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## Second Harmonic 527-GHz Gyrotron for DNP-NMR: Design and Experimental Results

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### Abstract

We report the design and experimental demonstration of a frequency tunable terahertz gyrotron at 527 GHz built for an 800 MHz Dynamic Nuclear Polarization enhanced Nuclear Magnetic Resonance (DNP-NMR) spectrometer. The gyrotron is designed at the second harmonic ( $\omega = 2\omega_c$ ) of the electron cyclotron frequency. It produces up to 9.3 W continuous microwave (CW) power at 527.2 GHz frequency using a diode type electron gun operating at  $V = 16.65$  kV,  $I_b = 110$  mA in a  $TE_{11,2,1}$  mode, corresponding to an efficiency of  $\sim 0.5\%$ . The gyrotron is tunable within  $\sim 0.4$  GHz by combining voltage and magnetic field tuning. The gyrotron has an internal mode converter that produces a Gaussian-like beam that couples to the  $HE_{11}$  mode of an internal 12 mm i.d. corrugated waveguide periscope assembly leading up to the output window. An external corrugated waveguide transmission line system is built including a corrugated taper from 12 mm to 16 mm i.d. waveguide followed by 3 m of the 16 mm i.d. waveguide. The microwave beam profile is measured using a pyroelectric camera showing  $\sim 84\%$   $HE_{11}$  mode content.

## Keywords

Terahertz; frequency tunable gyrotron; Dynamic Nuclear Polarization; electron beam; Gaussian beam; Corrugated waveguide

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## I. INTRODUCTION

NUCLEAR Magnetic Resonance spectroscopy at higher magnetic fields has become a unique tool for improving the spectral resolution to non-invasively probe atomic structures of biological macromolecules such as membrane proteins and non-protein molecular assemblies [1–4]. However, the problems associated with the poor polarization in conventional NMR spectroscopy, typically  $< 0.01\%$ , of the proton nuclei leads to low sensitivity and poor signal. Using microwaves matching the corresponding electron frequencies in the same electron-nuclei system one can dynamically transfer the polarization of unpaired electrons to the protons by redistributing the Boltzmann population of the spin states among the  $^1\text{H}$  protons, a process called Dynamic Nuclear Polarization or DNP-NMR. This increases the population difference between the nuclei spin-up and spin-down states compared to the thermal equilibrium leading to the increased signal [5–7]. For the current state of the art DNP-NMR systems, the microwave sources required for this purpose are typically in the frequency range of 250 – 700 GHz and power of the order of several watts to several tens of watts. Presently, gyrotrons are the most suitable microwave sources for this purpose.

Gyrotrons have historically been developed mainly for plasma heating in fusion tokamak reactors [8–13]. The required microwave power for this application is however in the megawatt range and the frequency requirements are below those for the DNP-NMR application [14, 15]. Other applications have also emerged in the past two decades where the required power is in the range of tens of watts to a few tens of kW [16–25]. A major application which has pushed the development of gyrotrons in the THz range is DNP-NMR spectroscopy. Since the required output power for this application is watts to tens of watts, therefore the development of gyrotrons is significantly relaxed compared to the megawatt gyrotrons in terms of technological challenges such as high-power handling, use of complex power supplies and significant reduction in cooling requirements. However, pushing the frequency into the sub-THz and THz range poses new challenges such as operation in the higher electron cyclotron harmonics, reduced efficiency, frequency tunability, high ohmic loss and higher operating mode competition etc. Despite these challenges, there is an intense interest in developing gyrotrons for DNP applications due to the vast potential of the DNP technique to revolutionize the field of NMR spectroscopy [26–44].

In addition to the requirement of watts of output power and excellent frequency stability, it is also desirable to have up to 1 GHz of frequency tunability to optimize the DNP matching conditions [45]. The frequency tunability is achieved by exciting successive axial cavity modes having very close eigenfrequencies [26, 27]. The work presented in this paper benefits from the previous successful work done by Torrezan et al. [28] at 460 GHz.

In this paper we report on the design, fabrication and testing of a 527 GHz gyrotron for use on an 800 MHz DNP-NMR spectrometer. This paper is organized as follows; Section-II describes the design of the gyrotron and the fabrication and experimental setup. Section-III describes the experimental results followed by conclusions in Section-IV.

## II. DESIGN AND EXPERIMENTAL SETUP

The schematic of the 527 GHz gyrotron showing different components is shown in Fig. 1. The design parameters of the gyrotron are listed in Table 1. The magnetic field is provided by a 10 T cryo-free superconducting magnet (SCM), identical to the one used by Alberti et al. [33] and Rozier et al. [43] for their 263.5 GHz DNP gyrotron. The electron beam is generated by a diode type electron gun designed using the electron optics code MICHELLE [46]. The electron gun is operated using a 4 kW CW power supply capable of providing a maximum operating voltage of 20 kV and a maximum beam current of 200 mA. The electron gun is designed to operate at  $\sim 17$  kV and 200 mA of beam current. The beam radius,  $r_e$  at the center of the cavity is 0.97 mm where it is compressed by a factor of the square root of 25.4 as compared to the emitter radius. The simulated electron velocity ratio,  $\alpha = v_{\perp}/v_{\parallel}$ , and perpendicular velocity spread,  $\Delta v_{\perp}$ , are shown in Fig. 2. The optimized values are  $\cong 1.8$  and  $<3\%$ , respectively, for the operating voltage in the range of 15 kV to 17 kV. An additional external Gun Coil electromagnet is used to fine tune the electron beam parameters, as shown in Fig. 1. The gun coil is typically used to reduce the magnetic field at the location of the electron gun, leading to negative values of the strength of the magnetic field of the gun coil.

The initial design of the resonator cavity and its operating mode selection were done by calculating the coupling coefficients of various modes at 527 GHz. The operating mode-  $TE_{11,2,q}$ ,  $q$  being the axial mode index, was chosen by optimizing various factors such as maximum coupling with the electron beam of radius 0.97 mm, separation from nearby competing fundamental and second harmonic modes, ohmic loss of the operating mode etc. The choice of the  $TE_{11,2}$  as the operating mode was also inspired by our past success of the 460 GHz gyrotron which successfully produced over 15 W [28]. The main operating mode is the  $TE_{11,2,1}$  mode ( $q = 1$ ), while the higher order modes, ( $q = 2,3,4$ , etc.) provide for frequency tuning. The interaction resonator cavity profile was optimized using cold cavity calculations. The total length of the cavity is  $\approx 78$  mm which consists of a straight resonator section of radius  $r_c = 1.593$  mm and length 25 mm and three consecutive up-tapers at  $0.3^\circ$ ,  $0.6^\circ$ , and  $1.9^\circ$  respectively. The length of the straight resonator section is  $\approx 45\lambda$  which is significantly longer compared to the length of cavities used in gyrotrons operating at power levels in the megawatts range for electron cyclotron resonance heating in fusion applications. The longer cavity was selected to consecutively excite the higher order axial modes,  $q$ , for continuous frequency tunability. Fig. 3 shows the cavity profile along with the normalized electric field amplitude of the  $TE_{11,2,1}$  mode. The calculated diffractive and ohmic quality factors of the operating modes are  $97500/q^2$  and 7500, respectively, assuming the electrical conductivity equal to one half of ideal copper. The diffractive quality factor is calculated using the cold-cavity approximation.

After the cold cavity axial field profile and the coupling coefficients were calculated, linear theory [48] was used to calculate the start oscillation current of the main operating mode-TE<sub>11,2,1</sub>, and the neighboring competing fundamental and second harmonic modes. This is shown in Fig. 4. The transverse electric field profile of the TE<sub>11,2</sub> mode and the circle representing the electron beam are also shown in the center of the figure. The starting current was calculated for an operating voltage of 16.7 kV,  $a = 1.8$ ,  $r_e = 0.97$  mm,  $v_{\perp}/v_{\parallel} = 3\%$  and for first axial mode  $q = 1$  of each mode shown in Fig. 4. It is clear that the main operating mode-TE<sub>11,2,1</sub> is well separated from the nearest fundamental mode TE<sub>7,1,1</sub> and the second harmonic modes TE<sub>5,4,1</sub> and TE<sub>-8,3,1</sub>. The predicted start oscillation current for the operating mode is  $\sim 20$  mA. The frequency tuning is accomplished by exciting a series of axial modes, TE<sub>11,2,q</sub>. The start oscillation current of these higher order axial modes ( $q$ ) was also calculated and is described below with the experimental results.

After the interaction, the microwaves were converted from the TE<sub>11,2</sub> mode to a Gaussian like beam using a quasi-optical mode converter assembly. The mode converter consists of a helical cut Vlasov type smooth wall mode launcher with waveguide radius of 2.35 mm and two curved mirrors, one parabolic and one spherical. The mode converter was designed using the electric-field integral equation code Surf3d [49]. After the second mirror the microwave beam was coupled into a corrugated waveguide of 12 mm diameter placed inside the tube as part of the periscope assembly with two miter-bends. Fig. 5(a) shows the fabricated and assembled mode converter assembly. This periscope assembly brings the microwaves in the form of an HE<sub>11</sub> mode to the window. The window is made of fused quartz and mounted on a conflat flange. Fig. 5(c) shows the electric field intensity in the cut plane along the axis of the tube calculated using Surf3D and radiated from mirror 1 to mirror 2 and then to the location of the corrugated waveguide aperture shown by the dashed line. The intensity profile at the waveguide aperture is also shown in the Fig. 5(b). The calculated Gaussian mode content of the beam just before the waveguide aperture is  $\sim 90\%$  and the HE<sub>11</sub> mode content coupled to the waveguide is  $\sim 84\%$ . The window used in the gyrotron tube is mounted on a CF flange and has a thickness of 3.352 mm and diameter of 42 mm.

The beam tunnel of the gyrotron consists of a section of axially slotted stainless-steel tube and a section of SiC and copper loaded disks to suppress any parasitic oscillations. The SiC is a standard absorbing material used in the gyrotron beam tunnel for low frequency oscillation suppression. The dielectric constant and loss tangent of different grades of the SiC used in gyrotron beam tunnels are  $\epsilon = 11.5$  to 13.5 and  $\tan \delta = 0.02$  to 0.06 [50, 51].

The collector is a simple undepressed water cooled collector at ground potential with respect to the cathode. The collector and cavity cooling requirements were calculated from the electron beam power used in the operation and the generated microwave power in the cavity from beam wave interaction. The cooling circuits were fabricated accordingly. All components were fabricated and tested in-house. A water chiller of cooling capacity of 5 kW is used in the experiment to meet the cooling requirements of the cavity and the collector. The cathode emitter was purchased from 3M Ceradyne Inc. A photograph of the installed gyrotron operating in the lab is shown in Fig. 6.

### III. EXPERIMENTAL RESULTS

The gyrotron was operated from 15 kV to 17 kV operating voltage and at magnetic field of 9.68 T to 9.75 T and the output power and frequency were measured. The internal corrugated waveguide positions the microwave beam at the center of the quartz window for coupling to the outside corrugated taper of 12 mm to 16 mm i.d. After the corrugated taper the beam is coupled into a 3 m long 16 mm i.d. corrugated waveguide. The corrugated waveguides were machined for operation at 527 GHz [52]. The output power was measured using a calibrated Scientech power meter. The calibration was verified using another THz power meter 3A-P-THz from Ophir Optonics Sol.

The start oscillation current of the operating modes,  $TE_{11,2,q}$  was first measured at an operating voltage of 16.7 kV and the gun coil magnet operating at  $\sim -37$  mT. The magnetic field was varied from 9.69 T to 9.75 T. The microwave output levels were measured using the 3A-P-THz power meter capable of measuring down to 0.1 mW of power. The measured start oscillation current is plotted in Fig. 7 along with the calculated start current of the axial modes  $TE_{11,2,q}$  for  $q = 1$  to 4, using the linear theory. The linear theory assumes a fixed axial field structure leading to separate starting current estimates for each axial mode  $q$ . In practice, these axial modes are coupled through the electron beam and form a continuously tunable range of output frequencies [53]. A minimum start current of  $\sim 21$  mA is observed experimentally. An excellent agreement is seen between the measured minimum start current and the theoretical start current for axial modes  $q = 1, 2, 3, 4$ .

The gyrotron was operated for measuring the power and frequency at a magnetic field of 9.708 T. The maximum output power of 9.3 W was observed for the operating voltage of 16.7 kV and beam current of 110 mA. The measured frequency at maximum power was 527.2 GHz. The measured power and frequency by varying the voltage at fixed magnetic field is shown in Fig. 8. In the voltage tuning measurement, the gun coil magnet was varied up to  $-35$  mT to maximize the output power. The lowest frequency of 527.12 GHz is observed with output power of 9.0 W. A total of  $\sim 0.35$  GHz continuous frequency tuning is observed. The output frequency is measured using an even harmonic mixer (EHM) in the WR1.9 frequency band from VDI Inc, and a heterodyne system with a local oscillator (LO) synthesizer operating at 18 GHz to 26 GHz. The output signal was verified by observing consecutive harmonics of the LO on a 2.5 GHz bandwidth oscilloscope. Fig. 9 shows the magnetic field tuning for output power and frequency at a fixed voltage of 16.7 kV. A total frequency tuning of  $\sim 0.25$  GHz is observed by varying the magnetic field. The multimode self-consistent nonlinear code MAGY [47] is used to theoretically calculate the output power and frequency by varying the magnetic field. The results are shown in Fig. 10 for the operating parameters used in Fig. 9. The operating voltage of 16.7 kV and beam current of 110 mA,  $\alpha = 1.85$  and  $v_{\perp}/v_{\parallel} = 5\%$  were used in simulations. The conductivity used in the calculations is one half of the ideal copper conductivity. The magnetic field center and the resonator cavity center are axially aligned in the simulations based on the tube position in the actual experiment for better comparison between the theory and experiments. The cavity radius used in the calculations is the actual measured cavity radius of  $r_c = 1.593$  mm. Although the velocity spread value calculated from MICHELLE is 2.5%, as shown in Fig. 2, in MAGY calculations we use 5% as it shows better agreement with the experiments. It is

not unusual that the velocity spread value in reality is higher than the calculations because there are additional contributors to it which are not fully considered in simulations in detail, such as misalignments in the gun assembly, emitter surface roughness and temperature nonuniformity. A theoretical frequency tuning of  $\sim 320$  MHz is observed by varying the magnetic field by successive excitation of axial modes,  $q$ . A very good qualitative agreement can be seen between the theory of Fig. 10 and the experimental results of Fig. 9 for both frequency tuning and output power considering the  $\sim 86\%$  efficiency of the Vlasov type mode converter. The overall frequency tunability of  $\sim 0.4$  GHz is observed by combining the voltage and magnetic tuning. An output power of more than 4 W is observed for  $\sim 0.35$  GHz of frequency variation. This is shown in Fig. 11. The voltage is varied from 15.5 kV to 17.0 kV and the magnetic field is varied from 9.69 T to 9.75 T and the gun coil field is varied from  $-25$  mT to  $-35$  mT. By careful alignment of the gyrotron tube in the magnet bore in transverse and axial directions it was verified that the gyrotron is operating in the second harmonic  $TE_{11,2}$  mode and no spurious modes, fundamental or second harmonic, were excited. The output power was also measured by varying the beam current at fixed voltage of 16.7 kV and magnetic field of 9.71 T. This is shown in Fig. 12. The gyrotron was operated continuously for 96 hours under computer control. The variation in power was less than 2% and the frequency was stable to within 2 MHz. No thermal issues were discovered during the operation of the gyrotron. For CW operation, the water temperature and flow rate at both the cavity and collector cooling channels are monitored in real time in LabView.

The output of the gyrotron is coupled into a 16 mm i.d. corrugated waveguide using a corrugated taper from 12 mm i.d. to 16 mm, with expected loss of less than 1% [52]. The corrugated waveguides of 30 cm-long sections are fabricated in house using a helical tap of nearly square shape corrugation [54] where depth and width of the corrugations are adapted for operation at  $\sim 527$  GHz. The waveguides are aligned and the beam profile at the end of the  $\sim 3$  m long section of corrugated waveguide is measured using a pyroelectric camera. The field profile was measured in a few different locations along the waveguide axis parallel to the aperture. The measured field profiles at 20 mm and 38 mm from the waveguide aperture are shown in Fig. 10. The calculated Gaussian content of the propagated field is  $\sim 92\%$ . The field measured in different locations along the waveguide axis was used to retrieve the phase of the beam and propagated back to the aperture [55]. The  $HE_{11}$  mode content of the field is thus measured to be  $\sim 84\%$ .

#### IV. CONCLUSIONS

We have designed, fabricated and experimentally tested a 527 GHz gyrotron operating at the second harmonic. The gyrotron has generated an output power of 9.3 Watts at 16.7 kV and 110 mA. This corresponds to an operating efficiency of  $\sim 0.5\%$ . Higher output power can be expected by further raising the operating beam current. This may be possible after additional processing of the gyrotron. The excitation of consecutive higher order axial modes by voltage and magnetic field variation led to a frequency tunability of 0.4 GHz. An output power of more than 4 W was observed for this frequency bandwidth. The output beam of the gyrotron is in the form of a corrugated waveguide  $HE_{11}$  mode with  $\sim 84\%$   $HE_{11}$  mode content.



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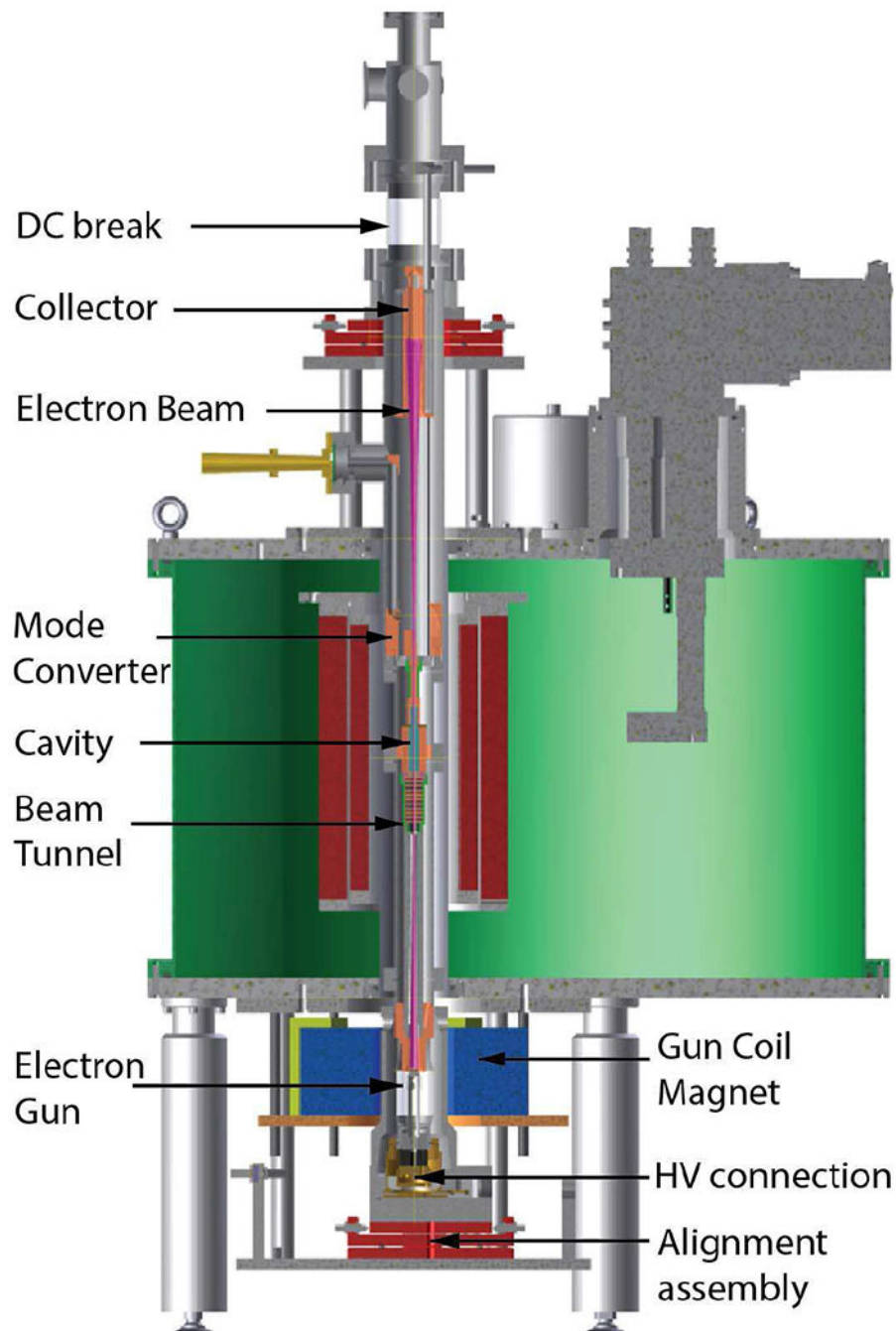
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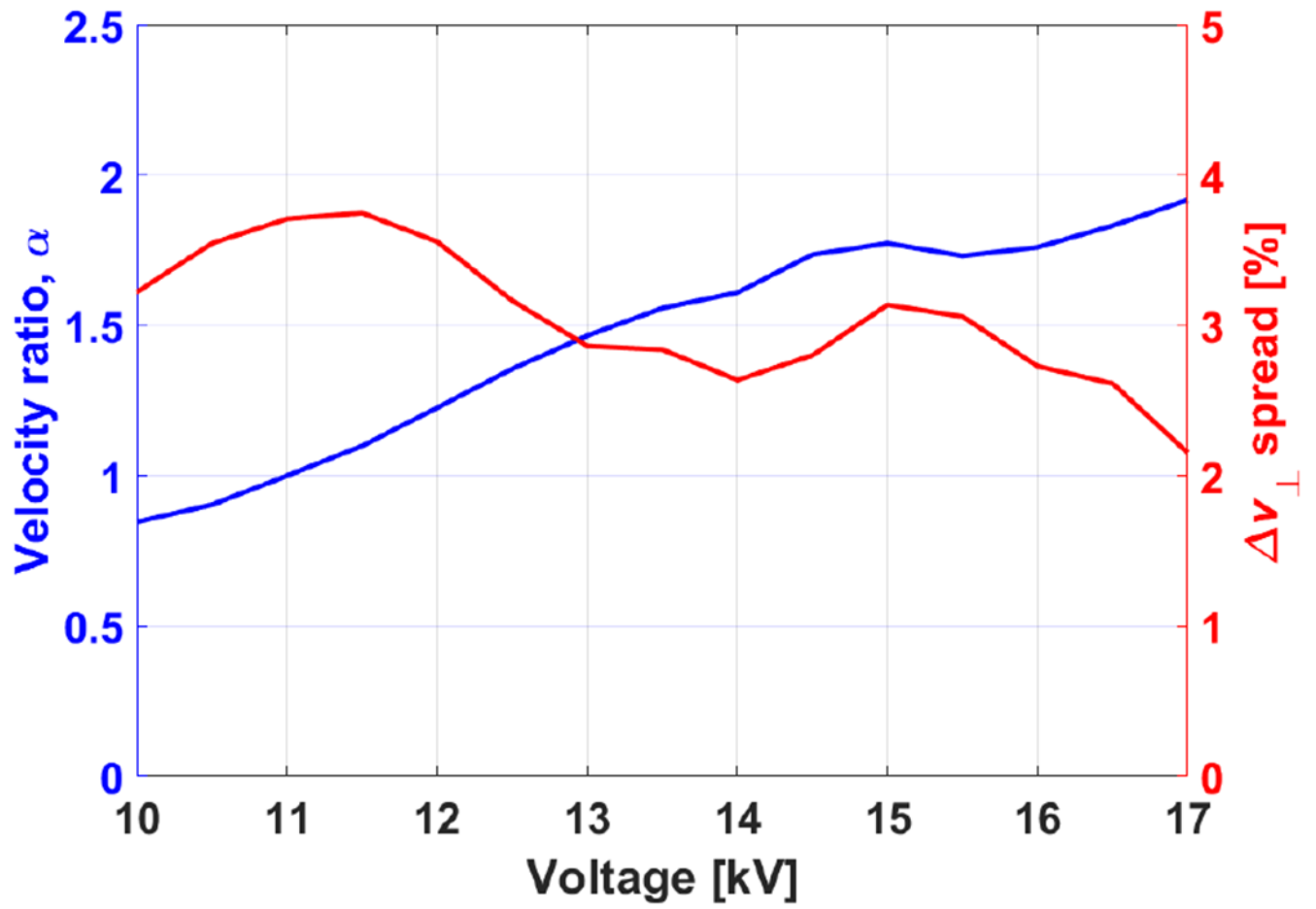
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**Fig. 1.** The schematic of the 527 GHz gyrotron in a 10 T SCM (shown in green) is shown. Different components of the tube assembly are labeled. The gyrotron tube is mounted on two assemblies at the top and bottom for alignment along the magnetic field axis.



**Fig. 2.** Calculated perpendicular to parallel velocity ratio,  $\alpha$ , and perpendicular velocity spread,  $v_{\perp}$  w.r.t. the operating voltage for the diode type electron gun.

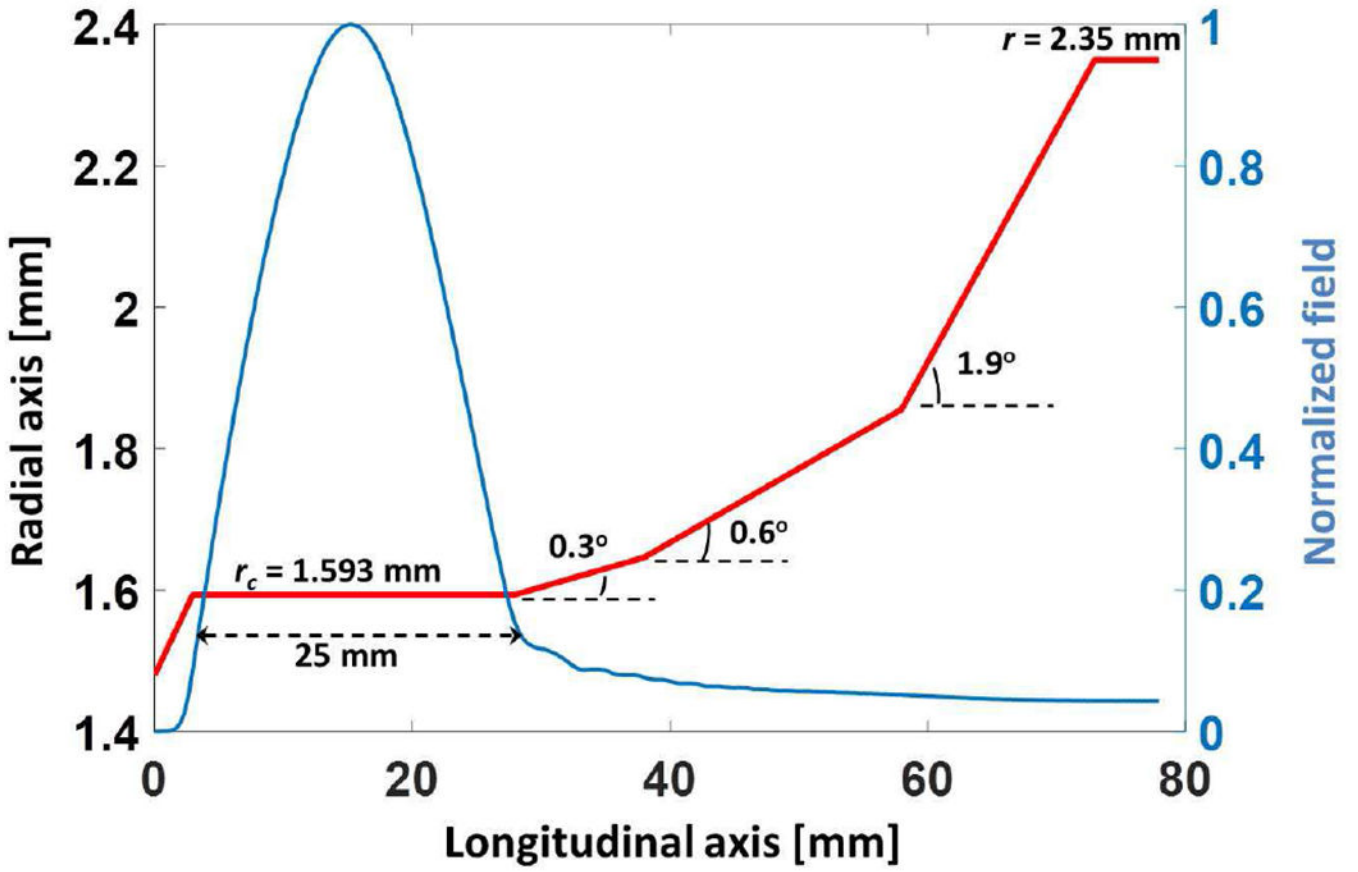
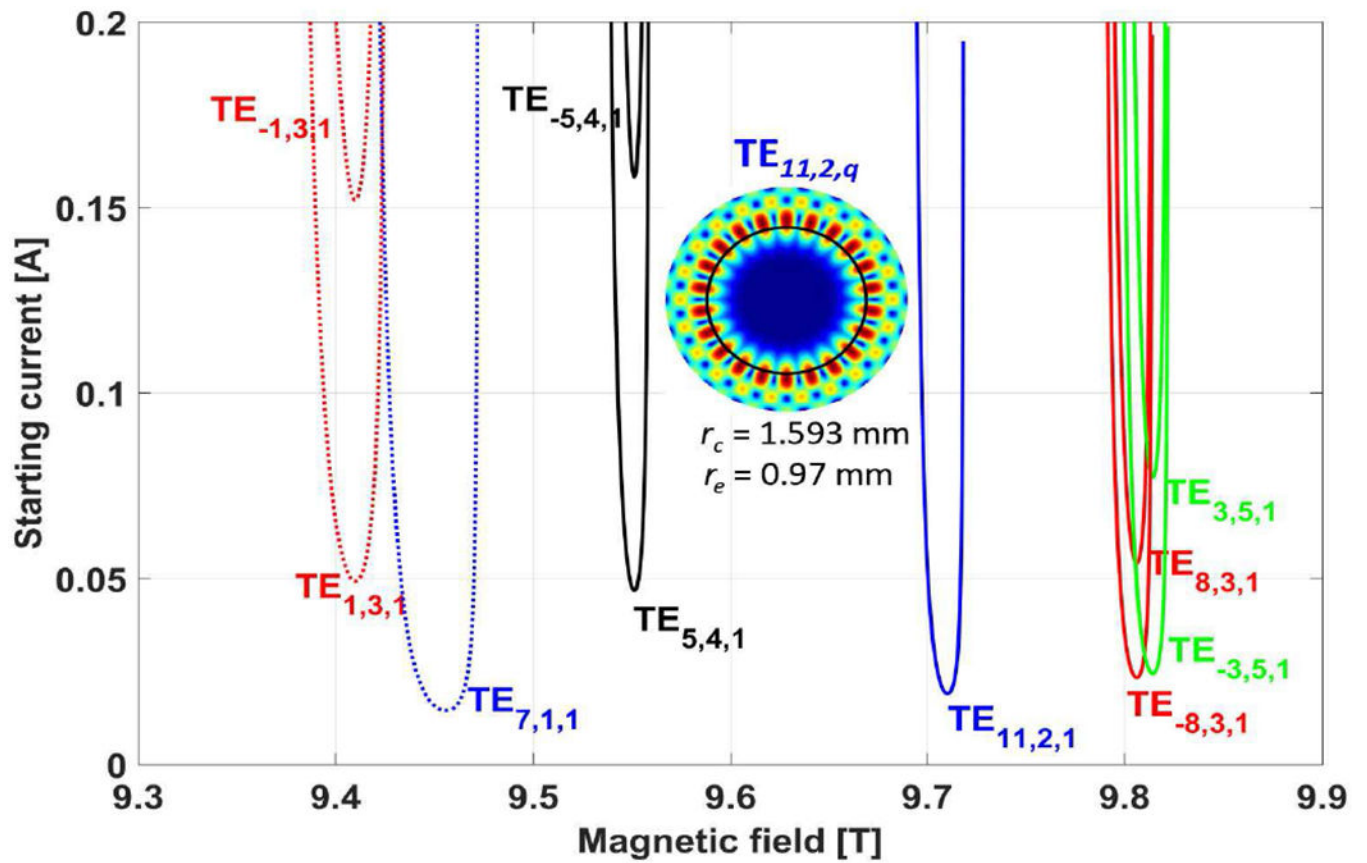
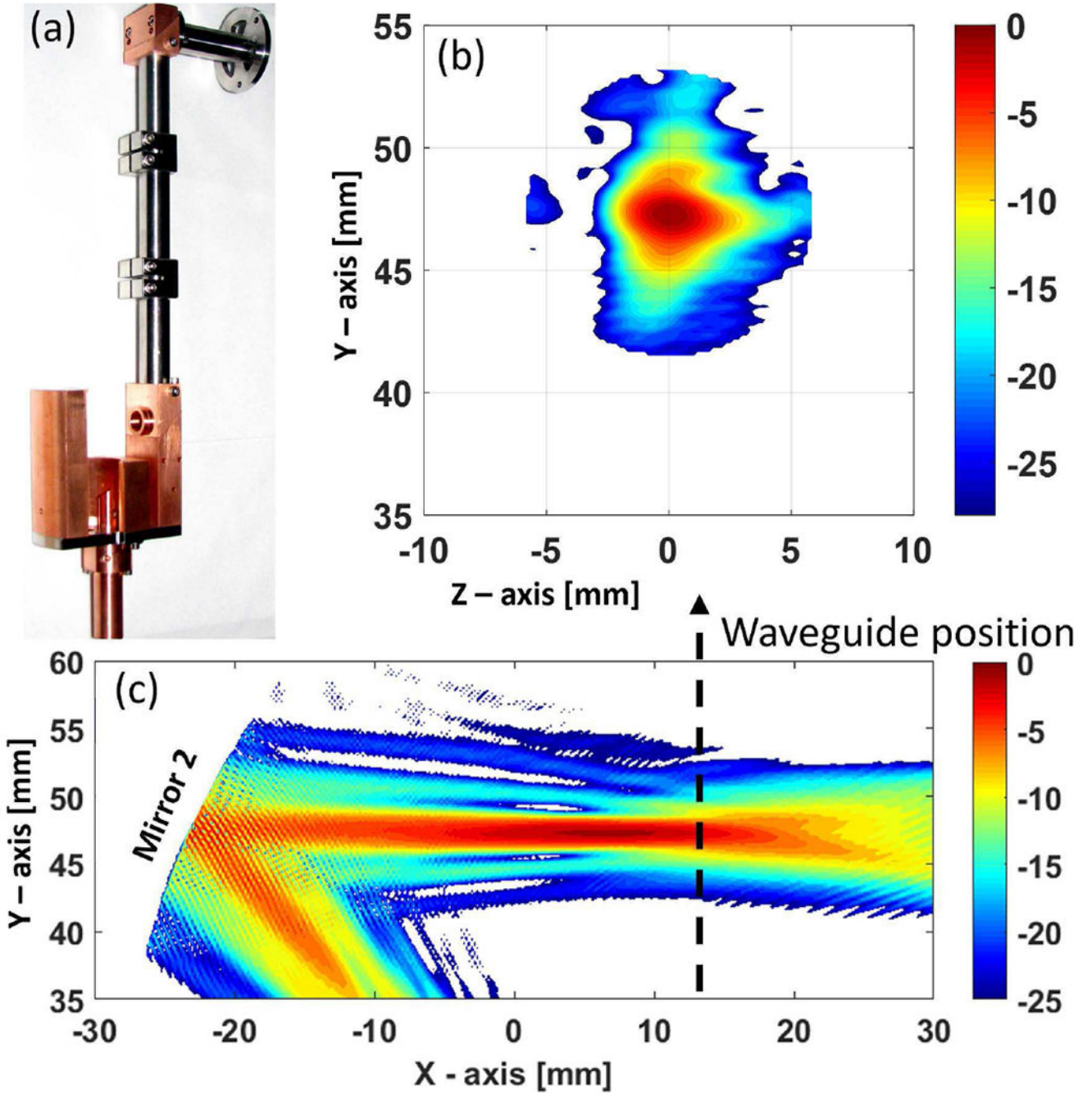


Fig. 3. The cavity profile showing 25 mm long resonator section and up tapers. The electric field profile of  $TE_{11,2,1}$  mode is also shown calculated using the self-consistent multimode code MAGY [47].



**Fig. 4.** Start oscillation current of the operating mode,  $TE_{11,2,1}$ , and the neighboring modes w.r.t. the operating magnetic field. Also shown in the center is the electric field profile of the operating mode  $TE_{11,2}$  in a resonator cavity of radius  $r_c = 1.593$  mm and the circle representing the electron beam of radius  $r_e = 0.97$  mm.

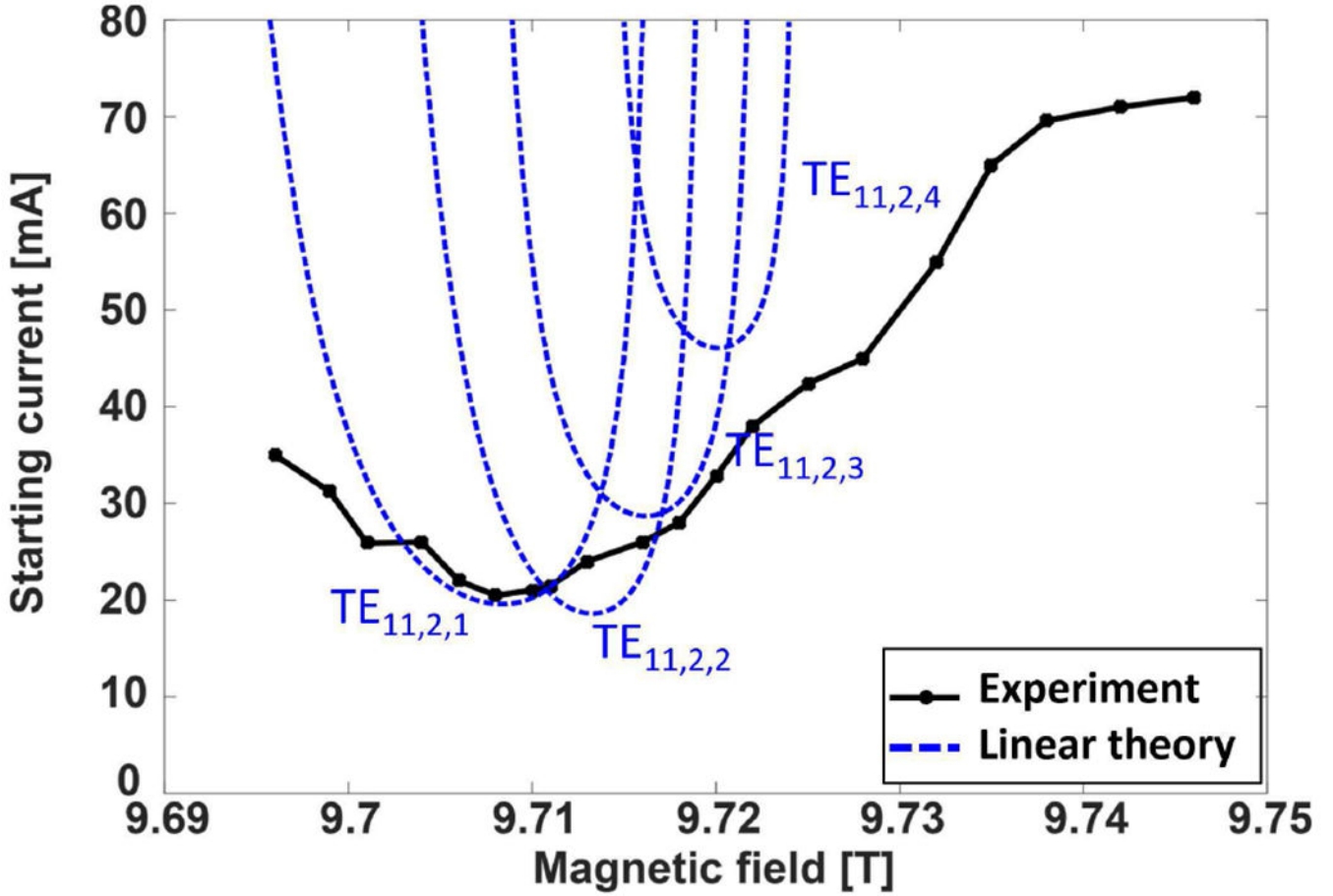




**Fig. 5.** (a) Fabricated mode converter assembly, (b) Intensity profile of the Gaussian beam calculated using Surf3D in a transverse plane before coupling into the corrugated waveguide. (c) Intensity profile for the radiated field after the 1<sup>st</sup> mirror and in a cut plane showing the position of 2<sup>nd</sup> mirror and the corrugated waveguide aperture.

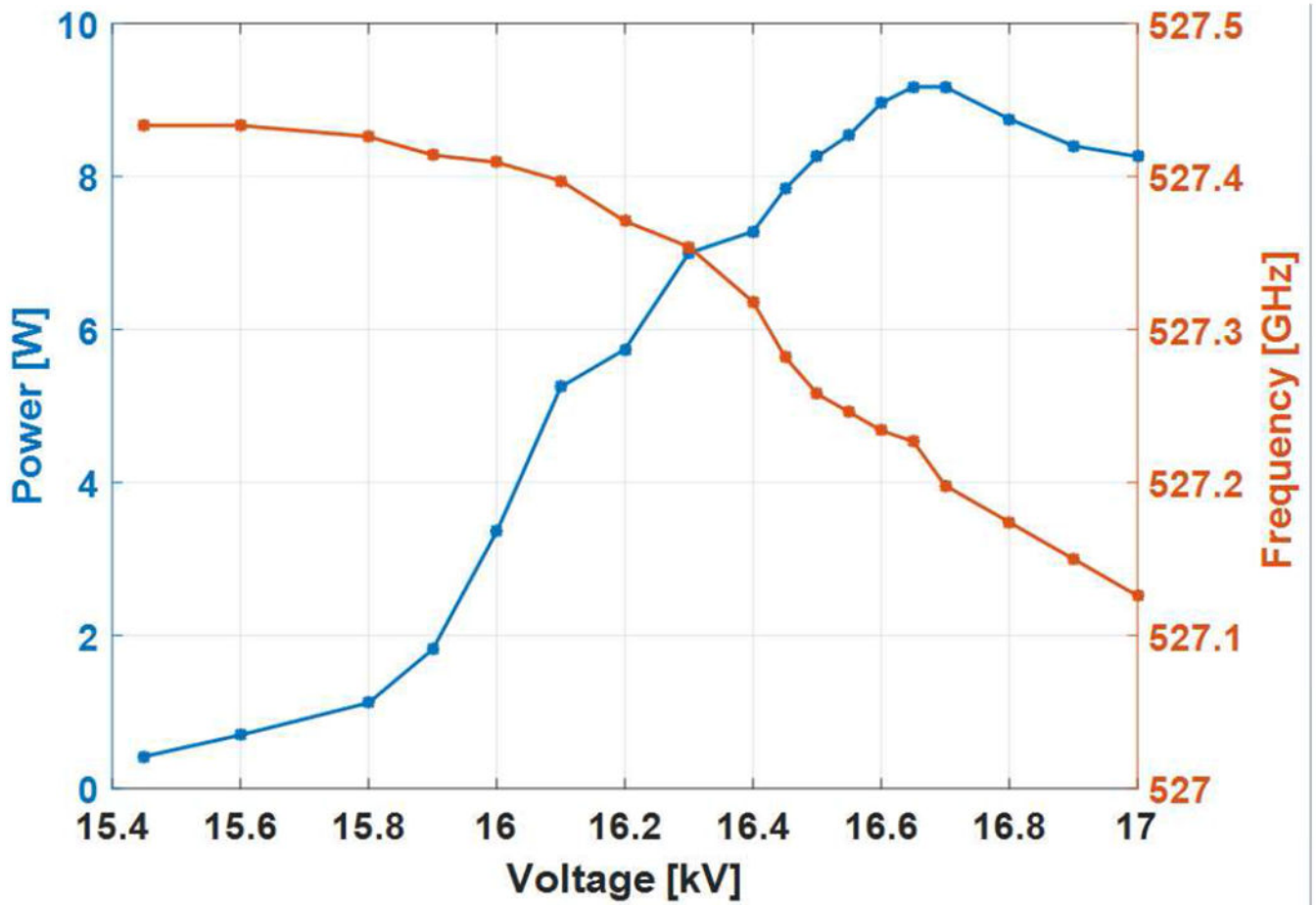


**Fig. 6.** Photograph of the 527 GHz gyrotron tube under test. A 10 T cryo-free superconducting magnet (green color) is seen in the picture along with the gyrotron tube and a small magnet coil, at the bottom of the main magnet, for fine tuning the electron beam parameters. A 16 mm i.d. corrugated waveguide transmission line made of brass is also seen in the picture.



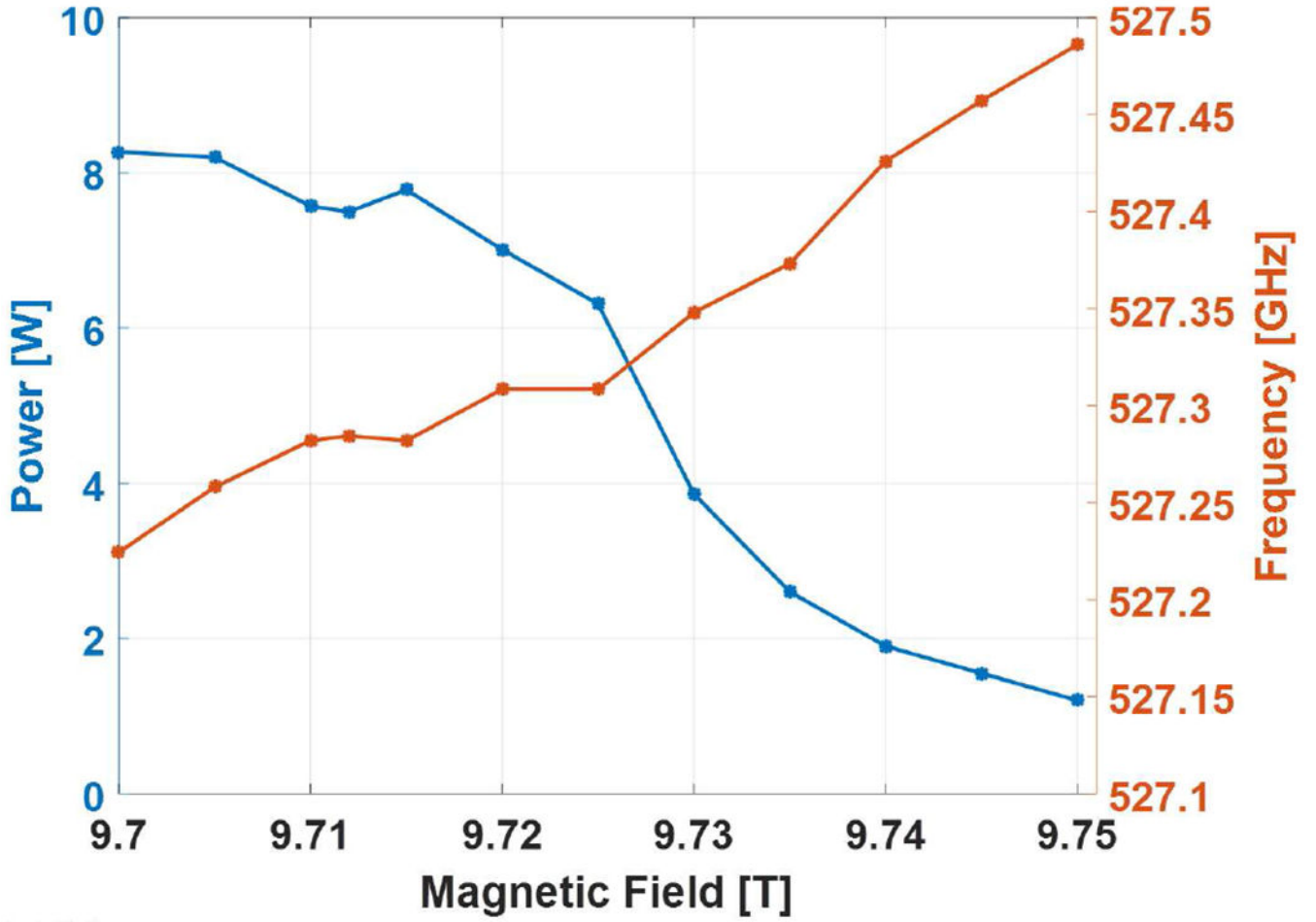
**Fig. 7.** The measured (solid black line with dots) and calculated (dashed blue lines) start oscillation current of the operating mode  $TE_{11,2,q}$ ,  $q = 1,2,3,4$  for different magnetic field values. A minimum start current of  $\sim 21$  mA is observed. The calculated start current assumed cold cavity electric field profile in the linear theory for cavity radius of  $r_c = 1.593$  and operating voltage of  $V = 16.7$  kV,  $\alpha = 1.8$  and  $v_{\perp} = 3\%$  with the gun coil magnetic field of  $\sim -37$  mT.



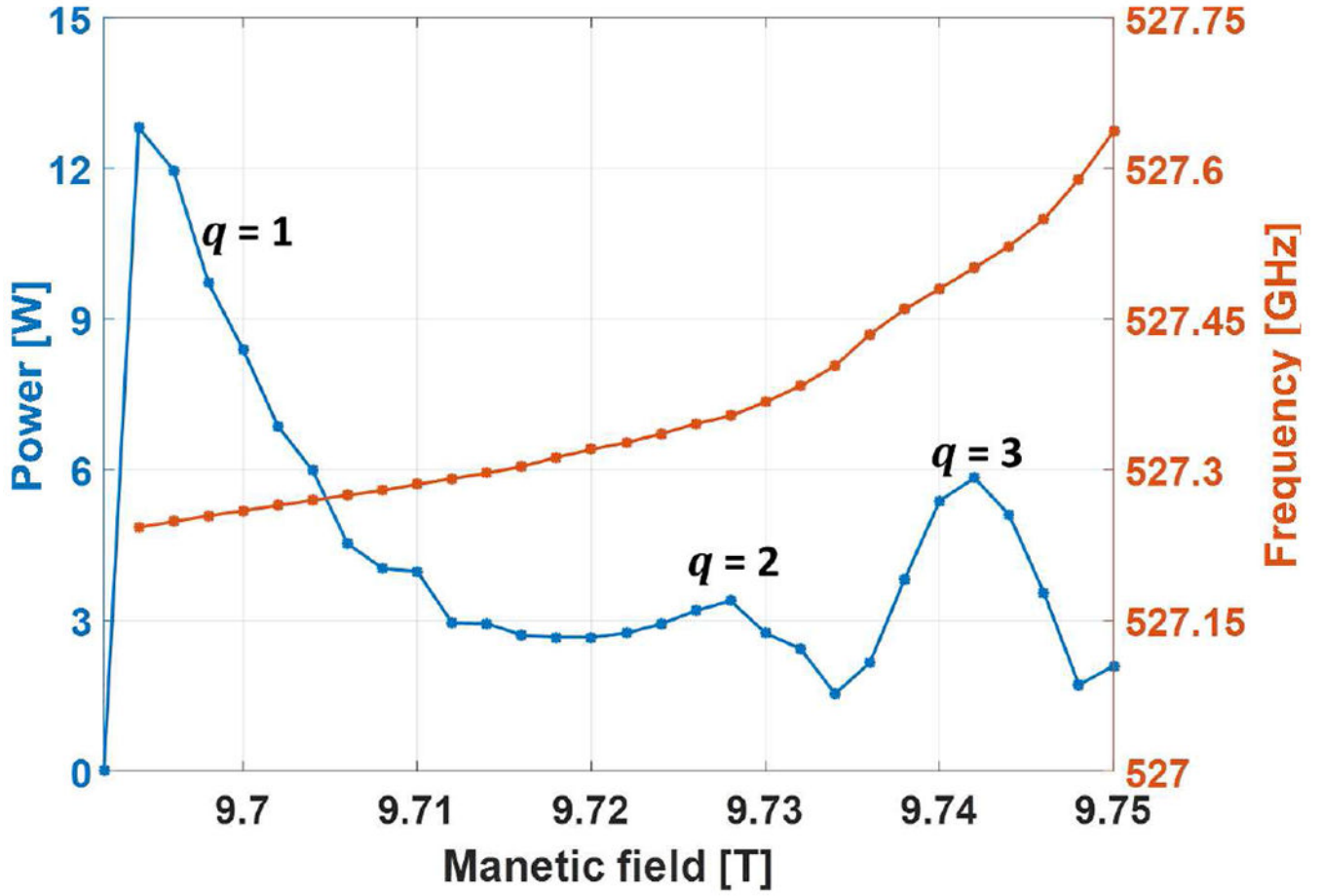


**Fig. 8.**

The measured output power and frequency with respect to the operating voltage at the magnetic field of 9.708 T and 110 mA of electron beam current. The gun coil magnet is operated up to  $\sim -35$  mT to maximize the output power.

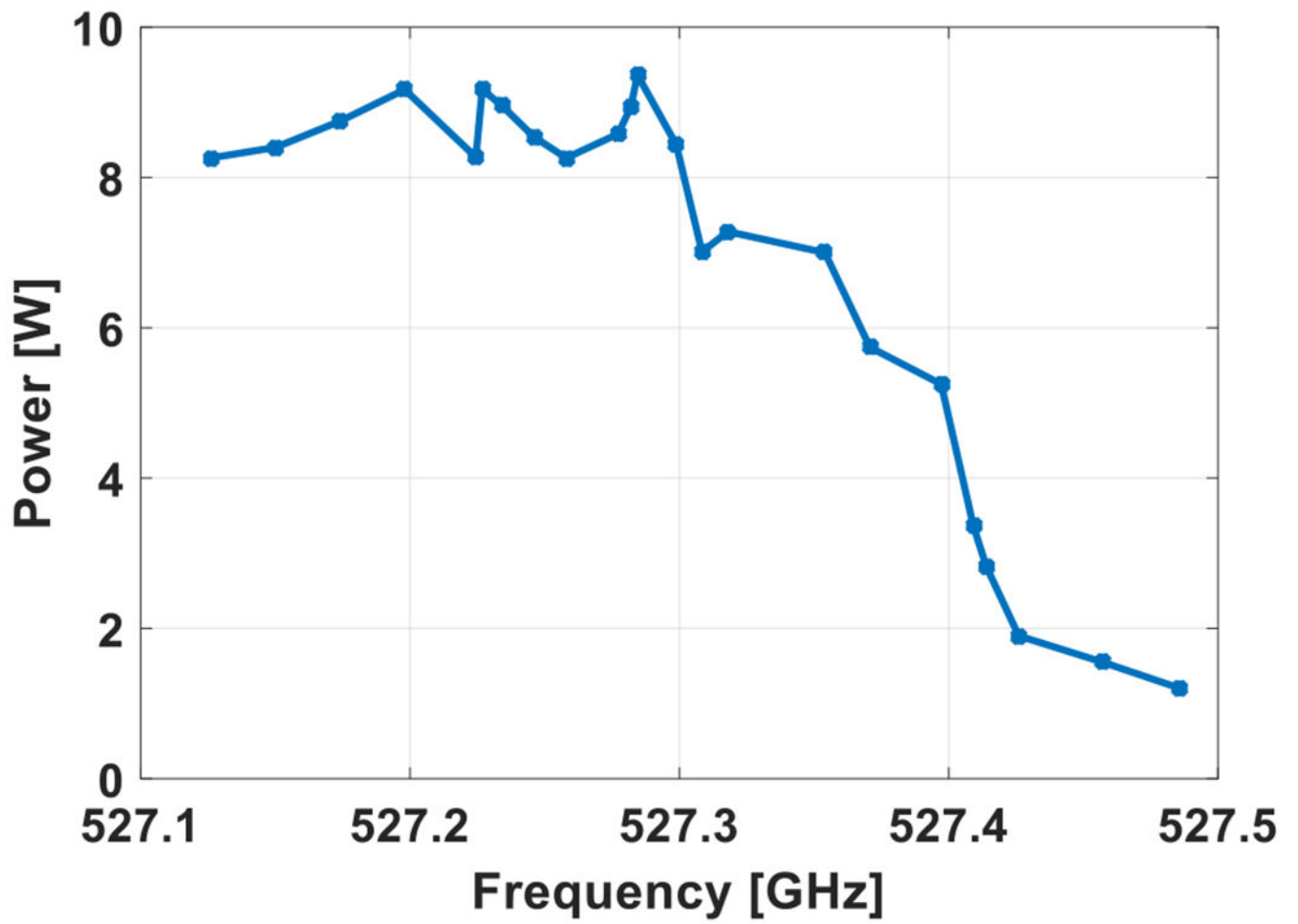


**Fig. 9.** The measured output power and frequency with respect to the operating magnetic field and fixed voltage of 16.7 kV and 90 mA of electron beam current. The gun coil magnet is operated up to  $\sim -35$  mT to maximize the output power.

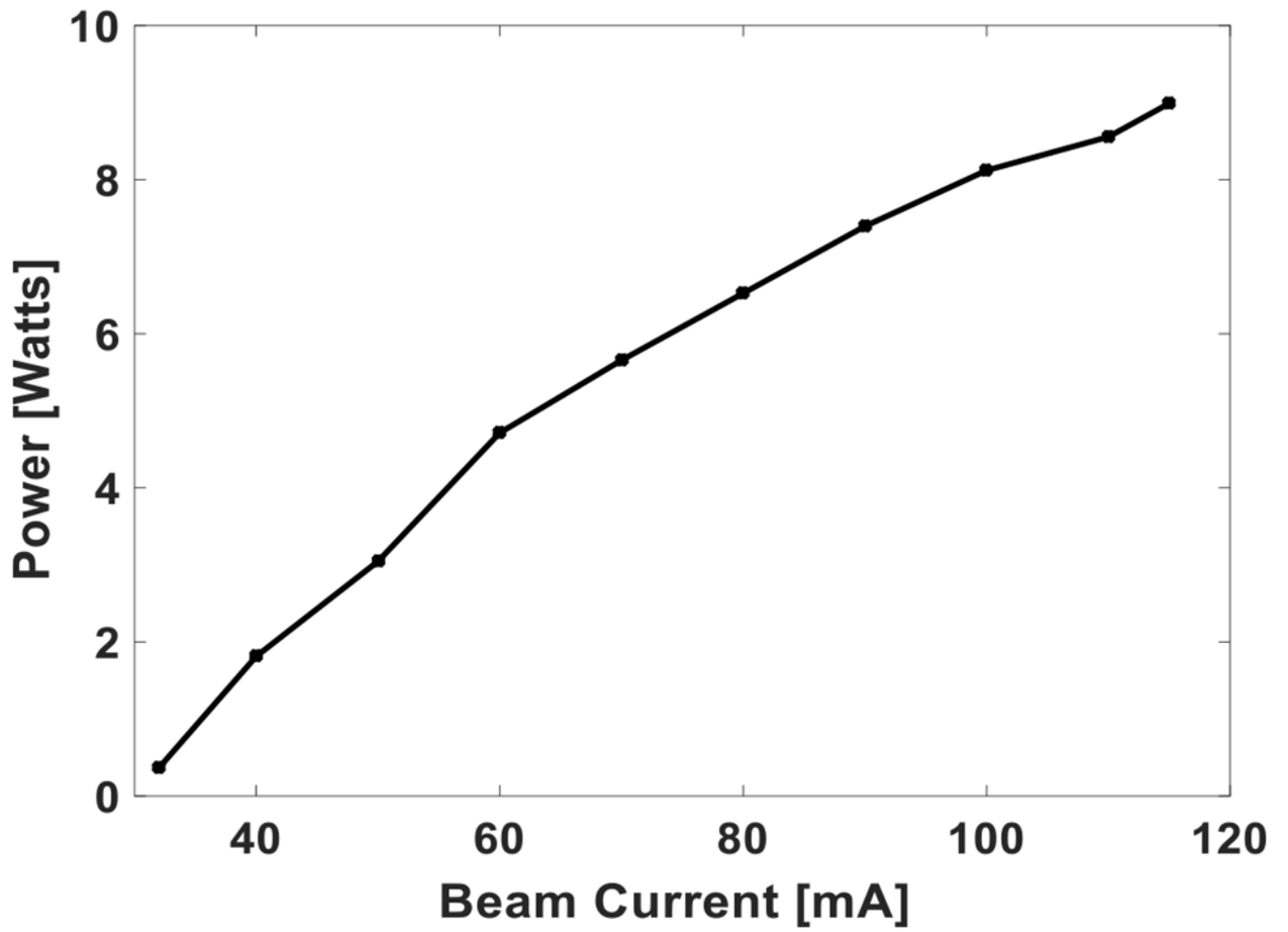


**Fig. 10.** The simulated output power and frequency with respect to the operating magnetic field at 16.7 kV voltage and 110 mA of electron beam current,  $\alpha = 1.85$  and  $v_{\perp}/v_{\parallel} = 5\%$ . The cavity radius used in the calculations is the measured radius of 1.593 mm.

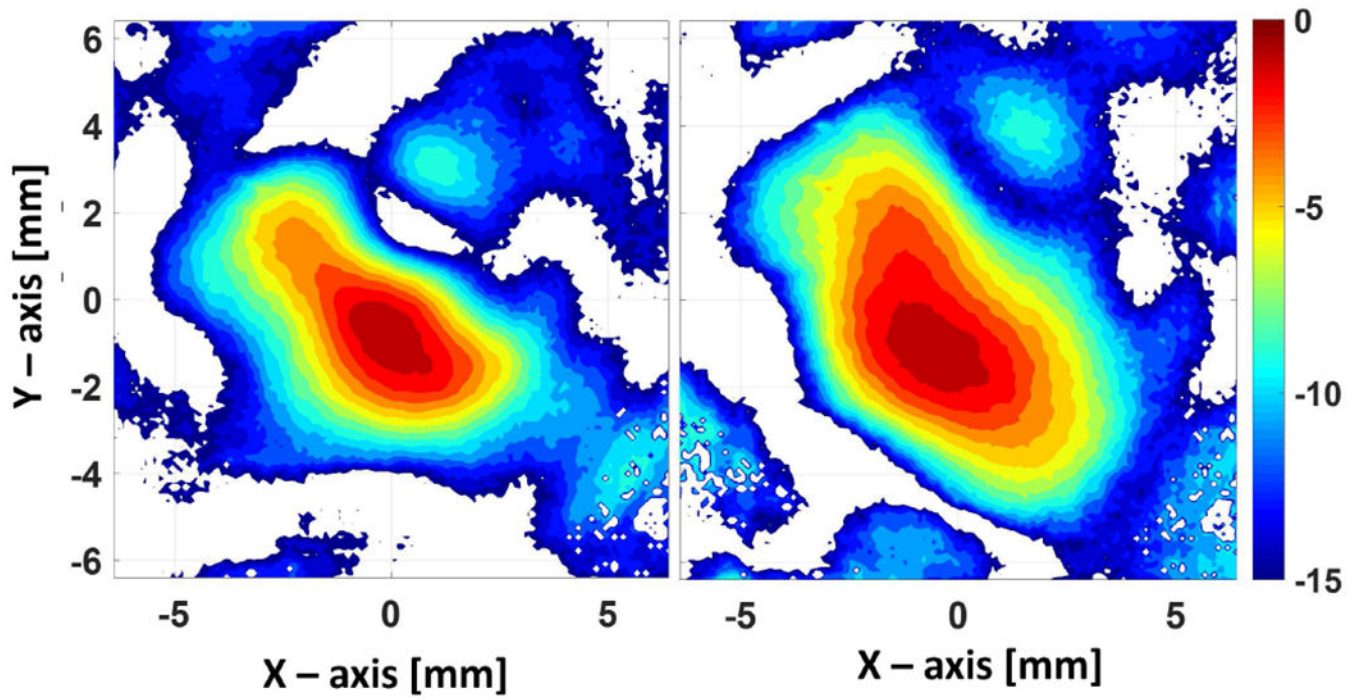




**Fig. 11.** The measured output power as a function of output frequency for varying magnetic field from 9.69 T to 9.75 T and voltage varying from 15.5 to 17 kV at a beam current of 110 mA.



**Fig. 12.**  
The measured output power with respect to the beam current at magnetic field of 9.71 T and 16.7 kV voltage.



**Fig. 13.**  
The measured intensity profile at 20 mm (left) and 38 mm (right) from the waveguide aperture after a 3 m long waveguide section using a pyroelectric camera.

**TABLE I**

## DESIGN PARAMETERS OF THE GYROTRON

Parameter	Value
Frequency	527 GHz
Cyclotron Harmonic	2 <sup>nd</sup>
Magnetic field	9.7 T
Operating Mode	TE <sub>11,2,q</sub>
Frequency Tuning	~ 1 GHz
Voltage	< 17 kV
Current	< 200 mA
Output Power	2 – 20 W
Cavity length	25 mm
Output mode	HE <sub>11</sub>

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