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# **Engineering Plasmonic Nanoparticles for Enhanced Photoacoustic Imaging**

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## **Abstract**

Photoacoustic (PA) imaging is an emerging imaging modality whereby pulsed laser illumination generates pressure transients that are detectable using conventional ultrasound. Plasmonic nanoparticles such as gold nanorods and nanostars are often used as PA contrast agents. The thermoelastic expansion model best describes the PA response from plasmonic nanoparticles: Light absorption causes a small increase in temperature leading to thermoelastic expansion. The conversion of optical energy into pressure waves  $(p_0)$  is dependent on several features: (i) the absorption coefficient  $(\mu_a)$ , (ii) the thermal expansion coefficient  $(\beta)$ , (iii) specific heat capacity  $(C<sub>n</sub>)$  of the absorbing material, (iv) speed of sound in the medium (c), and (v) the illumination fluence (F). Controlling the geometry, composition, coatings, and solvents around plasmonic nanostructures can help tune these variables to generate the optimum PA signal. The thermoelastic expansion model is not limited to plasmonic structures and holds true for all absorbing molecules. Here, we focus on ways to engineer these variables to enhance the PA response from plasmonic nanoparticles.

# **Graphical Abstract**



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#### **Keywords**

plasmonic nanoparticles; photoacoustic imaging; gold nanoparticles; coatings; enhancement; absorption coefficient; Grüneisen parameter; thermoelastic expansion

> Conventional optical imaging uses a light in–light out approach, which limits depth penetration due to scattering in tissue. Ultrasound imaging uses a sound in–sound out approach, which reduces scattering but lacks the spectral capabilities and high contrast of optical imaging. Photoacoustic (PA) imaging solves these limitations and provides spectral ultrasound imaging based on varying optical absorption. PA uses a light in–sound out approach, where pulsed light is absorbed to generate pressure transients that are detectable by ultrasound.<sup>1,2</sup>

PA imaging benefits from contrast agents that produce signal with different characteristics compared to the background. The most commonly used endogenous contrast agents are oxygenated and deoxygenated hemoglobin and melanin.<sup>2,3</sup> Oxy- and deoxyhemoglobin have distinct absorption spectra. Thus, PA signals obtained at different wavelengths can be unmixed to quantify blood oxygen saturation—a key parameter for diagnostics.<sup>4</sup> Exogenous contrast agents with different optical properties can increase contrast and can be engineered as activatable sensing and therapeutic platforms. These exogenous contrast agents range from dyes<sup>5–7</sup> and small molecules<sup>8–10</sup> to nano-<sup>11–13</sup> and microparticles.<sup>14–17</sup> Engineering highly efficient contrast agents can increase sensitivity, image quality, and early stage diagnostics; it can improve therapeutic efficiency.

To design efficient contrast agents, it is important to understand the method of signal generation, of which there are four primary mechanisms: thermal or thermoelastic expansion, vaporization,15,18,19 optical breakdown, and photochemical processes.20 For biomedical applications and contrast agent development, we are usually interested in the thermoelastic expansion route, because it does not change or break the properties of the biological sample under observation.<sup>1</sup> Herein, light absorption causes a small increase in temperature (in the milli-Kelvin range), leading to the brief expansion of the contrast agent. 1,20,21 This creates a pressure transient that can be detected acoustically. There is no bulk heating of the medium, because the expansion is spatially and thermally confined using light pulses shorter than the thermal relaxation time of the absorber.<sup>22</sup> The resulting PA signal is a function of the optical, thermal, and elastic properties of the sample.

The pressure difference generated due to thermal expansion  $(\rho_0)$  is given below in eq 1.<sup>1</sup> Here,  $\beta$  is the thermal expansion coefficient of the absorbing material in K<sup>-1</sup>; c is the speed of sound in ms<sup>-1</sup> in the medium;  $C_p$  is the specific heat capacity of the absorber in J/(K kg);  $\mu_a$  is the absorption coefficient in cm<sup>-1</sup>; and F is the irradiation fluence in J/cm<sup>2</sup>.  $\Gamma$  is the Grüneisen parameter (eq 2) characterizing the thermoacoustic conversion efficiency, and  $\vec{A}$ is the local energy deposition density in  $J/cm<sup>3</sup>$  (eq 3).

$$
\rho_0 = \left(\frac{\beta c^2}{C_p}\right) \mu_a F = \Gamma A \tag{1}
$$

$$
\Gamma = \left(\frac{\beta c^2}{C_p}\right) \tag{2}
$$

$$
A = \mu_{\rm a} F \tag{3}
$$

Plasmonic nanoparticles make good photoacoustic contrast agents because of their high absorption cross section, relative inertness, high stability, and a tunable localized surface plasmon resonance (LSPR).<sup>23,24</sup> Gold (spheres, rods, plates, shells, stars) and other metal nanostructures are particularly common as PA contrast agents. They are easily functionalized and can respond to specific chemical cues.25,26

To enhance PA signal generation from plasmonic nanoparticles, researchers have turned a few "knobs" as described in eq 1. The illumination fluence  $(F)$  is not a nanoparticle property and can be tuned separately, e.g., by laser power. The speed of sound  $(c)$  is a function of the tissue type and can be difficult to manipulate externally. However, the absorption coefficient  $(\mu_a)$ , thermal expansion coefficient  $(\beta)$ , and specific heat capacity of the absorber  $(C_p)$  can be tuned by modifying the nanoparticle shape, material, solvent, and so on. eq 1 can be applied to all absorbing molecules such as endogenous hemoglobin and melanin and is not just limited to plasmonic nanostructures.<sup>1</sup> Here, we review how the nanoparticle properties can be engineered to enhance PA signal  $(\rho_0)$ . We focus on the thermoelastic expansion model of PA generation and review ways to engineer the variables presented in eq 1.

# **ABSORPTION COEFFICIENT (** $μ_A$ **)**

The absorption coefficient describes how well a material absorbs incident photons at a defined wavelength—it is directly proportional to photoacoustic signal. As light traverses through a sample, it is either scattered or absorbed. Scattering and absorption is strong in tissues; hence, optical imaging techniques have limited depth penetration (approximately 100 μm). The reduced scattering coefficients for soft tissue and water from 600 to 900 nm range from 10 to 20 cm<sup>-1</sup> and to less than 0.003 cm<sup>-1</sup>, respectively. The absorption coefficient ranges from 0.1 to 0.5 cm<sup>-1</sup> and 0.0015–0.067 cm<sup>-1</sup>, respectively.<sup>27,28</sup> For PA signal generation, the excitation light must be absorbed to induce thermoelastic expansion (eq 1). A high  $\mu_a$  produces a high PA signal (eq 1), and there are various ways to increase the absorption coefficient of nanoparticles: increasing absorbance at a wavelength of interest, plasmon coupling leading to local-field enhancement and spectral shifts, and shape control of nanoparticles.

# **INCREASING ABSORBANCE**

Absorbance measures light attenuation through a sample. It combines the effects of photon absorption and scattering. Increased absorption and reduced scattering increase the absorption cross section of a particle. The absorption cross section is the probability that a photon of a particular wavelength is absorbed by a particle. Typical values for nanoparticles range from 2.9 to 33.5  $\times$  10<sup>11</sup> cm<sup>2</sup>.<sup>29</sup> Nanoparticles are employed as PA contrast agents

because of their high absorption cross section compared to that of other small molecules.<sup>30</sup> Generally increasing the absorbance leads to an increase in photoacoustic signal. $31-34$ Coatings like polydopamine,  $35,36$  melanin,  $37$  titanium dioxide,  $38$  graphene oxide,  $39$  and reduced graphene  $\alpha$ ide<sup>40,41</sup> on gold nanoparticles increase the absorbance and subsequent PA amplitude. Although uncoated nanoparticles and coatings alone showed no enhancement in all these reports, they do not show if a mixture of the core and shell materials leads to PA enhancement. An example was reported by Leng et al. in 2018; they engineered gold nanorod–copper sulfide heterostructures with enhanced photothermal conversion efficiency and photostability.<sup>42</sup> They showed that depositing  $Cu<sub>7</sub>S<sub>4</sub>$  on gold nanorod tips instead of a uniform shell around a nanorod core reduced scattering and increased absorption (Figure 1a–e). The increased absorption led to a 62% increase in photothermal efficiency. Its effect on PA signal generation remains to be seen.

Adding dyes like rhodamine  $B<sub>1</sub><sup>43</sup>$  indocyanine green,<sup>44–46</sup> cyanine 5.5, methylene blue, and quenchers like  $QSY_{21}$  to nanoparticles like gold or single walled carbon nanotubes increase absorption and enhance PA signals up to  $100$ -fold.<sup>47,48</sup> Upon irradiation, electrons excited to a higher energy state can relax back to their ground state by (i) a nonradiative pathway or vibrational relaxation, (ii) a radiative pathway or fluorescence, or (iii) intersystem crossing or phosphorescence. Maximizing nonradiative relaxation increases PA amplitude, as more energy is converted to heat and thermoelastic expansion.<sup>49</sup>

# **PLASMON COUPLING LEADING TO LOCAL-FIELD ENHANCEMENT AND SPECTRAL SHIFTS THAT CHANGE**  $μ_A$

Individual plasmonic oscillations on noble metal nanostructures can be coupled when nanoparticles are in close proximity to each other.<sup>50</sup> Plasmonic coupling can cause the electric field in the interparticle gap to be boosted by several orders of magnitude, which far exceeds the field enhancement for a single plasmonic nanoparticle (Figure 2a,b).<sup>51–55</sup> The plasmon modes of close particles can either be in phase or out of phase with each other. This causes attraction or repulsion leading to a decrease (red-shift) or an increase (blue-shift) in the hybridized plasmon resonance.<sup>56</sup> Usually, a decrease in nanoparticle size causes a blueshift, whereas an increase in size due to nucleation or aggregation causes a red-shift.<sup>57</sup> A red-shifted LSPR increases  $\mu_a$  in the near-infrared (NIR) range, which is favorable for PA imaging. The magnitude of the plasmon shift depends on the particle geometry, plasmon resonance of individual particles, and the interparticle spacing.52,58

PA enhancement due to plasmonic coupling has been extensively reported over the years. Gold nanoparticle strings,<sup>59</sup> biodegradable gold nanovesicles,<sup>60</sup> self-assembled semiconducting gold nanoparticles,  $^{61}$  aggregated gold–silica Janus nanoparticles,  $^{62}$  and many more examples show that plasmonic coupling between two or more nanoparticles shifts the LSPR into the biological optical window (650–900 nm). Increasing absorbance in the NIR range results in PA enhancement compared to uncoupled controls. Lu et al. reported a clever method for plasmon-coupling-related PA enhancement:63 They manipulated the LSPR band of AuNPs by enzymatic hydrolysis and induced self-assembly of nanoparticles. Nanoparticle assemblies responded to specific enzymatic cues (collagenase IV, Figure 2e)

Recently, our group reported iodide-doped silver-coated gold nanorods as an activatable PA contrast agent to measure oxidative stress.<sup>13</sup> We showed that the absorption spectra of gold nanorods (AuNRs) can be blue-shifted when coated with increasing amounts of silver (Figure 3a). Blue-shifting of the absorbance reduces  $\mu_a$  in the NIR range and the PA response at 680 nm. Selective etching of the silver iodide shell using reactive oxygen and nitrogen species can recover the absorption spectra and  $\mu_a$  of the bare AuNRs (Figure 3b). The increase in  $\mu$ <sub>a</sub> at 680 nm resulted in a 70% increase in the PA intensity in a mouse model of zymosan induced oxidative stress (Figure 3c,d).

# **SHAPE AND SIZE CONTROL**

Optimizing the geometry of plasmonic nanoparticles can significantly affect their optical and photoacoustic properties.64 Gold nanoparticles are the most commonly used plasmonic PA contrast agents.<sup>65</sup> They can be synthesized as spheres,<sup>66</sup> cubes,<sup>67</sup> plates,<sup>68</sup> rods,<sup>12,13,69–71</sup> cages,72–74 shells,75 bipyramids,76 and stars.77,78 The tips and edges lead to enhancement of the electromagnetic field as electron density increases at the tips. Simulations have shown that sharper tips have a higher  $\mu_a$  compared to rounded tips.<sup>79</sup> The emergence of hotspots where plasmons relax thermally via electron–phonon coupling can help generate the PA signal. Electron–phonon coupling describes the movement of atoms in a crystal lattice due to electron oscillations.80–82 Various shapes exhibit different ratios of absorption and scattering, so some shapes are better suited for PA than others. Gold nanospheres are easy to synthesize, but their LSPR (~530 nm) lies outside the biological optical window.

Gold nanorods are the most widely used PA contrast agents because of their easily tunable absorption spectrum; unfortunately, they are fairly unstable under high fluence laser illumination.83 Gold nanocages and nanoshells are also applied as biomedical PA contrast agents. Chigrin and co-workers modeled and showed that the most promising geometries are high aspect ratio gold nanorods and gold nanostars with high absorption and low scattering coefficients.<sup>84</sup> Knights *et al.* showed that 40–50 nm AuNRs produced the highest PA signal but were also cytotoxic; 10 nm AuNRs were efficient photoacoustic converters at equivalent total mass showing 2.5-fold PA enhancement.<sup>85</sup> There is a considerable amount of data on how PA is affected by shape and size of gold nanoparticles, but a comprehensive model that can predict PA response from other plasmonic geometries is still missing.

# **MOVING PLASMONICS INTO THE SECOND NEAR-INFRARED WINDOW (1000–1700 NM)**

Until recently, the field has devoted most of its efforts into engineering plasmonic nanostructures that operate within the first biological optical window (600–900 nm). There exists a second NIR window (1000–1700 nm) where biological tissues have low background. Designing contrast agents that operate within this range can help image deeper

by reducing scattering from tissue. Illumination fluence  $(F)$  can be increased when using longer (less energetic) wavelengths, which is governed by the maximum permissible exposure limit. Although water absorbs strongly at 1450 nm, researchers have shown promising results by moving plasmonics into the NIR-II window. Most of the work in this area is around semi-conductor-based polymeric nanoparticle assemblies with broadband absorption.86–88 Recently, Chen et al. reported a miniaturized gold nanorod for photoacoustic imaging in the NIR-II window.89 The LSPR peak of gold nanorods is controlled by the nanoparticle's aspect ratio. Increasing the aspect ratio red-shifts the LSPR of nanorods. Aspect ratios of 6 or higher can push the LSPR into the NIR-II range. The common seed-mediated synthesis of high aspect ratio gold nanorods results in particles that absorb in the NIR-II range but have large dimensions ( $120 \times 18$  nm).<sup>71,89</sup> But for *in vivo* applications, nanoparticles must be less than 100 nm in size to favor extravasation and accumulation in tumors. Chen and co-workers synthesized high aspect ratio gold nanorods  $(49 \times 8 \text{ nm})$  using a seedless synthesis process.<sup>89</sup> The resulting nanoparticles showed peak absorbance at 1064 nm and 3.5-fold photoacoustic enhancement. PA signal from miniaturized nanorods was higher due to a higher surface-area-to-volume ratio compared to their larger counterparts.

The absorption coefficient of plasmonic nanostructures can be easily tuned via the material properties or shape of the nanostructure.<sup>90</sup> For enhanced PA, the ideal nanoparticle should have high light absorbance and low scattering properties. Nanoparticle geometry and size have a significant impact on its optical properties. Geometries with sharp tips such as nanostars exhibit localized plasmon hotspots that result in a higher  $\mu_a$ . Moving plasmonics into the NIR-II window can help image deeper within tissues. The use of longer and lower energy wavelengths will aid in the clinical translation of PA by keeping exposure levels under the maximum permissible exposure limit while still providing sufficient imaging signal.

Another example was demonstrated by Wang *et al.* in 2014.<sup>91</sup> Using FDTD simulations, they modeled "rod-in-shell" nanoparticles with absorption in both the NIR-I and II windows (Figure 4). Their model nanoparticle had a gold core ( $60 \times 20$  nm), a composite Ag/Au (1:1) outer shell (4 nm), and a hollow gap in between (5 nm) (Figure 4c). Under longitudinal excitation mode, their nanoparticle showed two distinct absorption peaks at 681 and 1158 nm, respectively (Figure 4d). This was attributed to near-field enhancement and plasmon hybridization between the core and shell layers. With a higher absorption cross section compared to uncoated gold nanorods, these rod-in-shell particles could show high efficiency in both PA and photothermal therapy (PTT) applications.

Broadening the absorbance of gold nanoparticles into the NIR-II range via shape control, gold nanoechinus<sup>92</sup> or with coatings, or polydopamine blackbody<sup>36</sup> have been reported. As described earlier, plasmonic coupling of nanoparticles due to self-assembly or aggregation can also red-shift its absorbance spectra. The red-shift can be engineered into the NIR-II range. Overall, the high absorption cross section, tunability, and sensitivity of plasmonic nanoprobes along with the ability to image deeper and with higher fluence make their move into the second biological optical window highly desirable.<sup>49,93</sup> Since most of the work done

in this field is based on semiconducting and organic probes, plasmonics remains a relatively new but extremely promising area of study.

# **THERMAL EXPANSION COEFFICIENT (**β**)**

The thermal expansion coefficient describes how the size of an object changes as a function of temperature. The thermal expansion coefficient is a material property and hence very difficult to change in a nanoparticle. To the best of our knowledge, there have been no reports of engineering the thermal expansion coefficient of a nanoparticle. However, many groups have modeled and partially shown that the PA signal is a summed contribution of the thermal expansion of the nanoparticle and its solvent. During a laser pulse, heat cannot be completely confined within the nanoparticle—some of it diffuses into the solvent, leading to a shell-like layer of the solvent around the nanoparticle with elevated temperature. The resulting temperature increase induces thermal expansion in both the nanoparticle and the solvent (Figure 5). The amount of heat transferred into the solvent depends on the length of the laser pulse, interfacial thermal resistance (Kapitza resistance) at the nanoparticle–solvent interface, and the thermoelastic properties of the solvent itself.<sup>94</sup> Hence, tuning the solvent properties can enhance the PA signal.

The PA varying response of silica-coated gold nanospheres in water, silicone oil, and toluene was reported in 2012 by Chen *et al.*<sup>94</sup> They showed that the PA signal is proportional to the temperature-dependent Grüneisen parameter (eq 2) of the respective solvent (Figure 6). Ethanol has a much higher  $\beta$  and a much lower specific heat capacity ( $C_p$ ) compared to water at room temperature.<sup>95</sup> Hence, the Grüneisen parameter (eq 2) for ethanol is 11.7-fold greater than that of water at room temperature.<sup>95</sup> Shi *et al.* modeled and experimentally showed that changing the solvent from water to ethanol for gold nanospheres resulted in a 4.5-fold higher PA signal.<sup>95</sup> Furthermore, coating gold nanospheres with a shell having a higher thermoacoustic conversion efficiency than water could enhance the PA signal. The  $\beta$ for PDMS is 4.5-fold higher than that of water. By confining the heat distribution within the PDMS shell, a 3.5-fold PA enhancement was observed by Shi et al.<sup>95</sup>

Emelianov and others have modeled shown that over 90% of the PA signal is generated within the solvent compared to the absorbing nanoparticle.  $94-96$  The optical energy absorbed by the nanoparticle diffuses into the solvent as heat because of the small size of nanoparticles (high surface-area-to-volume ratio). The intensity of the PA signal is proportional to the outgoing heat flux.<sup>94</sup> Water is an ideal solvent to study the PA response of a nanoparticle. The thermal expansion of water disappears at 3.98 °C as water is densest at that temperature; hence, any PA signal generated from the heating solvent vanishes, but the PA generated from the nanoparticle can still be transmitted.<sup>94</sup>

Coatings such as silica,  $97-99$  reduced graphene oxide,  $40$  and zinc tera(4-pyridyl)porphyrin<sup>100</sup> can change the Kapitza resistance or thermal conductance at the nanoparticle–solvent interface, leading to PA enhancement. In the thermoelastic expansion model of PA, each nanoparticle has its own thermal profile: the particle itself and the heated solvent around it (Figure 5). When aggregated, the thermal profiles of individual nanoparticles can overlap and collectively add to increase the rate of heat flux into the solvent.101 This was

demonstrated by Bayer et. al in 2013.102 PA signal from aggregated silica-coated gold nanoparticles was up to 7-fold higher compared to dispersed silica-coated gold nanoparticles. A 25 nm thick silica shell around gold nanorods measuring  $65 \times 12$  nm in size ensured that there was no plasmonic coupling between aggregated nanoparticles (Figure 7a,b). In vitro tests showed that PA signal linearly increased as a function of particle concentration (Figure 7c,d).

Another example was demonstrated by Chen et al. in 2017; here, they synthesized temperature sensitive poly(N-isopropy-lacrylamide) (PNIPAM) nanoparticles loaded with gold nanorods or copper sulfide nanospheres.<sup>103</sup> PNIPAM shrinks above its lower critical solution temperature (LCST) at 32 °C. The gold nanorods or copper sulfide nanospheres were drawn closer together as the temperature increased above 32 °C—this increased the thermal flux from the cluster into the surrounding medium. Using this dynamic nanogel, they were able to enhance the PA signal from gold nanorods and copper sulfide nanospheres up to 22- and 30-fold, respectively (OD =  $3.33$  for 1 cm optical path length). Even below the LCST, PNIPAM–AuNR showed a 3.5-fold PA enhancement versus pure AuNR solution, which was attributed to random thermal overlap of nanoparticles within the gel matrix. PNIPAM–CuS showed no change in optical absorption due to aggregation but showed a 5.5 fold PA enhancement above the LCST. This confirmed that enhancement was due to thermodynamic and not plasmonic coupling effects.

The PA signal generated from core–shell nanoparticles, specifically gold–silica core–shell nanostructures, has been studied and modeled extensively. There are conflicting reports that show both PA enhancement (up to 3-fold for AuNRs)<sup>94,97,98,104,105</sup> and a PA decrease<sup>106</sup> upon silica coating. The most comprehensive model in this field was recently reported by Shahbazi *et al.* in 2019.<sup>96</sup> This group modeled the heat transfer and PA signal generation from a core–shell gold nanosphere using a nanosecond pulsed laser illumination. They considered the interfacial thermal resistance at the core–shell and shell–solvent interface and the temporal waveform of the laser pulse and showed that both parameters had a dramatic impact on PA amplitude. They also showed that only 7% of the PA signal is generated within the nanoparticle and that the fluid-generating region is a spherical shell almost twice as thick as the radius of the nanoparticle.

Other groups have looked at the ideal core–shell sizes and reported various values. Chen et al. observed the highest PA response from 13 nm gold cores coated with 18 nm silica shells. <sup>94</sup> Kumar et al. reported 6-fold enhancement using 35 nm gold cores coated with 20 nm silica shells.<sup>105</sup> Both groups reported that increasing the shell thickness above a certain value results in loss of PA signal as the thermal energy is confined within the shell and does not diffuse into the solvent. Using thin shells with short duration laser pulses result in the highest amplification.

On the contrary, Pang et al. reported that the Kapitza resistance has negligible effect on PA signal generation and silica coating causes a reduction in PA amplitude.<sup>106</sup> For uncoated gold nanospheres under 5 ns pulsed illumination, the Kapitza resistance had a negligible effect on PA amplitude. To see the effects of Kapitza resistance, they dropped the pulse duration to 0.1 ns. They assume that the Kapitza resistance at the shell–water interface is

Engineering the thermal expansion coefficient of a nanoparticle has yet to be reported. Most (>90%) of the PA response is generated by the solvent around the nanoparticle, and thus, choosing a solvent with a high  $\beta$  can significantly enhance the PA response; however, this is not always possible in biological applications. The Kapitza resistance can be tuned to maximize the thermal energy transfer from the nanoparticle to its solvent. Thin silica coatings that reduce the Kapitza resistance have shown promise in this direction.

# SPECIFIC HEAT CAPACITY (*C*<sub>P</sub>)

The specific heat capacity is the amount of thermal energy required to raise the temperature of a substance per unit of mass. Photoacoustic calorimetry is often used to measure the thermal properties of materials.108 To the best of our knowledge, there is no reported example of engineering the specific heat capacity of a nanoparticle to tune the PA response. Alloying metals can result in a new specific heat capacity compared to individual metals. Although there is no reported example of engineering the specific heat capacity of a nanoparticle, there is an example of alloyed Au–Ag nanoparticles that lead to PA enhancement. Hatef *et al.* reported that PA signals from Au–Ag alloyed nanoparticles were up to 2.5- and 10-fold higher compared to only gold or silver nanoparticles, respectively.<sup>109</sup> They attributed this enhancement to an increased absorption cross section and did not investigate the effects of alloying on the specific heat capacity of the particle. It would be interesting to study the effect of alloying and changing the specific heat capacity of nanoparticles to engineer a desired PA response.

# **LASER FLUENCE (***F***)**

Fluence is defined as the laser pulse energy (J) incident on the effective focal spot  $(cm<sup>2</sup>)$ . Laser pulse energy is a function of optical power and the pulse width (usually nanosecond scale). Laser fluence is an externally controlled variable as described in eq 1. Various groups have reported optimal laser configurations to obtain the maximum PA generation. The nonlinear increase in PA intensity as a function of laser fluence is widely reported for gold nanoparticles. The nonlinearity is usually attributed to changes in thermophysical parameters due to a localized increase in temperature at higher laser fluences.<sup>110</sup> Increasing the pulse width also increases the laser fluence, causing the PA signal to increase as long as thermal confinement is met. PA intensity usually plateaus at longer pulse widths.<sup>95</sup>

More recently, groups have reported femtosecond pulsed irradiation for PA imaging. Although laser fluence under fs pulses is lower, groups have found ingenious ways to enhance PA and thermal stability.<sup>111</sup> Femtosecond pulses can allow for two-photon photoacoustics using longer wavelengths and increasing depth penetration.<sup>112</sup> Masim *et al.* reported PA enhancement from gold nanoparticles using a pair of time-delayed femtosecond pulses (pulse width = 40 fs and  $E_p = 0.1$  mJ).<sup>113</sup> They studied photoacoustic intensity as a function of time delay between two pulses from 0 to 15 ns and showed that PA increased from 0 to 15 ns. They also showed that enhancement was maximized when the first pulse was 20–30% of the total energy delivered. The first pulse primed the solution for bubble generation (cavitation), while the second pulse was used to acquire the PA response (2-fold higher). Laser pulses with sharper temporal waveforms produce PA signals that are more sensitive to environmental changes (such as thermal resistance) and thus increase the probability of signal amplification.<sup>96</sup> Contrast agents that excite at the same wavelength but emit PA signal at various ultrasound frequencies could also help in multiplexing and sensing applications.

The laser fluence is a key parameter to obtain relevant PA signal. It is important to keep the laser fluence under the maximum permissible exposure limit (typically under  $20 \text{ mJ/cm}^2$ ). The use of femtosecond pulses can generate stronger PA signal via other mechanisms such as cavitation. Although increasing laser fluence can lead to PA enhancement, high fluence illumination can also lead to nanoparticle degradation especially in nonspherical structures. 99

# **FUTURE OUTLOOKS**

The thermal expansion model of nanoparticle-based photoacoustic signal generation still remains the primary mode to create contrast in biomedical applications. The absorption coefficient  $(\mu_a)$  is the easiest knob to tune to enhance PA contrast. Increasing the absorbance and reducing scattering using plasmonically coupled heterostructures or by incorporating NIR absorbing dyes can help design more efficient PA contrast agents. Shape control can engineer the local electromagnetic-field enhancement around a nanoparticle, but the field lacks a more comprehensive comparison between different shapes. Nanoparticles that absorb in the second optical window (1000–1700 nm) can help image deeper and reduce scattering. <sup>89</sup> Challenges like the absorption of water (1450 nm) still need to be overcome, but there are significant advances in this area already.<sup>114</sup> Engineering dispersion and shapes of chiral nanoparticles can lead to higher sensitivity in the optical properties (refractive index) of the nanoparticle. Polarization-dependent excitation could offer strong contrast even in highly absorbing media and biological tissues.<sup>115–117</sup> Nanoprobes that change their optical properties in different suspension environments can be excellent sensing platforms. Recently, CuS nanoplates that responded to refractive index changes (160–600 nm RIU−1) in their suspension media with changes in particle diameter and operation wavelength (1100– 1250 nm) were reported.<sup>118</sup>

Tuning the thermal expansion coefficient  $(\beta)$  of a nanoparticle is challenging. While phase changing hybrid plasmonic nanoparticles could be promising,  $119$  there is no such report for PA applications. Most of the PA signal is generated within the solvent, and tuning the rate of heat flux into the solvent is a key parameter to optimize the PA response. Silica-coated gold nanoparticles are the most studied nanosystem in the field, and it may be useful to study other plasmonic-semiconducting nanosystems.

The specific heat capacity  $(C_p)$  of a nanoparticle is difficult to engineer, but alloyed nanoparticles with tunable thermophysical properties can further enhance the PA contrast. The laser fluence  $(F)$  is a key parameter for good contrast. Anisotropic nanoparticles tend to be unstable under high fluence illumination, because they deform into a more thermodynamically favorable shape. Using femtosecond pulses that enhance contrast via other methods of PA generation is encouraging.

Other ways to enhance PA such as cavitation are promising.120 Nanoparticle-coated microbubbles have shown up to a 10-fold increase in PA amplitude due to cavitation even at low laser fluences.<sup>15</sup>

Deep-learning tools are usually used to increase PA contrast during image processing.<sup>121</sup> Tools that can model and predict the photoacoustic behavior of various old and novel nanoparticles may benefit the field in the future.

Gold is the most commonly used contrast agent due to its highly tunable optical properties, high photoacoustic conversion efficiency, and relative biological inertness. Of course, other materials can also generate PA signal: polypyrrole-coated iron–platinum- $122$  and platinumchelated bilirubin have also been reported<sup>123</sup> with excellent biocompatibility and a strong PA response under 808 nm illumination. Homan et al. and others have shown that functionalized silver nanoparticles can be used as a PA sensing platform for a host of biomedical imaging applications.<sup>124,125</sup> Ku *et al.* reported copper sulfide nanoparticles operating in the NIR-II window (1064 nm) as a PA contrast agent. Operating in the NIR-II window, they were able to image  $\sim$ 5 cm deep in chicken breast.<sup>126</sup> Although PA contrast agents other than gold have been reported, an empirical study to compare different nanoplatforms is still missing.

It is important to note that the thermoelastic expansion model of PA generation in eq 1 holds true for all absorbing molecules such as melanin, hemoglobin, and small molecules, not just for plasmonic nanoparticles.<sup>1</sup> A substantial amount of work has been done with nonplasmonic PA contrast agents, $127-129$  but more work is required with other plasmonic nanoplatforms.

## **PLASMONIC STABILITY**

Although gold nanorods are widely used as photoacoustic contrast agents, they are plagued with low photothermal stability. The rod structure is thermodynamically unfavorable, leading to nanoparticle reshaping under high fluence illumination. Reducing illumination fluence by using light emitting diodes (LEDs) instead of lasers can offer a cheap and easy solution. LEDs operate well below the MPE, which is the currently holding back the clinical translation of laser-based PA systems. But eq 1 clearly shows that reducing  $F$  will also reduce  $p_0$ . This may be acceptable if  $\mu_a$ ,  $\beta$ , or  $C_p$  are tuned to make a highly efficient contrast agent. Furthermore, two-photon excitation using less energetic and longer wavelengths can improve photostability and depth penetration. Some work on organic polymer–plasmonic and dye–plasmonic hybrid nanostructures has been reported in this space.<sup>43,130,131</sup>

In the advent that low-fluence illumination is not feasible, nanoprobes can also be shielded using coatings to improve their photothermal stability. For gold nanorods, coatings like silica,  $99$  reduced graphene oxide,  $40$  and polydopamine  $35$  have increased photothermal stability. Increasing the surface-area-to-volume ratio of nanoparticles allows faster dissipation of heat into the solvent, which increases PA amplitude and photothermal stability. Miniaturized gold nanorods  $(48 \times 8 \text{ nm})$  show 3-fold higher photothermal stability compared to larger rods with the same aspect ratio ( $120 \times 18$  nm).<sup>89</sup> Other materials and morphologies such as palladium nanosheets have also shown enhanced stability, but the tests in this study were carried out at a relatively low fluence of 4 mJ/cm<sup>2</sup>.<sup>132</sup> Improving photothermal stability will always be highly sought after, as it allows contrast agents to be imaged for longer time periods. Higher stability can also increase sensitivity and allow longer illumination times, which can improve the efficacy for PTT applications.

# **MULTIFUNCTIONAL PHOTOACOUSTIC PLASMONIC NANOPROBES**

Multifunctional nanoprobes that provide contrast in photoacoustics as well as other imaging modalities are especially helpful. Using proven and standard imaging techniques like magnetic resonance imaging (MRI) and ultrasound, multimodal contrast agents can independently validate the PA response.<sup>12,133,134</sup> This is particularly helpful, since there is no clinically approved photoacoustic imaging system on the market yet. One such triple modal contrast agent was reported by Vijayaraghvan et al. in 2016.<sup>92</sup> Here, they synthesized plasmonic gold nanoechinus and used them as upconverting, fluorescent, and photoacoustic contrast agents operating in the NIR-II range. Regura *et al.* reported a Janus plasmonic– magnetic gold–iron oxide nanoparticle that provided contrast with X-ray, CT,  $T_2$ - weighted MRI, PA, and dark-field optical microscopy.<sup>135</sup> The ability to multiplex various techniques can help extract maximum information from the probe.

Theranostic nanoprobes that can diagnose and treat a condition are extremely desirable. The ability to sense, activate, and treat can help with targeted drug delivery and improve treatment efficacy. PTT is the most common theranostic application for photoacoustic contrast agents.136–138 PA contrast agents are ideal for PTT, since they have a high absorption cross section and high photothermal conversion efficiency. Plasmonic nanoparticles can be engineered with molecular beacons that seek out specific diseased sites such as tumors where they are retained through the enhanced permeability and retention (EPR) effect. The increasing concentration of nanoparticles within the tumor microenvironment not only increases PA contrast but also helps in ablating the tumor with PTT.<sup>139–142</sup> Theranostic applications beyond PTT are a relatively new field of study, which is very encouraging. Hybrid plasmonic structures have shown promising results. Recently, our group and others reported a gold–silver hybrid nanoparticle for treatment and photoacoustic imaging of drug resistant bacterial infections.<sup>68,70</sup>

## **SAFETY AND CLINICAL TRANSLATION**

Reducing the cytotoxicity of contrast agents is paramount for successful clinical translation. Gold is a commonly used contrast agent for PA because of its relative biological inertness. However, the both the surface chemistry and the underlying substrate are important. Indeed,

gold nanorods are plagued with cytotoxicity issues because of the cetyltrimethy-lammonium bromide (CTAB) surfactant used during gold nanorod synthesis. CTAB forms a bilayer around the gold nanorod, giving it a high positive charge, keeping it colloidally stable. CTAB can cause defects in the cell membrane, leading to cell death.<sup>143,144</sup> High doses of CTAB are often used as a positive control for cell death in cytotoxicity assays.13 A substantial body of work has been done to reduce the cytotoxic effects of gold nanoparticles through surface modifications using polyethylene glycol,<sup>145</sup> phosphatidylcholine,<sup>146</sup> bovine serum albumin,<sup>147</sup> poly(acrylic acid), and poly(allylamine) hydrochloride.<sup>144</sup> A comparative study on shape-dependent cytotoxic effects of gold nanoparticles (spheres, rods, and stars) revealed that cytotoxicity was highest in rods followed by stars and nanospheres, respectively.148 Gold nanoparticles were also found to be more biocompatible (72%) than silver nanoparticles (58%) at 100  $\mu$ M using HepG2 cells.

The cytotoxicity of plasmonic contrast agents is the biggest hurdle for their clinical translation. Although designing multifunction probes is beneficial for complex sensing and therapeutic applications, it is often accompanied by complex synthesis processes and cytotoxic nanoparticles. Designing simpler yet safer nanoparticles would aid in the quick translation of PA contrast agents into clinics. While others like Halas and West have used plasmonics for photothermal therapy,  $149$  we are unaware of research groups using plasmonics and PA in humans, but the use of endogenous contrast agents such as melanin<sup>150</sup> and hemoglobin<sup>151</sup> has been reported. Unlike plasmonics, the endogenous nature of these contrast agents means that they cannot be tuned to maximize their PA response. The ability to carefully engineer plasmonic probes to specifically optimize the PA response for various biomedical applications makes them highly desirable.<sup>152,153</sup>

## **CONCLUSION**

Overall, plasmonic nanoparticles can be engineered for photoacoustic imaging by tuning one of the many "knobs" presented in eq 1. The acquired PA signal is a complex response from the nanoparticle, its shell, and the solvent around it. Although other methods of PA enhancement like cavitation are not properly understood, they are still promising. The field will continue to strive toward increasing the biocompatibility and photothermal stability of plasmonic nanostructures. This will facilitate more sensitive contrast agents for clinical translation. Multimodal and multifunctional nanoprobes are an important future research focus for photoacoustics.

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#### **VOCABULARY**

Absorption coefficient,  $\mu_a$  (cm<sup>-1</sup>), describes how well a material absorbs incident photons at a defined wavelength; Thermal expansion coefficient,  $\beta$  (K<sup>-1</sup>), describes how the size of an object changes as a function of temperature; Specific heat capacity,  $C_p$  (J/(K kg)), describes the amount of thermal energy required to raise the temperature of a substance per unit of mass; Fluence,  $F$ , describes the laser pulse energy  $(J)$  incident on the effective focal spot (cm<sup>2</sup>); Localized surface plasmon resonance (LSPR), describes the locally confined resonant oscillation of conduction electrons upon light irradiation



#### **Figure 1.**

Increasing absorbance enhances PA. (a–c) Calculated extinction (Ext.), absorption (Abs.), and scattering (Sca.) spectra of AuNRs, core–shell, and dumbbell-like AuNRs– $Cu<sub>7</sub>S<sub>4</sub>$ heterostructures, respectively. Adding Cu<sub>7</sub>S<sub>4</sub> on the tips versus a shell increases absorption and reduces scattering. (d,e) FDTD simulations of the localized electric field of (d) core– shell and (e) dumbbell-like AuNRs–Cu<sub>7</sub>S<sub>4</sub> heterostructures show electric-field enhancement when  $Cu<sub>7</sub>S<sub>4</sub>$  is deposited on AuNR tips. (f) Wrapping with reduced graphene oxides increases the absorbance at 700 and 800 nm. (g) PA amplitude acquired with increasing laser power inputs for samples presented in (f). Wrapping with reduced graphene oxide enhances the PA signal by 2-fold compared to unwrapped controls. (h) PA images obtained at 700 and 800 nm (laser power:  $6.2 \text{ mJ/cm}^2$ ; pulse rate: 10 Hz). Panels (a-e) were reprinted with permission from ref 42. Copyright 2018 John Wiley and Sons. Panels (f–h) were adapted with permission from ref 40. Copyright 2015 American Chemical Society.



#### **Figure 2.**

Plasmonic coupling causing local-field enhancement and spectral shifts. (a) Electric-field intensity distribution on top surface and (b) within the interparticle gap as a function of distance  $(g)$  between two gold nanorings. As  $g$  reduces, the electric-field intensity between the two particles is enhanced due to plasmonic coupling. (c) Induced aggregation and plasmon coupling causes a red-shift in the absorption spectrum. As a result, the absorbance at 710 nm increases. (d,e) Aggregation-based enhancement of the PA signal in vitro. (d) PA intensity curves of AuNPs and AuNPs@P3—a positive control that self-aggregates showing that increased NIR absorbance leads to a higher PA signal compared to the unaggregated control. (e) PA signal of AuNPs@P1 and AuNPs@P2 (negative control) treated with collagenase IV for 3 h in PBS buffer. Nanoparticles aggregated in the presence of collagenase IV results in a higher PA signal. Panels (a,b) were reprinted with permission from ref 53. Copyright 2012 American Chemical Society. Panels (c–e) were reprinted with permission from ref 63. Copyright 2020 John Wiley and Sons.



#### **Figure 3.**

Spectral shifts affect the absorption coefficient. (a) Absorption spectra of silver-coated AuNRs with increasing shell thickness (right to left). Increasing shell thickness causes a blue-shift in the absorption spectra due to a reduction in the AuNR aspect ratio. The  $\mu_a$  at 680 nm reduces because of the blue-shift. (b) Selective etching of the silver shell with an increasing concentration of reactive oxygen species  $(H_2O_2)$  restores the AuNR absorption and  $\mu$ <sub>a</sub> at 680 nm. (c) PA intensity at 680 nm increases as a function of  $\text{H}_2\text{O}_2$  concentration as  $\mu_a$  increases in (b). (d) *In vivo* detection of oxidative stress as a function of time. Iodidedoped silver-coated particles are more sensitive to oxidative stress and show a 70% increase in PA intensity over 90 min (680 nm). Adapted with permission from ref 13. Copyright 2020 The Royal Society of Chemistry.



#### **Figure 4.**

Moving plasmonics into the NIR-II window. (a,c) Schematic illustrations and (b,d) calculated absorption spectra of (a,d) a Au nanorod ( $D = 20$  nm,  $L = 60$  nm) or (c,d) a Au rod-in-shell particle ( $D = 20$  nm,  $L = 60$  nm,  $T = 4$  nm, and  $G = 5$  nm) when illuminated by the (black) transversely or (red) longitudinally polarized light. The blue and yellow shaded areas indicate the NIR-I and -II biological optical windows, respectively. (e) Near-field intensity distribution of the resonance modes (i–iv). Reprinted with permission from ref 91. Copyright 2014 The Royal Society of Chemistry.



#### **Figure 5.**

Solvent-generated PA signal. During a laser pulse, heat is transferred into the solvent, which undergoes thermoelastic expansion and produces up to 93% of the PA response (yellow curve) compared to the absorbing nanoparticle (purple curve). PA amplitude is directly proportional to the rate of heat flux from the nanoparticle to the solvent and depends on the laser pulse width, Kapitza resistance (thermal resistance at the nanoparticle–solvent interface), and the thermoelastic properties of the solvent itself. The red region around the nanoparticle depicts the heated region of the solvent.



#### **Figure 6.**

PA is proportional to the thermal expansion coefficient of the solvent. Thermal expansion coefficient and the Grüneisen parameter are directly proportional to each other. The PA amplitude of gold nanospheres with varying shell thicknesses as a function of temperature in various solvents: (a) water, (b) silicon oil, (c) toluene. The silica thicknesses are 0 nm (circle), 6 nm (triangle), 18 nm (square), and 38 nm (rhombus). (d) Temperature dependence of the Grüneisen parameter (left axis) for water (open circles) and toluene (open triangles) as well as the measured photoacoustic signal (right axis) of water (solid circle) and silicon oil (solid square). The error bar represents the standard error of 16 signals of 8 pulses each. Reprinted with permission from ref 94. Copyright 2011 John Wiley and Sons.

Mantri and Jokerst Page 29



#### **Figure 7.**

Overlapping thermal profiles increase the rate of thermal flux into the solvent. (a) Absorption spectra of dispersed and aggregated silica-coated gold nanorods show no plasmon coupling between aggregates. (b) Transmission electron microscopy images confirming aggregation of particles. PA signal from aggregated silica-coated AuNRs is higher than that of the dispersed controls. (c) PA images corresponding to varying concentrations of nanoparticles. (d) Average PA signal as a function of nanorod concentration for aggregated and dispersed particles. Aggregation causes the thermal profiles of the nanoparticles to overlap, increasing the rate of heat flux into the solvent and hence enhancing PA. Reprinted with permission under a Creative Commons (BY 4.0) – Gold open access license from ref 102. Copyright 2013 International Society of Optics and Photonics (SPIE).