# **Author Manuscript**

Accepted for publication in a peer-reviewed journal

# N**LS** National Institute of Standards and Technology • U.S. Department of Commerce

Published in final edited form as:

J Phys Chem C Nanomater Interfaces. 2020 ; 124(16): .

# **Comparable Enhancement of TERS Signals from WSe2 on Chromium and Gold**

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# **Abstract**

Plasmonic tip-sample junctions, at which the incident and scattered optical fields are localized and optimally enhanced, are often exploited to achieve ultrasensitive and highly spatially localized tip-

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enhanced Raman scattering (TERS). Recent work has demonstrated that the sensitivity and spatial resolution that are required to probe single molecules are attainable in such platforms. In this work, we observe and rationalize comparable TERS from few-layer  $WSe<sub>2</sub>$  single crystals exfoliated onto Au- and Cr-coated Si substrates, using a plasmonic TERS probe excited with a 638 nm laser. Our experimental observations are supported by finite-difference time-domain simulations that illustrate that the attainable field enhancement factors at the Au–Au and Au–Cr tip-sample junctions are comparable in magnitude. Through a combined experimental and theoretical analysis, we propose that besides Au/Ag, several metallic substrates may be used to record bright TERS spectral images.

#### **Graphical Abstract**



# **INTRODUCTION**

Tip-enhanced Raman scattering  $(TERS)^{1-4}$ -based nanoscale chemical imaging is ideally suited for the characterization of low-dimensional molecular and material systems.<sup>5-9</sup> This is especially the case for ultrathin 2D materials: a classical example of which is graphene sandwiched between a plasmonic TERS tip and a plasmonic substrate.<sup>10–17</sup> Beyond carbonbased material systems, relatively recent initial observations of  $TERS<sup>18</sup>$  as well as 'tipenhanced photoluminescence' (TEPL)<sup>19</sup> from  $MoS<sub>2</sub>$  were immediately followed by numerous demonstrations in which nanoscale heterogeneities (down to 10 nm) in  $\text{WS}_2$ <sup>20</sup>  $MoSe<sub>2</sub>,<sup>21</sup>$  and  $WSe<sub>2</sub><sup>22</sup>$  were directly visualized using these two complementary techniques. These results bolster the utility and promise of TERS and TEPL to advance the existing understanding of the fundamental properties of van-der-Waals layered materials.

In most of the existing TERS studies, both the probes and the substrates are made of single crystalline or polycrystalline gold or silver.<sup>5–16</sup> The premise is that visible light irradiation excites a hybrid tip–substrate junction plasmon, which effectively leads to the extreme localization and enhancement of the incident and scattered optical fields.<sup>5–9</sup> Indeed, a signal decrease of  $1-2$  orders of magnitude accompanies the switch from "gap mode"<sup>21,22</sup> to tiponly TERS20 from well-defined 2D material systems, the dimensions of which are gauged through simultaneous/sequential topographic atomic force microscopy (AFM) measurements. In this regard, bright TERS requires highly optimized samples, which greatly limits the generality of the technique and necessitates transferring the sample, typically onto a gold or silver substrate. At the same time, unavoidable interactions between the sample and a substrate may unnecessarily obfuscate the information content in TERS spectra and images.23–25 In this work, we demonstrate that other metallic substrates may be used to achieve high TERS enhancement factors for 2D semiconductors and, as a result, high signalto-noise TERS spectral images. The case is made through a comparison of the signal levels in hyperspectral TERS maps recorded from WSe<sub>2</sub> single crystals exfoliated on Cr-on-Si vs Au-on-Si substrates.

# **EXPERIMENTAL SECTION**

#### **Substrate/Sample Preparation.**

WSe<sub>2</sub> flakes were grown using a chemical vapor transport method with  $\text{SeBr}_4$  as a transport agent. First, polycrystalline  $WSe<sub>2</sub>$  powder was synthesized by reacting stoichiometric amounts of W (99.9%, Strem Chemicals) and Se (99.999%, Strem Chemicals) powders in a vacuum-sealed quartz ampule at 750 °C for 2 days followed by annealing at 850 °C for 3 days in the same ampule.<sup>1</sup> Then, a quartz ampule containing ≈1 g of WSe<sub>2</sub> charge and ≈0.1 g (5 mg/cm<sup>3</sup>) of SeBr<sub>4</sub> transport agent (99%, Strem Chemicals) was sealed under a vacuum and placed in a horizontal quartz tube furnace. The growth lasted 90 h with the charge and growth zone temperatures set to 825 and 700 °C, respectively.

For TERS measurements,  $WSe<sub>2</sub>$  flakes were mechanically exfoliated onto chromium and gold-coated silicon chips. The 30 nm thick Cr and Au films were deposited onto Si(100) substrates using the 4Wave dual chamber ion beam (Cr) and bias target (Au) cluster deposition system at the base pressure of  $2.6 \times 10^{-6}$  Pa ( $2 \times 10^{-8}$  Torr). Prior to the Au film deposition, a 20 nm layer of Ti was deposited onto Si surface to promote good adhesion of Au.

#### **TERS Measurements.**

AFM characterization and TERS imaging were carried out in air on an OmegaScope-R system (Horiba Scientific), coupled to XploRA confocal Raman spectrometer (Horiba Scientific). A linearly polarized 638 nm laser (100  $\mu$ W/ $\mu$ m<sup>2</sup>) was incident onto the sample through a 100x objective (Mitutoyo,  $NA = 0.7$ ) at a 65 $^{\circ}$  with respect to the surface normal.

<sup>&</sup>lt;sup>1</sup>Disclaimer: Certain commercial equipment, instruments, or materials are identified in this paper in order to specify the experimental procedure adequately. Such identification is not intended to imply recommendation or endorsement by the National Institute of Standards and Technology nor is it intended to imply that the materials or equipment identified are necessarily the best available for the purpose.

J Phys Chem C Nanomater Interfaces. Author manuscript; available in PMC 2021 January 01.

An OMNI-TERS AFM probe (silver coated, with a thin protective gold overlayer;  $k = 3$ )  $N/m$ ,  $f = 73.5$  kHz, APPNano, Horiba Scientific) was used to record the TERS maps shown herein, and the laser polarization was set along the long axis of the tip. A dedicated TERS imaging mode (SpecTop) was employed for fast and efficient TERS mapping. Using this mode, TERS signals are collected when the tip is in direct contact with the surface, with a typical force in the 2–10 nN range. A semicontact mode is used to move the sample relative to the tip (pixel to pixel) to preserve the sharpness and optical properties of the tip and to minimize the lateral forces that otherwise perturb the sample or contaminate the tip in conventional constant contact mode feedback.

#### **Numerical Simulations.**

Simulations were performed using a commercial FDTD package (Lumerical FDTD Solutions). The interaction of electromagnetic plane waves with the tip–surface construct (vide infra) is evaluated by iteratively solving finite-difference analogues of the timedependent Maxwell equations. The calculations model a solid silver tip with a 50 nm radius that is coated with a 20 nm gold overlayer. The tip is positioned 4 nm above a surface composed of 30 nm-thick metallic films supported by a silicon substrate. For reference, we also consider the case of tip (same as above) atop an uncoated glass substrate. A Gaussian source centered at 638 nm illuminates the tip/surface structure at a 65° angle with respect to the surface normal through an air objective  $(NA = 0.7)$ . Much like in the experiment, the incident light source is nominally polarized along the tip axis. The calculation outputs a three-dimensional surface of field enhancement which is then compared to the experiment. The dielectric permittivity of silver is obtained from Yang et al.,  $^{26}$  gold constants are obtained from Olmon et al.<sup>27</sup> Other constants are otherwise sourced from Palik.<sup>28–30</sup>

# **RESULTS AND DISCUSSION**

Our first task is to establish that we can directly compare TERS signals recorded from  $WSe<sub>2</sub>$ flakes exfoliated onto Cr-on-Si and Au-on-Si substrates. To this end, freshly prepared samples were imaged. To set a baseline for our discussion, we begin with a brief exposition of the TERS spectral images of WSe<sub>2</sub> on Au.

Few-layer WSe<sub>2</sub> flakes ( $\approx$ 7 nm in height as measured through topographic AFM, Figure 1A,B) exhibit strong resonant TERS scattering following 638 nm excitation. It is important to note that the far field/diffraction-limited response is weak, approximately an order of magnitude lower than the near field signal, as evidenced through the observation of welldefined sharp edges of the  $WSe<sub>2</sub>$  flake in the TERS map that correctly reproduce their analogues in the topographic image (Figure 1A and the inset of Figure 1C). If the far-field signal contribution was significant, the TERS map would exhibit more gradual, diffractionlimited (≈500 nm) transition at the edges of the flake. A detailed comparison between the far and near field components of the overall signal are provided in Figure S1. The enhanced resonant Raman spectra shown in Figure 1C exhibit the out-of-plane/in-plane  $A_{1g}/E_{2g}$  peak at around 250 cm−1 overlapping with the peak of the second order longitudinal acoustic phonon at M point 2LA(M) at 260 cm<sup>-1</sup>, as well as two weaker features at ≈130 and 395 cm

 $-1$ , which correspond to the first and third order of longitudinal acoustic mode at M point  $LA(M)$  and  $3LA(M)$  respectively.<sup>31</sup>

WSe<sub>2</sub> flakes that are comparable in their overall topographic features to the ones showed in Figure 1 were subsequently located on the Cr-on-Si substrate, see Figure 2, using the same TERS probe. Somewhat surprisingly, the TERS activity of these flakes were comparable to their Au-supported analogues. In terms of the characteristic spectral features, the longitudinal acoustic phonon peaks LA(M) (~130 cm<sup>-1</sup>), 2LA(M) (~260 cm<sup>-1</sup>), and  $3LA(M)$  (~395 cm<sup>-1</sup>) are still observable, albeit exhibiting slightly less prominent relative intensities in these spectra (Figure 2D) as compared to their analogues on Au (Figure 1C). Note that the predominant  $A^1_{\rm g}/E^2_{\rm g}$  modes at 250 cm<sup>-1</sup> are otherwise unscathed. These two observations suggest that the attainable TERS activities in the Au and Cr cases are not a result of significantly modified local chemical environments at the  $WSe_2/Au$  interface in these two samples.

A more direct comparison between the TERS signals that were recorded from the two samples emphasizes the statement above; see Figure 3. Namely, the overall similarities between the (spatially averaged) TERS spectra of  $WSe<sub>2</sub>$  on Au (red spectrum) and Cr (black trace) are evident. Interestingly, when the tip is in direct contact with the underlying metals 100–200 nm away from the edges of WSe2 crystals, (Au/Cr in green/blue, respectively), a much more pronounced TERS "background" is observed for the case of the plasmonic support. On the basis of prior analyses, we tentatively attribute this observation to electronic Raman scattering from nanometrically rough  $Au<sup>32</sup>$  Its absence in the case of Cr is intriguing and suggests that either (i) the difference in the surface roughness of the Cr and gold substrates (RMS 0.2 and 0.44 nm, correspondingly) is sufficient to affect the intensity of the background signal, or (ii) plasmonic modes in the visible (in the vicinity of the excitation wavelength) are prerequisite to the observation of backgrounds from imperfect mirrors/ nanometrically roughened metals.

Despite the fact that the measured heights of the crystals on gold and chromium were close, 7 and 5.5 nm respectively, we expect that the mere difference in the thickness of the crystals might affect the intensity of the gap-mode TERS signal. In order to understand how the crystal thickness affects the intensity of TERS signal, we measured TERS response from WSe<sub>2</sub> crystals of different heights exfoliated on gold (see Figure S2 in Supporting Information). As expected, the intensity of the TERS signal decreases with increasing thickness. Extrapolating the data from the graph in Figure S2, we can estimate that the thickness-adjusted intensity of  $WSe<sub>2</sub>$  on chromium was about 70% of the signal on gold.

To ensure the generality of our observations, several WSe<sub>2</sub> crystals on a Cr-on-Si substrate were imaged. TERS spectral images recorded from a strategically selected flake are shown in Figure 4. The crystal of interest features two distinct areas that exhibit apparent heights of ≈3 nm and ≈7 nm, respectively. Although both parts of the flake exhibit a pronounced TERS response, the thinner one also shows a PL peak, which is rigorously indicative of a  $WSe<sub>2</sub>$ monolayer.<sup>19,22</sup> This observation suggests that  $\approx$  2.3 nm of the measured apparent height of the flakes in the samples consists of a contamination layer that isolates the substrate from the WSe<sub>2</sub> monolayer, an observation eluded to in a prior report.<sup>33</sup> On this basis, we can re-

estimate the number of layers in the imaged  $WSe<sub>2</sub>$  flakes, where 3 and 7 nm apparent heights correspond to 1 and 6 WSe<sub>2</sub> layers, respectively, taking into account an approximate thickness of a monolayer of 0.7 nm.

To further ensure reproducibility of the results obtained in the first round of measurements, we repeated the experiments on new set of WSe<sub>2</sub> crystals exfoliated to gold and chromium. Even though in these samples we have not been able to locate crystals of similar height, the thickness-adjusted intensity of the TERS signal on chromium was found to be about 50% of corresponding signal on gold (see Figure S3 and related text in Supporting Information), which qualitatively agrees with the above-described results of our first round of measurements.

To rationalize our experimental observation of comparable TERS intensities from few-layer WSe2 single crystals exfoliated onto Au-on-Si and Cr-on-Si substrates, we performed finitedifference time-domain simulations; see Figure 5B–C. For reference, we also simulated the fields sustained at the apex when our plasmonic Au TERS probe is atop a glass substrate; see Figure 5A. We additionally performed predictive simulations that consider several other metallic substrates to assess the generality of our observations and analysis; see Figure 5D,E. Note that the general asymmetry in the field profiles shown in Figure 5 arises from the low-angle-of incidence of the incident microscopic laser field. For all metals tested, approaching the plasmonic tip toward the substrate leads to the confinement and enhancement of the electromagnetic field (Figure 5B–F). On the basis of our calculations, it is useful to comment on the attainable field enhancement factors. Assuming that all 4 electric fields involved in the TERS process are enhanced,  $34$  the total TERS signal enhancement follows the  $(E/E_0)^4$  law. In this case, all of the simulated plasmonic tipmetallic substrate constructs support enhancement factors that only differ by a single order of magnitude, i.e., in the  $\approx$ (1–9)  $\times$  10<sup>5</sup> range. This translates into comparable TERS signals for measurements performed on Au-on-Si vs Cr-on-Si substrates. In the absence of the underlying metallic support (Figure 5A), modest tip-only TERS would be expected, which is again consistent with prior observations.<sup>35</sup> More generally, our simulations suggest that the choice of underlying metal (Cr, Au, Mg, Al, Ag) does not affect the overall TERS signals in practice. In this regard, our results are consistent with a previous report on TERS signals from aromatic molecules on metallic substrates.17 Ultimately, both the real and imaginary parts of the dielectric determine the overall field enhancement. As the real dielectric becomes less negative, the field enhancement increases until the imaginary part becomes appreciable.36 At the red excitation wavelength used, Cr has a larger imaginary dielectric constant but smaller real dielectric constant relative to gold, resulting in a similar field enhancement and therefore TERS signal.

# **CONCLUSIONS**

The experimental results and theoretical analysis presented in this work reveal that bright TERS signals are attainable from analytes sandwiched between Au/Ag TERS probes and various visible light-reflecting metals. Namely, the choice of metallic substrate is not necessarily restricted to gold, silver, and other noble metals. Experimentally, this was demonstrated through TERS-based nanoscale chemical imaging of few-layer  $WSe<sub>2</sub>$  single

crystals exfoliated onto Au-on-Si and Cr-on-Si substrates. Unlike molecular reporters, where a plethora of effects can contribute to the variations in the recorded signal magnitudes, 2D materials are ideal targets for fundamental TERS studies. Finite-difference time-domain simulations corroborate our observables and further suggest that other metal substrates (including Mg, Al, and potentially other metallic mirrors) are expected to support TERS signals that are similar in magnitude to those reported herein. The prospect of utilizing several metallic substrates in TERS spectral imaging measurements should contribute to the growth and generality of this powerful nanoscale chemical imaging technique. The ability to utilize substrates other than Au and Ag may simplify general TERS sample preparation and expand the choice of substrates, which for some applications such as bioimaging, are currently restricted to Au.

#### **Supplementary Material**

Refer to Web version on PubMed Central for supplementary material.

#### **ACKNOWLEDGMENTS**

A.B. was supported in part by the Department of Energy's (DOE) Office of Biological and Environmental Research Bioimaging Technology project #69212. A.K. acknowledges support through Horiba Instruments Inc. A.G.J. and P.Z.E. acknowledge support from the Laboratory Directed Research and Development Program at Pacific Northwest National Laboratory (PNNL). S.K. acknowledges support from the U.S. Department of Commerce, National Institute of Standards and Technology under the financial assistance award 70NANB18H155

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#### **Figure 1.**

(A) Topographic AFM image of  $WSe<sub>2</sub>$  flakes on an Au-on-Si substrate. The measured height of these flakes is about 7 nm and is uniform across the sheets, as illustrated by a crosssectional height profile in panel B. Panel A also shows a 1.2  $\mu$ m × 0.6  $\mu$ m TERS image (88 pixels/line, 13.6 nm/pixel) where the intensity of the peaks within the 235–270 cm−1 range is overlaid over the topographic AFM image. The TERS image is expanded in the inset of panel C, which also shows the spatially averaged TERS spectra of  $WSe<sub>2</sub>$  that exhibit the familiar resonant Raman response of this system following red-laser excitation. Rather sharp edges of the flake are evident in the TERS map, which is indicative of strong TERS enhancement and nearly pixel-limited spatial resolution.

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#### **Figure 2.**

(A) Topographic AFM image of WSe<sub>2</sub> flakes on a Cr-on-Si substrate. The measured height of these flakes is about 5.5 nm, and is uniform across the sheets, as illustrated by a crosssectional height profile in panel B. (C) 1.5  $\mu$ m × 0.75  $\mu$ m TERS image (100 pixels/line, 15 nm/pixel) where the intensity of the peaks within the 235–270 cm<sup>-1</sup> range is overlaid over the topographic image. The same image is expanded in the inset of panel D, which also shows a spatially averaged TERS spectrum of WSe<sub>2</sub> that again exhibits the familiar resonant Raman response of this system following red-laser excitation. Much like in Figure 1, rather sharp edges of the flake are evident in the TERS map, which is indicative of strong TERS enhancement, pixel-limited spatial resolution, and a dim far field response.

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#### **Figure 3.**

Comparison between the TERS signals averaged over a  $150 \times 150$  nm area, obtained with the same TERS probe from WSe<sub>2</sub> on Au (red) and Cr (black). The similarity between the two spectra in terms of the observable resonances and signal levels is noteworthy. Background signals observed when the tip was in direct contact with the metals in two cases are also shown. A significantly more intense gold (green) compared to chromium (blue) background is observed. See text for more details. Ripples in the spectra, particularly pronounced in the background signal on gold are an artifact caused by optical interference on the back-illuminated CCD sensor.

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#### **Figure 4.**

(A) Topographic AFM image of the  $WSe<sub>2</sub>$  flake on Cr surface. (B) Corresponding profiles along the red and blue lines in panel A, featuring two distinct parts of the flake with approximate heights of 3 and 7 nm, respectively. (C) TERS/TEPL spectra from the 3 nm (red line) and 7 nm (blue line) thick areas, offset vertically by 50 counts for clarity. TERS (Panels D, E) and TEPL (panel F) maps from the square area, indicated by a dotted square in panel A. The spectral window for each map is: (D) from 195 to 295 cm−1 (green band in panel C), (E) from 360 to 420 cm−1 (blue band in panel C), (F) from 1700 to 1850 cm−1 (red band in panel C).

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#### **Figure 5.**

Electric field (E/E<sub>0</sub>) image slices in the XZ plane are shown in panels A (SiO<sub>2</sub> substrate), B (Cr-on-Si), C (Au-on-Si), D (Mg-on-Si), E (Al-on-Si), and F (Ag-on-Si). As discussed in the experimental and theoretical section of this manuscript, the simulations emulate all aspects of the experiments, including the laser color/polarization/angle of incidence. Note that the laser is incident from left to right, which leads to the asymmetry in the field maps for all constructs used. See text for more details. Note the differences in the scale bar ranges in A– F.