# Efficient and Stable Planar n–i–p Sb<sub>2</sub>Se<sub>3</sub> Solar Cells Enabled by Oriented 1D Trigonal Selenium Structures

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Environmentally benign and potentially cost-effective Sb<sub>2</sub>Se<sub>3</sub> solar cells have drawn much attention by continuously achieving new efficiency records. This article reports a compatible strategy to enhance the efficiency of planar n-i-p Sb<sub>2</sub>Se<sub>2</sub> solar cells through Sb<sub>2</sub>Se<sub>2</sub> surface modification and an architecture with oriented 1D van der Waals material, trigonal selenium (t-Se). A seed layer assisted successive close spaced sublimation (CSS) is developed to fabricate highly crystalline Sb<sub>2</sub>Se<sub>3</sub> absorbers. It is found that the Sb<sub>2</sub>Se<sub>3</sub> absorber exhibits a Se-deficient surface and negative surface band bending. Reactive Se is innovatively introduced to compensate the surface Se deficiency and form an (101) oriented 1D t-Se interlayer. The p-type t-Se layer promotes a favored band alignment and band bending at the Sb<sub>2</sub>Se<sub>3</sub>/t-Se interface, and functionally works as a surface passivation and hole transport material, which significantly suppresses interface recombination and enhances carrier extraction efficiency. An efficiency of 7.45% is obtained in a planar Sb<sub>2</sub>Se<sub>3</sub> solar cell in superstrate n-i-p configuration, which is the highest efficiency for planar Sb<sub>2</sub>Se<sub>3</sub> solar cells prepared by CSS. The all-inorganic Sb<sub>2</sub>Se<sub>3</sub> solar cell with t-Se shows superb stability, retaining  $\approx$ 98% of the initial efficiency after 40 days storage in open air without encapsulation.

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### DOI: 10.1002/advs.202001013

# 1. Introduction

Quasi-1D(Q1D) V–VI antimonv chalcogenide-based solar cells have drawn intensive attention as a result of the meteoric rise in power conversion efficiencies over the past few years.[1-14] In particular, antimony selenide (Sb<sub>2</sub>Se<sub>3</sub>) has recently emerged as a highly attractive light-harvesting material arising from its extraordinary properties such as high absorption coefficient (> $10^5$  cm<sup>-1</sup>), suitable optical bandgap (1.1-1.3 eV), good carrier mobility ( $\approx 10 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ ), and simple single phase structure.<sup>[1-5,14,15]</sup> The earth-abundant low-toxic materials, together with scalable and cost-effective fabrication processes promise Sb<sub>2</sub>Se<sub>3</sub> solar cells a competitive technology for high-efficiency, low-cost photovoltaics. The unique Q1D structure of Sb<sub>2</sub>Se<sub>3</sub> enables intrinsically benign grain boundaries if the crystals are properly aligned.<sup>[2]</sup> To date, the highest reported power conversion efficiencies of both superstrate-planar and substrate-nanorod type Sb<sub>2</sub>Se<sub>3</sub> solar cells are 7.6% and 9.2%, respectively.<sup>[4,5]</sup>

Nevertheless, the low carrier density of Sb<sub>2</sub>Se<sub>3</sub> and high interface recombination loss represent the two most urgent obstacles to further improvement of the device efficiency.<sup>[16,17]</sup> The intrinsic defects in the low-symmetry Sb<sub>2</sub>Se<sub>3</sub> semiconductors are complicated and limit their p-type doping level.<sup>[18]</sup> To obtain high-quality Sb<sub>2</sub>Se<sub>3</sub> absorbers, several approaches have been developed, such as rapid thermal evaporation,<sup>[1,2]</sup> close spaced sublimation (CSS),<sup>[5,19,20]</sup> vapor transport deposition.<sup>[3,4,21-23]</sup> Early work mainly focused on improving the quality of the absorber (orientation control, deep defect passivation) and on engineering junction interfaces.<sup>[1-4]</sup> To further improve device performance, increasing attention has been paid to the trapping and carrier transport at the Sb<sub>2</sub>Se<sub>3</sub> rear surface and related interface. Tang and co-workers reported that, due to the electrical loss induced by rear surface/interface defects and the presence of a contact barrier, the additional deficit of open-circuit voltage of Sb<sub>2</sub>Se<sub>3</sub> solar cells can exceed 0.2 V.<sup>[16]</sup> Nevertheless, little attention has been paid to the properties of the Sb<sub>2</sub>Se<sub>3</sub> rear surface, despite their role in limiting device performance. In this regard, effective surface treatment or passivation methods are yet to be developed to further advance the device efficiency.

In addition to the lack of well-established surface modification methods, most of the existing superstrate Sb<sub>2</sub>Se<sub>3</sub> cell structures do not have a wide-bandgap barrier/passivating laver at the absorber interface to provide carrier confinement and reduce interface recombination velocity. Furthermore, considering the quasi-intrinsic nature of the Sb<sub>2</sub>Se<sub>3</sub> absorber layer, incorporation of a p-type interfacial layer to construct n-i-p configuration is a feasible way to overcome the bottleneck of the low doping concentration (1013 cm-3, as i-layer) of the absorber layer and improve the carrier extraction efficiency.<sup>[17,24]</sup> Exploration of efficient charge-extraction material or hole-transporting material (HTM) is the key to construct n-i-p Sb<sub>2</sub>Se<sub>3</sub> solar cells. Recently, organic materials (Spiro-OMeTAD, CZ-TA, PCDTBT) and inorganic materials (PbS quantum dot, CuSCN, NiO) have been investigated as HTMs to improve device performance (Figure S1, Supporting Information).<sup>[16,17,20,22,24,25]</sup> Especially, PbS quantum dots have been successfully used to prepare record-efficiency (7.62%) Sb<sub>2</sub>Se<sub>3</sub> solar cells.<sup>[4]</sup> However, the toxicity of Pb, the highmobility Cu impurity and the instability of organic materials limit their practical applications. An ideal interfacial structure should feature compatible material composition, matched energy band level, low interface defect density, good conformal coverage, and low-contact barrier. Meeting all these requirements is not easy but essential to realize efficient Sb<sub>2</sub>Se<sub>3</sub> devices.

In this work, we present a strategy that uses oriented trigonal selenium (t-Se) to modify the Sb<sub>2</sub>Se<sub>3</sub> surface and construct n-i-p Sb<sub>2</sub>Se<sub>3</sub> solar cells. Different from the ring structure of selenium clusters, trigonal selenium is a 1D van der Waals solid with helical atomic chains.<sup>[26]</sup> The unique properties of wide bandgap (1.8-2.0 eV), high work function (>5.0 eV), p-type conductivity ( $\approx 10^{14}$ cm<sup>-3</sup>), and high mobility ( $\approx 60 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ ) make t-Se a promising p-type interfacial layer for constructing efficient Sb<sub>2</sub>Se<sub>3</sub> solar cells.<sup>[26,27]</sup> We developed a seed layer assisted successive CSS process to fabricate high-crystalline Sb<sub>2</sub>Se<sub>3</sub> films. The chemical and electrical properties of the Sb<sub>2</sub>Se<sub>3</sub> surface were carefully analyzed. To reduce surface Se deficiency and the related downward band bending, reactive Se was introduced to compensate the surface Se deficiency and form an oriented 1D trigonal Se interlayer on the Sb<sub>2</sub>Se<sub>3</sub> surface. It is found that the prepared t-Se, which has (101) preferred orientation matches well with the crystal orientation of the Sb<sub>2</sub>Se<sub>3</sub> absorber. Further, the p-type t-Se functionally worked as surface passivation and p-type hole transport layer, facilitating the formation of n–i–p structured Sb<sub>2</sub>Se<sub>3</sub> solar cells. As a consequence, an efficiency of 7.45% was achieved, which is the highest efficiency ever reported for superstrate Sb<sub>2</sub>Se<sub>3</sub> solar cells prepared by CSS.

# 2. Results and Discussion

# 2.1. Fabrication of $Sb_2Se_3$ Films via Seed Layer Assisted Successive CSS Process

In this study,  $Sb_2Se_3$  films were prepared by the CSS method according to our previous work with some modifications.<sup>[19]</sup> A ceramic light visor was inserted between source and substrate to prevent overheating the graphite substrate holder by interactive light radiation (**Figure 1**a). This optimization promoted more independent control of substrate and source temperatures. The maximum steady-state temperature difference between the source and substrate was enlarged from  $\approx$ 120 to  $\approx$ 150 °C (the source temperature was 510 °C).

We developed a seed layer assisted successive CSS process to enhance preferred orientation and grain growth. As shown in Figure 1b, a thin Sb<sub>2</sub>Se<sub>3</sub> seed layer (≈10 nm) was first deposited onto the CdS surface, then the main Sb<sub>2</sub>Se<sub>3</sub> absorber layer (≈500 nm) was deposited on the seed layer. The two deposition processes were conducted in a successive manner without breaking the vacuum while keeping the substrate temperature higher than 270 °C. This successive high-temperature process without low-temperature stage can eliminate repeated nucleation on the grain surface and favor continuous grain growth.<sup>[28,29]</sup> Figure 1c,d shows the cross-sectional and top-view scanning electron microscopy (SEM) images of the Sb<sub>2</sub>Se<sub>3</sub> film. It can be seen that the Sb<sub>2</sub>Se<sub>3</sub> film has an excellent crystallinity with Sb<sub>2</sub>Se<sub>3</sub> grains as tall as the thickness of the film (through-thickness grains). The grains are densely packed and have a lateral size of 0.5-1 µm, which is about two times larger than that of the grains in films made by the conventional single-step CSS process.<sup>[19]</sup> The XRD patterns of the Sb<sub>2</sub>Se<sub>3</sub> films deposited by seed layer assisted successive CSS process and single-step CSS process are shown in Figure 1e. The two films had the same thickness and were measured under identical conditions. The XRD peak intensity of the Sb<sub>2</sub>Se<sub>3</sub> film prepared by seed layer assisted successive CSS is much stronger than that of the film by single-step CSS, indicating that seed layer assisted successive CSS process could strongly enhances the crystallinity. Furthermore, the film by seed layer assisted successive CSS shows a more (221) preferred orientation. As  $Sb_2Se_3$  is composed of 1D  $(Sb_4Se_6)_n$  ribbons stacked together via weak van der Waals interaction, strong (221) preferred orientation growth with through-thickness grains is highly favored for efficient carrier transport and benign grain boundaries.<sup>[2]</sup>

## 2.2. Surface Se Deficiency and Band Bending

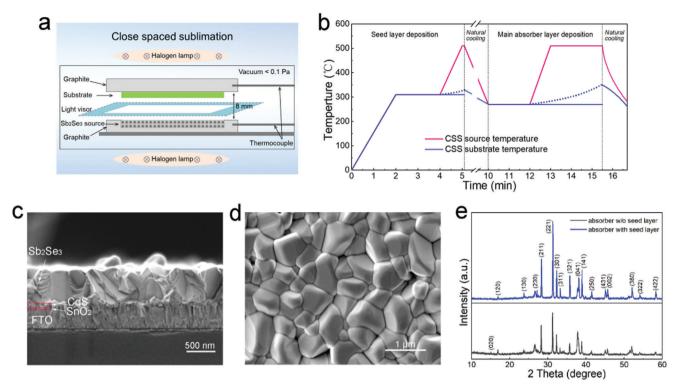
In addition to structural properties, the chemical composition was characterized by high-resolution X-ray photoelectron spectroscopy (XPS). Figure 2a shows the Sb 3d spectrum of the asdeposited Sb<sub>2</sub>Se<sub>2</sub> film. It was observed that the shoulder peak with lower binding energy at 528.7 eV was identified, which is representative of the reduced Sb elemental state or Se-deficient component at the Sb<sub>2</sub>Se<sub>3</sub> surface.<sup>[30,31]</sup> It has been reported that Sb<sub>2</sub>Se<sub>3</sub> is slightly decomposed during high-temperature process, leading to the formation of Se-deficient Sb<sub>2</sub>Se<sub>3-x</sub> because of the high vapor pressure of Se.<sup>[32-34]</sup> We further conducted XPS depthprofile measurements to investigate the chemical distribution of Sb<sub>2</sub>Se<sub>3</sub> absorber. Figure S2 of the Supporting Information shows XPS scans of the Sb 3d at various depths. The intensity of the shoulder peak decreased and disappeared with increasing depth, indicating that the detectable Se deficiency only exists within several nanometers from the surface. Because of the high vapor pressure of Se, the high-temperature deposited Sb<sub>2</sub>Se<sub>3</sub> film is prone to deviate from the stoichiometric ratio.[18,32-34] Theoretically, Se-related defects produce a series of intrinsic doping levels in the bandgap and make Sb<sub>2</sub>Se<sub>3</sub> p-type under the Se-rich condition and n-type under the Se-poor condition.<sup>[18,32]</sup> We assume



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**Figure 1.** Fabrication and characterization of Sb<sub>2</sub>Se<sub>3</sub> layer. a) Schematic diagram of the close spaced sublimation (CSS) apparatus for Sb<sub>2</sub>Se<sub>3</sub> film deposition. b) The temperature profile for seed layer assisted successive CSS process. c) Cross-sectional SEM image and d) top-view SEM image of Sb<sub>2</sub>Se<sub>3</sub> film. e) XRD spectra of the Sb<sub>2</sub>Se<sub>3</sub> films deposited by seed layer assisted successive CSS process and single CSS process.

that the surface Se deficiency is related to the change of chemical environment during the whole CSS process. When the source temperature is high (510 °C), due to the higher vapor pressure of Se, source vapor provides a Se-rich environment around the substrate, leading to a less Se-deficient condition. A sufficiently high Se content is critical to guarantee intrinsic p-type conductivity in Sb<sub>2</sub>Se<sub>3</sub> films. During the cooling process, however, the Se vapor around the hot substrate tends to diminish with decreasing source temperature. Under such circumstances, the deposited Sb<sub>2</sub>Se<sub>3</sub> film may decompose or re-evaporate, which lead to surface Se deficiency because of the high vapor pressure of Se.

In order to tackle the Se deficiency at the rear surface of Sb<sub>2</sub>Se<sub>3</sub>, we deposited an ultrathin Se layer (2 nm, monitored by crystal oscillator) on the Sb<sub>2</sub>Se<sub>3</sub> surface by high-vacuum thermal evaporation. The reactivity of Se on the Sb<sub>2</sub>Se<sub>3</sub> surface was regulated by altering the substrate temperature. It is found that, when the substrate temperature was lower than 80 °C, the device performance changes were negligible. When the substrate temperature was elevated to 100 °C, enhanced spectral response at long wavelengths and enhanced device performance could be observed (Figure S3, Supporting Information). However, when the substrate temperature was higher than 110 °C, it was difficult to deposit Se on the Sb<sub>2</sub>Se<sub>3</sub> surface. To understand the mechanism of the performance enhancement, the Se-modified Sb<sub>2</sub>Se<sub>3</sub> surface was characterized. In order to characterize the change in chemical composition with respect to the pure Sb<sub>2</sub>Se<sub>3</sub> surface, the samples with Se coating were cleaned to make sure that no superfluous Se residue was present on the surface (Figure S4, Supporting Information). As shown in Figure 2d, the shoulder peak at 528.7 eV

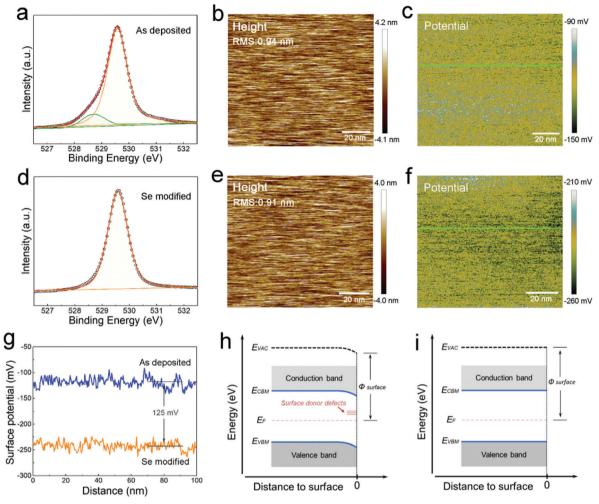
disappeared, and the Sb 3d spectrum shows only one symmetric peak at 529.6 eV corresponding to Sb<sup>3+</sup> in Sb<sub>2</sub>Se<sub>3</sub>.<sup>[31,35]</sup> These observations provide evidence that the Se modification process is fairly effective in compensating surface Se deficiency of the Sb<sub>2</sub>Se<sub>3</sub> absorber.

To further characterize the surface electrical properties of Sb<sub>2</sub>Se<sub>3</sub> films, Kelvin probe force microscopy (KPFM) was employed to map the surface morphology and potential of the asdeposited Sb<sub>2</sub>Se<sub>3</sub> surface and the Se-modified Sb<sub>2</sub>Se<sub>3</sub> surface. To calibrate the potential value of the probe, an Au sample was measured together with the Sb<sub>2</sub>Se<sub>3</sub> samples (Figure S5, Supporting Information). The atomic force microscopy (AFM) images (Figure 2b,e) demonstrate similar height morphology with rootmean-square (RMS) roughness in the range of 0.91-0.94 nm. Figure 2c,f,g shows the KPFM images and the corresponding line profiles of the surface potentials. It is observed that the average surface potential of the Sb<sub>2</sub>Se<sub>3</sub> film decreased by 125 mV after Se modification. The surface potential difference relates quantitatively to the difference in surface work functions  $(\Phi)$ .<sup>[36,37]</sup> The surface potential measurements indicate a work function increase of 125 mV for the Se-modified Sb<sub>2</sub>Se<sub>3</sub> surface as compared to the as-deposited Sb<sub>2</sub>Se<sub>3</sub> surface. This is a considerable change for a 1.2 eV bandgap semiconductor. The reason for the work function variation of the Sb<sub>2</sub>Se<sub>3</sub> surface with different chemical composition can be explained by means of the schematic band diagram in Figure 2h,i. In general, the work function of a semiconductor can be written as  $\Phi = E^a + \Delta E_{fn}$ , where  $E^a$  is the electron affinity and  $\Delta E_{\rm fn}$  is the Fermi level position referenced to the bottom of the conduction band.<sup>[36]</sup> In case of the as-deposited

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**Figure 2.** Surface chemical and electrical analysis. a) XPS peak of Sb  $3d_{5/2}$  obtained from the surface of as-deposited  $Sb_2Se_3$ . b) AFM image of the as-deposited  $Sb_2Se_3$  surface, and c) is the corresponding KPFM image obtained by scanning the same area. d) XPS peak of Sb  $3d_{5/2}$  obtained from the surface of Se-modified  $Sb_2Se_3$ . e) The AFM image of the Se-modified  $Sb_2Se_3$  surface, and f) the corresponding KPFM image of the Se-modified  $Sb_2Se_3$  surface, and f) the corresponding KPFM image obtained by scanning the same area. g) Potential profiles obtained along the two lines indicated in (c) and (f). Schematic diagrams of surface band bending of h) an as-deposited  $Sb_2Se_3$  film and i) a Se-modified  $Sb_2Se_3$  film.

Sb<sub>2</sub>Se<sub>3</sub> surface, the Se-deficiency related defects, such as Se vacancies, usually produce deep-donor levels in the bandgap which may act as carrier recombination centers.<sup>[18,32]</sup> The deep-donor states near the conduction band uplift the Fermi energy, resulting in a downward band bending, smaller  $\Delta E_{\rm fn}$ , and decreased  $\Phi$ at the surface (Figure 2h).<sup>[36–38]</sup> For the Se-modified Sb<sub>2</sub>Se<sub>3</sub> surface, due to the fact that the Se deficiency was compensated, the position of the Fermi energy at the surface is consistent with that of bulk Sb<sub>2</sub>Se<sub>3</sub>, leading to a surface under more flat band conditions (Figure 2i).

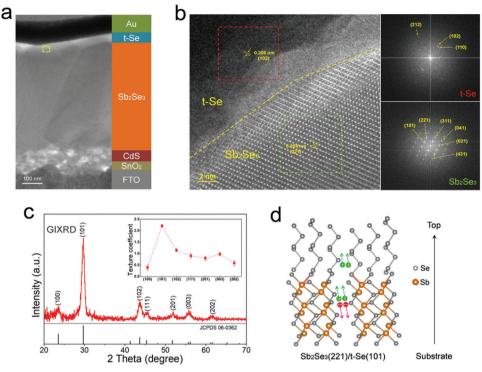
It is worth pointing out that surface band bending or even band inversion has been found in several kinds of semiconductors.<sup>[37–40]</sup> Experimental results have indicated that, in the region of p–n junction interface, an inverted surface band bending in the absorber layer is beneficial for device performance due to the reduction of the recombination rate by shifting the electrical junction away from the defective physical junction interface.<sup>[39,40]</sup> However, at the back contact interface, an inverted surface band bending in the absorber layer is detrimental to the device performance.<sup>[41]</sup> On the one hand, the downward valence band bending forms a hole barrier, which impedes holes transport to the back electrode, and on the other hand, due to the negative conduction band bending, a portion of the photogenerated electrons would diffuse backward and recombine with holes at the back surface. Hence, modification of surface Se-deficiency is essential not only to passivate the Se-deficiency related deep-level defects but also to construct a favorable band structure for carrier transport at the surface/interface.

## 2.3. Interfacial Structure and Orientation

Considering the compatible chemical and electrical properties of Se with  $Sb_2Se_3$ , Se layers with different thicknesses were deposited as interfacial layer between  $Sb_2Se_3$  and Au electrode. Figure S6 of the Supporting Information shows the device



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**Figure 3.** Interface structure and orientation. a) Cross-sectional TEM image of the  $Sb_2Se_3$  solar cell. b) HRTEM images of  $Sb_2Se_3/t$ -Se interface and the Fourier transforms of corresponding lattice fringe. c) XRD pattern of t-Se layer on the  $Sb_2Se_3$  substrate using GIXRD mode. d) Atomistic model of [101]-oriented 1D t-Se on the (221) plane of  $Sb_2Se_3$ .

performance of the Sb<sub>2</sub>Se<sub>3</sub> solar cells with Se layer thicknesses of 0, 10, 15, 20, and 30 nm. All devices with a Se layer modification show performance improvements. The solar cell with a 15 nm Se layer shows the highest power conversion efficiency with all device parameters improved. To analyze the interface structure, high-resolution transmission electron microscopy (HRTEM) was conducted. The sample for cross-sectional imaging was prepared by focused ion beam and a TEM image of the device is shown in Figure 3a. Figure 3b shows the HRTEM image of Sb<sub>2</sub>Se<sub>3</sub> interface taken from the dashed rectangle of Figure 3a. As expected, the Sb<sub>2</sub>Se<sub>3</sub> displays an orthorhombic phase, judging from multiple interplane distances. On the other hand, the Se layer displays a trigonal phase with multiple typical interplane distances. The Fourier transforms of the corresponding lattice fringe of the orthorhombic Sb<sub>2</sub>Se<sub>3</sub> and trigonal Se, and the major diffraction spots are indexed. Furthermore, grazing incidence X-ray diffraction (GIXRD) with incidence angle of 0.5° was performed to confirm the crystal phases and lattice orientations of the Se layer on the Sb<sub>2</sub>Se<sub>3</sub> surface (Figure 3c). All diffraction peaks are identified and indexed to the 1D trigonal Se (JCPDS 06-0362). The t-Se layer shows preferential growth with an intensity of the (101) peak much higher than that of all other peaks. The orientation parameter of a crystalline plane (hkl) was quantitatively described by the texture coefficient  $(TC_{hkl})$ .<sup>[42]</sup> As can be seen in the inset of Figure 3c, TC(101) has a large value of 2.2, and all the other six TC values are close to or smaller than 1, indicating a strong (101) preferred orientation growth. Due to the distinct Q1D crystal characteristics, the disordered orientations at the Sb<sub>2</sub>Se<sub>3</sub>/t-Se interface easily introduce trap centers.<sup>[1,2,8]</sup> Especially for 1D

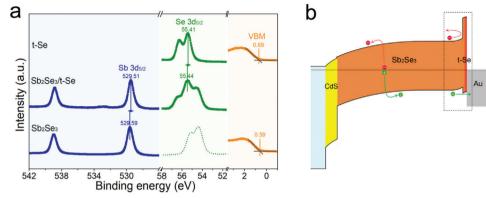
t-Se, it has been shown that the conductivity of t-Se crystals is  $\approx 10$  times higher when measured parallel to the Se chains than in the perpendicular direction.<sup>[43]</sup> The (101) preferred orientation of t-Se indicates that the Se chains are perpendicular to the substrate and almost parallel to the Sb<sub>2</sub>Se<sub>3</sub> ribbons with (221) preferred orientation (Figure 3d).<sup>[44]</sup> The ordered and matched orientations at the Sb<sub>2</sub>Se<sub>3</sub>/t-Se interface can reduce the possible interfacial defects and form a low-resistance path for efficient carrier transport.

#### 2.4. Interface Band Alignment

To provide insight into the impact of t-Se layer on the interfacial band structure, XPS measurements including core-level spectroscopy and valence band (VB) spectroscopy were conducted at the Sb<sub>2</sub>Se<sub>3</sub>/t-Se interface. The quantitative valence band offset (VBO) can be obtained by VBO =  $E_{VB}^b - E_{VB}^a + V_{bb}$ , where  $E_{VB}^b$  and  $E_{VB}^a$  are the energy positions of the valence band edges of bulk Sb<sub>2</sub>Se<sub>3</sub> and bulk t-Se, respectively, and  $V_{bb}$  is the band bending at the interface.<sup>[45]</sup> The  $V_{bb}$  can be estimated from the core level energy shift between the bulk area and the interface, and can be calculated by  $V_{bb} = (E_{cl}^a - E_{cl}^a(i)) - (E_{cl}^b(i) - E_{cl}^b)$ , where  $E_{cl}^b$  and  $E_{cl}^a$  are the core level energies of Sb 3d<sub>5/2</sub> and Se 3d<sub>5/2</sub> in the bulk Sb<sub>2</sub>Se<sub>3</sub> and bulk t-Se, respectively, and  $E_{cl}^b(i)$  and  $E_{cl}^a(i)$  are the core level energies of Sb 3d<sub>5/2</sub> at the Sb<sub>2</sub>Se<sub>3</sub>/t-Se interface, respectively. Note that the values in the equation are calculated under the assumption that binding energy below Fermi edge is negative. The XPS curves and the binding



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**Figure 4.** Band alignment at the  $Sb_2Se_3/t$ -Se interface. a) The XPS spectra of Sb 3d, Se 3d and valence band in the bulk t-Se,  $Sb_2Se_3/t$ -Se interface, and bulk  $Sb_2Se_3$ . b) Schematic energy band diagram of the  $Sb_2Se_3$  device drawn by XPS measurements.

energies are shown in Figure 4a. Considering that the Se element is included in both Sb<sub>2</sub>Se<sub>3</sub> and t-Se, we chose higher peak position ( $\approx$ 55.4 eV) as the core level energy of Se 3d<sub>5/2</sub>, which corresponds to the binding energy of Se  $3d_{5/2}$  in t-Se (Figure S4, Supporting Information). The difference of valence band maxima (VBM) of Sb<sub>2</sub>Se<sub>3</sub> layer and t-Se layer is 0.10 eV. The total  $V_{\rm bb}$ and VBO at the Sb<sub>2</sub>Se<sub>3</sub>/t-Se interface are 0.11 and 0.21 eV, respectively. To confirm the valence band alignment and band bending at the interface, ultraviolet photoelectric spectroscopy (UPS) was used to examine the energy levels (Figure S7, Supporting Information). The calculated VBM and work functions are 0.56 and 4.55 eV for the Sb<sub>2</sub>Se<sub>3</sub> film, and 0.70 and 5.18 eV for the t-Se film, respectively. The VBM values are in good agreement with those shown in Figure 4a. Besides, the higher work function of t-Se favors an interface electrical field directed from Sb<sub>2</sub>Se<sub>3</sub> to t-Se, which corresponds to an upward band bending at the interface of Sb<sub>2</sub>Se<sub>3</sub>/t-Se. Given that the bandgaps of bulk Sb<sub>2</sub>Se<sub>3</sub> and t-Se remain unchanged, the conduction band offset (CBO) can be calculated by CBO =  $(E_{gap}^b - E_{gap}^a)$  + VBO. The bandgaps of the Sb<sub>2</sub>Se<sub>3</sub> and t-Se were measured to be 1.20 and 1.98 eV, respectively (Figure S8, Supporting Information), and the CBO was calculated to be -0.57 eV.

Figure 4b shows the band diagram of the solar cell with a layer stack of fluorine tin oxide (FTO)/CdS/Sb<sub>2</sub>Se<sub>3</sub>/t-Se/Au. The Sb<sub>2</sub>Se<sub>3</sub>/t-Se forms a favorable band alignment at the contact interface. The upward band bending at the Sb<sub>2</sub>Se<sub>3</sub> surface along with the close VBM levels between the Sb<sub>2</sub>Se<sub>3</sub> and t-Se will facilitate the efficient extraction of holes from Sb<sub>2</sub>Se<sub>3</sub> to electrode. At the same time, due to the higher conduction band maximum, the t-Se layer acts as a back field reflector, which blocks photogenerated electrons from injecting into the back electrode and consequently reduces carrier recombination at the back contact. Furthermore, the p-type conductivity ( $\approx 10^{14}$  cm<sup>-3</sup>, about one order of magnitude higher than Sb<sub>2</sub>Se<sub>3</sub>) and the preferred orientation of 1D t-Se favors hole injection and transport in the t-Se layer. Based on the discussions above, it can be concluded that the t-Se layer functionally works as a p-type hole transport material in the ni-p device, by passivating the rear surface, facilitating both hole extraction and electron reflection. The matching interfacial properties enable to effectively suppress the interface recombination and reduce the series resistance of the devices, which is expected to deliver a significantly enhanced photovoltaic performance.

Table 1. Photovoltaic performance parameters for the  $Sb_2Se_3$  solar cells with and without t-Se layer.

Device	$J_{\rm SC}$ [mA cm <sup>-2</sup> ]	V <sub>OC</sub> [V]	FF [%]	PCE [%]	$R_{\rm S} ~[\Omega~{\rm cm}^2]$	$R_{\rm SH} ~[\Omega~{ m cm}^2]$
w/o t-Se	25.4	0.380	55.1	5.32	3.1	158
with t-Se	28.9	0.413	62.3	7.45	2.2	249

## 2.5. Device Performance and Analysis

As mentioned earlier, we examined the effect of t-Se growth temperatures and layer thicknesses. Solar cells with 15 nm t-Se layer deposited at 100 °C gave the best and most reproducible performance. Figure 5a shows the statistic device performance of Sb<sub>2</sub>Se<sub>3</sub> solar cells with and without 15 nm thick t-Se layer. It is obvious that, compared to the control device, solar cells modified with t-Se layers exhibit noticeably higher overall performance in terms of all photovoltaic parameters. The average power conversion efficiency (PCE) dramatically increases from ≈5.2% to over 7%. Figure 5b,c presents the current density-voltage (J-V) and external quantum efficiency (EQE) curves of the best devices with and without t-Se measured under standard test conditions. The detailed device parameters are summarized in Table 1. The parasitic resistances  $R_{\rm S}$  and  $R_{\rm SH}$  were extracted from the J–V curves using Site's method.<sup>[46]</sup> Compared with the control devices, the open-circuit voltage ( $V_{OC}$ ), short-circuit current density ( $J_{SC}$ ), and fill factor (FF) of the device with t-Se layer are enhanced by 8.7%, 13.8%, and 13.1%, respectively, resulting in a 40% improvement in PCE, from 5.32% to 7.45%. To the best of our knowledge, the obtained efficiency of 7.45% in this study is the highest efficiency value among all reported superstrate Sb<sub>2</sub>Se<sub>3</sub> solar cells fabricated by CSS, and very close to the record efficiency (7.62%) of the superstrate Sb<sub>2</sub>Se<sub>3</sub> solar cells, which is employed with toxic PbS hole transport layer.<sup>[4]</sup> In addition, no hysteresis between forward and reverse scans was observed in the Sb<sub>2</sub>Se<sub>3</sub> solar cell with t-Se layer (Figure S9, Supporting Information). The EQE characterizations were carried out to verify the reliability of the J-V measurement. The integrated  $J_{\rm SC}$  values from EQE curves are 25.2 and 28.6 mA cm<sup>-2</sup> for Sb<sub>2</sub>Se<sub>3</sub> solar cells without and with t-Se layer, respectively, which are in good agreement with the values derived from the *J*–*V* measurements (Figure 5c).





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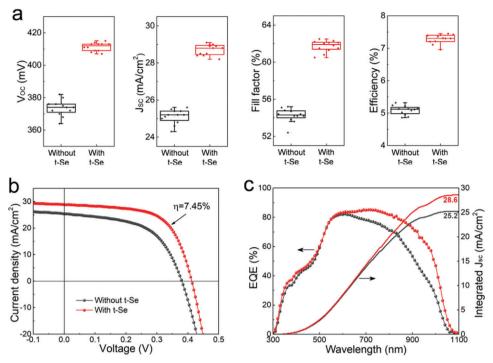


Figure 5. Device performance. a) Performance statistics of  $Sb_2Se_3$  devices with and without t-Se layer, 12  $Sb_2Se_3$  solar cells were fabricated per experimental condition. b) The *J*–V curves and c) EQE curves of the best performing  $Sb_2Se_3$  devices with and without t-Se layer.

In order to elucidate the origin of the performance improvement, specific device characterizations were performed. First, it is obviously seen from Figure 5c that the higher  $I_{SC}$  of solar cells with t-Se layer is mainly results from spectral response enhancement in the long-wavelength region (>600 nm). The EQE at the wavelength of 1000 nm is still about 50%, which is higher than most previously reported values (≈30%).[16,17,20,24] The EQE improvement in the long wavelength region (600-1100 nm) contributes to an integrated photocurrent increase of  $\approx 3$  mA cm<sup>-2</sup>. Due to the photons at the short-wavelength region (<625 nm) have smaller absorption depth, the t-Se layer does not absorb the photons passing through the Sb<sub>2</sub>Se<sub>3</sub> absorber and contribute to the EQE enhancement (Figure S10, Supporting Information). It is known that the photons with long wavelength are generally absorbed at a deep position in the absorber layer, the increased EQE response at long wavelengths is mainly due to the sufficient generation and collection of carriers close to the deep or back side.<sup>[46]</sup> Considering that the limited thickness of the Sb<sub>2</sub>Se<sub>3</sub> layer ( $\approx$ 500 nm), the transport and recombination of carriers generated at deeply in the absorber layer can be easily affected by the rear surface or interface properties. With the introduction of t-Se layer, the upward band bending in the surface region promotes holes drifting as well as injection via the back field. The passivated deep-donor surface defects and the blocking effect of the t-Se layer for electrons reduce the carrier recombination at the back interface. In addition, the depletion width  $(x_d)$  of the main junction is found to be increased from 223 to 238 nm for the cell with the t-Se layer, which strengthens the charge carrier separation and transport deeply in the Sb<sub>2</sub>Se<sub>3</sub> absorber layer (Figure 6d). Overall, the favored band structure and suppressed recombination are responsible for the EQE

enhancement at long wavelengths and the corresponding  $J_{\rm SC}$  increase.

The carrier transport behavior was analyzed by temperaturedependent  $V_{\rm OC}$  and light intensity-dependent  $J_{\rm SC}$  measurements. As shown in Figure 6a, the plot of  $V_{\rm OC}$  versus temperature provides information about the activation energy  $(E_A)$  of the main recombination mechanism.<sup>[47]</sup> The  $E_A$  linearly extrapolated from the curve of  $V_{\rm OC}$  versus temperature at 0 K shows a high value of 1.08 eV for the solar cells with a t-Se layer, which is close to the bandgap of the Sb<sub>2</sub>Se<sub>3</sub> absorber (1.20 eV), suggesting a dominant Schottky-Read-Hall recombination process in the space charge region. On the other hand, for the device without t-Se, the  $E_{A}$  is calculated to be 0.99 eV, which is about 0.21 eV lower than the bandgap, indicating additional interface recombination and thus the  $V_{\rm OC}$  is also limited as well.<sup>[47]</sup> Figure 6b shows a linear dependence of the  $J_{SC}$  on the light intensity for both types of solar cells, in which both slopes of the lines are close to 1, indicating that both devices suffer negligible from space-charge limited photocurrent. The relatively higher value of the slope of the device with t-Se layer indicates that a better back contact leads to a much easier charge transport.[16]

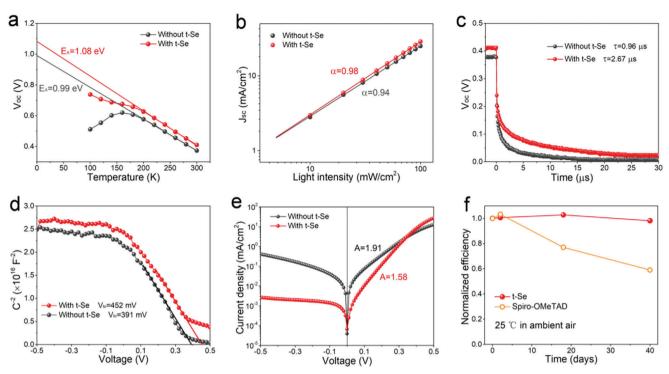
The carrier recombination dynamics was quantitatively analyzed by transient  $V_{\rm OC}$  decay measurements. In a typical solar cell under open-circuit conditions, the decay time of  $V_{\rm OC}$  upon removal of light is representative of the recombination time constant of the solar cell.<sup>[5,17]</sup> Figure 6c displays the  $V_{\rm OC}$  decay curves. The calculated decay times are determined to be 2.67 and 0.96 µs for Sb<sub>2</sub>Se<sub>3</sub> solar cells with and without t-Se layer, respectively. The longer  $V_{\rm OC}$  decay time and the better retardation of carrier recombination confirm the reduction of interface defects and the related electronic recombination loss for the device with t-Se layer.

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**Figure 6.** Device analysis. a) Temperature-dependent  $V_{OC}$ , and b)  $J_{SC}$  versus light intensity, and c)  $V_{OC}$  decay curves, and d)  $C^{-2}-V$  curves, and e) J-V curves under dark conditions of Sb<sub>2</sub>Se<sub>3</sub> devices with and without t-Se layer, respectively. f) Long-term stability of Sb<sub>2</sub>Se<sub>3</sub> devices with t-Se layer and with Spiro-OMeTAD layer.

Mott–Schottky ( $C^{-2}-V$ ) measurements were conducted to understand the origin of the  $V_{\rm OC}$  improvement from the built-in potential  $(V_{\rm bi})$  variation. As shown in Figure 6d, compared to the device without t-Se layer, the  $V_{\rm bi}$  of the device with t-Se layer increased from 0.391 to 0.452 V. As  $V_{\rm bi}$  is positively correlated to  $V_{\rm OC}$ , the increased  $V_{\rm bi}$  theoretically contributes to the improvement of V<sub>OC</sub>.<sup>[17]</sup> Considering the identical CdS/Sb<sub>2</sub>Se<sub>3</sub> junction structure and the same slopes of the linear regions, the increased  $V_{\rm bi}$  can be attributed mainly to the band bending at the Sb<sub>2</sub>Se<sub>3</sub>/t-Se interface. Further, the diode properties are determined from the corresponding dark *I*–*V* curves of the two cells, as shown in Figure 6e. The device with t-Se layer demonstrates a better diode rectification, which is consistent with the increased shunt resistance  $(R_{SH})$  and the decreased series resistance  $(R_S)$  (Table 1). The optimized parasitic resistances contribute to the higher FF. The reduced leakage current can be attributed to the reduced recombination current and the possible shunt paths. The smaller  $R_{\rm s}$ of the t-Se based device is due to the enhanced carrier extraction efficiency and the reduced Schottky barrier at the back interface (Figure S11, Supporting Information). The diode ideality factors (A) are calculated to be 1.58 and 1.91 for the cells with and without t-Se layer. The variation in A between 1 and 2 depends on the energies of the deep defects that act as dominant trap states or the tunneling contribution.<sup>[46]</sup> The lower value of A affirms the reduced trap-assisted recombination in the device with t-Se layer.

In addition, we studied the long-term stability of the Sb<sub>2</sub>Se<sub>3</sub> solar cell, and compared this to the cell with the commonly used organic Spiro-OMeTAD HTL. The two devices have the same structure except the HTL materials. The details of the Sb<sub>2</sub>Se<sub>3</sub> solar cell with Spiro-OMeTAD HTL used in the present stability studies can be found in our previous work.<sup>[48]</sup> Both devices were stored and measured in ambient environment without encapsulation. As shown in Figure 6f, the normalized efficiency of the Sb<sub>2</sub>Se<sub>3</sub> solar cell with t-Se retains  $\approx$ 98% of its initial value after storage for 40 days. By contrast, the efficiency shows about 40% decline for the device with the Spiro-OMeTAD HTL after the same storage duration. These results suggest that the chemically compatible t-Se based all-inorganic structure is an effective approach to fabricate highly efficient and stable Sb<sub>2</sub>Se<sub>3</sub> solar cells.

## 3. Conclusion

In summary, we have demonstrated a compatible approach to modify the Sb<sub>2</sub>Se<sub>3</sub> surface and construct efficient n-i-p Sb<sub>2</sub>Se<sub>3</sub> solar cells. A seed layer assisted successive CSS process was developed to fabricate high-crystalline Sb<sub>2</sub>Se<sub>3</sub> film. The surface Se deficiency and negative surface band bending were discovered. Reactive Se was innovatively introduced to compensate surface Se deficiency and form an oriented 1D t-Se interlayer on the Sb<sub>2</sub>Se<sub>3</sub> surface. It is found that the t-Se layer has a matched (101) preferred orientation, and functionally works as surface passivation and p-type hole transport material. Benefiting from the suppressed recombination and enhanced band structure, the device performance was significantly improved. An efficiency as high as 7.45% was achieved, which is the highest efficiency reported for superstrate Sb<sub>2</sub>Se<sub>3</sub> solar cells prepared by CSS. Besides, the allinorganic Sb<sub>2</sub>Se<sub>3</sub> solar cell shows excellent device stability. This work provides a compatible component and device architecture design for highly efficient and stable Sb<sub>2</sub>Se<sub>3</sub> solar cells.

# 4. Experimental Section

Solar Cell Fabrication: FTO-coated glass with sheet resistance of 8  $\Omega$  $sq^{-1}$  was cleaned by sonication with by acetone, ethanol, and deionized water for 10 min, respectively. A resistive 30 nm SnO<sub>2</sub> buffer layer was deposited on the FTO surface by the RF magnetron sputtering technique at the substrate temperature of 200 °C. CdS layers with a thickness of  $\approx$ 60 nm were prepared on glass/FTO/SnO<sub>2</sub> substrates by chemical bath deposition at 85 °C from a solution composed of deionized water, cadmium acetate, ammonium acetate, and thiourea. The substrate with deposited CdS layer was then immersed into saturated CdCl<sub>2</sub> methanol solution to dipcoat a thin CdCl<sub>2</sub> layer, and heat treated in a furnace in air atmosphere at a temperature of 380 °C for 10 min. The Sb<sub>2</sub>Se<sub>3</sub> absorber layers were deposited by a seed layer assisted successive CSS process in a modified CSS facility, as sketched in Figure 1a,b. The source material is commercial Sb<sub>2</sub>Se<sub>3</sub> powder with 99.999% purity. The deposition of Sb<sub>2</sub>Se<sub>3</sub> layer includes two processes: seed layer deposition and main absorber layer deposition. The CSS process was conducted at a vacuum pressure of  $\approx$ 0.1 Pa. There is no inert gas or reactive gas in the chamber. First, seed layer deposition: set source and substrate temperatures at 310 °C for 120 s, raise source temperature to 510 °C in 60 s, keep the substrate temperature at 310 °C and actually the temperature increases to  $\approx$ 330 °C after 60 s, keep the source temperature at 510 °C for 10 s; then the source and substrate temperatures naturally cooling to 270 °C and begin the main absorber layer deposition: keep the source and substrate temperatures at 270 °C for 120 s, raise source temperature to 510 °C in 60 s, and then keep at 510 °C for 150 s, the substrate temperature is set at 270 °C the whole process and actually increases to  $\approx$ 350 °C when the process is over, then turn off the power to stop the deposition and finally take the film out when it is cooled to 150 °C. After that, the Sb<sub>2</sub>Se<sub>3</sub> film was rinsed with carbon disulfide (CS<sub>2</sub>, Aladdin, AR) for 1 min and then rinsed by deionized water. The Se films were deposited by thermal evaporation with 99.999% pure Se particles. The chamber was pumped to a pressure of less than  $5\times10^{-4}$  Pa before evaporation. The substrates were rear heated at an average rate of 10 °C min<sup>-1</sup> from room temperature to the designed temperature. The deposition rate was monitored by crystal oscillator. The Deposition rate was kept at  $\approx 2$  Å s<sup>-1</sup>. Subsequently, gold back contacts ( $\approx 100$  nm) were deposited by thermal evaporation in a vacuum chamber and the active area of individual device was defined as 0.126 cm<sup>2</sup>.

Sample Preparation: Samples for Sb<sub>2</sub>Se<sub>3</sub> surface measurements: the as-deposited Sb<sub>2</sub>Se<sub>3</sub> after rinsed by CS<sub>2</sub> solution for 1 min was defined as as-deposited Sb<sub>2</sub>Se<sub>3</sub>; the Sb<sub>2</sub>Se<sub>3</sub>/t-Se (2 nm) was first re-evaporated in an RTP furnace for 60 s at 150 °C under a vacuum pressure below 0.1 Pa, then immersed in CS<sub>2</sub> solution for 5 min and rinsed by deionized water, the obtained sample with pure Sb<sub>2</sub>Se<sub>3</sub> surface without Se residue was defined as Se-modified Sb<sub>2</sub>Se<sub>3</sub>. Samples for interface band alignment measurements: a 50 nm thick t-Se film deposited on top of the Sb<sub>2</sub>Se<sub>3</sub> substrate was defined as Sb<sub>2</sub>Se<sub>3</sub>/t-Se interface; the Sb<sub>2</sub>Se<sub>3</sub>/t-Se(3 nm) was re-evaporated at 150 °C and immersed in CS<sub>2</sub> solution, the sample with clean Sb<sub>2</sub>Se<sub>3</sub> surface without Se residue was defined as Sufface in CS<sub>2</sub> solution, the sample with clean Sb<sub>2</sub>Se<sub>3</sub> surface without Se residual was defined as bulk Sb<sub>2</sub>Se<sub>3</sub>.

Material Characterization: The film morphologies were characterized by SEM (FEI Apreo LoVac). The optical properties were measured by UVvis spectrophotometer (Agilent Tech, Cary 5000). XRD and GIXRD: XRD measurements were performed by a Bruker D8 Advance X-ray diffractometer with Cu K $\alpha$  as the radiation source. GIXRD pattern was performed with incidence angle of 0.5° and five consecutive measurements were collected and averaged into single spectra. TEM and HRTEM: Samples for TEM were prepared by ablating the device using a focused ion beam microscope. Pt layers were deposited on top of the device for protection. TEM and HRTEM images were taken on a JEOL JEM 3200FS TEM at 300 kV. AFM and KPFM: The AFM and KPFM measurements were performed on a Bruker Dimension Icon with Nanoscope V controller. All the KPFM images were collected in lift mode by using SCM–PIT probes with a spring constant of 2.8 N  $m^{-1}$ and a resonant frequency of 78.3 kHz. During all KPFM measurements, the optimal lift height is 30 nm. UPS and XPS: UPS and XPS characterizations were obtained with a Thermo Fisher Scientific K-Alpha<sup>+</sup>, using the

HeI (21.22 eV) emission line and Al K $\alpha$  radiation (1486.6 eV). X-ray photoelectron depth profiling spectra were recorded by using Ar<sup>+</sup> sputtering gun operated at 3 keV, rastering a 2.0 × 2.0 mm<sup>2</sup> area with 5 s postetch delays.

Device Characterization: The J-V curves were measured using a Keithley 2400 source meter. The illumination was provided by a Newport Oriel 92192 solar simulator with an AM1.5G spectrum, operating at 100 mW cm<sup>-2</sup>, which was calibrated by a standard silicon solar cell (91150 V) from Newport. The EQE spectra were measured in AC mode on a spectrum corresponding system (Enlitech QE-R), calibrated by Si reference solar cell. Capacitance-voltage (C-V) was measured with the bias voltage varying from -0.5 to 0.5 V at room temperature under near-dark condition. Besides, the measurement frequency was set to 100 kHz and the AC modulation voltage was set to 20 mV. The voltage dependent depletion width was calculated according the formulation  $x_d = \epsilon A/C(V)$ , where  $x_d$  is the depletion width,  $\varepsilon$  is the permittivity of Sb<sub>2</sub>Se<sub>3</sub>, A is the device area, and C(V) is voltage dependent capacitance, respectively. For low temperature J-V measurements, the devices were placed in a closed-cycle nitrogen cryostat, and the temperature was varied from 100 to 300 K with a step size of 20 K. A LabVIEW control program was used to operate the temperature controller and the Keithley 2400 source measurement unit for current and voltage data acquisition. Light intensity /-V measurements and transient photovoltage V<sub>OC</sub> decay curves were measured by an electrochemical workstation (ZAHNER CIMPS, Germany). The light source is the standard wideband white LED light source for ZAHNER CIMPS. The light intensity changed from 1000 to 100 W m<sup>-2</sup> in light intensity J-V measurements. The light intensity used in transient photovoltage  $V_{OC}$  decay measurements was 1000 W m<sup>-2</sup>.

# **Supporting Information**

Supporting Information is available from the Wiley Online Library or from the author.

# Acknowledgements

This research was financially supported by the National Natural Science Foundation of China (Grant No. 61804064), National Key Research and Development Program of China (Grant No. 2019YFB1503400), Natural Science Foundation of Guangdong Province (Grant Nos. 2017A030310543 and 2019A1515011616), and Guangzhou Science and Technology Program key projects (Grant No. 201707010254).

# **Conflict of Interest**

The authors declare no conflict of interest.

## **Keywords**

high efficiency, n–i–p structure, orientation,  $\mathsf{Sb}_2\mathsf{Se}_3$  solar cells, trigonal selenium

Received: March 18, 2020 Revised: May 7, 2020 Published online: June 30, 2020

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