

Since January 2020 Elsevier has created a COVID-19 resource centre with free information in English and Mandarin on the novel coronavirus COVID-19. The COVID-19 resource centre is hosted on Elsevier Connect, the company's public news and information website.

Elsevier hereby grants permission to make all its COVID-19-related research that is available on the COVID-19 resource centre - including this research content - immediately available in PubMed Central and other publicly funded repositories, such as the WHO COVID database with rights for unrestricted research re-use and analyses in any form or by any means with acknowledgement of the original source. These permissions are granted for free by Elsevier for as long as the COVID-19 resource centre remains active.





www.elsevier.com/locate/physa

**PHYSICA** 

Physica A 379 (2007) 607-614

# On fractional order differential equations model for nonlocal epidemics

E. Ahmed<sup>a</sup>, A.S. Elgazzar<sup>b,c,\*</sup>

<sup>a</sup>Mathematics Department, Faculty of Science, Mansoura University, 35516 Mansoura, Egypt

<sup>b</sup>Mathematics Department, Faculty of Science, Al-Jouf University, P.O. Box 2014, Sakaka, Al-Jouf, Saudi Arabia

<sup>c</sup>Mathematics Department, Faculty of Education, 45111 El-Arish, Egypt

Received 17 June 2006 Available online 16 February 2007

#### Abstract

A fractional order model for nonlocal epidemics is given. Stability of fractional order equations is studied. The results are expected to be relevant to foot-and-mouth disease, SARS and avian flu.
© 2007 Elsevier B.V. All rights reserved.

Keywords: Fractional order differential equations; Stability; Applications to nonlocal epidemic models

## 1. Introduction

Recently three major epidemic diseases have occurred namely foot-and-mouth disease, severe acute respiratory syndrome (SARS) and avian (bird's) flu. Hopefully this will increase the awareness of modeling infectious diseases spreading that is an important topic in mathematical biology [1]. There are different approaches to this topic, e.g. ordinary differential equations, difference equations, partial differential equations and coupled map lattice. Here we use fractional order differential equations (FOD). The reason is that FOD are naturally related to systems with memory which exists in most biological systems. Also they are closely related to fractals which are abundant in biological systems. Consider the following evolution equation [2]:

$$\frac{\mathrm{d}f(t)}{\mathrm{d}t} = -\lambda^2 \int_0^t k(t-t')f(t')\,\mathrm{d}t'.$$

If the system has no memory then  $k(t-t') = \delta(t-t')$ , and one gets  $f(t) = f_0 \exp(-\lambda^2 t)$ . If the system has an ideal memory, then

$$k(t - t') = \begin{cases} 1 & \text{if} \quad t \ge t' \\ 0 & \text{if} \quad t < t' \end{cases},$$

<sup>\*</sup>Corresponding author. Mathematics Department, Al-Jouf University, P.O. Box 2014, Sakaka, Al-Jouf, Saudi Arabia. *E-mail address:* elgazzar@mans.edu.eg (A.S. Elgazzar).

hence  $f \approx f_0 \cos \lambda t$ . Using Laplace transform  $L[f] = \int_0^\infty f(t) \exp(-st) dt$ , one gets L[f] = 1 if there is no memory and L[f] = 1/s if there is ideal memory hence the case of non-ideal memory is expected to be given by  $L[f] = 1/s^{\alpha}$ ,  $0 < \alpha < 1$ . In this case the above equation becomes

$$\frac{\mathrm{d}f(t)}{\mathrm{d}t} = \frac{1}{\Gamma(\alpha)} \int_0^t (t - t')^{\alpha - 1} f(t') \, \mathrm{d}t',$$

where  $\Gamma(\alpha)$  is the Gamma function. This system has the following solution:

$$f(t) = f_0 E_{\alpha+1}(-\lambda^2 t^{\alpha+1}),$$

where  $E_{\alpha}(z)$  is the Mittag-Leffler function given by

$$E_{\alpha}(z) = \sum_{k=0}^{\infty} \frac{z^k}{\Gamma(\alpha k + 1)}.$$

It is direct to see that  $E_1(z) = \exp(z)$ ,  $E_2(z) = \cos z$ .

Following a similar procedure to study a random process with memory, one obtains the following fractional evolution equation:

$$\frac{\partial^{\alpha+1} P(x,t)}{\partial t^{\alpha+1}} = \sum_{n} \frac{(-1)^n}{n!} \frac{\partial^n [K_n(x) P(x,t)]}{\partial x^n}, \quad 0 < \alpha < 1,$$

where P(x, t) is a measure of the probability to find a particle at time t at position x. We expect that the above result will be relevant to many complex adaptive systems and to systems where fractal structures are relevant since it is argued that there is a strong relevance between fractals and fractional differentiation [3].

For the case of fractional diffusion equation the results are

$$\frac{\partial^{\alpha+1} P(x,t)}{\partial t^{\alpha+1}} = D \frac{\partial^2 P(x,t)}{\partial x^2}, \quad P(x,0) = \delta(x), \quad \frac{\partial P(x,0)}{\partial t} = 0,$$

then

$$P = \frac{1}{2\sqrt{D}t^{\beta}} M\left(\left(\frac{|x|}{\sqrt{D}t^{\beta}}\right); \beta\right), \quad \beta = \frac{\alpha + 1}{2},$$

$$M(z;\beta) = \sum_{n=0}^{\infty} \frac{(-1)^n z^n}{n!\Gamma(-\beta n + 1 - \beta)}.$$

For the case of no memory  $\alpha = 0 \Rightarrow M(z; 1/2) = \exp(-z^2/4)$ .

The paper is organized as follows: in Section 2, we study the stability of FOD. The stability conditions are derived and several examples are given. The stability conditions for some fractional order differential coupled map lattices are concluded in Section 3. Applications to nonlocal epidemics is introduced in Section 4. The stability conditions for the disease-free state are discussed. Some conclusions are presented in Section 5.

## 2. Stability of fractional order differential equations [4–6]

Consider the following system:

$$D^{\alpha}x(t) = f(x, y), D^{\alpha}y(t) = g(x, y), \quad \alpha \in [0, 1),$$
 (1)

where the fractional derivative in Eq. (1) is in the sense of Caputo. The equilibrium solutions are defined by  $f(x_{eq}, y_{eq}) = 0$ ,  $g(x_{eq}, y_{eq}) = 0$  and it is locally asymptotically stable if all the eigenvalues  $\lambda$  of the Jacobian

matrix  $A = \begin{bmatrix} \partial f/\partial x & \partial f/\partial y \\ \partial g/\partial x & \partial g/\partial y \end{bmatrix}$  evaluated at the equilibrium point satisfies the following condition [4,5]:

$$\left|\arg(\lambda)\right| > \frac{\alpha\pi}{2}$$
. (2)

The condition in Eq. (2) poses an interesting question namely:

What are the conditions that all the roots of the polynomial equation

$$P(\lambda) = 0, P(\lambda) = \lambda^n + a_1 \lambda^{n-1} + a_2 \lambda^{n-2} + \dots + a_n$$
(3)

satisfy Eq. (2) where all the coefficients in Eq. (3) are real?

For  $\alpha = 1$ , the solution is the Routh–Hurwitz conditions [7]

$$a_1 > 0, \begin{vmatrix} a_1 & 1 \\ a_3 & a_2 \end{vmatrix} > 0, \begin{vmatrix} a_1 & 1 & 0 \\ a_3 & a_2 & a_1 \\ a_5 & a_4 & a_3 \end{vmatrix} > 0, \dots$$
 (4)

For  $\alpha \in [0, 1)$ , these conditions are sufficient but not necessary. Since most biologically interesting systems are 1,2 and 3-dimensions, we will study the problem (3) for n = 1, 2, 3.

**Definition 1.** The discriminant D(f) of a polynomial

$$f(x) = x^{n} + a_{1}x^{n-1} + a_{2}x^{n-2} + \dots + a_{n}$$

is defined by  $D(f) = (-1)^{n(n-1)/2}R(f, f')$ , where f' is the derivative of f, if  $g(x) = x^n + b_1x^{l-1} + b_2x^{l-2} + \cdots + b_l$ , R(f, g) is the determinant of the corresponding Sylvester  $(n+l)\otimes(n+l)$  matrix. The Sylvester matrix is formed by filling the matrix beginning with the upper left corner with the coefficients of f(x), then shifting down one row and one column to the right and filling in the coefficients starting there until they hit the right side. The process is then repeated for the coefficients of g(x).

Using the results of Ref. [3], if D(f) > 0 < 0 then there is an even (odd) number of pairs of complex roots for the equation f(x) = 0. For n = 3, this implies that D(f) > 0, and all the roots are real while D(f) < 0 implies that there is only one real root and one complex root and its complex conjugate. For n = 3, we have

$$D(f) = 18a_1a_2a_3 + (a_1a_2)^2 - 4a_3a_1^3 - 4a_2^3 - 27a_3^2.$$

# Proposition 1.

- (i) For n = 1, the condition for (3) is  $a_1 > 0$ .
- (ii) For n = 2, the conditions for (3) are either Routh-Hurwitz conditions or

$$|a_1| < 0, 4a_2 > (a_1)^2, \left| \tan^{-1} \left( \frac{\sqrt{4a_2 - (a_1)^2}}{a_1} \right) \right| > \frac{\alpha \pi}{2}.$$
 (5)

(iii) For n = 3, if the discriminant of  $P(\lambda)$ , D(P) is positive, then Routh–Hurwitz conditions are the necessary and sufficient conditions for (3) i.e.

$$a_1 > 0, a_3 > 0, a_1 a_2 > a_3$$
 if  $D(P) > 0$ . (6)

- (iv) If D(P) < 0,  $a_1 \ge 0$ ,  $a_2 \ge 0$ ,  $a_3 > 0$ ,  $\alpha < 2/3$ , then condition (3) is satisfied. Also if D(p) < 0,  $a_1 < 0$ ,  $a_2 < 0$ ,  $\alpha > 2/3$ , then all roots of  $P(\lambda) = 0$  satisfies  $|\arg(\lambda)| < \alpha \pi/2$ .
- (v) If D(P) < 0,  $a_1 > 0$ ,  $a_2 > 0$ ,  $a_1 a_2 = a_3$  then condition (3) is satisfied for all  $\alpha \in [0,1)$ .
- (vi) For general n,  $a_n > 0$  is a necessary condition for condition (3) to be satisfied.
- (vii) If  $\forall \lambda$ ,  $P(\lambda) = P(-\lambda)$  then define  $x = \lambda^2$  and Routh–Hurwitz conditions for the resulting polynomial in x are necessary conditions for (3) for all  $\alpha \in [0,1)$ .

(viii) For n > 1, the necessary and sufficient condition for (3) is  $\int_{-\infty}^{0} dz/P(z)|_{C1} + \int_{0}^{\infty} dz/P(z)|_{C2} = 0$ , where  $C_1$  is the curve  $z = x(1-i\tan\alpha\pi/2)$  and  $C_2$  is the curve  $z = x(1+i\tan\alpha\pi/2)$ ,  $i = \sqrt{-1}$ .

**Proof.** Case (i) is obvious. For (ii), the two roots are  $\lambda_{\pm} = \left[ -a_1 \pm \sqrt{(a_1^2 - 4a_2)} \right] / 2$ . If both roots are real or complex conjugates with negative real parts then condition (3) is equivalent to the Routh-Hurwitz case. If the two roots are complex conjugate with positive real parts, then the two roots become  $\lambda_{\pm} = \left[ -a_1 \pm i \sqrt{4a_2 - (a_1)^2} \right] / 2$  and one gets Eq. (5).

To prove (iii) not that if n = 3, D(p) > 0, then all the roots of  $p(\lambda) = 0$  are real hence Routh–Hurwitz conditions are both necessary and sufficient for (3).

To prove (iv) not that if n = 3, D(P) > 0, then the roots of  $P(\lambda) = 0$  are one real and a complex conjugate pair thus

$$p(\lambda) = (\lambda + b)(\lambda - \beta - i\gamma)(\lambda - \beta + i\gamma) \Rightarrow$$

$$a_1 = b - 2\beta, \ a_2 = \beta^2 + \gamma^2 - 2b\beta, \ a_3 = b(\beta^2 + \gamma^2), \ b \geqslant 0$$
(7)

and  $a_1 > 0 \Rightarrow b > 2\beta$ ,  $a_2 > 0 \Rightarrow \beta^2 \sec^2 \theta > 2b\beta > 4\beta^2 \Rightarrow \theta > \pi/3$ , where  $\theta = |\arg(\lambda)|$ . The second part is proved similarly.

To prove (v) if n = 3, D(P) > 0, then Eq. (7) is valid. Now  $a_1 a_2 = a_3 \Rightarrow b^2 \beta + \beta(\beta^2 + \gamma^2) = 2b\beta^2 \Rightarrow \beta = 0$  or  $\beta^2 + \gamma^2 + b^2 = 2b\beta$ 

The last equality is not valid if both  $a_1 > 0$ ,  $a_2 > 0$ .

To prove (vi) use the fact that for general n

$$P(\lambda) = \left[ \prod_{j} (\lambda + b_{j}) \right] \left[ \prod_{k} (\lambda - \beta_{k} - i\gamma_{k})(\lambda - \beta_{k} + i\gamma_{k}) \right] \Rightarrow$$

$$a_{n} = \left[ \prod_{j} b_{j} \right] \left[ \prod_{k} (\beta_{k}^{2} + \gamma_{k}^{2}) \right]. \tag{8}$$

To prove (vii) use Eq. (8) and  $P(\lambda) = P(-\lambda) \forall \lambda \Rightarrow P(\lambda)$  contains only even power of

$$\lambda \Rightarrow b_j = 0 \quad \forall j, \quad \beta_k = 0 \forall k \Rightarrow$$

$$P(\lambda) = \prod_k (\lambda^2 + \gamma_k^2) = \prod_k (x + \gamma_k^2).$$

And all the roots x should be negative.

To prove (viii) not that if P(z) has no roots in the region  $|\arg(\lambda)| < \alpha \pi/2$ , hence the function 1/P(z) will be analytic in this region. Using Cauchy theorem  $\oint_C f(z) dz = 0$  for all f(z) analytic within and on the curve C, and that P(z) is polynomial of degree >1 this completes the proof.  $\square$ 

One can give explicit examples where the Routh-Hurwitz conditions are not valid yet Eq. (3) is satisfied for explicit  $\alpha$ , e.g.  $a = \frac{1}{2}$ . The first example is

$$\lambda^{3} - \frac{\lambda^{2}}{2} + \frac{\lambda}{2} + \frac{1}{2} = 0,$$
  
$$\lambda^{3} - \left(1 - \frac{1}{\sqrt{2}}\right)\lambda^{2} + \left(\frac{1}{2} - \frac{1}{\sqrt{2}}\right)\lambda + \frac{1}{2\sqrt{2}} = 0.$$

For n = 3, one can solve the polynomial equation explicitly and apply Eq. (3) to the solutions. The solution method is:

Define  $y = \lambda - a_1/2$ , then the polynomial becomes

$$y^3 + py + q = 0$$
,  $p = -a_1^2/3 + a_2$ ,  $q = 2a_1^3/27 - a_1a_2/3 + a_3$ ,

$$y = \rho + \sigma, \ \rho = \sqrt[3]{-(q/2) + \sqrt{(q/2)^2 + (p/3)^3}}, \ \ \sigma = \sqrt[3]{-(q/2) - \sqrt{(q/2)^2 + (p/3)^3}},$$

choose the roots such that  $\rho \sigma = -p/3$ .

**Conjecture 1.** For all n > 3, if  $\Delta_1, \Delta_2, \dots, \Delta_n$  are Routh–Hurwitz determinants

$$\Delta_1 = a_1, \, \Delta_2 = \begin{vmatrix} a_1 & 1 \\ a_3 & a_2 \end{vmatrix}, \, \Delta_3 = \begin{vmatrix} a_1 & 1 & 0 \\ a_3 & a_2 & a_1 \\ a_5 & a_4 & a_3 \end{vmatrix}, \, \ldots,$$

then the conditions

$$\Delta_i > 0, i = 1, 2, \dots, n-2, \quad a_n > 0, \quad \Delta_{n-1} = 0,$$
 (9)

are sufficient conditions that Eq. (3) is valid for all  $\alpha \in [0, 1)$ .

This conjecture can be proved for n = 4. The case n = 3 is proved in Proposition 1.

Now we apply Proposition 1 to derive the value of  $\alpha$  at which chaos or instability disappears at some models. Fractional order Lotka-Volterra (FLV) predator–rey model is given by

$$D^{\alpha}x = \beta x - xy$$
,  $D^{\alpha}y = -\gamma y + xy$ ,  $\alpha \in [0, 1]$ ,

where  $\beta$ ,  $\gamma$  are positive constants. The equilibrium solutions of FLV are (0, 0) and  $(\gamma, \beta)$ . The eigenvalues corresponding to (0, 0) are  $\beta$ ,  $-\gamma$  respectively hence (0, 0) is unstable for all  $\alpha \in [0, 1]$ . The eigenvalues corresponding to  $(\gamma, \beta)$  are  $\pm i\sqrt{\beta\gamma}$  hence  $|\arg(\lambda)| = \pi/2$ , then it is locally asymptotically stable for all  $\alpha \in [0, 1)$ . This agrees with the numerical results of Ref. [5].

Fractional order Chen (FOC) model is given by [7]

$$D^{\alpha}x = a(y - x), D^{\alpha}y = (c - a)x - xz + cy, D^{\alpha}z = xy - bz,$$

where a=35, b=3, c=28. The equilibrium solutions for FOC are (0, 0, 0) and  $(x^*, x^*, x^{*2}/b)$ , where  $x^*=\pm\sqrt{b(2c-a)}$ . The eigenvalues of (0, 0, 0) are  $\lambda=-b$ ,  $\lambda^2+(a-c)\lambda-a(2c-a)=0$ . Using Proposition 1 part (vi), then (0, 0, 0) is unstable for all  $\alpha \in [0, 1]$ . The eigenvalues of the internal equilibrium solution is  $\lambda^3+(a-c+b)\lambda^2+bc\lambda+2ab(2c-a)=0$ . Since all the coefficients of this polynomial are positive and it is easy to check that D(P)<0, using Proposition 1 part (iv) then the internal solution of FOC will be stable for  $a<\frac{2}{3}$  which is in excellent agreement with the numerical result obtained by and Chen [8].

A similar result is obtained for fractional order Lorenz system

$$D^{\alpha}x = c(v-x)$$
,  $D^{\alpha}v = rx - xz - v$ ,  $D^{\alpha}z = xv - bz$ ,

where r, b, c are positive constants and r > 1.

The fractional order Rossler system (FOR) is defined by [9]

$$D^{\alpha}x = -(v+z)$$
,  $D^{\alpha}v = x + av$ ,  $D^{\alpha}z = 0.2 + xz - 10z$ ,  $a = 0.6$ 

There are one equilibrium z = -y, x = -ay,  $y = [-10 + \sqrt{100 - 0.8a}]/2a$ , and the characteristic polynomial is

$$\lambda^3 + (ay - a + 10)\lambda^2 + (1 - ya^2 - 10a - y)\lambda + 10 + 2ay = 0.$$

It is direct to see that  $a_2 < 0$ ,  $a_1 > 0$  hence by solving the characteristic equation directly one gets the following roots:

$$\lambda = -9.986, \ \lambda = 0.3 - 0.954i, \ \lambda = 0.3 + 0.954i.$$

This implies that the equilibrium solution is locally asymptotically stable if  $\alpha < 0.8$  which is in excellent agreement with the result obtained numerically in Ref. [9]:  $\alpha \in (0.7, 0.8)$ .

Considering the fractional order Chua model

$$D^{\alpha}x = a(y - x) - g(x), D^{\alpha}y = s[-a(y - x) + z],$$
  
 $D^{\alpha}z = -c(y + rz),$ 

$$g(x) = \begin{cases} -x & \text{if } |x| < 1, \\ [-1 + b(|x| - 1)] \operatorname{sgn}(x) & \text{if } 1 < |x| < 10, \\ [10(|x| - 10) + 9b - 1] \operatorname{sgn}(x) & \text{if } |x| > 10, \end{cases}$$

where a = 0.923 or 1, b = 0.636, r = 0.071 or 0, s = 0.066. The equilibrium solution is  $x \approx \pm (1-b)/(a-b)$ ,  $y \approx 0$ ,  $z \approx -ax$ . Linearizing about the quilibrium solution one gets

$$\begin{bmatrix} -a+b & a & 0 \\ sa & -sa & s \\ 0 & -c & -rc \end{bmatrix},$$

whose eigenvalues for a=1, b=0.636, c=0.779, s=0.066, r=0.071 are  $\lambda=-0.486$ ,  $\lambda=5\times 10^{-4}\pm 0.184i$  which implies that this equilibrium is locally asymptotically stable for  $\alpha\in[0,1)$ . The eigenvalues for a=0.923, b=0.636, c=0.779, s=0.066, r=0 are  $\lambda=-0.406$ ,  $\lambda=0.029\pm.188i$  which implies that this equilibrium is locally asymptotically stable for  $\alpha<0.9$ .

# 3. Stability conditions for some fractional order differential coupled map lattices

Spatial effects are important in many biological systems. Thus generalizing fractional order systems (FOS) to include them is important. The standard approach is fractional order partial differential equations. However since most biologically interesting systems are nonlinear [10], one gets fractional order nonlinear partial differential equations whose existence and uniqueness has not been established yet. Therefore, we use coupled map lattices (CML) [11] as an alternative approach to include spatial effects in FOS. Consider the 1-system CML

$$D^{\alpha}u_{i}(t) = (1 - D)f(u_{i}) + \frac{D}{2}[f(u_{i+1}) + f(u_{i-1})],$$
  

$$i = 1, 2, \dots, n,$$
(10)

where D is a positive constant. The homogeneous equilibrium solution of Eq. (10) satisfies  $f(u_{eq}) = 0$  and it is stable if all the eigenvalues of the circulant matrix [12]

$$\begin{bmatrix} (1-D)f' & (D/2)f' & 0 & \dots & 0 & (D/2)f' \\ (D/2)f & (1-D)f' & (D/2)f & 0 & \dots & 0 \\ 0 & (D/2)f' & (1-D)f' & (D/2)f' & \dots & 0 \\ \dots & \dots & \dots & \dots & \dots & \dots \\ 0 & \dots & 0 & (D/2)f' & (1-D)f' & (D/2)f' \end{bmatrix}$$

satisfies Eq. (2). Since the eigenvalues of the circulant matrices [12] are known, the local asymptotic stability conditions become

$$\left| \arg \left[ (1 - D)f'(u_{eq}) + Df'(u_{eq}) \cos \frac{2\pi k}{n} \right] \right| > \frac{\alpha \pi}{2},$$

$$k = 0, 1, \dots, n - 1$$

$$(11)$$

Generalizing to 2-system CML

$$D^{\alpha}x_{i} = (1 - D)f(x_{j}, y_{j}) + \frac{D}{2} \Big[ f(x_{j+1}, y_{j+1}) + f(x_{j-1}, y_{j-1}) \Big], D^{\alpha}y_{i} = (1 - D)g(x_{j}, y_{j}) + \frac{D}{2} \Big[ g(x_{j+1}, y_{j+1}) + g(x_{j-1}, y_{j-1}) \Big],$$
(12)

j = 1, 2, ..., n. The homogeneous equilibrium solutions are given by  $f(x_{eq}, y_{eq}) = 0$ ,  $g(x_{eq}, y_{eq}) = 0$  and is locally asymptotically stable if all the eigenvalues of the following matrix B satisfies Eq. (2):

$$B = \left[ (1 - D) + \frac{D}{2} \cos\left(\frac{2\pi k}{n}\right) \right] \begin{bmatrix} \partial f/\partial x & \partial f/\partial y \\ \partial g/\partial x & \partial g/\partial y \end{bmatrix},\tag{13}$$

where the Jacobian matrix is evaluated at the equilibrium point.

## 4. Application to a nonlocal epidemic model

Now we present a nonlocal interacting epidemic model. The nonlocal interactions in the epidemic spreading is widely observed in many outbreaks, especially the recent ones of FMD, SARS and bird's flu. The Lajmanovich-Yorke model [13,14] is a special case of coupled map lattices. It is a globally interacting model that emphasizes the diffusion of a disease from other infected farms beside its spread within the same farm. We have generalized it to the inhomogeneous case [15] and introduced it as an approximation for the epidemic spreading on small-world networks [15]. Consider *n* patches, where each one contains a certain number of individuals (say animals). In general these patches are not identical. Infection spreads from infected animals within the patch and due to those diffusing from other patches. The fractional order form of the model is given by

$$D^{\alpha}y_{i} = \lambda_{i}y_{i}(1 - y_{i}) + \mu_{i}(1 - y_{i}) \sum_{i \neq i} y_{j} - \gamma_{i}y_{i}, i = 1, 2, \dots, n,$$
(14)

where  $y_i$  is the number of infected individuals in the *i*th patch. The first term represents the infection within the patch with a rate  $\lambda_i$ . The second term represents the effect of other patches both nearby and far away at a rate  $\mu_i$ . The recovery rate is represented by  $\gamma_i$ . Since the effect of other patches (second term in Eq. (14)) is significantly smaller than the first one, we expect that  $\mu_i \ll \lambda_i$ . The disease is eradicated if  $y_i = 0$ ,  $\forall i = 1, 2, ..., n$ . So the stability of this solution is studied. Using the same procedures in the previous section, the stability of the zero solution (disease free state) is that all the eigenvalues x of the matrix A satisfy  $|\arg(x)| > \alpha \pi/2$  where

For n = 2, the solutions  $y_i = 0$ , i = 1, 2 are locally asymptotically stable if

$$\left| \tan^{-1} \left( \frac{\sqrt{(\lambda_1 - \gamma_1 + \lambda_2 - \gamma_2)^2 - 4[(\lambda_1 - \gamma_1)(\lambda_1 - \gamma_1) - \mu_1 \mu_2]}}{\lambda_1 - \gamma_1 + \lambda_2 - \gamma_2} \right) \right|$$

$$> \frac{\alpha \pi}{2}.$$

For n = 3, the characteristic polynomial of A is

$$x^3 + a_1 x^2 + a_2 x + a_3 = 0,$$

where

$$a_1 = -[\lambda_1 - \gamma_1 + \lambda_2 - \gamma_2 + \lambda_3 - \gamma_3],$$

$$a_2 = (\lambda_1 - \gamma_1)(\lambda_2 - \gamma_2) - \mu_1\mu_2 + (\lambda_1 - \gamma_1)(\lambda_3 - \gamma_3) - \mu_1\mu_3 + (\lambda_3 - \gamma_3)(\lambda_2 - \gamma_2) - \mu_3\mu_2$$

and

$$a_3 = -[(\lambda_1 - \gamma_1)(\lambda_2 - \gamma_2)(\lambda_3 - \gamma_3) + 2\mu_1\mu_2\mu_3 - (\lambda_1 - \gamma_1)\mu_2\mu_3 - (\lambda_3 - \gamma_3)\mu_2\mu_1 - (\lambda_2 - \gamma_2)\mu_1\mu_3].$$

Using part (vi) in Proposition 1, the disease free state is unstable if  $a_3 < 0$ .

Finally vaccination may affect the parameters  $\lambda_i$  but not  $\mu_i$ . Also it is important to know whether the vaccine works once administered or it takes time till it becomes effective.

#### 5. Conclusions

A fractional order evolution equation describing a random process with memory is concluded. This will be relevant to many complex adaptive systems. The stability conditions for both fractional order differential equations and fractional order differential coupled map lattices are derived. A fractional order form of Lajmanovich-Yorke model is introduced. This model emphasis the nonlocal interaction beside the local interaction in the epidemic spreading. Both types of interaction are widely observed in the recent FMD, SARS and bird's flu outbreaks.

## Acknowledgment

One of us (E. Ahmed) wishes to thank Prof. H. M. Yehia for discussions and help.

#### References

- [1] O. Diekmann, J.A.P. Heesterbeek, Mathematical Epidemiology, Wiley, New York, 2000.
- [2] A.A. Stanislavsky, Memory effects and macroscopic manifestation of randomness, Phys. Rev. E 61 (2000) 4752.
- [3] A. Rocco, B.J. West, Fractional calculus and the evolution of fractal phenomena, Physica A 265 (1999) 535.
- [4] D. Matignon, Stability results for fractional differential equations with applications to control processing, Computat. Eng. Syst. 2 (1996) (Lille France 963).
- [5] E. Ahmed, A.M.A. El-Sayed, H.A.A. Elsaka, Equilibrium points, stability and numerical solutions of fractional-order predator-prey and rabies models, Jour. Math. Anal. Appl. 325 (2007) 542.
- [6] E. Ahmed, A.M.A. El-Sayed, H.A.A. Elsaka, (2005), Submitted.
- [7] A.P. Mishina, I.V. Proskuryakov, Higher algebra, Nauka (1965).
- [8] C. Li, G. Chen, Chaos in the fractional order Chen model and its control, Chaos Soliton. Fract. 22 (2004) 549.
- [9] C. Li, G. Chen, Chaos and hyperchaos in fractional order Rossler equations, Physica A 341 (2004) 55.
- [10] L. Edelstein-Keshet, Mathematical models in biology, Siam Classic.Appl. Math. 46 (2004).
- [11] K. Kaneko, Theory and Applications of Coupled Map Lattices, Wiley, New York, 1993.
- [12] S. Barnett, Matrices, Cambridge Univ. Press., Cambridge, 1990.
- [13] A. Lajmanovich, J.A. Yorke, A deterministic model for gonorrhea in a non-homogeneous population, Math. Biosci. 28 (1976) 221.
- [14] F. Ball, Stochastic and deterministic models for SIS epidemics among a population partitioned into households, Math. Biosci. 156 (1999) 41.
- [15] E. Ahmed, A.S. Hegazi, A.S. Elgazzar, An epidemic model on small-world networks and ring vaccination, Int. J. Mod. Phys. C 13 (2002) 189.