# **Supplementary information**

# CFTR function, pathology and pharmacology at single-molecule resolution

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1	CFTR function, pathology and pharmacology at single-molecule resolution
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### 8 Supplementary Methods

## 9 Calculations of the steady state population fractions

Here, we demonstrate the kinetic scheme formulation at physiologically relevant saturating ATP concentrations (> 1 mM) denoted with [ATP] starting with the state occupation probabilities when the system is at steady state. The population fractions of the states are labeled with  $S_l$  for l = 1, 2, ..., 8. Temporal evolution of these population fractions is governed by the following set of equations where the transition rates from state l to l' are shown with  $k_{l \rightarrow l'}$  for all  $l \neq l' = 1, 2, ..., 8$ .

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$$\frac{dS_1}{dt} = -(k_{1 \to 2} [ATP] + k_{1 \to 8})S_1 + k_{2 \to 1}S_2 + k_{8 \to 1}S_8,$$

16 
$$\frac{dS_2}{dt} = -(k_{2 \to 1} + k_{2 \to 3})S_2 + k_{3 \to 2}S_3 + k_{1 \to 2}[ATP]S_1,$$

17 
$$\frac{dS_3}{dt} = -(k_{3\to 8} + k_{3\to 2} + k_{3\to 4})S_3 + k_{2\to 3}S_2 + k_{4\to 3}S_4 + k_{8\to 3}[ATP]S_8,$$

18 
$$\frac{dS_4}{dt} = -(k_{4\to 3}+k_{4\to 5})S_4+k_{3\to 4}S_3+k_{5\to 4}S_5,$$

19 
$$\frac{dS_5}{dt} = -(k_{5 \to 4} + k_{5 \to 6})S_5 + k_{4 \to 5}S_4 + k_{6 \to 5}S_6$$

20 
$$\frac{dS_6}{dt} = -(k_{6\to 5} + k_{6\to 7})S_6 + k_{7\to 6}S_7 + k_{5\to 6}S_5,$$

21 
$$\frac{dS_7}{dt} = -(k_{7 \to 6} + k_{7 \to 8})S_7 + k_{6 \to 7}S_6 + k_{8 \to 7}S_8,$$

22 
$$\frac{dS_8}{dt} = -(k_{8\to 3} [ATP] + k_{8\to 7} + k_{8\to 1})S_8 + k_{1\to 8}S_1 + k_{8\to 1}S_8,$$

23 
$$\sum_{k=1}^{8} S_k = 1.$$

Here, the last equation imposes conservation of mass on the population fractions. We find the steady state solution for the state population fractions. steady state population fractions is used to initialize our data simulation. Below we provide the details of the data simulation.

#### 27 Simulation of simultaneous electrophysiology and smFRET data

In our simulation setup, we imagine visualizing the single molecule conformations with smFRET and electrophysiology, simultaneously. As the single molecule transitions amongst its conformational states labeled with  $S_l$  for l = 1, 2, ..., 8 where some of these states give rise the same signal in smFRET or electrophysiology measurements as summarized below in **Supplementary Table 1** with kinetic rates provided in **Supplementary Table 2**.

33 Supplementary Table 1. Conformational states and their associated signals in smFRET and electrophysiology

34 measurements.

occupied conformational	smFRET signal	electrophysiology signal
state		
<i>S</i> <sub>1</sub>	low FRET	no current
<i>S</i> <sub>2</sub>	low FRET	no current
S <sub>3</sub>	high FRET	no current
$S_4$	high FRET	no current
<i>S</i> <sub>5</sub>	high FRET	current
<i>S</i> <sub>6</sub>	high FRET	current
<i>S</i> <sub>7</sub>	high FRET	no current
<i>S</i> <sub>8</sub>	high FRET	no current

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36 Here, we simulate the trajectory reporting on the conformational dynamics of the single molecule in 37 continuous time as a Gillespie<sup>60</sup> trajectory where these conformational dynamics are governed by the 38 following generator matrix labeled with  $\overline{\bar{G}}$ 

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#### 40 where the \* notation shows the negative of the sum of each row.

Subsequently, this trajectory of the single molecule gives rise to smFRET and electrophysiology measurements. For example, if the molecule has dimerized NBDs and an open pore then this conformation of the molecule leads to high FRET and high current measurements in smFRET and electrophysiology, respectively.

In a prototypical smFRET experiment, FRET efficiency (FRET) is computed based on the photons emitted from donor  $(n_D)$  and acceptor dyes  $(n_A)$  that are attached to the single molecule where  $FRET = \frac{n_A}{n_A + n_D}$ .

47 In our simulations, we adjusted the photon emission rates of donor and acceptor dyes according to known 48 total experimentally observed photon counts such that the low and high FRET signals are equivalent to ~20 % and ~55 % FRET, respectively. For the electrophysiology measurements, based on the occupied 49 states of the single molecule's trajectory, we generate measurements. Specifically, low (closed channel 50 51 conformation) and high (open channel conformation) current levels are adjusted according to the 52 experimentally observed conformational state levels in electrophysiology experiments where low and high 53 current levels are  $\sim 0$  pA and  $\sim 2$  pA. Both smFRET and electrophysiology measurements are acquired with a camera-based detector<sup>61</sup> thereby the measurements are subject to temporal averaging provided that the 54 conformational dynamics of the molecule are faster than the data acquisition rate labeled with  $\frac{1}{\Delta t}$  where  $\Delta t$ 55 denotes the data acquisition period. We use  $\Delta t^F$  and  $\Delta t^E$  for smFRET and electrophysiology experiments, 56 respectively. Here, we denote the trajectory of the molecule with  $\tau(t)$ , for all  $t \ge 0$ , where photon emission 57

rates for donor and acceptor dyes that lead to low and high FRET levels are denoted with  $\bar{\mu}_D = (\mu_1^D, \mu_2^D)$  and  $\bar{\mu}_A = (\mu_1^A, \mu_2^A)$ . We also note that for a given time  $t, \tau(t) = S_l$  for some l = 1, 2, ..., 8. Low and high current state levels for the electrophysiology measurements are denoted with  $\bar{C} = (C_1, C_2)$ . We assumed shot noise for the smFRET measurements and normally distributed measurement noise for the electrophysiology experiments.

Now, we summarize how simulated measurements in smFRET and electrophysiology experiments aregenerated.

Based on Gillespie simulation, we generate  $\tau(.) \sim Gillespie(\overline{\bar{G}}, \tau(0))$ , where '~' stands for sampling of  $\tau(.)$ 65 66 as a Gillespie trajectory with a given starting point denoted with  $\tau(0)$  at time t = 0 that is determined based 67 on the steady state probability of state occupation explained above. Subsequently, given the photon emission rates  $\bar{\mu}_D$  and  $\bar{\mu}_A$  of donor and acceptor dyeswe can now formulate the collected photons associated 68 with donor and acceptor dyes. The measurements are labeled with  $\overline{w}_D = (w_1^D, w_2^D, \dots, w_N^D)$  and  $\overline{w}_A =$ 69  $(w_1^A, w_2^A, \dots, w_N^A)$  on the donor and acceptor dyes associated detectors in smFRET experimental setup 70 where  $w_n^D, w_n^A$  denote the n-th measurement for all n = 1, 2, ..., N with N labeling the last smFRET 71 measurement. Then, we now formulate the n-th measurement during the data acquisition period  $\Delta t^F$  within 72 time window of  $[t_{n-1}, t_{n-1} + \Delta t^F]$  such that  $w_n^D \sim Poisson(\int_{t_{n-1}}^{t_{n-1}+\Delta t^F} dt \, \mu_{\tau(t)}^D)$ 73 and the  $w_n^A \sim Poisson(\int_{t_{n-1}}^{t_{n-1}+\Delta t^F} dt \, \mu_{\tau(t)}^A)$  for all n = 1, ..., N. Here, "Poisson" denotes the distribution to 74 75 formulate the shot noise statistics in the measurements. Next, we move to the electrophysiology measurements labeled with  $\bar{y} = (y_1, y_2, \dots, y_M)$  where M labels the last measurement in 76 77 electrophysiology experiments. In the case of acquiring electrophysiological data, we assume that the 78 measurement noise is governed by "Normal" distribution and thereby the m-th measurement labeled with  $y_m$  collected with data acquisition rate  $\frac{1}{\Delta t^E}$  within the time window  $[t_{m-1}, t_{m-1} + \Delta t^E]$  shown as 79

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$$y_m \sim Normal(\frac{1}{\Delta t^E} \int_{t_{m-1}}^{t_{m-1} + \Delta t^E} dt C_{\tau(t)}), \sigma^2)$$
 for all  $m = 1, 2, ..., M$  where  $\sigma^2 (pA^2)$  denotes the measurement  
81 noise.

In our current study, the only noise sources are the shot noise due to photon emissions for smFRET and normally distributed measurement noise for electrophysiology. We have not considered the read-out or any jittering noises dictated by the type of detector used in the experiments. Those additional improvements related to the measurement models in our simulations that are left for future studies.

86 Supplementary Table 2. Kinetic rate parameters used in simulations. \*For simulation of E1371Q gating

 $k_{5 \rightarrow 6}$  was set to zero.

rates	units	values
$k_{1 \rightarrow 2}$	μΜ <sup>-1</sup> s <sup>-1</sup>	1
$k_{2 \rightarrow 1}$	s <sup>-1</sup>	50
$k_{2\rightarrow 3}$	s <sup>-1</sup>	10
$k_{3\rightarrow 2}$	s <sup>-1</sup>	0.3
$k_{1 \rightarrow 8}$	s <sup>-1</sup>	0.3
$k_{8 \rightarrow 1}$	s <sup>-1</sup>	0.3
$k_{3 \rightarrow 8}$	s <sup>-1</sup>	0.6
$k_{8 \rightarrow 3}$	$\mu M^{-1}s^{-1}$	0.012
$k_{3\rightarrow 4}$	s <sup>-1</sup>	1.3
$k_{4\rightarrow 3}$	s <sup>-1</sup>	0.002
$k_{5 \rightarrow 4}$	s <sup>-1</sup>	20
$k_{4\rightarrow 5}$	s <sup>-1</sup>	100
$k_{5 \rightarrow 6}$	s <sup>-1</sup>	2*
$k_{6 \rightarrow 5}$	s <sup>-1</sup>	0
$k_{6 \rightarrow 7}$	s <sup>-1</sup>	4
$k_{7 \rightarrow 6}$	s <sup>-1</sup>	0
$k_{7 \rightarrow 8}$	s <sup>-1</sup>	4.5
$k_{8 \rightarrow 7}$	s <sup>-1</sup>	0

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#### 101 Supplementary Figure 1 | Gel source data

102 Uncropped gel scans of the gel imaged in Extended Data Figure 2a. The gel was first imaged for LD555

103 and LD655 fluorescence, and then Coomassie-stained. The dashed boxes indicate the regions of the scans

104 that are presented in Extended Data Figure 2a.