

Published in final edited form as:

*Appl Phys Lett.* 2019 ; 115(7): . doi:10.1063/1.5100154.

## Improved contacts to p-type MoS<sub>2</sub> transistors by charge-transfer doping and contact engineering

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### Abstract

MoS<sub>2</sub> is known to show stubborn n-type behavior due to its intrinsic band structure and Fermi level pinning. Here, we investigate the combined effects of molecular doping and contact engineering on the transport and contact properties of monolayer (ML) MoS<sub>2</sub> devices. Significant p-type (hole-transport) behavior was only observed for chemically doped MoS<sub>2</sub> devices with high work function palladium (Pd) contacts, while MoS<sub>2</sub> devices with low work function metal contacts made from titanium showed ambipolar behavior with electron transport favored even after prolonged p-doping treatment. ML MoS<sub>2</sub> transistors with Pd contacts exhibit effective hole mobilities of  $(2.3 \pm 0.7) \text{ cm}^2 \text{ V}^{-1} \text{ S}^{-1}$  and an on/off ratio exceeding  $10^6$ . We also show that p-doping can help to improve electrical contacts in p-type field-effect transistors: relatively low contact resistances of  $(482 \pm 40) \text{ k}\Omega \mu\text{m}$  and a Schottky barrier height of  $\approx 156 \text{ meV}$  were obtained for ML MoS<sub>2</sub> transistors. To demonstrate the potential application of 2D-based complementary electronic devices, a MoS<sub>2</sub> inverter based on pristine (n-type) and p-doped monolayer MoS<sub>2</sub> was fabricated. This work presents a simple and effective route for contact engineering, which enables the exploration and development of high-efficiency 2D-based semiconductor devices.

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While transition metal dichalcogenides (TMDs) have been studied for decades, recent advances in nanoscale materials characterization and device fabrication have opened new opportunities for 2D materials in nanoelectronics and optoelectronics. To enable complementary device operation, both n- and p-type field-effect transistors (n-FETs and p-FETs) are needed. Using different TMDs for n-FETs and p-FETs is possible but increases the complexity of the device fabrication. In TMD-based devices, contacts and interfaces of 2D materials can significantly impact device performance and are viewed as limiting factors in their performance.<sup>1-3</sup> Since the pristine surface of 2D materials does not tend to form covalent bonds, a van der Waals (vdW) gap is formed at the interfaces between the metal and 2D materials. This vdW gap acts as a tunnel barrier for carriers. Researchers have pointed out that the three criteria for efficient charge injection are a strong orbital overlap between the contacts and channel materials, a low Schottky barrier height (SBH), and a narrow tunnel

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See the supplementary material for additional results on the optical characterization of MoS<sub>2</sub>, Schottky barrier extraction, optical image of the MoS<sub>2</sub> inverter, and doping results on SiO<sub>2</sub> covered transistors.

barrier.<sup>4</sup> MoS<sub>2</sub> typically exhibits strong n-type semiconductor properties because of Fermi level pinning close to the conduction band edge, regardless of the contact metal used.<sup>5</sup> Several approaches have been tested for achieving p-type transport properties of MoS<sub>2</sub>, including contact work function engineering,<sup>6</sup> implantation of foreign atoms,<sup>7–9</sup> plasma treatment,<sup>10</sup> and ionic gating.<sup>11</sup> However, most of these chemical doping approaches suffer from thermal and chemical stability problems and introduce mobile ions into the system. These ions can easily diffuse and disrupt the device performance, resulting in low hole mobility and/or high contact resistance.

Recently, we have demonstrated effective doping of monolayer TMDs by using molecular reductants and oxidants.<sup>12</sup> The molecular reductants and oxidants are introduced onto the surface of 2D materials from solution. P-type transport behavior of MoS<sub>2</sub> was demonstrated by using tris(4-bromophenyl)ammoniumyl hexachloroantimonate (“Magic Blue,” MB, [N(C<sub>6</sub>H<sub>4</sub>-*p*-Br)<sub>3</sub>]SbCl<sub>6</sub>), as the molecular oxidant (p-dopant).<sup>13,14</sup> In this study, we investigated the effects of molecular doping and contact engineering on the charge transport properties and 2D/metal interfaces of MoS<sub>2</sub> devices. We have shown that MoS<sub>2</sub> p-FETs can only be achieved through combining molecular doping with high-work function palladium (Pd) ( $\Phi = 5.6$  eV) contacts. Only ambipolar behavior can be achieved when using the low work function metal titanium (Ti) ( $\Phi = 4.3$  eV).<sup>15,16</sup> The p-doped monolayer MoS<sub>2</sub> transistor with a Pd contact exhibits effective hole mobilities of  $(2.3 \pm 0.7)$  cm<sup>2</sup> V<sup>-1</sup> S<sup>-1</sup>, an on/off ratio exceeding 10<sup>6</sup>, and a relatively low contact resistance of  $(482 \pm 40)$  k $\Omega$   $\mu$ m. SBHs were extracted by measuring the activation energy in the thermionic emission regime. A relatively low hole injection barrier of  $\approx 156$  meV was obtained for p-doped MoS<sub>2</sub> with Pd contacts. A MoS<sub>2</sub> inverter based on pristine (n-type) and p-doped monolayer MoS<sub>2</sub> was fabricated, demonstrating the potential of charge-transfer doping in the application of 2D-based complementary electronic devices.

Monolayer (ML) MoS<sub>2</sub> flakes were exfoliated onto 300 nm thermally grown SiO<sub>2</sub> on a low-resistivity silicon wafer according to the literature-reported gold-mediated exfoliation method from the bulk crystal.<sup>17</sup> The number of layers was confirmed by Raman and photoluminescence (PL) spectroscopy (shown in supplementary material Fig. S1). The source/drain contacts were defined by e-beam lithography, followed by electron beam metal deposition and a lift-off process. The deposited source/drain contacts were 30 nm Pd or Ti with 70 nm Au on top to protect from oxidation. As shown in Fig. 1(c), the channel length (L) varies between 1.5  $\mu$ m and 9  $\mu$ m. The active area of the channel was patterned using a reactive ion etcher (RIE) by using a gas mixture of SF<sub>6</sub> and O<sub>2</sub>. A bottom-gate top-contact configuration was used with a 300 nm SiO<sub>2</sub> dielectric. Figure 1(a) shows the p-dopant, MB, used in this study. MB is a strong p-dopant, and the reduction potential of the functional group [N(C<sub>6</sub>H<sub>4</sub>-*p*-Br)<sub>3</sub>]<sup>+</sup> is +0.70 V vs ferrocene.<sup>18</sup> During the doping process, the tris(4-bromophenyl)-ammoniumyl radical cation accepts an electron from MoS<sub>2</sub> and forms the neutral tris(4-bromophenyl)amine which largely remains in solution and/or is washed away in the following rinsing step. The hexachloroantimonate counterions (SbCl<sub>6</sub><sup>-</sup>) are left on the surface to balance the holes introduced to the film. MB was dissolved in anhydrous dichloromethane (DCM) to make a 5 mmol L<sup>-1</sup> solution. The doping treatments were carried out inside a nitrogen glovebox. The devices were immersed in the dopant solutions for various times and were then rinsed with fresh solvent three to five times to remove

physisorbed dopants. The electrical characterization of the fabricated devices was collected inside an inert-atmosphere glovebox using a commercial parameter analyzer at room temperature. Unless otherwise noted, electrical parameters reported in the text are the mean and standard deviation.

Figures 1(d) and 1(e) show the experimental transfer characteristics of typical pristine bottom-gated top-contact ML MoS<sub>2</sub> with Ti and Pd metal contacts measured at 300 K. The threshold voltage ( $V_T$ ) is extracted by using the linear extrapolation method with the drain current measured as a function of gate voltage. The  $V_T$  values for Ti and Pd contacts was found to be  $(-54 \pm 3)$  V and  $(-25 \pm 4)$  V, respectively, averaged over tens of devices. The devices exhibit n-FET characteristics regardless of the work function of the metals. The on-current of the high work function metal Pd is smaller than that of the low work function Ti, suggesting that Ti is better aligned with the conduction band of MoS<sub>2</sub> and has better electron injection than Pd. These trends are consistent with the previously reported studies of the metal-to-MoS<sub>2</sub> interface.<sup>19</sup>

The MoS<sub>2</sub> FET devices were first immersed in the 5 mmol L<sup>-1</sup> dopant solutions for 10 min. As shown in Fig. 2(a), for both Ti- and Pd-contact FETs, the p-doping treatments shifted the transfer curve to a more positive gate bias after 10 min of treatment [Fig. 2(a), solid square]. An increase in the hole current at negative gate voltage was observed for both contacts to MoS<sub>2</sub>. Devices with Ti contacts exhibit ambipolar properties with electron transport favored, while Pd contacted MoS<sub>2</sub> shows preferred hole transport. The same samples were retreated with an additional 1 h molecular exposure causing  $V_T$  to shift to more positive back-gate voltage ( $V_{GS}$ ), consistent with stronger p-doping. However, for the Ti-contact FETs, a decrease is observed in both the electron and hole current and the MoS<sub>2</sub> FETs still appear to be ambipolar. For the Pd-contact FETs, large  $V_T$  shifts (from 40 V to 75 V) and saturation of the ON current at  $10^{-7}$  A/ $\mu$ m were observed. Significant  $V_T$  shifts ( $V_T$ ) up to 64 V and 100 V were achieved for Ti- and Pd-contact devices, respectively. The change in hole carrier density ( $p$ ) can be estimated from the  $p = C_{bg} V_T/q$  expression, where back-gate capacitance  $C_{bg}$  is estimated from a parallel-plate capacitor model,  $C_{bg} = \epsilon_0 \epsilon_r/d$ , with  $\epsilon_r = 3.9$  for SiO<sub>2</sub> and the thickness of the gate dielectric  $d = 300$  nm, and the change of hole sheet density ( $p$ ) is estimated to be  $4.6 \times 10^{12}$  cm<sup>-2</sup> and  $7.2 \times 10^{12}$  cm<sup>-2</sup> for Ti- and Pd-contact FETs. Since the MoS<sub>2</sub> devices were exposed to the same molecular doping conditions, it is reasonable to assume that the same number of charge carriers was introduced for both Ti- and Pd-contact devices. The smaller change in hole sheet density for the Ti-contacted devices may be attributed to the misalignment of the metal work function and MoS<sub>2</sub> valence band, where the doping-induced charge carriers have little impact on the current. The output curves of the pristine n-FET with Pd contacts and doped p-FET are shown in Fig. 2(b). The device performance is summarized in Table I, and the results are the average of 10–15 devices. The pristine Ti-contacted MoS<sub>2</sub> shows a higher electron mobility ( $12.2 \pm 1.7$  cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup>) than Pd-contact FETs ( $8.3 \pm 0.9$ ) cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup>, while 10 min doped Pd-contacted MoS<sub>2</sub> devices exhibit a hole mobility of  $(2.3 \pm 0.7)$  cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup>. This trend is consistent with a smaller electron injection barrier between Ti-MoS<sub>2</sub> favoring electron-transport for both the pristine and doped devices, while in the Pd case, doping has a more dramatic impact, suggesting that the barrier for hole injection is substantially decreased. According to theoretical calculations and experimental results, there are several causes for the Fermi level

pinning at the metal contact of MoS<sub>2</sub>: the van der Waals gap formed at the interface which acts as a tunnel barrier,<sup>19</sup> the lack of orbital overlaps between metal and TMDs,<sup>20,21</sup> and possible sulfur vacancy defect sites which can act as donors.<sup>22</sup> The fact that p-FETs were not observed in Ti-contact devices may be attributed to the lack of orbital overlaps between Ti and p-doped MoS<sub>2</sub>. For the Pd-contact FETs, the longer treatment time (1 h) decreases the mobility to  $(0.82 \pm 0.12 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1})$ . This decrease in the hole mobility is attributed to the increased concentration of ionized scattering centers formed by the doping product,  $\text{SbCl}_6^-$  anions. This observation is consistent with previous reports on molecular doping of organic semiconductors<sup>23,24</sup> and doping of inorganic semiconductors.<sup>25</sup>

Figure 3 shows qualitative energy band diagrams for MoS<sub>2</sub> FETs with Ti and Pd contacts at the metal-MoS<sub>2</sub> interface. Both experimental and theoretical results show that the SBH is correlated with the work function of the metal contacts. Lower work function metals (Ti in this case) show improved contacts with pristine MoS<sub>2</sub> flakes, as indicated by higher electron injection efficiency and lower contact resistance. Before the molecular doping, the Fermi level of MoS<sub>2</sub> is pinned close to the conduction band.<sup>22</sup> The interface is strongly impacted by this pinning which occurs close to the conduction band of MoS<sub>2</sub>, making electron transport dominating. Because of the complicated metal-MoS<sub>2</sub> interface, this barrier height cannot be modified effectively by varying the work function of the metal contacts alone.<sup>26</sup> The molecular p-dopant molecule interacts with the MoS<sub>2</sub> surface via a charge-transfer reaction where electrons are transferred from MoS<sub>2</sub> to the p-dopants, making the MoS<sub>2</sub> Fermi level shift closer to the valence band edge.<sup>27</sup> Thus, doping induces a much higher hole carrier concentration. For the doped low work function Ti-contacted devices, ambipolar behavior was observed with electron transport favored, demonstrating that both high hole concentrations and a low injection barrier are needed for efficient hole transport.

The contact resistance,  $R_c$ , before and after doping with different contacts was extracted by using the transfer length method (TLM). The channel width ( $W$ ) of the lithographically defined MoS<sub>2</sub> strip (device channel) is 25  $\mu\text{m}$ . The measured resistance is the aggregate value calculated from the linear I-V characteristics for large voltage sweeps. The  $V_{GS}$  is 0 V. Figure 4(a) shows the measured total resistance for different channel lengths. The contact and sheet resistances were extracted from the linear fit of the total resistance, and the contact resistance was normalized for the device width. Before doping, the contact resistances of MoS<sub>2</sub> n-FETs were calculated to be  $(805 \pm 65) \text{ k}\Omega \mu\text{m}$  with Ti contacts and  $(1430 \pm 107) \text{ k}\Omega \mu\text{m}$  with Pd contacts. These numbers are comparable to previously reported contact resistance for ML MoS<sub>2</sub> FETs.<sup>28</sup> The Pd contact resistance was reduced to  $(482 \pm 40) \text{ k}\Omega \mu\text{m}$  after p-doping. Two notable observations are (1) the total resistance increases for the Ti-contact device and decreases for the Pd-contact device after 1 h treatment; (2) contact resistance increases by more than two orders of magnitude for Ti-contact FETs, indicating a larger SBH, and decreases by a factor of three for Pd-contact FETs, indicating improved carrier injection. p-Doping shifts the Fermi level closer to the valence band. Thus, it is expected that higher  $R_c$  and higher charge injection barrier are observed with Ti contacts since the work function is closer to the conduction band. In the case of Pd contacts, the work function is closer to the midgap of MoS<sub>2</sub>, making these devices more sensitive to doping, showing lower  $R_c$  and more efficient hole injection. Even though only the channel area was exposed to the dopant treatment, the change in the contact property may be due to the doped

one-dimensional edge between the contacts and ML MoS<sub>2</sub>. It was reported in the literature that graphene with a 1D edge contact can have remarkably low contact resistance.<sup>1</sup> It is also possible that the dopants and the injected charge carriers diffuse beneath the metals during immersion in the solution and alter the interfaces. While data such as shown in Fig. S5 suggest such a mechanism, further experimental and theoretical studies are needed to fully understand this effect.

To further evaluate the impacts of doping and different metals on contacts, we also extracted the SBH by using the following equation:

$$I_{DS} = AT^2 \exp \left[ \frac{q}{k_B T} \left( \Phi_B - \frac{V_{DS}}{n} \right) \right], \quad (1)$$

where  $A$  is Richardson's constant,  $k_B$  is the Boltzmann constant,  $q$  is the electronic charge,  $T$  is the temperature,  $\Phi_B$  is the effective barrier height, and  $n$  is the ideality factor.<sup>5</sup> Figures 4(b) and 4(c) show an example of flatband Schottky barrier extraction for the 1 h p-doped MoS<sub>2</sub> transistor with Pd contacts. Figure 4(b) shows transfer characteristics at different operating temperatures. These data were used to create the Arrhenius plot in Fig. 4(c). As explained in detail in the literature, thermally assisted tunneling determines the inverse subthreshold slope in Schottky barrier devices.<sup>29,30</sup> Figure 4(d) shows that, at sufficiently positive gate voltages, the effective barrier height,  $\Phi_B$ , responds linearly to the gate voltage,  $V_{GS}$ , when only thermionic current flows through the device. The flatband voltage,  $V_{FB}$ , is reached when tunneling starts to play a role and  $\Phi_B$  deviates from the linear trend mentioned previously. A p-type SBH ( $\Phi_{SB}$ ) of  $\approx 156$  meV is determined from the Arrhenius plot. The SBH extractions for the pristine MoS<sub>2</sub> transistor with Ti and Pd contacts are shown in supplementary material Fig. S2. In the n-type transistor,  $\Phi_B$  responds linearly to  $V_{GS}$  at sufficiently negative gate voltages (thermionic regime). The Schottky barrier heights were determined to be  $\approx 100$  meV and  $\approx 230$  meV for Ti- and Pd-contact pristine MoS<sub>2</sub> transistors, respectively. A SBH for the p-doped MoS<sub>2</sub> transistor with Ti contacts could not be extracted because the tunneling regime was not reached (Fig. S3). A relatively low hole injection barrier can only be obtained for p-doped MoS<sub>2</sub> with Pd contacts. The tunneling barrier for holes is too high in Ti-contact devices even after p-doping. These results are consistent with the contact resistance results.

To demonstrate the use of chemical doping for device applications, a complementary inverter was fabricated by wire bonding two Pd-contact MoS<sub>2</sub> FETs on the same chip. One of the transistors was immersed in the dopant solution. Figures 5(a) and 5(b) show the device schematic and the electronic circuit. A photograph of the completed device is shown in Fig. S4. The back gate was applied as the input voltage ( $V_{IN}$ ), and a supply voltage ( $V_{DD}$ ) was applied to the p-FET. Figure 5(c) shows the voltage transfer characteristics of the inverter at various  $V_{DD}$  values. The device operates at a negative input voltage and the switching threshold voltage of the inverter located in the negative region. The slope of the voltage transfer curves in Fig. 5(c) represents the gain of the inverter, which is  $\approx 2.73$ . Integrating the transistor with a thin high- $k$  dielectric in a top-gate configuration should enhance the voltage gain greatly.

In 2D-based electronics, the contacts and interfacial properties play important roles in determining device performance. In this work, high-performance p-type field-effect (p-FET) transistors were demonstrated by using chemically doped monolayer MoS<sub>2</sub> and high work function metal contacts (Pd). The impact of molecular doping and the metal work function on the charge transport and contact properties of MoS<sub>2</sub> devices was presented. Molecular doped p-FETs show an effective hole mobility of 2.3 cm<sup>2</sup> V<sup>-1</sup> S<sup>-1</sup>, an on/off ratio exceeding 10<sup>6</sup>, and an improved contact resistance of ≈482 kΩ μm. However, when a low work function metal (Ti) is used to contact MoS<sub>2</sub>, the device only exhibits ambipolar transport with electron transport favored. It was shown that the SBH can be effectively modulated through doping of the channel materials and the contact engineering by changing the contact metal. A relatively low hole injection barrier of ≈ 156 meV was obtained for p-doped MoS<sub>2</sub> with Pd contacts. The utility of MoS<sub>2</sub> p-FETs and n-FETs was demonstrated with a complementary inverter. In summary, controlled doping is an important tool for tuning the electronic properties of 2D devices. Both controllable doping and careful contact engineering are important to realize efficient charge injection/extraction and for improving overall device performance. Dopants and contact engineering of TMDs and other emerging materials are powerful tools to fuel the further development of nanoelectronics and optoelectronics.

## Supplementary Material

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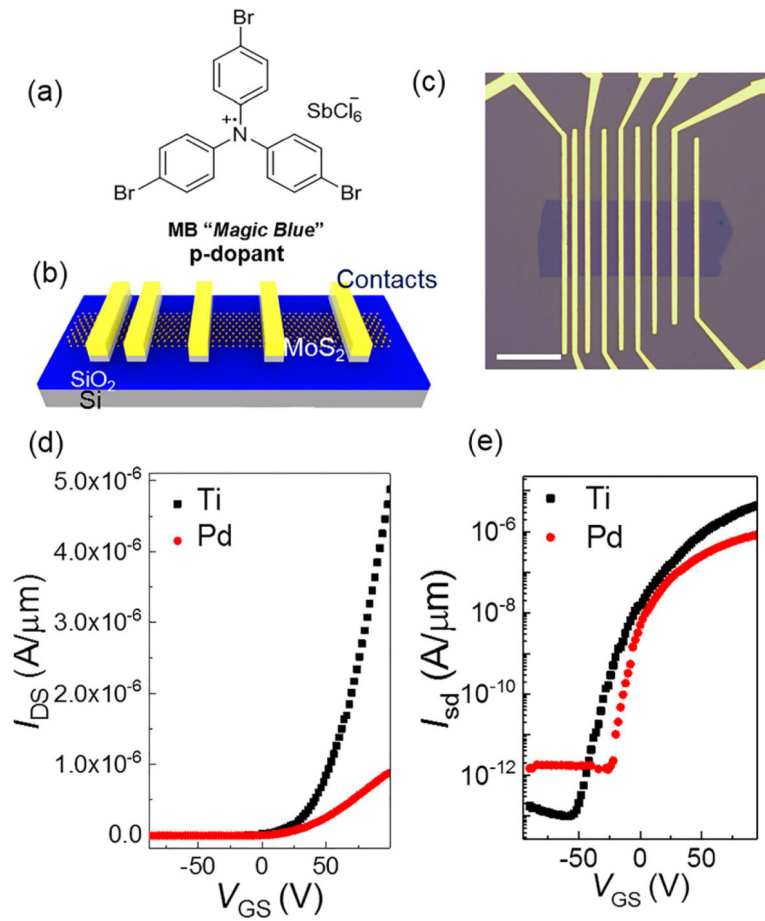
## Acknowledgments

S.Z. acknowledges support from the National Institute of Standards and Technology (NIST) Financial Assistance Award with Federal Award ID No. 70NANB16H228. The authors are thankful to Dr. David J. Gundlach and Dr. Emily G. Bittle for their assistance with the electrical measurements.

## REFERENCES

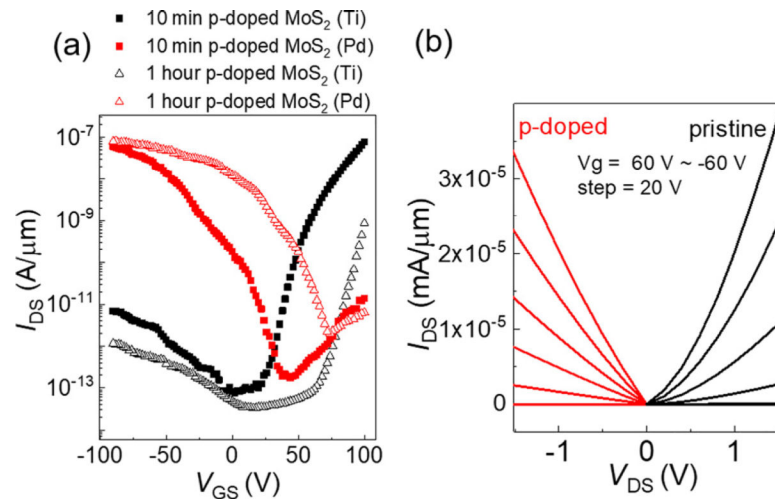
1. Wang L, Meric I, Huang P, Gao Q, Gao Y, Tran H, Taniguchi T, Watanabe K, Campos L, and Muller D, *Science* 342(6158), 614–617 (2013). [PubMed: 24179223]
2. Schulman DS, Arnold AJ, and Das S, *Chem. Soc. Rev* 47(9), 3037–3058 (2018). [PubMed: 29498729]
3. Guros NB, Le ST, Zhang S, Sperling BA, Klauda JB, Richter CA, and Balijepalli A, *ACS Appl. Mater. Interfaces* 11(18), 16683–16692 (2019). [PubMed: 30990006]
4. Kang J, Liu W, Sarkar D, Jena D, and Banerjee K, *Phys. Rev X* 4, 031005 (2014).
5. Das S, Chen H-Y, Penumatcha AV, and Appenzeller J, *Nano Lett.* 13(1), 100–105 (2013). [PubMed: 23240655]
6. Kim G-S, Kim S-H, Park J, Han KH, Kim J, and Yu H-Y, *ACS Nano* 12(6), 6292–6300 (2018). [PubMed: 29851473]
7. Laskar MR, Nath DN, Ma L, Lee EW, Lee CH, Kent T, Yang Z, Mishra R, Roldan MA, and Idrobo J-C, *Appl. Phys. Lett.* 104(9), 092104 (2014).
8. Liu X, Qu D, Ryu J, Ahmed F, Yang Z, Lee D, and Yoo WJ, *Adv. Mater.* 28(12), 2345–2351 (2016). [PubMed: 26808483]
9. Nipane A, Karmakar D, Kaushik N, Karande S, and Lodha S, *ACS Nano* 10(2), 2128–2137 (2016). [PubMed: 26789206]
10. Neal AT, Pachter R, and Mou S, *Appl. Phys. Lett.* 110(19), 193103 (2017).

11. Zhang YJ, Ye JT, Yomogida Y, Takenobu T, and Iwasa Y, *Nano Lett.* 13(7), 3023–3028 (2013). [PubMed: 23795701]
12. Zhang S, Hill HM, Moudgil K, Richter CA, Walker ARH, Barlow S, Marder SR, Hacker CA, and Pookpanratana SJ, *Adv. Mater.* 30(36), 1802991 (2018).
13. Gusev OV, Denisovich LI, Peterleitner MG, Rubezhov AZ, Ustynyuk NA, and Maitlis PM, *J. Organomet. Chem.* 452(1–2), 219–222 (1993).
14. Bell F, Ledwith A, and Sherrington D, *J. Chem. Soc. C* 1969, 2719–2720.
15. Hajzus JR, Biacchi AJ, Le ST, Richter CA, Walker ARH, and Porter LM, *Nanoscale* 10(1), 319–327 (2018).
16. Hölzl J and Schulze F, *Springer Tracts in Modern Physics* (Springer-Verlag, 1979).
17. Desai SB, Madhvapathy SR, Amani M, Kiriya D, Hettick M, Tosun M, Zhou Y, Dubey M, Ager JW, Chrzan D, and Javey A, *Adv. Mater* 28(21), 4053–4058 (2016). [PubMed: 27007751]
18. Connelly NG and Geiger WE, *Chem. Rev.* 96(2), 877–910 (1996). [PubMed: 11848774]
19. Kim C, Moon I, Lee D, Choi MS, Ahmed F, Nam S, Cho Y, Shin H-J, Park S, and Yoo WJ, *ACS Nano* 11(2), 1588–1596 (2017). [PubMed: 28088846]
20. Hu Z, Wu Z, Han C, He J, Ni Z, and Chen W, *Chem. Soc. Rev.* 47(9), 3100 (2018). [PubMed: 29509206]
21. Gong C, Colombo L, Wallace RM, and Cho K, *Nano Lett.* 14(4), 1714–1720 (2014). [PubMed: 24660782]
22. Liu D, Guo Y, Fang L, and Robertson J, *Appl. Phys. Lett.* 103(18), 183113 (2013).
23. Kröger M, Hamwi S, Meyer J, Riedl T, Kowalsky W, and Kahn A, *Org. Electron.* 10(5), 932–938 (2009).
24. Olthof S, Mehraeen S, Mohapatra SK, Barlow S, Coropceanu V, Brédas J-L, Marder SR, and Kahn A, *Phys. Rev. Lett.* 109(17), 176601 (2012). [PubMed: 23215211]
25. Van Zeghbroeck B, *Principles of Electronic Devices* (University of Colorado, 2011).
26. Yang L, Majumdar K, Liu H, Du Y, Wu H, Hatzistergos M, Hung P, Tieckelmann R, Tsai W, and Hobbs C, *Nano Lett.* 14(11), 6275–6280 (2014). [PubMed: 25310177]
27. Tarasov A, Zhang S, Tsai M-Y, Campbell PM, Graham S, Barlow S, Marder SR, and Vogel EM, *Adv. Mater.* 27(7), 1175–1181 (2015). [PubMed: 25580926]
28. Li S-L, Komatsu K, Nakaharai S, Lin Y-F, Yamamoto M, Duan X, and Tsukagoshi K, *ACS Nano* 8(12), 12836–12842 (2014). [PubMed: 25470503]
29. Appenzeller J, Radosavljević M, Knoch J, and Avouris P, *Phys. Rev. Lett.* 92(4), 048301 (2004). [PubMed: 14995412]
30. Appenzeller J, Knoch J, Derycke V, Martel R, Wind S, and Avouris P, *Phys. Rev. Lett.* 89(12), 126801 (2002). [PubMed: 12225112]
31. The I<sub>ON</sub>/off and V<sub>T</sub> values are extracted from the hole side for 10 min and 1 h p-doped MoS<sub>2</sub> FETs with Pd contacts.

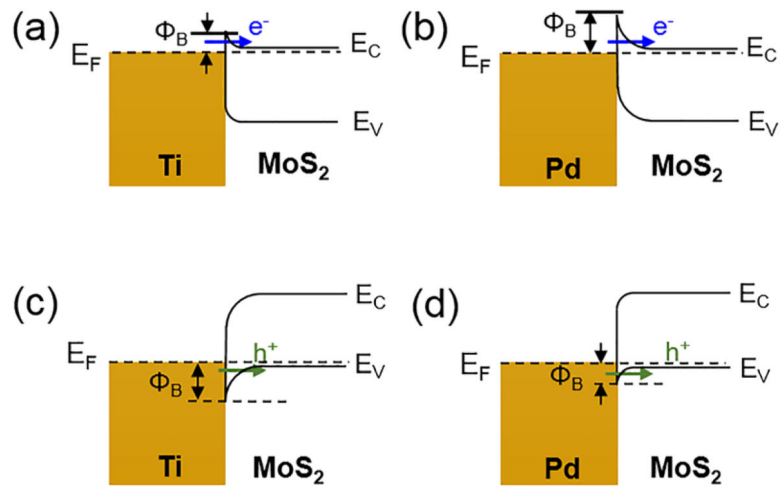
**FIG. 1.**

(a) Chemical structures of the p-dopant used in this study. (b) Schematics of a top-contact back-gate monolayer MoS<sub>2</sub> transistor with a TLM test structure. (c) Optical image of the MoS<sub>2</sub> transistor (scale bar is 25 μm). Representative transfer characteristics of a pristine MoS<sub>2</sub> FET at 300 K with either Ti or Pd as the contacts are shown on linear (d) and log (e) scales,  $V_{DS} = 5$  V and  $L = 5$  μm.

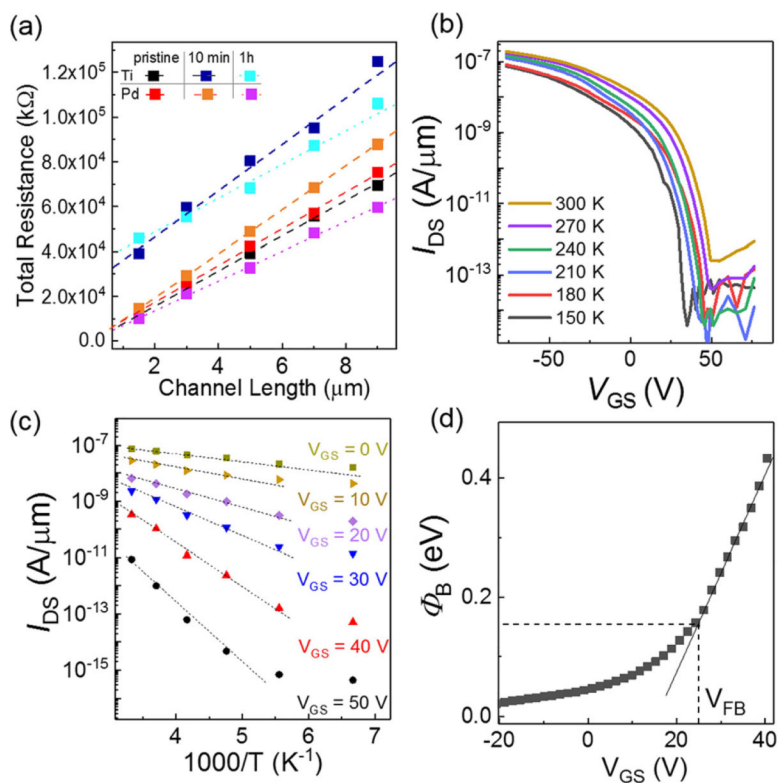




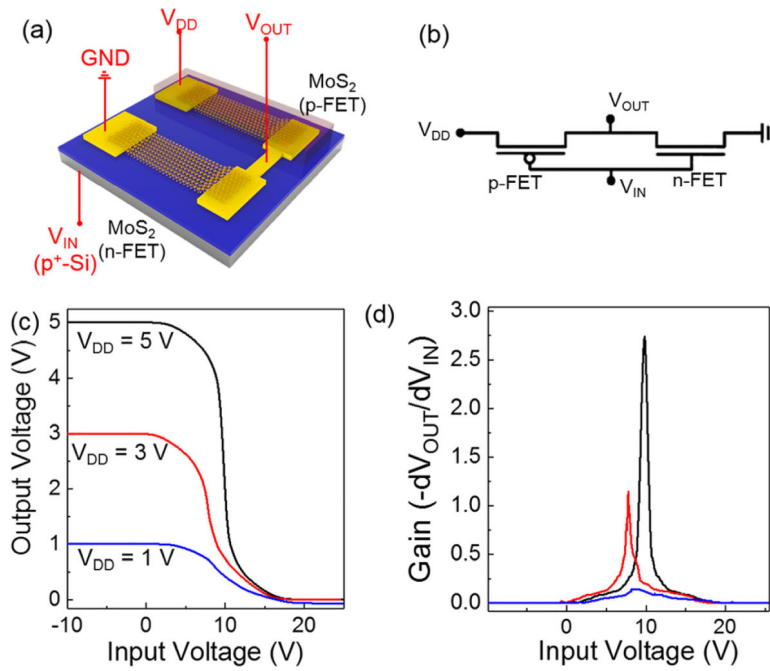
**FIG. 2.** (a) Representative transfer characteristics of a MoS<sub>2</sub> FET with Ti and Pd metal contacts after continuously p-doping with MB with two different exposure times.  $V_{DS} = 5$  V and  $L = 5$   $\mu\text{m}$ . (b) Output curves collected from the same Pd-contact device before and after 1 h of MB doping.



**FIG. 3.** Qualitative energy band diagrams for Ti (a) and (c) and Pd (b) and (d) contacted MoS<sub>2</sub> FETs in the ON-state showing the SBH height for the electron (a) and (b) and hole (c) and (d) injection at the interface after the p-doping.



**FIG. 4.** Electrical characteristics of p-doped MoS<sub>2</sub> devices. (a) The total resistance of MoS<sub>2</sub> FETs with different channel lengths. Channel width  $W = 25 \mu\text{m}$ . (b) Experimental transfer characteristics of the p-doped MoS<sub>2</sub> transistor with Pd contacts at different operating temperatures. The bias voltage is fixed at 5 V, and the channel length is  $5 \mu\text{m}$ . (c) Arrhenius-type plot constructed at different gate voltages. (d) Extracted barrier height as a function of applied gate voltages.



**FIG. 5.** (a) Schematic diagram of a MoS<sub>2</sub> inverter, (b) The electric circuit of the MoS<sub>2</sub> inverter. (c) Voltage-transfer characteristics of the MoS<sub>2</sub> inverter at different  $V_{DD}$  values. (d) The direct-current voltage gain of the inverter at  $V_{DD} = 1$  V, 3 V, and 5 V. The channels for both n-FET and p-FET are 5  $\mu\text{m}$ .

Summary of the comparative transistor performance of ML MoS<sub>2</sub> FETs with Ti and Pd contacts before and after the p-doping. Average of ca. 10 samples.<sup>31</sup>

**TABLE I.**

Contacts	Treatment	Mobility (cm <sup>2</sup> V <sup>-1</sup> s <sup>-1</sup> )			$I_{\text{ON/OFF}}$	$V_{\text{T,avg}}$ (V)	$\Phi_{\text{SB}}$ (meV)
		Electron	Hole	Hole			
Ti	Pristine	12.2 ± 1.7	...	...	10 <sup>7</sup>	-54 ± 3	100 <sup>a</sup>
	10 min p-doped	0.84 ± 0.10	6.0 × 10 <sup>-5</sup>	...	10 <sup>6</sup>	0 ± 7	...
	1 h p-doped	0.017 ± 0.009	9.3 × 10 <sup>-6</sup>	...	10 <sup>4</sup>	...	...
Pd	Pristine	8.3 ± 0.9	...	...	10 <sup>6</sup>	-25 ± 4	230 <sup>a</sup>
	10 min p-doped	1.4 × 10 <sup>-3</sup>	2.3 ± 0.7	...	10 <sup>6</sup>	40 ± 5	...
	1 h p-doped	7.4 × 10 <sup>-4</sup>	0.82 ± 0.12	...	10 <sup>5</sup>	75 ± 3	156 <sup>b</sup>

<sup>a</sup>For electron transportation.

<sup>b</sup>For hole transportation.