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## Differences in the properties of porcine cortical and nuclear fiber cell plasma membranes revealed by saturation recovery EPR spin labeling measurements

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### Abstract

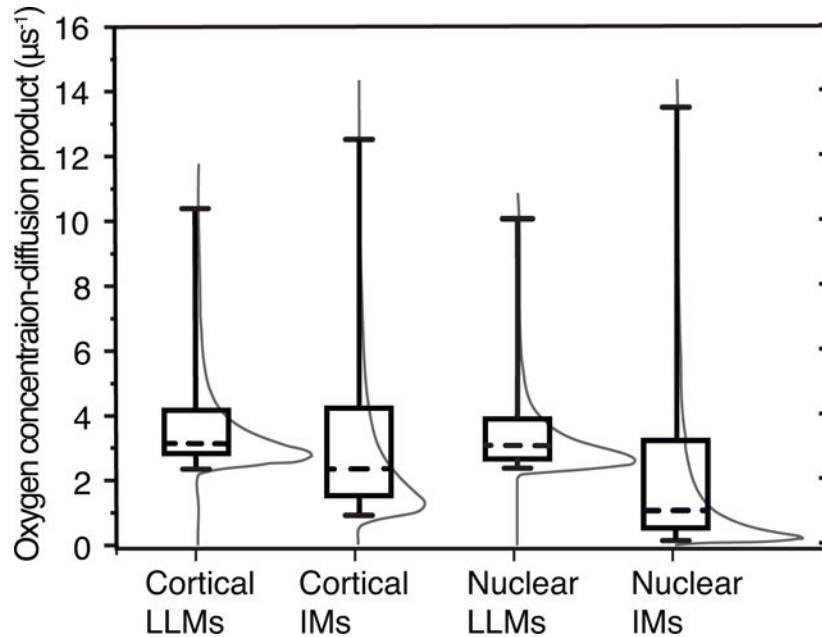
Eye lens membranes are complex biological samples. They consist of a variety of lipids that form the lipid bilayer matrix, integral proteins embedded into the lipid bilayer, and peripheral proteins. This molecular diversity in membrane composition induces formation of lipid domains with particular physical properties that are responsible for the maintenance of proper membrane functions. These domains can be, and have been, effectively described in terms of the rotational diffusion of lipid spin labels and oxygen collision with spin labels using the saturation recovery (SR) electron paramagnetic resonance method and, now, using stretched exponential function for the analysis of SR signals. Here, we report the application of the stretched exponential function analysis of SR electron paramagnetic resonance signals coming from cholesterol analog, androstane spin label (ASL) in the lipid bilayer portion of intact fiber cell plasma membranes (IMs) isolated from the cortex and nucleus of porcine eye lenses. Further, we compare the properties of these IMs with model lens lipid membranes (LLMs) derived from the total lipids extracted from cortical and nuclear IMs. With this approach, the IM can be characterized by the continuous probability density distribution of the spin-lattice relaxation rates associated with the rotational diffusion of a spin label, and by the distribution of the oxygen transport parameter within the IM (*i.e.*, the collision rate of molecular oxygen with the spin label). We found that the cortical and nuclear LLMs possess very different, albeit homogenous, spin lattice relaxation rates due to the rotational diffusion of ASL, indicating that the local rigidity around the spin label in nuclear LLMs is considerably greater than that in cortical LLMs. However, the oxygen transport parameter around the spin label is very similar and slightly heterogenous for LLMs from both sources. This heterogeneity was previously missed when distinct exponential analysis was used. The spin lattice relaxation rates due to either the rotational diffusion of ASL or the oxygen collision with the spin label in nuclear IMs have slower values and wider distributions compared with those of cortical IMs. From this evidence, we conclude that lipids in nuclear IMs are less fluid and more heterogeneous than those in cortical membranes. Additionally, a comparison of

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properties of IMs with corresponding LLMs, and lipid and protein composition analysis, allow us to conclude that the decreased lipid-to-protein ratio not only induces greater rigidity of nuclear IMs, but also creates domains with the considerably decreased and variable oxygen accessibility. The advantages and disadvantages of this method, as well as its use for the cluster analysis, are discussed.

## Graphical Abstract



## Keywords

stretched exponential; membrane fluidity; membrane heterogeneity; saturation recovery EPR; spin label; eye lens

## 1. Introduction

The application of the saturation recovery (SR) electron paramagnetic resonance (EPR) technique, with the use of lipid-analog nitroxide spin labels, was focused on studies of the organization and dynamics of lipid bilayer membranes (Altenbach et al., 1989; Kawasaki et al., 2001; Mainali et al., 2017a, 2013b, 2011b; Marsh, 2018; Raguz et al., 2011a, 2011b; Subczynski et al., 2007a, 2007b, 1991; Yin et al., 1987). By fitting the SR signal to exponential function, the spin-lattice relaxation rate ( $T_1^{-1}$ , inverse of the spin-lattice relaxation time) can be obtained. The  $T_1^{-1}$ , in the absence of other paramagnetic molecules, is determined primarily by the rotational diffusion of the nitroxide moiety of spin labels (Mailer et al., 2005; Marsh, 2018; Robinson et al., 1994), thus giving information about the membrane fluidity sensed by this nitroxide moiety. Using phospholipid (PL) and cholesterol (Chol) analog spin labels (*i.e.*, probe molecules that mimic the behavior of the parent molecules), the information about rotational motion of PLs and Chol can be obtained.

Introduction of the relaxation paramagnetic agent, molecular oxygen, into investigated membranes (by equilibrating them with the controlled partial pressure of oxygen) tremendously increased the application of the SR EPR methods to study membrane organization and dynamics. Use of this approach allowed us, for the first time, to evaluate the oxygen permeability coefficient across model (Mainali et al., 2013b; Raguz et al., 2008; Subczynski et al., 1991, 1989; Widomska et al., 2007) and biological membranes (Raguz et al., 2015a; Subczynski et al., 1992; W. K. Subczynski et al., 2017). It also allowed discrimination of membrane domains in model (Ashikawa et al., 1994; Kawasaki et al., 2001; Mainali et al., 2015, 2013c, 2013b; Raguz et al., 2011a, 2011b, 2008; Subczynski et al., 2010, 2007a; Widomska et al., 2007; Wisniewska and Subczynski, 2008) and biological membranes (Raguz et al., 2015a, 2014), which cannot be discriminated by spin labels alone. The spin-label oximetry approach was used to monitor membrane fluidity the diffusion of small probe molecules within the membrane, namely, diffusion of molecular oxygen. The profiles of the oxygen diffusion-concentration product across membranes that we obtained are new and very informative; they provide a sensitive way to describe membrane fluidity and, importantly, the neighboring lipid non-conformability.

Recently, we advanced the SR EPR approach for membrane studies by introducing the new way to analyze the SR signals using stretched exponential function (SEF) (Stein and Subczynski, 2020). To the best of our knowledge, this was the first time SEF was used to analyze SR EPR signals of spin-labeled biological membranes. This approach removes the requirement for a specific number of exponential components in the SR signal, each of which is characterized by the strict spin lattice relaxation rate and assigned to a distinct homogeneous membrane domain. With only two fitting parameters (the characteristic spin-lattice relaxation rate and the heterogeneity parameter), the analysis of SR signals with SEF offers a way to construct a continuous probability distribution of all possible spin lattice relaxation rates rather than indicating two or three strict distinct exponentials. The probability of finding a range of spin lattice relaxation rates associated with particular domains or neighboring lipids can be evaluated.

In our first paper (Stein et al., 2019), we applied the SEF to analyze SR signals for spin-labeled complex biological membranes (cortical and nuclear fiber cell plasma membranes of the porcine eye lenses) when the probability distributions of spin-lattice relaxation rates were determined only by the rotational diffusion of spin labels. These conditions were achieved for deoxygenated membrane samples. We demonstrated that the rates and the rate distributions sensed by PL analog spin labels (12-SASL) and Chol analog androstane spin labels (ASL) are faster and narrower in cortical membranes as compared with those measured in nuclear membranes. Constructed probability distribution functions and cumulative distribution functions allowed us to conclude that ~40% of PL molecules and ~25% of Chol molecules in nuclear membranes have rotational diffusions slower than those encountered in cortical membranes. Thus, in nuclear membranes, a very rigid environment exists that is not seen in cortical membranes. In that paper, we introduced the rescaling procedure that allowed construction of probability distribution functions and cumulative distribution functions with the x-axis expressed as  $T_1^{-1}$ , which allowed a simple and straightforward analysis of the data obtained for each membrane sample with its individual characteristic relaxation times and stretching parameters. Finally, we showed that the

stretched exponential parameters can be used to determine the significance of experimental differences and to cluster the samples based on origin using a K-means clustering analysis that is useful in exploratory research.

In our second publication (Stein and Subczynski, 2020), we advanced the previously established spin label SR EPR methods with the use of the SEF analysis to the case where the second relaxation process, namely the Heisenberg exchange between the spin label and molecular oxygen that occurs during bimolecular collisions, contributes to the decay of SR signals. We developed the theory and formulated basic equations for the analysis of SR EPR signals using SEF in membranes equilibrated with air, which allowed separate construction of probability distributions of spin-lattice relaxation rates determined by the rotational diffusion of spin labels, and the distribution of relaxations induced strictly by collisions with molecular oxygen. The latter distribution is determined by the distribution of the oxygen diffusion-concentration product within the membrane, which forms a new way to describe membrane fluidity and heterogeneity. We demonstrated the application of these new approaches, using intact fiber cell plasma membranes (IMs) of the cortex and nucleus of porcine eye lenses and lens lipid membranes (LLMs) prepared from the total lipid extract from these IMs, both labeled with ASL. These investigations clearly showed the extent to which integral and peripheral membrane proteins decrease the fluidity and increase the heterogeneity of membrane fluidity sensed by the Chol analog ASL alone and sensed by the movement and solubility of molecular oxygen. They also demonstrated that molecular oxygen is a very sensitive probe molecule for such investigations.

Here, we applied the above described approaches for the analysis of the SR EPR data with the application of the SEF to compare properties of membranes from different origins, namely IMs isolated from the cortex and nucleus of two-year-old porcine eye lenses. We also compared properties of these complex IMs with appropriate LLMs prepared from the total lipid extracts from IMs. For these investigations, we chose ASL. This work, in which we further developed and tested the method of analysis SR EPR signals with SEF approach, completes the research discussed our previous two papers (Stein et al., 2019; Stein and Subczynski, 2020).

## 2. Materials and Methods

Methods developed in (Stein et al., 2019; Stein and Subczynski, 2020) are applied to study differences between IMs isolated from the cortex and nucleus of porcine eye lenses labeled with the Chol analog ASL. These membranes were chosen because they were previously investigated in detail using both continuous wave (CW) and SR EPR approaches (Mainali et al., 2018, 2012a; Raguz et al., 2015b, 2008). Also, new methodologies for applying the SEF to analyze SR signals were first demonstrated on these membranes. Because of that, in the sections below, we only briefly describe all necessary procedures for samples preparations and handlings, referring readers to appropriate references. We provide the detailed procedures and indicated equations needed for the analysis of SR signals with SEF.

## 2.1. Materials

Spin-labeled Chol analog (*i.e.*, ASL) was purchased from Molecular Probes (Eugene, OR) (see the structures in Fig. 2 of Ref. (Witold K Subczynski et al., 2017)). Other chemicals, of at least reagent grade, were purchased from Sigma-Aldrich (St. Louis, MO).

## 2.2. Isolation of IMs and total lipids from cortical and nuclear samples

These procedures are described in (Stein and Subczynski, 2020). Porcine eyes were obtained on the day of slaughter from Johnsonville Sausage, LLC (Watertown, WI). Lenses were dissected and kept at  $-80^{\circ}\text{C}$  until needed. For the presented experiments, 36 lenses were thawed and decapsulated, and the cortex, the soft outer portion of the lens, was separated from the nucleus, the firm inner portion of the lens (Estrada and Yappert, 2004; Rujoi et al., 2003). Cortical and nuclear IMs were isolated separately (Bloemendal et al., 1972; Cenedella and Fleschner, 1992; Chandrasekher and Cenedella, 1995). First, the fiber cells were homogenized, and the soluble component was removed by centrifugation. Special care was taken to produce a uniform suspension by repeatedly aspirating the insoluble membrane suspension through a syringe fitted with an 18-gauge needle. Half of each sample was aliquoted into 18 different tubes each and stored in  $-80^{\circ}\text{C}$  as IMs. The other half of the sample was lyophilized. The dry weights were recorded, and lipids were extracted using Folch methods. The Chol-to-PLs ratio were determined by mass spectrometry at Medical College of Wisconsin the Pharmacology and Toxicology Mass Spectrometer Facility (Milwaukee, WI).

## 2.3. Spin labeling of IMs

Spin labeling from the dry film of ASL on the bottom of a test tube was performed as described previously (Raguz et al., 2015b). The ASL was introduced as 1 mol% of the total lipid in each sample. IM suspensions were added to the test tubes and shaken for about 2 h at room temperature. Finally, spin-labeled membrane suspensions were centrifuged for a short time, and the loose pellet was transferred to a 0.6 mm i.d. capillary made of gas-permeable methylpentene polymer and used for EPR measurements (Subczynski et al., 2005).

## 2.4. Preparation of spin-labeled LLMs

Multilamellar dispersions made from the total lipids extracted either from the cortical or the nuclear samples and containing 1 mol% of ASL (*i.e.*, LLMs) were prepared using the rapid solvent exchange method (Mainali et al., 2013b). The final membrane dispersions containing 1 to 2 mg of total lipids were centrifuged briefly (12000 g, 15 min,  $4^{\circ}\text{C}$ ), and the loose pellet was used for EPR measurements.

## 2.7. SR EPR measurements

SR EPR signals were obtained on the central line of the ASL CW EPR spectrum, as indicated by the arrows in Fig. 1, using short-pulse SR EPR at X-band with a loop-gap resonator (Mainali et al., 2017a). The X-band EPR spectrometer with SR capacity was developed at the National Biomedical EPR Center (Milwaukee, WI). Recently, the apparatus received two major hardware improvements: One improvement allows a short (300 ns duration) saturating pulse of 1 W of power to be delivered to the loop-gap resonator, and the

other reduces the dead time after the pump pulse from 300 ns to 100 ns. All measurements were carried out at  $\sim 36^\circ\text{C}$  for samples equilibrated with the same gas that was used for temperature control (*i.e.*, a controlled mixture of nitrogen and dry air adjusted with flowmeters [Matheson Gas Products; model 7631H-604]).

## 2.8. Analysis of SR signals using SEF

The first 120 ns from all recorded signals was removed before the further analysis because of the irregularity of the data. All SR signals were normalized and, using the Levenberg-Marquardt algorithm in the Origin 2020 (Northampton, MA) software package, fitted to Eq. (1):

$$I(t) = \exp\left(-\left(tT_{1\text{strobs}}^{-1}\right)^{\beta_{\text{obs}}}\right), \quad (1)$$

where  $I(t)$  is the signal at time  $t$ ,  $T_{1\text{strobs}}^{-1}$  is the fitted rate parameter at a particular air fraction, and  $\beta_{\text{obs}}$  is the corresponding fitted heterogeneity parameter. The zero time was taken at 220 ns after the pump pulse. Fitted parameters,  $T_{1\text{strobs}}^{-1}$  and  $\beta_{\text{obs}}$ , were obtained for each record. At least three records were used to calculate the averages and standard deviations per each air fraction.

**2.8.1. Analysis of SR signals coming from IMs and LLMs**—The average  $T_{1\text{strobs}}^{-1}$  values for each sample at various air fractions ( $f_{\text{Air}}$ ) were plotted versus the air fraction at which they were obtained and fitted to Eq. (2):

$$T_{1\text{strobs}}^{-1} = W_{\text{str}} f_{\text{Air}} + T_{1\text{strN}_2}^{-1}, \quad (2)$$

where  $W_{\text{str}}$  is the SOTP rate parameter and  $T_{1\text{strN}_2}^{-1}$  is the spin-lattice relaxation rate due to rotational diffusion and should be the same as the observed value obtained experimentally under nitrogen.

The averaged  $\beta_{\text{obs}}$  values were also plotted versus the air fraction at which they were obtained. The plots for each sample were fitted to Eq. (3):

$$\beta_{\text{obs}} = \beta_{\text{N}_2} \left(T_{1\text{strN}_2}^{-1} / T_{1\text{strobs}}^{-1}\right) + \beta_{\text{W}} \left(W_{\text{str}} f_{\text{Air}} / T_{1\text{strobs}}^{-1}\right), \quad (3)$$

where  $\beta_{\text{N}_2}$  is the heterogeneity parameter associated with  $T_{1\text{strN}_2}^{-1}$ , and  $\beta_{\text{W}}$  is the SOTP heterogeneity parameter.

The probability density distributions and the corresponding cumulative distributions were obtained as described previously (Stein and Subczynski, 2020). The probability density distributions for the given  $\beta$  parameters were computed in MATLAB R2018a (MathWorks, Natick, MA) using Eq. 11 from (Johnston, 2006). The  $s$ -values, as defined in (Johnston, 2006), for each distribution were multiplied by the corresponding stretched rate parameter to produce rate values. To maintain normalization of the probability distribution, the density values had to be divided by the same values as the corresponding stretched rates. The cumulative distributions were obtained by integrating the area under the density distribution

from zero to the particular rate. The values of the 5<sup>th</sup>, 25<sup>th</sup>, 50<sup>th</sup> (median), 75<sup>th</sup>, and 95<sup>th</sup> percentiles were obtained from the cumulative distributions.

### 3. Results and Discussion

When we introduce new methods to study lens fiber cell plasma membranes, we test them on IMs from the cortex and nucleus of two-year-old porcine eye lenses before transferring them to studies of human lenses. Here, we show broad applications of the recently developed SR EPR spin labeling approaches with SEF analysis (Stein et al., 2019; Stein and Subczynski, 2020) to compare the properties of IMs isolated from the cortex and nucleus of porcine eye lenses as well as the properties of these IMs with their models, namely LLMs prepared from the total lipid extracts from IMs. The Chol analog ASL is a good choice because this spin label can be easily incorporated into the IMs. Additionally, Chol is the major lipid component of lens fiber cell plasma membranes (Borchman and Yappert, 2010; Deeley et al., 2008; Rujoi et al., 2003).

#### 3.1. Characterization of cortical and nuclear IMs and LLMs

The mean weights and standard deviations of the cortical and nuclear portions from a single lens were roughly the same,  $0.432 \pm 0.0911$  g and  $0.461 \pm 0.0813$  g, respectively. The total wet weights from 36 lenses were 15.6 g for cortical and 16.6 g for nuclear samples. Half of each IM sample was lyophilized, and the total lipids were extracted providing us with the weights of proteins and lipids in the IMs. The dry weights of cortical and nuclear lyophilized IMs were 0.0637 g and 0.199 g, respectively, whereas the total lipids from the cortical and nuclear IMs weighed 12.1 mg and 7.50 mg, respectively. The dry membrane component of fiber cells, which includes peripheral and integral membrane proteins and the bilayer lipids, comprises 0.816% and 2.40% of the total cortical and nuclear fiber cell wet weights, respectively. This is consistent with the well-documented observation that fiber cells undergo compaction and lose water as they age, with the least water found in the center of the lens (Bassnett and Costello, 2017). Additionally, the proteins account for 81% and 96% of the cortical membrane and nuclear membrane dry weights, respectively. For comparison, the protein in the red blood cell membranes comprises only 44% of dry weight (Dupuy and Engelman, 2008). This increase is consistent with the previously published data, indicating that as fiber cells age certain phospholipids are depleted (Estrada and Yappert, 2004; Yappert et al., 2003) and previously soluble proteins bind to the membranes with higher affinity (Chandrasekher and Cenedella, 1995; Su et al., 2011; Truscott et al., 2011; Wang et al., 2013).

The Chol-to-PL ratios were found to be 0.60 and 2.7 in cortical and nuclear membranes, with an average PL weight of 744 Da. Given that these ratios and the amounts of lipids extracted from each subsection, the cortical and nuclear membranes contained 12.4  $\mu$ moles and 4.19  $\mu$ moles, respectively, of PLs for the total of 16.59  $\mu$ moles in the whole lens. Chol accounts for 7.42  $\mu$ moles and 11.3  $\mu$ moles in cortical and nuclear membranes, respectively, for the total of 19.02  $\mu$ moles in a whole lens. The Chol-to-PL ratio in the whole lens is 1.14. These results are slightly lower than those published by (Deeley et al., 2008), where the approximate Chol-to-PL mole ratio for porcine membranes can be calculated to be

approximately 1.4 for the whole lens. In addition, we can conclude that lipid-to-protein mole ratio decreases from 11.5 in cortical membranes to 2.42 in nuclear membranes, assuming the average protein weight is 30 kDa. The PL-to-protein mole ratio changes from 7.21 in cortical membranes to 0.657 in nuclear membranes. Similar trend, although not as pronounced, is observed in for cholesterol 4.31 in cortical and 1.77 in nuclear IMs.

It is important to note that the assumed average protein molecular weight above is a rough estimate. The intact eye lens membranes are complex in their protein composition (Bassnett and Costello, 2017; Song et al., 2009). For instance, aquaporin-0 is considered to be the major integral membrane protein with a full molecular weight of about 28 kDa. However, this protein is prone to cleavage and truncations that generates smaller fragments (Korlimbinis et al., 2009; Wenke et al., 2015). While it comprises a detectable amount in membranes isolated from younger human lenses, aquaporin-0 is barely detectable in membranes isolated from older lenses partly due to truncation and partly due to comprising smaller fraction due to crystallin binding (Truscott et al., 2011). Connexins 46 and 50 are 46 kDa and 50 kDa, respectively, form another major type of integral membrane protein. These proteins seem to be depleted with age and are missing in senile cataractous lenses (Buzhynskyy et al., 2011). But the majority of the protein as isolated in native or in partially urea stripped membranes is crystallin protein that becomes water insoluble with age (Chandrasekher and Cenedella, 1995; Truscott et al., 2011).

### 3.2. CW spectra and SR signals of ASL in cortical and nuclear IMs and LLMs

For brevity, the maximum splitting between the outermost peaks obtained directly from the CW EPR spectra of ASL (see inserts in Fig. 1) was used as a convenient parameter to monitor the order parameter of anisotropic motion of ASL (Hubbell and McConnell, 1971). The wobbling motion generates an angular distribution of the long axis of ASL molecules within the confines of a cone imposed by the membrane environment and the order parameter indicates the amplitude of that wobbling motion. However, it is often used to monitor membrane fluidity as was done by Kusumi et al. (Kusumi et al., 1986) (a greater maximum splitting indicates lower membrane fluidity). Maximum splitting in cortical (Fig. 1A) and nuclear (Fig. 1B) LLMs are 37.7 G and 39 G, respectively, which reflects the greater Chol content and the greater ordering effect of Chol in nuclear LLMs. The maximum splitting in the EPR spectrum of cortical IMs (Fig. 1A') increases only slightly, by 0.3 G, as compared with that in cortical LLM (Fig. 1A), indicating that the presence of proteins and their organization in these membranes does not affect ASL order as detected by CW EPR. It can be partly explained by the fact that Chol molecules as well as their analog ASL are substantially excluded from boundary lipids surrounding integral membrane proteins (Bieri and Hoelzl Wallach, 1975; Warren et al., 1975), or that the residence time of the ASL at the protein boundary is so short that no specific interactions are detected in cortical IMs (Páli and Kóta, 2019)

The CW EPR spectrum of ASL in nuclear IMs (Fig. 1B') is very different than that in cortical IMs (Fig. 1A'). It shows two clear components indicated by two maximum splitting. As noted in our previous publications (Mainali et al., 2012a; Raguz et al., 2015b), the greater maximal splitting is the characteristic of the strongly immobilized ASL and indicates the



presence of very rigid environment for Chol molecules existing in nuclear IMs with the exchange rate with bulk lipids slower than the time scale of CW method. The values of two maximum splitting in nuclear IMs are 40 G and 61 G, respectively, close to values estimated earlier for similar membranes (Mainali et al., 2012a; Raguz et al., 2015b). We assigned the strongly immobilized component as coming from ASL molecules trapped within aggregates and arrays of integral membrane proteins or otherwise immobilized by the presence of side chains on the partially inserted peripheral proteins. Our recent unpublished data indicate that the immobilized component in the EPR spectrum of ASL observed in nuclear IMs disappears upon treatment with 8 M urea. We concluded that the presence of peripheral proteins creates a highly immobile lipid environment, with the slow ASL exchange rate with other domains, or that the partially inserted, unfolded peripheral proteins or truncated peptide (Grami et al., 2005; Tang and Borchman, 1998) form rigid complexes with ASL. Interestingly, that similar pattern is observed for human IMs.

The immobilized component in the EPR spectra of ASL coming from cortical IMs appears only for membranes isolated from the lenses of 60- to 80-year-old donors (Raguz et al., 2015a). The immobilized component is always present in nuclear IMs from lenses of donors of all ages. These introductory data indicate that both porcine eye lens membranes and ASL as a monitoring molecule are good system to present the abilities of the SR EPR SEF methods.

### 3.3. Application of SEF to fit SR signals and obtain SEF parameters

**3.3.1. Rotational diffusion in LLMs**—Fig. 1 presents the procedure of fitting the representative SR signals coming from ASL in cortical and nuclear IMs and LLMs, recorded after deoxygenation, and after equilibration with 50% air, to the SEF using Eq. (1). Based on the residuals, both signals obtained under nitrogen and 50% air fit well to Eq. (1). Such experiments were performed for membranes equilibrated with various air fractions from 0% to 100%. These procedures allow the SEF fitting parameters  $T_{1\text{strob}}^{-1}$  and  $\beta_{\text{obs}}$  to be obtained for all SR signals analyzed with the SEF. The cumulative data for the obtained parameters are presented in Fig. 2 as a function of the air fraction in the equilibrating gas mixture.

It is important to note that, in the magnetically diluted deoxygenated samples, the spin lattice relaxation rates were determined primarily by the rotational diffusion of spin labels, and the  $T_1^{-1}$ s are directly proportional to the rotational diffusion rates of spin labels (Mainali et al., 2013a, 2011a). For this reason, the fitting parameters  $T_{1\text{strob}}^{-1}$  and  $\beta_{\text{obs}}$  obtained under nitrogen are denoted as  $T_{1\text{strN}_2}^{-1}$  and  $\beta_{\text{N}_2}$ . The averages of these parameters and their standard deviations obtained from at least three signals were found to be  $0.564 \pm 2.25 \times 10^{-3} \mu\text{s}^{-1}$  and  $0.991 \pm 6.15 \times 10^{-3}$  for cortical LLMs and  $0.367 \pm 1.98 \times 10^{-3} \mu\text{s}^{-1}$  and  $0.993 \pm 4.15 \times 10^{-3}$  for nuclear LLMs. Because  $\beta_{\text{N}_2}$  parameters are practically equal to 1, these SR signals represent single exponential decays and  $T_{1\text{strN}_2}^{-1}$ s are the same as  $T_{1\text{N}_2}^{-1}$ s, single spin lattice relaxation rates that characterize the rotational diffusion of ASL in LLMs. The ASL rotational diffusion is slower in nuclear LLMs as compared with that in cortical LLMs, reflecting the more rigid environment induced by a high Chol-to-PL ratio and the depletion of unsaturated PLs from the nuclear region of the eye lens (Estrada and Yappert,

2004). However, for similar Chol-to-PL ratios, these membranes are significantly more fluid than model membranes made of sphingomyelin-Chol at the same temperature (Mainali et al., 2012b). This may be because the model membranes were made of a single PL, whereas the eye lens contains a mixture of PLs with mixed acyl chains lengths and a degree of unsaturation.

**3.3.2 Rotational diffusion in IMs**—For IMs, the SEF parameters  $T_{1\text{strN}_2}^{-1}$  and  $\beta_{\text{N}_2}$  determined for deoxygenated samples are  $0.348 \pm 2.14 \times 10^{-3} \mu\text{s}^{-1}$  and  $0.947 \pm 5.75 \times 10^{-3}$  for cortical IMs, and  $0.252 \pm 2.86 \times 10^{-3} \mu\text{s}^{-1}$  and  $0.904 \pm 0.011$  for nuclear IMs, confirming our previous results (Stein et al., 2019) that, on average, the ASL probes in nuclear membranes are in more rigid environments than those in cortical membranes. The  $\beta_{\text{N}_2}$  parameters are sufficiently lower than 1, which indicates that the membrane fluidities sensed by ASL in IMs are heterogeneous and should be characterized by the probability distribution of  $T_{1\text{N}_2}^{-1}$ s. Additionally, the  $\beta_{\text{N}_2}$  value obtained for the nuclear membranes is lower than that obtained for the cortical membranes, confirming a higher heterogeneity of the ASL environment in nuclear membranes than in cortical membranes.

**3.3.3.  $T_{1\text{str}}$  and the stretching parameter  $\beta$ —effect of molecular oxygen in LLMs**—Cumulative data for  $T_{1\text{strobs}}^{-1}$  and  $\beta_{\text{obs}}$  for cortical and nuclear LLMs for samples equilibrated with different fractions of air are presented in Fig. 2A and B. Recall from Sect. 3.3.1 that the  $T_{1\text{strN}_2}^{-1}$  for cortical LLMs is almost twice as fast as that for nuclear LLMs, which indicates that the rotational diffusion rate of ASL in cortical LLMs is significantly greater than that in nuclear LLMs. As shown in Fig. 2A, in both cortical and nuclear LLMs, the  $T_{1\text{strobs}}^{-1}$  increases linearly upon addition of air in an equilibrating gas mixture. Despite having such different  $T_{1\text{strN}_2}^{-1}$ , the slopes of these linear dependences are practically the same for both LLMs. The values of  $T_{1\text{strobs}}^{-1}$  as a function of air fraction were fitted with the linear equation, with an intercept and a slope of  $0.529 \pm 0.0295$  and  $3.62 \pm 0.136 \mu\text{s}^{-1}$  for cortical LLMs, respectively, and of  $0.367 \pm 1.94 \times 10^{-3}$  and  $3.43 \pm 0.0123 \mu\text{s}^{-1}$  for nuclear LLMs, respectively. The values following the plus/minus sign indicate the error of fitting. The intercepts are the fitted  $T_{1\text{strN}_2}^{-1}$ s and are very close to those obtained experimentally in Sect. 3.3.1. The slopes are measures of the stretched oxygen transport parameters  $W_{\text{str}s}$ , which depend on the collision rates of oxygen with ASL and are very similar to each other despite the difference in  $T_{1\text{strN}_2}^{-1}$ s and Chol-to-PL ratios.

The  $\beta$  parameter is essential for stretched exponential analysis of SR signals because it determines the heterogeneity of the distribution of rates. The  $\beta_{\text{obs}}$  parameters for cortical and nuclear LLMs are plotted versus the air fractions in Fig. 2B. These values represent the sum of  $\beta_{\text{N}_2}$  and  $\beta_{\text{W}}$  parameters weighted by their respective rates at a given air fraction (Stein and Subczynski, 2020). The  $\beta_{\text{W}}$ s were determined by fitting the plots of  $\beta_{\text{obs}}$ s for cortical and nuclear LLMs versus the air fractions to Eq. 3 while holding the  $T_{1\text{strN}_2}^{-1}$ s and  $W_{\text{str}s}$  constant.  $\beta_{\text{W}}$ s were found to be  $0.914 \pm 0.033$  for cortical LLMs and  $0.937 \pm 0.012$  for nuclear LLMs.

**3.3.4.  $T_{1\text{str}}$  and the stretching parameter  $\beta$ —effect of molecular oxygen in IMs**—Cumulative data for  $T_{1\text{strobs}}^{-1}$  and  $\beta_{\text{obs}}$  for cortical and nuclear IMs for samples equilibrated with different fractions of air are presented in Fig. 2C and D. Upon the increase

of the fraction of air in the equilibrating gas mixture, the  $T_{1\text{strobs}}^{-1}$ s exhibit amazing linearity with the increase of air fraction in both cortical and nuclear membranes (Fig. 2C). The values of  $T_{1\text{strobs}}^{-1}$  as a function of air fraction were fitted with the linear equation with a slope and a zero intercept of  $2.73 \pm 0.0581 \mu\text{s}^{-1}$  and  $0.348 \pm 0.034 \mu\text{s}^{-1}$  for cortical IMs, respectively, and of  $1.34 \pm 0.0345 \mu\text{s}^{-1}$  and  $0.252 \pm 3.7 \times 10^{-3} \mu\text{s}^{-1}$  for nuclear IMs, respectively. The intercept values reflect the  $T_{1\text{strN}_2}^{-1}$ s and are measures of the membrane fluidity sensed by the ASL (the rotational diffusion of ASL). As with LLMs, the zero-intercept values match those measured experimentally.

In Fig. 2D, the  $\beta_{\text{obs}}$  parameters obtained for cortical and nuclear LLMs are plotted versus the air fractions. This analysis clearly indicates that the heterogeneity of the membrane fluidity sensed by molecular oxygen is significantly greater in nuclear IMs than in cortical IMs. The  $\beta_{\text{WS}}$  obtained by fitting  $\beta_{\text{obs}}$  data points to Eq. (3) are  $0.724 \pm 0.0125$  and  $0.407 \pm 0.0143$  for cortical IMs and nuclear IMs, respectively.

### 3.4. Distribution of spin-lattice relaxation rates of ASL in cortical and nuclear LLMs and IMs

The probability density distributions for the  $T_{1\text{N}_2}^{-1}$ s of ASL in the investigated membranes were constructed as described in [methods] using fitted parameter from Sec. 3.3. (Fig. 2). In Fig. 3, in addition to the probability density distributions, we indicated regions of  $T_{1\text{N}_2}^{-1}$ s where appropriate percentiles of possible rates are located using box-and-whiskers plots. The LLMs are represented by single lines because no distributions of  $T_{1\text{N}_2}^{-1}$  exist in these samples. The lipid exchange within the ASL environment seems to be sufficiently faster than the relaxation via rotational diffusion; thus, only a single relaxation rate is visible. Nonetheless, cortical LLMs are more fluid than nuclear LLMs, as expected from Chol-to-PL ratios and age-related depletion of unsaturated PLs in nuclear membranes (Estrada and Yappert, 2004).

The major conclusion from these results is that, because LLMs are characterized by a single  $T_{1\text{N}_2}^{-1}$  that are faster than their corresponding IMs, the integral and peripheral proteins in IMs create a variety of rigid environments that have a low lipid exchange constant between domains, which contributes to the detected heterogeneity of the ASL rotational diffusion. Additionally, the  $T_{1\text{N}_2}^{-1}$ s in intact membranes of the successive layers of fiber cells may be different enough to contribute to the observed heterogeneity. The extent of each effect will be investigated in future work. The  $T_{1\text{N}_2}^{-1}$ s for cortical and nuclear LLMs are faster than at least 75% of  $T_{1\text{N}_2}^{-1}$ s that can be found in their corresponding IMs. (Very fluid environments should be associated with such fast rates in IMs.) The increased protein content in nuclear IMs compared with that in cortical LLMs corresponds to the wider and more rigid distribution of  $T_{1\text{N}_2}^{-1}$ s. In IMs, 70% of the of cortical  $T_{1\text{N}_2}^{-1}$ s density has values between  $0.27 \mu\text{s}^{-1}$  and  $0.37 \mu\text{s}^{-1}$ , and another 20% is distributed as a long tail up to  $0.72 \mu\text{s}^{-1}$ . With the decrease of the lipid-to-protein ratio from 11.5 in cortical IMs to 2.24 in nuclear IMs, the width of the bulk 70% of the  $T_{1\text{N}_2}^{-1}$ s density has values between  $0.16 \mu\text{s}^{-1}$  and  $0.30 \mu\text{s}^{-1}$ . The long 20% tail distribution also widened up to  $0.79 \mu\text{s}^{-1}$ .

### 3.5. Distribution of oxygen transport parameters in cortical and nuclear LLMs and IMs

The probability density distribution and subsequent cumulative probability distributions of  $Ws$  in cortical and nuclear LLMs and IMs were constructed to examine the difference in values of the oxygen transport parameter among the samples. Cumulative results are presented in Fig. 4 as box-and-whiskers plots with probability density distribution functions for each sample for reference.

The distributions of  $Ws$  in cortical and nuclear LLMs are surprisingly similar despite their considerable differences in Chol-to-PL ratios and  $T_{1N_2}^{-1}$  rates. In both samples 75% of the density is distributed between approximately  $2 \mu s^{-1}$  and  $4 \mu s^{-1}$ , and another 20% extends to  $10 \mu s^{-1}$ . Because  $W$  is a function of the oxygen diffusion-concentration product, it is possible that the difference in PL types and Chol compositions lands ASL in similar environments regardless of the Chol-to-PL ratio. In our previous publications, we detected no heterogeneity of  $Ws$  in membranes made from the total lipids extracted from whole porcine lens (Raguz et al., 2008) or in separated cortical porcine LLMs (Mainali et al., 2012a). This slight heterogeneity might have been missed because the goodness of fit was estimated visually based on the residual. Using the SEF allows us to detect heterogeneity where the double exponential fit might have missed it. Mainali 2012 (Mainali et al., 2012a) reported the presence of some amount of  $W$  with the value of  $0.5 \mu s^{-1}$  in porcine nuclear LLMs; no such low value is present in the current distributions or when fitted to double-exponential function.

Comparison of results obtained for IMs with those obtained for LLMs clearly indicates that the presence of integral and peripheral proteins creates lipid environments in IMs with low ASL  $Ws$ . In cortical IMs, 70% of the  $W$  distribution extends from  $0.95 \mu s^{-1}$  to  $4.23 \mu s^{-1}$ , and another 20% stretches all the way to  $12.5 \mu s^{-1}$ . The effect is even more pronounced in nuclear IMs where the slowest 5% of the density has rates below  $0.15 \mu s^{-1}$ , and the slowest 75% below  $3.25 \mu s^{-1}$ . Comparison of the distribution of  $Ws$  between cortical and nuclear IMs indicates that oxygen diffusion-concentration products in nuclear IMs are suppressed by the presence of integral and peripheral membrane proteins much stronger than in cortical IMs. This can be explained by the fact that nuclear IMs are composed of a greater fraction of integral and peripheral proteins (Grami et al., 2005; Tang and Borchman, 1998) which induce formation of lipid domains with very low oxygen accessibility. All the presented data were obtained with spin labels located in different membrane environments. This probe approach permits measurement of collision rates of molecular oxygen with the nitroxide moiety of spin labels, and calculation of the product of the local oxygen diffusion coefficient and the local oxygen concentration. Separation of this product into component factors, diffusion and solubility, is not possible.

$T_{1N_2}^{-1}$  measurements show membrane fluidity and the heterogeneity sensed by lipid spin labels based on their rotational diffusion, and  $Ws$  measurements report membrane fluidity and the heterogeneity through the oxygen diffusion-concentration product in the membrane (by movement and solubility of small probe molecules). Both methods complement each other. A comparison of Figs. 3 and 4 indicates that the  $Ws$  method is more than an order of magnitude more sensitive than the  $T_{1N_2}^{-1}$ s (in both methods, the measured values are spin-lattice relaxation rates).

We previously demonstrated that the  $T_{1\text{str}}^{-1}$  and  $\beta$  parameters obtained from the SEF analysis of SR EPR signals coming from spin-labeled deoxygenated cortical and nuclear fiber cell plasma membranes from porcine eye lenses can be used for a multivariate K-means cluster analysis. This analysis is based on the differences in the separation of indicated parameters coming from samples of different origins. Data presented in Fig. 5 indicate that the separation of these data points increases enormously when  $T_{1\text{strobs}}^{-1}$ s and  $\beta_{\text{obs}}$ s obtained for samples equilibrated with air are used to separate them. For deoxygenated samples,  $T_{1\text{strN}_2}^{-1}$  and  $\beta_{\text{N}_2}$  parameters are relatively close to each other, with a Euclidean distance of 0.104. For samples saturated with air, the separation of  $T_{1\text{strobs}}^{-1}$ s and  $\beta_{\text{obs}}$ s parameters increased ~16 times to the value of 1.66.

The separation of SEF parameters reflects differences in the fluidity and heterogeneity of lipids in investigated membranes. Presented in Fig. 5 separation of data points was obtained for membranes of relatively young animals. We expect that parameters obtained from SEF analysis of SR EPR signals coming from human lenses will be clearly separated reflecting age related and health related differences in the organization of lipids in fiber cell membranes.

#### 4. Concluding Remarks and Future Directions

In this paper, we describe the fluidity of the lipid bilayer portion of IMs of the porcine eye lens as detected by ASL using SR EPR spectroscopy with the SEF analysis. The membrane fluidity was reported by the rotational diffusion of ASL and by the oxygen transport parameter (oxygen diffusion-concentration product). Two papers, previously published in the specific methodological journal, introduced this new method to analyze the organization, fluidity, and heterogeneity of complex biological membranes when only the rotational diffusion mode of spin lattice relaxation is considered (Stein et al., 2019), and when rotational diffusion and oxygen collisions contribute to the observed SR signal (Stein and Subczynski, 2020).

The conventional CW EPR technique with the use of lipid spin labels was more broadly used to study the structure of eye lens membranes (Babizhayev et al., 1992; Puskin and Wiese, 1982). Our group has also significant achievements in this direction. We applied PL analog spin labels (n-PC and n-SASL) and Chol analogs ASL and CSL to obtain profiles of the physical properties across membranes and to quantitate the amount of lipids with low and high mobility in cortical and nuclear bovine (Raguz et al., 2009; Widomska et al., 2007), porcine (Mainali et al., 2018, 2012a, 2011b; Raguz et al., 2015b), and human (Mainali et al., 2017b, 2015, 2013b; Raguz et al., 2015a, 2014) eye lens IMs as well as across corresponding LLMs. The CW EPR spectrum of ASL in nuclear IMs (recorded as a first derivative of the absorption EPR spectrum) consistently show two components indicated by two clearly separated maximum splitting. The EPR spectra of ASL in other investigated membranes, namely cortical and nuclear LLMs and cortical IMs, show only one component (one maximum splitting is observed). The appearance of the highly immobilized component in the CW EPR spectrum correlates with the aging of fiber cell membranes (Mainali et al., 2015; Raguz et al., 2015a). Aging also correlates with the increased of fraction of protein in the membrane, and associated with the membrane as reported here and elsewhere

(Borchman et al., 1989; Chandrasekher and Cenedella, 1995), breakdown of gap junction and formation of square arrays and microplacia (Bassnett et al., 2011; Costello et al., 1989) possibly caused by truncation and other modifications of integral membrane proteins (Korlimbinis et al., 2009; Wenke et al., 2015). These processes are insidious and may contribute to the formation of trapped and highly immobilized lipids in a variety of domains throughout the layers of cells represented in a sample. Diversity of membrane organization suggests that each component may contain the less abundant signals that get lost in the total broad-spectrum analysis. We hypothesize that the SR EPR method described here can provide this detailed analysis.

SR EPR in magnetically diluted samples and in the absence of oxygen provides information regarding the rotational diffusion of spin labels that depends on the local membrane fluidity (Mainali et al., 2013a, 2011a). Introduction of oxygen into a sample further informs about the membrane fluidity in terms to concentration/diffusion product of the small hydrophobic molecule (Kusumi et al., 1982). The fluidity of biological membranes affects the function of integral proteins, and membrane barrier properties (Almeida et al., 2005; Sezgin et al., 2017; Spector and Yorek, 1985; Subczynski et al., 2009) that is required to maintain membrane and cell homeostasis. The SR signal is recorded as an exponential-like decay. Observed with aging removal of unsaturated phospholipids (Borchman et al., 2004; Estrada and Yappert, 2004) and increase of Chol-to-PL ratios decreases membrane fluidity and creates formation of more rigid lipid environment (Almeida et al., 2005; Subczynski et al., 2009). Simultaneously, with aging the protein mole fraction is increased by the depletion of phospholipids and the increased binding of peripheral proteins. These processes create gradual rigidification of membranes towards the center of the lens. Because components of the multi-exponential decay are non-orthogonal, their numbers and values are difficult to determine (Smith et al., 1976). It is a common practice to fit such signals to double- or triple-exponential decay function and to assign specific values of decay times to select two or three domains, as we did in our previous publications using similar samples. Another method is to report the weighted averages of the two components, as is done in many similar spectroscopies and where stretched exponential application function was also used such as fluorescence life time imaging (Benny Lee et al., 2001) nuclear magnetic imaging (Bennett et al., 2003).

The use of the SEF to analyze a multi-exponential decay SR signal that consists of an undefined number of relaxation constants provides two fitting parameters, the characteristic or stretched spin-lattice relaxation rate  $T_{1\text{str}}^{-1}$  and the heterogeneity parameter  $\beta$ . If  $\beta$  is 1, then the signal comes from a signal exponential decay. If  $\beta$  is less than 1, this parameter is used to construct the continuous density probability distribution of all possible rates (see schematic explanation in Fig. 1 of (Stein and Subczynski, 2020)). Such distribution removes the necessity of estimating the number of rate constants and allows the rate to be treated as a continuous variable. The clear advantage of this treatment over double exponential is that instead of having only two fixed values a probability of finding any a range of values may be determined. This is convenient for complex biological specimens such as the eye lens membranes, where various domains from an individual fiber cell exist within a membrane, and the domains from different layers of fiber cells may also be different in physical properties measured by SR EPR. Of course, in a continuous distribution, the probability of

any single rate is nil. Instead, the probability of a range of rates must be determined. A range of a certain width under the mode of the distribution will have a higher probability and, therefore, is the larger component of the signal than a range of the same width under the long tail of the distribution. Another important factor to consider is that the distribution is determined by the inverse Laplace transform of the SEF, which is unimodal and has a long tail; this idealized shape may conceal finer details such as the existence of smaller modes.

The exciting application of the SEF is to use the fitting parameters ( $T_{1\text{str}}^{-1}$ ,  $\beta$ ) for a multivariate cluster analysis of stretched-exponential data from complex IMs as is demonstrated in (Stein et al., 2019). As demonstrated in Fig. 5, the separation of data points used for cluster analysis increases enormously for samples equilibrated with air. For deoxygenated samples, the  $T_{1\text{strN}_2}^{-1}$  and  $\beta_{\text{N}_2}$  parameters are relatively close to each other, with a Euclidean distance of 0.11. For samples saturated with air, the separation of the  $T_{1\text{strob}}^{-1}$  and  $\beta_{\text{obs}}$  parameters increases by an order of magnitude to the value of a Euclidean distance of 1.37. The oxygen-enhanced sensitivity of the advanced SR EPR technique, together with the SEF analysis, makes it possible to obtain measurements on membranes separated from a single human lens samples volume as small as 30 nL to 150 nL. Also, it permits the health history of the donor to be considered in the data analysis. We believe that the air-enhanced multivariate K-means cluster analysis of samples from different human populations and with different health histories should help elucidate predispositions for cataract formation.

Here, we show the general abilities of the SR EPR spin-labeling method with the SEF analysis in studies of IMs isolated from the cortex and nucleus of porcine eye lenses. We also indicate future directions in development and applications of this method. The next step will be to apply of this method to study IMs from human eye lenses.

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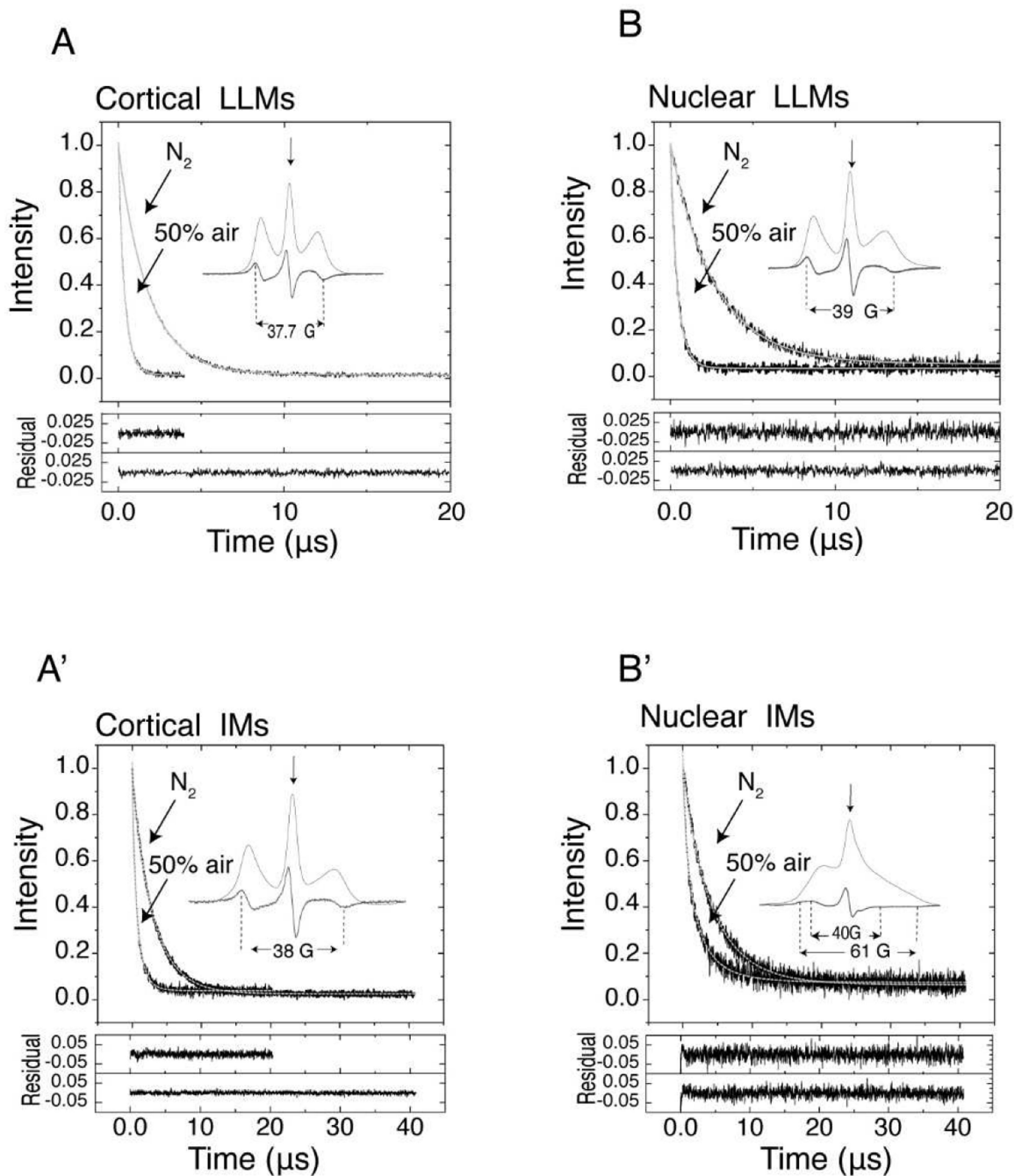
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**Highlights:**

Distribution of the rotational diffusion of cholesterol in eye lens membranes

Distribution of the oxygen transport parameter in eye lens membranes

Application of stretched exponential analysis to study membrane dynamics



**Fig. 1.**

Representative SR signals with fitted curves and residuals (the experimental signal minus the fitted curve) for ASL in deoxygenated cortical and nuclear LLMs and in membranes equilibrated with 50% air. SR signals presented in Fig. 1, which were obtained at different air fractions in the equilibrating gas mixture, were fitted to Eq. (1). The fitting parameters  $T_{1\text{strobs}}^{-1}$  and  $\beta_{\text{obs}}$  are displayed in Fig. 2. Inserts show the CW EPR spectra of ASL in deoxygenated membranes (thick lines) also displayed as absorption spectra (thin lines). The position of certain peaks in the EPR spectra was evaluated by monitoring the

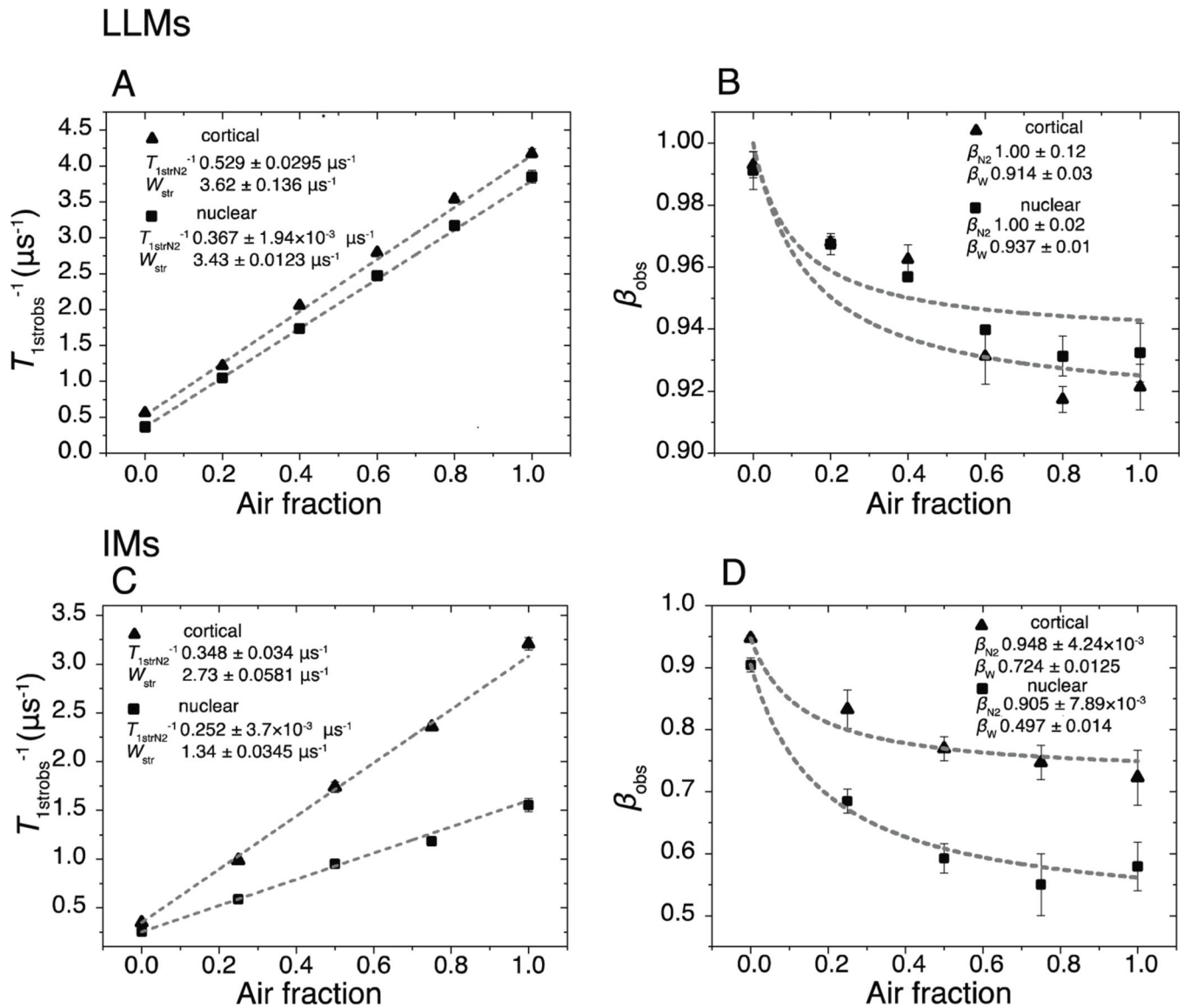
spectra at higher receiver gains and higher modulation amplitudes. The saturating microwave pump pulse is located at the field position of maximum intensity in the EPR absorption spectrum (indicated by an arrow). The recovery of the signal was recorded at the same field position with a low microwave power. All necessary details are indicated in the appropriate figures.

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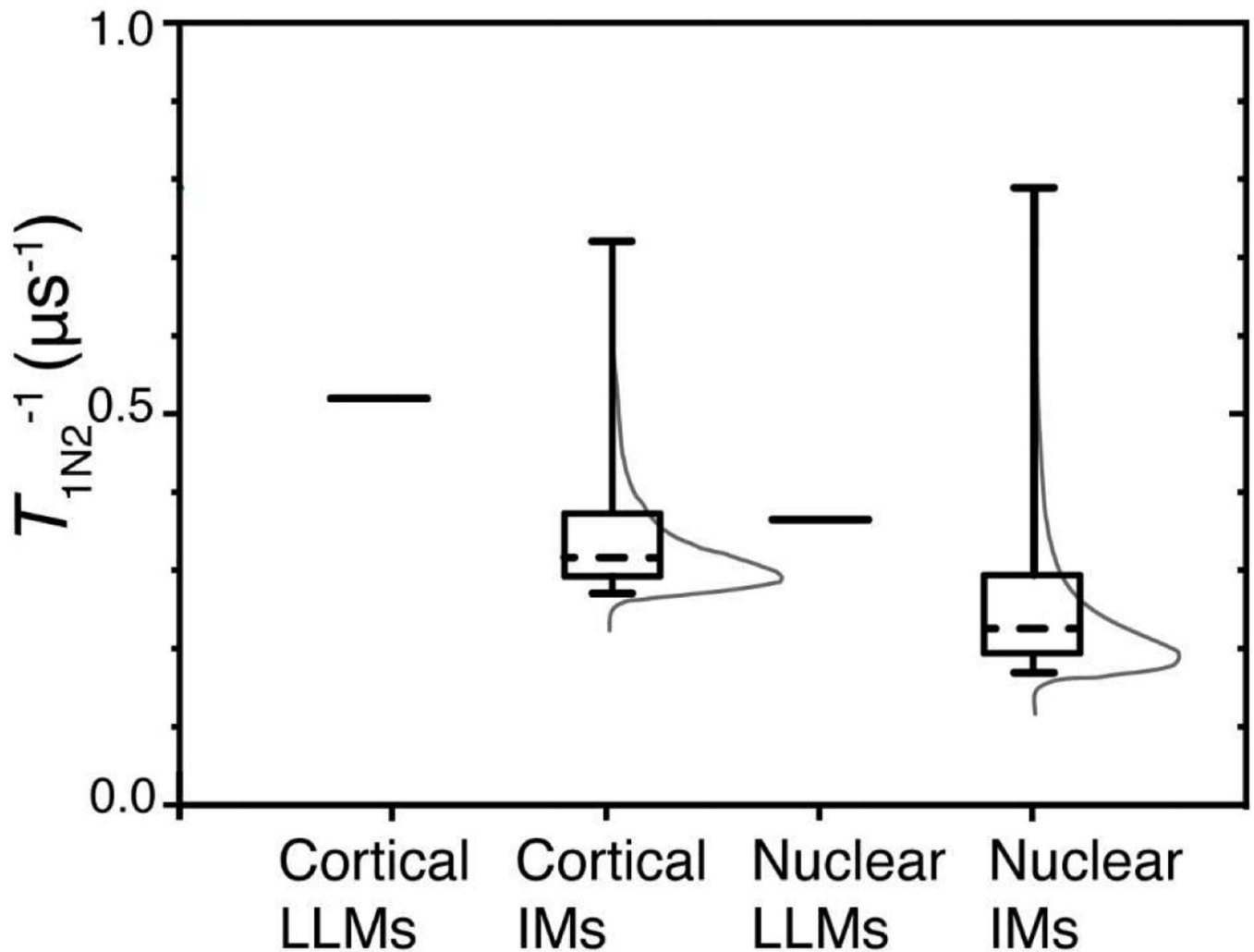
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**Fig. 2.**

In A and B, fitting parameters  $T_{1\text{stombs}}^{-1}$  and  $\beta_{\text{obs}}$ , which were obtained from fitting the SR signals of the cortical and nuclear LLMs using Eq. (1), are plotted versus the air fractions at which samples were equilibrated. **A.** The average  $T_{1\text{stombs}}^{-1}$  values from fitting at least three SR signals using Eq. (1) are indicated by black triangles and squares, respectively, for cortical and nuclear LLMs, with vertical bars that represent the standard deviations. The lines are linear fits of these data points. The values following the  $\pm$  sign are standard errors of the parameters. **B.** The black triangles and black squares are the average values of  $\beta_{\text{obs}}$  obtained from at least three SR signals, and the vertical lines correspond to standard deviations. Fits to Eq. 3 are indicated by lines. In C and D, fitting parameters  $T_{1\text{stombs}}^{-1}$  and  $\beta_{\text{obs}}$ , which were obtained from fitting the SR signals of the cortical and nuclear IMs using Eq. (1), are plotted versus the air fractions at which samples were equilibrated. **C.** The average  $T_{1\text{stombs}}^{-1}$  values from fitting at least three SR signals using Eq. (1) are indicated by black triangles and squares for cortical and nuclear IMs, respectively, with vertical bars that

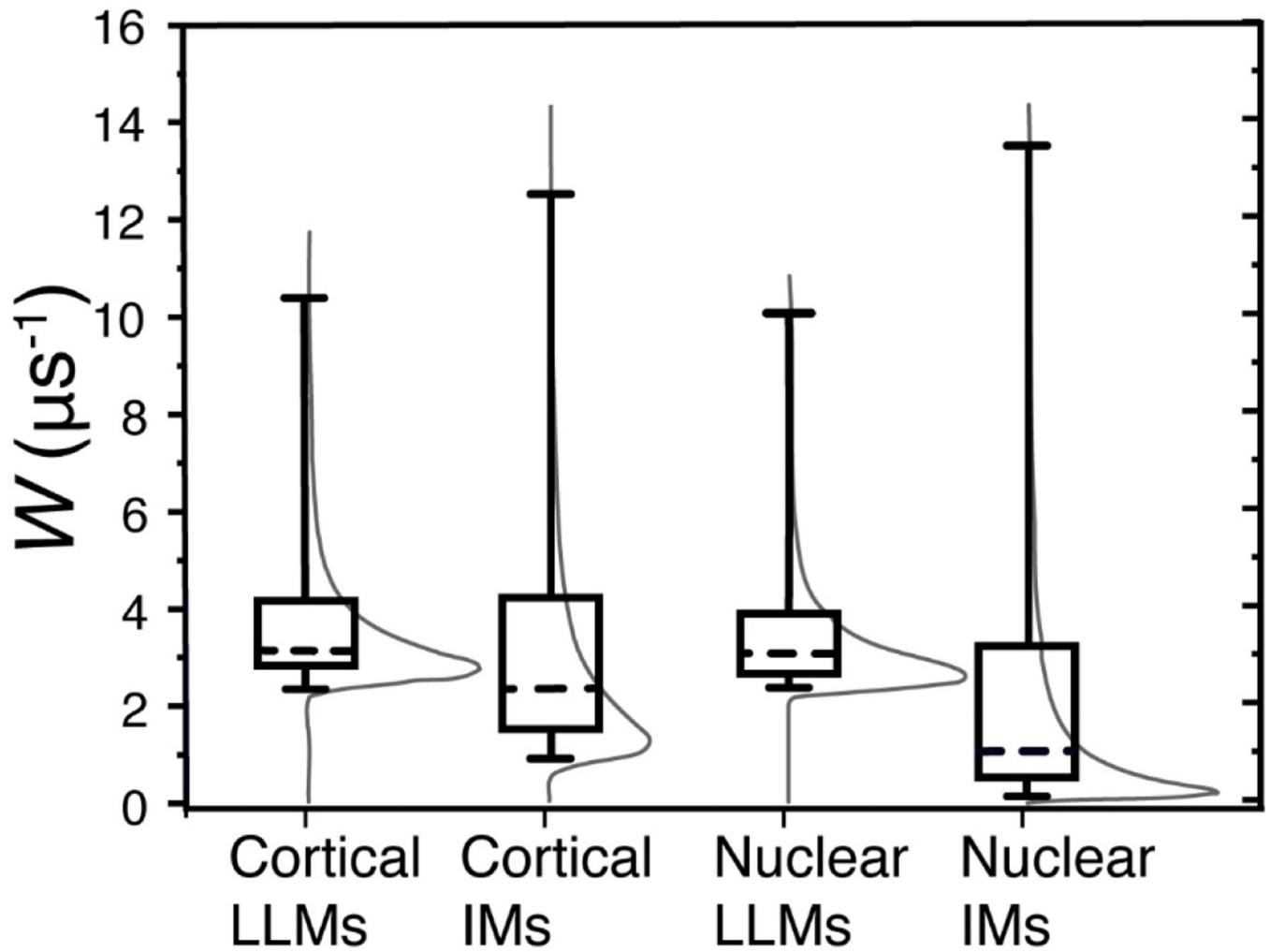


represent the standard deviations. The linear fits of these data points with the air fraction are also shown. The values following the  $\pm$  sign are standard errors of the parameters. **D.** The black triangles and black squares are the average values of  $\beta_{\text{obs}}$  obtained from at least three SR signals, and the vertical lines correspond to standard deviations. Fits to Eq. 3 are indicated by lines.

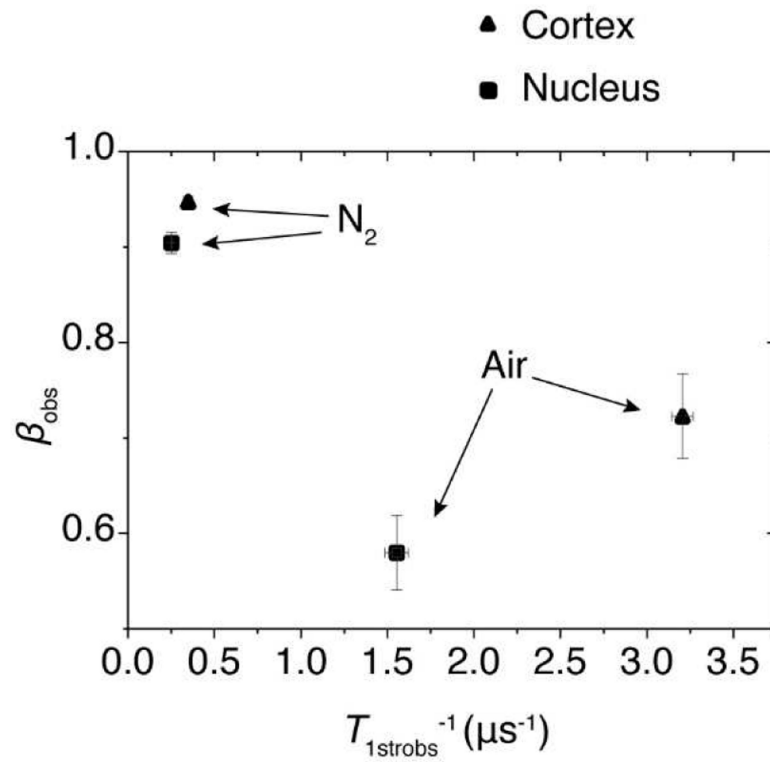


**Fig. 3.**

Using Laplace transform (Johnston, 2006; Stein et al., 2019) and the fitted parameters  $T_{1strN2}^{-1}$  and  $\beta_{N2}$ , probability distribution densities of  $T_{1N2}^{-1}$ s were constructed for each sample. These distributions, which are associated with the rotational diffusion of ASL in IMs, are represented by thin gray lines. Rotational diffusions of spin-lattice relaxation rates of ASL in cortical and nuclear LLMs are characterized by a single  $T_{1N2}^{-1}$  and, for both types of membrane, these distributions are depicted by single lines (delta functions) at  $T_{1N2}^{-1}$ , which are equal to  $0.529 \mu s^{-1}$  and  $0.367 \mu s^{-1}$  for cortical and nuclear LLMs, respectively. The rate percentiles were determined from cumulative probability distributions by integrating probability density distributions [First]. The whiskers in Fig. 3 extend from the 5<sup>th</sup> to 95<sup>th</sup> percentiles of all possible rates in a given distribution, the boxes extend from the 25<sup>th</sup> to 75<sup>th</sup> percentiles, and the dashed lines represent the 50<sup>th</sup> percentiles of cumulative probability distributions.



**Fig. 4.** The distributions of  $W$ s in cortical and nuclear IMs and in the LLMs associated with them, they are indicated by thin gray lines. The whiskers (cumulative probability distributions) extend from the 5<sup>th</sup> to 95<sup>th</sup> percentile of all possible collision rates in a given sample. The boxes extend from the 25<sup>th</sup> to 75<sup>th</sup> percentiles, and the dashed lines represent the 50<sup>th</sup> percentiles of cumulative probability distributions.



**Fig. 5.** Multivariable plot of  $T_{1strobe}^{-1}$  and  $\beta_{obs}$  parameters obtained with the SEF analysis of ASL SR signals from IMs for deoxygenated samples and samples equilibrated with air. Data from (▲) cortical IMs and (■) nuclear IMs. The error bars represent the standard deviations.