critical domains are required to confirm the allosteric relay proposed by MD simulations

the process. Analysis of coupled motions through a generalized correlation (GC) method revealed that PAM binding significantly strengthens the motional correlations of Cas9, particularly at HNH, inducing a shift in its conformational dynamics consistent with the function of other allosteric effectors (Guo and Zhou 2016). Notably, by introducing a per-domain correlation score ( $C_{s_i}$ ) matrix that accumulates (and normalizes) the  $GC_{ij}$  coefficients over each domain, high inter-domain correlations are observed for the motions of HNH and RuvC. This dynamic cross-talk is not observed in the absence of PAM, indicating that PAM binding induces the coupled dynamics of HNH and RuvC (Sternberg et al. 2015). In the presence of PAM, HNH also correlates strongly with the  $\alpha$ -helical REC lobe, supporting a direct signal transfer (Chen et al. 2017) and suggesting that PAM acts as a long-range allosteric effector (Sternberg et al. 2014).

To further understand the role of PAM in the allosteric network, community network analysis (Rivalta et al. 2012; Sethi et al. 2009) was applied to structure Cas9 correlations (Fig. 9a, described above in computational approaches). This analysis shows that in the absence of PAM (Fig. 9a, top panel), a high fragmentation of the community structure is observed, which is then reduced upon PAM binding (Fig. 9a, bottom panel). Thus, in the absence of PAM, communication between regions of Cas9 is weak, hampering facile signal transduction. Indeed, upon PAM binding, the communities reorganize (evidenced by changes in the edge betweenness, *i.e.*, thicker bonds) to strongly couple REC2 (community 4), HNH (8), and RuvC (1), a sign of increased correlation or signal transfer and supporting the role of PAM as an effector of inter-domain cross-talk (Sternberg et al. 2014). These community networks were analyzed to determine the route(s) of information transfer between HNH and RuvC. The allosteric “pathways” between residues computed from the dynamic network as a sum of their edge lengths, identifying the most likely signaling routes (*i.e.*, the pathways exhibiting the shortest edge lengths) (Rivalta et al. 2012; Sethi et al. 2009). This analysis revealed that the PAM-mediated allosteric signaling is likely to be transduced through the L1/L2 flexible loops, connecting HNH and RuvC (Fig. 9b). This is consistent with the origin of a dynamic pathway throughout HNH determined by MD (Fig. 9c), and also agrees well with previous experimental evidence, suggesting an allosteric signal transport across REC–HNH–RuvC (Chen et al. 2017). This conclusion is supported by previous structural studies suggesting that conformational changes in L1 and L2 could transmit information about HNH activation to RuvC (Jiang et al. 2016). The importance of single amino acids within these loops was quantified through an analysis of the node betweenness (*i.e.*, the number of shortest pathways that cross the node). This study indicated that residues displaying the highest node betweenness, Q771 and E584 (within L1 and L2, respectively) engage in interactions with K775 and R905, forming essential edges



**Fig. 9** **a** Community network analysis of the Cas9 structure in the absence (top) and presence (bottom) of PAM (DNA). Bond thickness between communities represents the strength of the intercommunication. **b** Allosteric pathway captured from the dynamic network, showing biological information flow between the HNH (green) and RuvC (blue) nucleases utilizes the L1 and L2 loops. Residues displaying the highest node betweenness, Q771 and E584, engage in interactions with K775 and R905, forming essential edges of communication between HNH and RuvC. **c** Proposed network of flexible residues computed from dynamical network models of Cas9 after optimizing motional correlation between residues adjacent to RuvC and REC2. Portions of this figure are reproduced with permission from Palermo et al. *J. Am. Chem. Soc.* 2017, 139, 16028. Copyright American Chemical Society

of communication between HNH and RuvC. K775A or R905A mutations were shown to alter Cas9 selectivity (Chen et al. 2017; Slaymaker et al. 2016), opening an avenue for controlling nuclease activity through its allosteric signaling. Overall, these collective studies indicate that PAM binding induces a tight cross-talk between HNH and RuvC (Palermo et al. 2017b; Sternberg et al. 2014, 2015). However, data collected to this point do not yet clarify how the effect of PAM binding propagates to the nuclease domains. This open question is currently driving additional integrative NMR and computational research, with the goal of deciphering the mechanism of information transfer between the PAM-binding site and the HNH and RuvC domains.

## Summary

Allostery in biomolecular assemblies provides a high level of functional control over vital biological processes. The combination of experimental solution NMR with in silico techniques can elucidate the mechanisms by which allosteric control is maintained in nucleoprotein assemblies. The highlighted case studies illustrate a broad range of molecular size and structural complexity that can be explored with these methodologies.

**Acknowledgments** G.P. acknowledges funding from the National Science Foundation (NSF), as this material is based upon work supported by the National Science Foundation under Grant No. CHE-1905374. GPL acknowledges funding from the COBRE Center for Computational Biology of Human Disease (NIGMS P20GM109035).

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