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Development of a fast calibration method for image mapping spectrometry

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Abstract

An image mapping spectrometer (IMS) is a snapshot hyperspectral imager that simultaneously captures both the spatial (x, y) and spectral (λ) information of incoming light. The IMS maps a three-dimensional (3D) datacube (x, y, λ) to a two-dimensional (2D) detector array (x, y) for parallel measurement. To reconstruct the original 3D datacube, one must construct a lookup table that connects voxels in the datacube and pixels in the raw image. Previous calibration methods suffer from either low speed or poor image quality. We herein present a slit-scan calibration method that can significantly reduce the calibration time while maintaining high accuracy. Moreover, we quantitatively analyzed the major artifact in the IMS, the striped image, and developed three numerical methods to correct for it.

1. INTRODUCTION

Hyperspectral imaging (HSI) has been extensively employed in myriad applications such as biomedical imaging $[1-4]$, remote sensing $[5-8]$, and machine vision $[9-11]$. HSI is a functional combination of a two-dimensional (2D) camera and a spectrometer, acquiring both the spatial (x, y) and spectral (λ) information of a scene. While capable of acquiring a hyperspectral datacube (x, y, λ) , most existing HSI techniques require scanning either in the

Data Availability

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Disclosures. The authors declare no conflicts of interest.

The data that support the plots within this paper and other findings of this study are available from the corresponding author upon reasonable request.

Code Availability

Codes used for this work are available from the corresponding author upon reasonable request.

spatial [12–14] or spectral domain [15–18]. Because only a small portion of the hyperspectral datacube can be seen at a time by these imagers [19], the scanning mechanism significantly jeopardizes the light throughput and hinders the imaging of dynamics.

In contrast, snapshot HSI techniques overcome the above limitations by acquiring all hyperspectral datacube voxels in parallel [20,21]. Representative techniques encompass coded aperture snapshot spectral imaging (CASSI) [22–24], computed tomography imaging spectrometry (CTIS) [25,26], and image mapping spectrometry (IMS) [27–34]. While all these techniques operate in a snapshot format, only IMS features 100% light throughput while maintaining a compact form factor and high computational efficiency, making it suitable for real-time field imaging applications.

The IMS operates by mapping a three-dimensional (3D) hyperspectral datacube to a 2D camera through an angled mirror array, referred to as an image mapper. The underlying principle is detailed in Ref. [35]. In brief, the image mapper splits an image into strips and redirects them to different locations on the 2D camera. Because the mirror facets on the array have varied tilts, a blank region is created between adjacent sliced images. A dispersion element, such as a prism, then disperses the sliced image and fills this blank region with spectrum. In this way, each pixel on the 2D camera is encoded with the unique spatial and spectral information of the original hyperspectral datacube, and remapping generates the image.

The image remapping during reconstruction requires a lookup table, which connects each voxel in the hyperspectral datacube to a pixel in the raw image. The calibration of this lookup table is nontrivial. Previous calibration methods, such as edge alignment [27] and point scan [35], suffer from either low image quality or slow calibration speed. For example, the edge alignment method images a target with a sharp edge and constructs the lookup table by aligning the image slices with respect to this feature. Because image slices experience different levels of distortion, the lookup table built upon local feature alignment cannot be faithfully applied to the global image, leading to a reduced image quality and spectrum accuracy [36]. In contrast, the point scan method builds the table by illuminating one datacube voxel at a time, followed by pinpointing the centroid of the impulse response image on the 2D camera. Because the calibration is performed on each datacube voxel, the reconstructed image quality and spectral accuracy are superior to those achieved by edge alignment. However, due to the reliance on scanning, the calibration process is time consuming, typically taking hundreds of hours to complete.

To overcome the above limitations, we herein present a fast and accurate calibration method, which we refer to as slit scan. Slit scan can correct for the same image slice distortion as with point scan; however, it does not have the need for prolonged 2D scanning, thereby significantly reducing the calibration time to tens of minutes. Moreover, we quantitatively analyzed the primary artifact in the IMS, the striped image, and we provide several solutions to correct for it. The radiometric calibration was detailed in Ref. [35].

2. OPTICAL SETUP OF THE IMS

A schematic of the IMS is shown in Fig. 1. The mapping mirror is placed at a plane that is conjugate with both the object and the 2D camera. The object is imaged by an objective lens and a collecting lens, and an intermediate image is formed at the image mapper. As a result, each mirror facet on the image mapper reflects a slice of the intermediate image. Because the angles of the mirror facets are grouped into periodic blocks [29], the reflected image slices are directed to the correspondent subfields. In our experiment, a total of 480 image slices are grouped into 5×8 subfields, each containing 12 image slices. Within a subfield, there is a blank region between adjacent image slices. After passing through a dispersion prism, the spectrum of each image slice occupies this blank region. The dispersion prism array in our system works in a visible light range (400–700 nm) as detailed in Ref. [31]. A bandpass filter is used to limit the input spectral range, thereby avoiding spectral overlap between adjacent sliced images. A microlens array with 5×8 microlenses is used to focus the final image on the camera. Under monochromatic illumination, the width of each image slice is one pixel on the camera, and the width of a blank region between two adjacent image slices is 40 pixels. Under white-light illumination, these 40 pixels represent 40 different color channels in the visible spectral range. The parameters of the IMS are detailed in Refs. [31,32].

3. IMS CALIBRATION METHODS

The IMS maps a 3D hyperspectral datacube (x, y, λ) to a 2D image (x, y) on the camera. If the point spread function (PSF) of the IMS is a delta function, the mapping between a voxel in the datacube and a pixel in the image would be one to one. The imaging process can be described by the following equation:

$$
I(x, y) = T \cdot O(x, y, \lambda),\tag{1}
$$

where *I* is the captured 2D raw image, O is the 3D hyperspectral datacube, and *T* is the mapping matrix of the IMS. The goal of calibration is to determine the inverse of T. With that being known, the original datacube can be easily reconstructed based on

$$
O(x, y, \lambda) = T^{-1}I(x, y).
$$
 (2)

In practice, rather than calculating the mapping matrix, we experimentally build a lookup table, which contains the index of each voxel in the datacube and locations of the corresponding pixel in the image.

As mentioned, two methods have been developed for constructing the lookup table: edge alignment [27] and point scan [35]. Edge alignment simultaneously extracts all image slices in the raw image and rearranges them to form a spectral channel. In this case, the lookup table contains the locations of image slices both in the datacube and in the raw image. The sequence of image slices in the raw image is determined by the design of the image mapper, and such knowledge is known as priors. Due to distortion, the image slices in the raw image are tilted and exhibit different magnification ratios. A common approach to extracting these

image slices is to capture an empty field under uniform illumination, followed by image binarization. After this procedure, each image slice is isolated, and a simple curve fitting reveals the shape contour. Next, a target with a sharp edge that is perpendicular to the mirror facet direction is imaged, and all the resultant images slices are aligned with respect to this edge. The image processing pipeline is illustrated in Fig. 2(a). Finally, the starting coordinates of the aligned image slices in the raw image are recorded to fill in the lookup table.

Despite being fast in implementation, the major drawback of edge alignment is that the reconstructed image suffers from distorted edges. This is because the image slices experience varied magnification ratios when they pass through different lenslets in the array. As illustrated in Fig. 2(b), an object consisting of two straight lines is divided into three image slices. These three adjacent image slices are imaged by different lenslets with varied magnification ratios. If we extract these imaged slices from the raw image and align them with respect to line 2, the edge of line 1 would appear distorted. We further show a realworld example when imaging a US air force (USAF) resolution target. The reconstructed image through edge alignment is shown in Fig. 2(c). Area 3 shows the bars of group 2 element 4, indicating 5.66 lp/mm. All image slices were aligned with respect to a horizontal line that crosses the central field of view (FOV).As shown in the zoomed inset, Area 2 is close to the reference edge and thus has a high reconstruction quality. In contrast, Area 1 and Area 3 located at the upper and bottom FOV exhibit distorted edges.

In comparison, the point scan method constructs the lookup table by illuminating hypercube voxels one at a time and recording the location of the impulse response on the 2D camera. To calibrate the entire hypercube, one must raster scan a monochromatic point light source —such as a pinhole illuminated by narrow band light—across the 2D FOV. Because this method experimentally identifies the mapping relation for each voxel in the hypercube, the errors induced by the image distortion and ununiform magnification are corrected. However, the implementation of the point scan method is time consuming. For example, to calibrate an image of 500×500 pixels at a given wavelength, a total of 139 his needed, using a dwell time of 2 sat each scanning step.

4. SLIT-SCAN CALIBRATION

To overcome previous limitations and enable fast and accurate IMS calibration, we herein present slit scan, a method that maintains the accuracy of point scan but at a speed orders of magnitude faster. Simply put, rather than scanning a point source across the 2D FOV, we use a uniformly illuminated slit and scan only the axis parallel to the mirror facet to construct the lookup table.

Our method has been made possible due to the unique optical architecture of the IMS, where the image mapper slices the incident field into strips and projects them to different locations of the detector array. Upon illumination by a thin slit source that is perpendicular to the slice direction, the image of the slit is divided into a series of points, which create impulse responses separable in space in the raw image. Pinpointing the centroids of these impulse responses then generates the lookup table for all the voxels illuminated in parallel. In

essence, this method uses the image mapper to perform point scans on all the points of the slit simultaneously.

The experimental setup is shown in Fig. 3(a). A 10 μm optical slit (Lenox Laser, G-SLIT-1- DISC-10) was placed at the object plane of the IMS, where the intermediate image of the slit at the image mapper is perpendicular to the mirror facet. The slit is mounted on a motorized translation stage (Thorlabs, MTS50-Z8) for scanning. We chose the width of the slit to match with the size of a resolution cell at the object plane. The image of the slit is then sliced by the mirror facets. Because the length of the slit is longer than the width of the FOV, each mirror facet creates a unique response on the detector array. A representative raw image is shown in Fig. 3(b). By scanning the slit along the direction of mirror facets and recording the locations of all responses, we obtained a complete lookup table.

Compared with the edge alignment and point scan methods, slit scan offers two prominent advantages. First, slit scan calibrates the mapping for each voxel in the hypercube, and therefore it automatically corrects for the image distortion and varied magnification in the IMS. A reconstructed image of a USAF resolution target at 532 nm is shown in Fig. 4. Compared with the result obtained through edge alignment in Fig. 3(d), the reconstructed image shows no distorted edges. To quantitatively compare the reconstruction quality, we calculated the standard deviation of the edge in Area 1 from both results. Here we define the standard deviation of a horizontal edge as the variation in its pixel index along the vertical axis. The standard deviation (STD) of the edge reconstructed through edge alignment is 1.47 pixels, compared with 0.55 pixels obtained through slit scan, the reconstruction accuracy has been imporved by 60%. The STD of the edge in Area 2 is 0.31 pixel, and the STD of the edge of one bar in Area 3 is 0.56 pixel in Fig. 4. Second, slit scan can be implemented and performed rapidly. In our experiment, a wavelength layer of the hypercube (i.e., a spectral channel image) has 445×480 voxels, where 480 is the total number of mirror facets on the image mapper. Because the scanning pitch is one voxel, only 445 steps are required. Again with a 2 s dwell time for each step, calibration takes \sim 30 min to build the entire lookup table. In contrast, the point scan method needs to scan a total of 445×480 steps. Therefore, our method increases the calibration speed by a factor of 480, which is equal to the total number of mirror facets on the image mapper. This dramatic speed improvement, which we refer to as the parallelization advantage, becomes more critical for high-resolution IMSs with a large number of mirror facets.

Note that it is sufficient to calibrate the lookup table at only one wavelength. As shown in Fig. 5(a), the white stripe represents an image slice on the camera, and it disperses horizontally, as denoted by the rainbow area. The lookup table at other wavelengths can be obtained by shifting the recorded locations of responses horizontally, and the shifted distance in pixels can be deduced from the chromatic dispersion curve of the prism as shown in Fig. 5(b). In order to measure the dispersion curve, an optical pinhole was placed at the object plane, and six color filters (Thorlabs, FB470–10, FB510–10, FB550–10, FB590–10, FB630–10, and FB670–10) were used sequentially to find the relationship between the wavelength and pixel locations. Considering the spectral resolution of the current system (-5) nm), a spectral bandwidth of 10 nm occupies only two pixels, leading to a negligible localization error. To further minimize the localization error of the wavelength responses, we

performed multiple measurements and averaging. The final curve was fitted with a secondorder polynomial based on following approxiamtion:

$$
\delta(\lambda) \approx [n(\lambda) - 1]\alpha,\tag{3}
$$

where λ is wavelength of light, $\delta(\lambda)$ is the wavelength-dependent dispersion, $n(\lambda)$ is the refractive index of the prism glass, and α is the angle of the prism. The relation between reflective index and wavelength follows the equation

$$
n(\lambda) = A + \frac{B}{\lambda^2},\tag{4}
$$

where A and B are constant coefficients.

5. STRIPE ARTIFACT

Although the slit-scan method can generate an accurate lookup table, the process itself does not correct for the image slices' intensity variations due to the reflectivity difference of the mirror facets and the defocus of the microlens array. The subsequently induced artifact, referred to as the striped image, degrades the image quality in the IMS. Here we develop a mathematical model to quantitatively analyze this artifact.

We describe the effect of mirror reflectivity and microlens defocus on the image quality using the line spread function (LSF). Although other optical aberrations may also exist in the system, we found that the defocus is the dominant factor that accounts for the stripped image artifacts in the IMS [27]. The image of a mirror facet can be written as

$$
I_m(x, y) = L_m(x) * I'_m(x, y),
$$
\n(5)

where $L_m(x)$ is the LSF of the *m*th mirror facet, denotes convolution, and $I'_m(x, y)$ depicts the geometric image of the mth mirror facet. Two representative LSFs are shown in Fig. 6(a). The intensity profile along the dashed line is plotted in Fig. 6(b). The full width at halfmaxima (FWHM) of these two LSFs are different. To reconstruct a spectral channel image, our calibration method extracts the peak light intensity from only one pixel and maps it back to the hypercube. The peak intensity is sensitive to both the microlens defocus and mirror reflectivity on the image mapper: the low mirror reflectivity reduces the total intensity of the LSF; on the other hand, the defocus broadens the LSF and, therefore, decreases the peak intensity. The nonuniform peak intensity introduces the stripe artifact. As an example, we show an image of a sharp edge under uniform illumination in Fig. 6(c).

To remove the stripe artifact, we developed three methods: Fourier domain filtering, deconvolution, and datacube normalization. Fourier domain filtering is a fast method to remove the stripes of specific frequencies. The image processing pipeline is shown in Fig. 7(a). The dominant striping frequency is associated with the periodic mirror facet blocks on the image mapper as illustrated in Fig. 7(b). For simplicity, we assume the image mapper consists of three mirror facet blocks, each containing two mirror facets. The incident field is then divided into six image slices. Slices 1, 3, and 5 are imaged by microlens 1, while slices

2, 4, and 6 are imaged by microlens 2. If the two microlenses have different levels of defocus, the intensities of these two groups of image slices would be different. Therefore, the dominant frequency of stripes in the reconstructed image is determined by the total number of mirror facets within one block. By applying notch filters in the Fourier domain, we can readily remove the correspondent stripes. However, this method cannot eliminate stripes at other frequencies, which are induced by the nonuniform reflectivity of the mirror facets and the nonuniform defocus of the microlenses. Additionally, the spatial frequency information at the same frequency range as that of the stripes is lost. As an example, we applied the Fourier domain filtering method to the image in Fig. 6(c), and the result is shown in Fig. 7(c). The stripes within the white area are effectively removed.

In contrast, deconvolution can suppress the stripe artifacts at all frequencies. By imaging an empty field illuminated by a monochromatic light source, we can simultaneously measure all combined LSFs in the raw image. This knowledge can then be used for deconvolution, compensating for the intensity variation of image slices. Additionally, deconvolution increases the spectral resolution by reducing the width of LSFs, which is also the impulse response of monochromatic light. However, because deconvolution is sensitive to noise, the performance of this method is highly dependent on the signal-to-noise ratio (SNR) of the image.

Lastly, datacube normalization captures an empty field under uniform illumination and constructs a reference hypercube. Dividing the real measurement by this reference cube yields the corrected results. The image processing pipeline is shown in Fig. 8. This method is essentially a simplified version of deconvolution, provided that the LSF of an image slice is a delta function. Therefore, the results are accurate only when the defocus is negligible.

To quantitively evaluate the stripe artifacts in the IMS, we define the striping factor as:

$$
c = \frac{\text{std}}{\text{mean}}\tag{6}
$$

for an image corresponding to an empty field under uniform illumination. Here mean is the mean intensity of the image and std is the standard deviation. The quotient c is zero if the image has no stripe artifact. In order to compare the performance of the proposed methods, we simulated an IMS with different levels of Gaussian noise and defocus. The Gaussian noise has a zero mean, and its standard deviation is normalized by the peak intensity in the image slice without any defocus or noise. Other simulation parameters and representative results are summarized in Table 1. We simulated a transparent object that alters the spectrum of the original illumination, and we generated three raw images: a reference image corresponding to an empty field uniformly illuminated by a light source with a rectangular spectrum centered at 532 nm and with a 50 nm bandwidth, an image that contains the LSFs under monochromatic illumination, and an image with the transparent object with a Gaussian transmission spectrum centered at 532 nm and with a 12 nm FWHM. Next we reconstructed a spectral channel image at 532 nm and corrected for the stripe artifact using the three methods (Fourier domain filter, deconvolution, and datacube normalization). We calculated the metric c before and after we applied the correction. As shown in Table 1, Fourier domain filtering becomes more effective along with an increase in defocus, as it

removes the dominant stripes introduced by the microlens array. Deconvolution outperforms other methods when the noise level is low, and datacube normalization fails when defocus is large.

Finally, we performed a tolerance analysis for the microlens array defocus based on the criterion defined. We use $c < 0.05$ as a criterion for acceptable image quality. Here we consider only the defocus and neglect other lens aberrations. To establish the relationship between the defocus of the microlens and the corresponding peak intensity in the image slice, we simulated our system using a wave propagation method, and the result is shown in Fig. 9. The wavelength used for simulation is 532 nm. All intensities were normalized by the peak intensity in the image slice without defocus. Note that the resultant curve is dependent on both the detector pixel size and the numerical aperture (NA) of the microlenses. If defocus exists, the Strehl ratio of the system decreases; however, because the signals measured by a pixel is ensquared energy, the peak intensity of the LSFs may not change as much as the Strehl ratio. On the other hand, for a fixed defocus, a smaller NA microlens has a larger depth of focus. Therefore, the image slice has a higher peak intensity compared to using a large-NA microlens. For our system, the detector pixel size is 7.4 μm, and the NA of the microlens array is 0.125.

Next we ran a Monte Carlo simulation. The parameters used for simulation are listed in Table 2. A group of microlenses with different levels of defocus was generated, and the probability density function (PDF) of the defocus of each microlens was assumed as a Gaussian distribution with a zero mean and a σ standard deviation. The number of simulated microlenses in a group is the same as the number of microlenses used in our system. The striping factor c was then calculated for this group of microlenses based on the peak intensities, which can be deduced from the result in Fig. 9. For each σ , we generated 1000 groups of microlenses and calculated corresponding striping factor c. The goal of the simulation is to find a threshold σ , so that 90% of microlens groups meet the acceptable image quality criterion (i.e., $c < 0.05$). The results are shown in Fig. 10. In our case, the threshold σ is 0.073 mm, indicating that 95% of microlenses should have a defocus of less than 2σ or 0.146 mm if operating the IMS at 532 nm.

6. CONCLUSION

In this paper, we developed a slit-scan method for constructing a lookup table for the IMS. By using an optical slit that is perpendicular to the mirror facets and scanning the FOV along one axis, we increased the calibration speed by 2 orders of magnitude. Also, we improved image reconstruction accuracy by 60% when imaging a resolution target as demonstrated in Section 4. Moreover, we proposed three methods to reduce the stripe artifacts, and we compared their pros and cons. We expect this work to lay the foundation for future IMS development.

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Fig. 2.

Edge alignment method. (a) Image slices alignment. (b) Distorted edge due to the variation in the magnification ratio of adjacent image slices. (c) Reconstructed image at 532 nm. The three insets show the zoomed view of the boxed areas. Area 3 shows the bars of group 2 element 4, indicating 5.66 lp/mm.

Fig. 4.

Reconstructed image of a USAF resolution target through slit-scan calibration. Compared to the result in Fig. 2(c), the standard deviation of the edge in Area 1 decreases from 1.47 pixels to 0.55 pixels.

Wavelength calibration. (a) Illustration of a dispersed image slice on the detector. Area 1 is the image slice, and Area 2 is the spectrum. (b) Dispersion curve of the prism.

Line spread functions and stripe artifact. (a) Two example line spread functions (LSFs). (b) Intensity plot along the orange dashed line in (a). (c) Image of a sharp edge with stripe artifacts.

Fig. 7.

Image processing pipeline of Fourier domain filtering. (a) Schematic pipeline. FFT, Fourier transform; iFFT, inverse Fourier transform; red square, notch filter. The frequencies within the notch filter are blocked. (b) The dominant frequency of stripes is determined by the total number of mirror facets within one block. (c) Fourier domain filtering result of the image shown in Fig. 6(c). The stripes within the white area are effectively removed.

Fig. 8.

Image processing pipeline of datacube normalization.

Fig. 10.

Percent of microlens groups that satisfy $c < 0.05$ versus standard deviation of the defocus of a microlens.

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Table 2.

Parameters for Monte Carlo Simulation

