

# Motor-driven advection competes with crowding to drive spatiotemporally heterogeneous transport in cytoskeleton composites

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## *Supplementary Material*

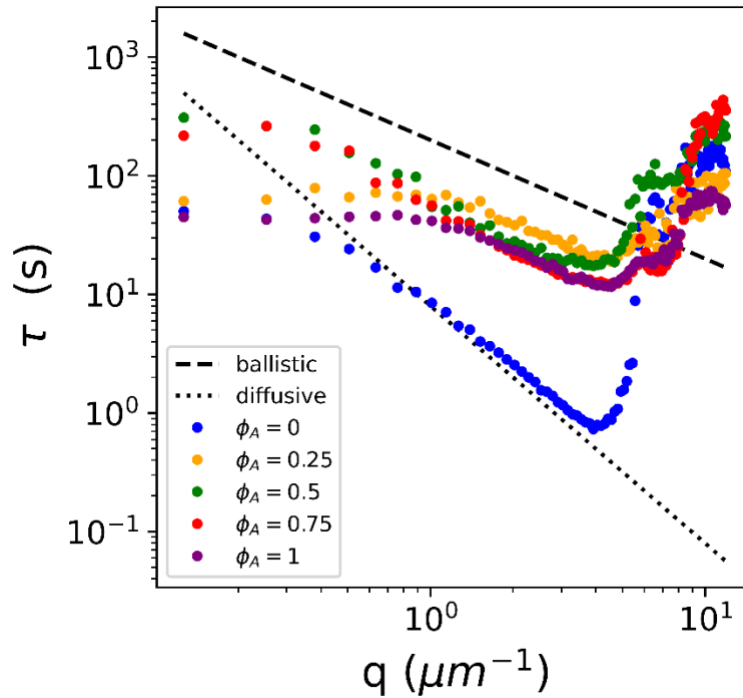
### 1 Supplementary Figures and Tables

**Figure S1:**  $\tau(q)$  versus wavevector  $q$  evaluated over the entire  $q$  range over which corresponding  $D(q, \Delta t)$  curves are fit.

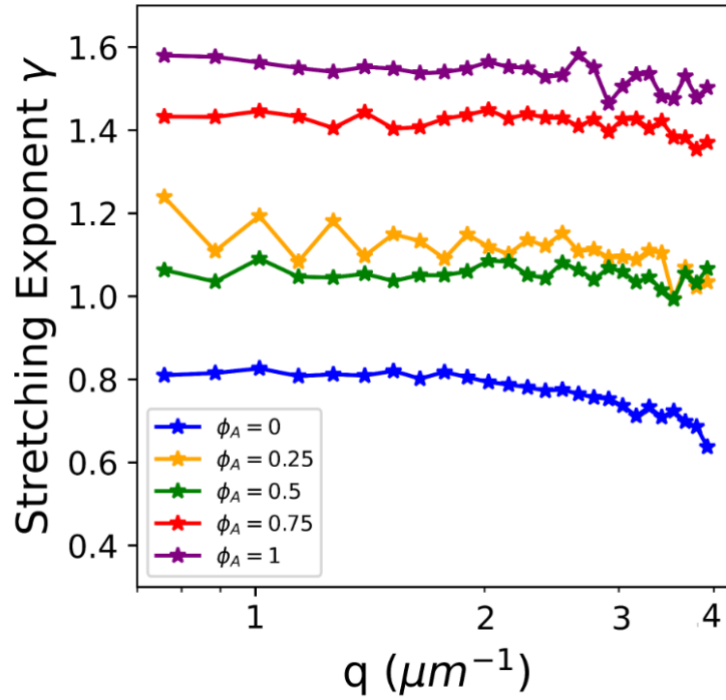
**Figure S2:** Stretching exponents  $\gamma(q)$  determined from fits to  $D(q, \Delta t)$  and plotted for all  $q$  values used to determine the power-law dependence of  $\tau(q)$ .

**Figure S3:** van Hove distributions of particle displacements in the  $x$ - and  $y$ - directions for all composites.

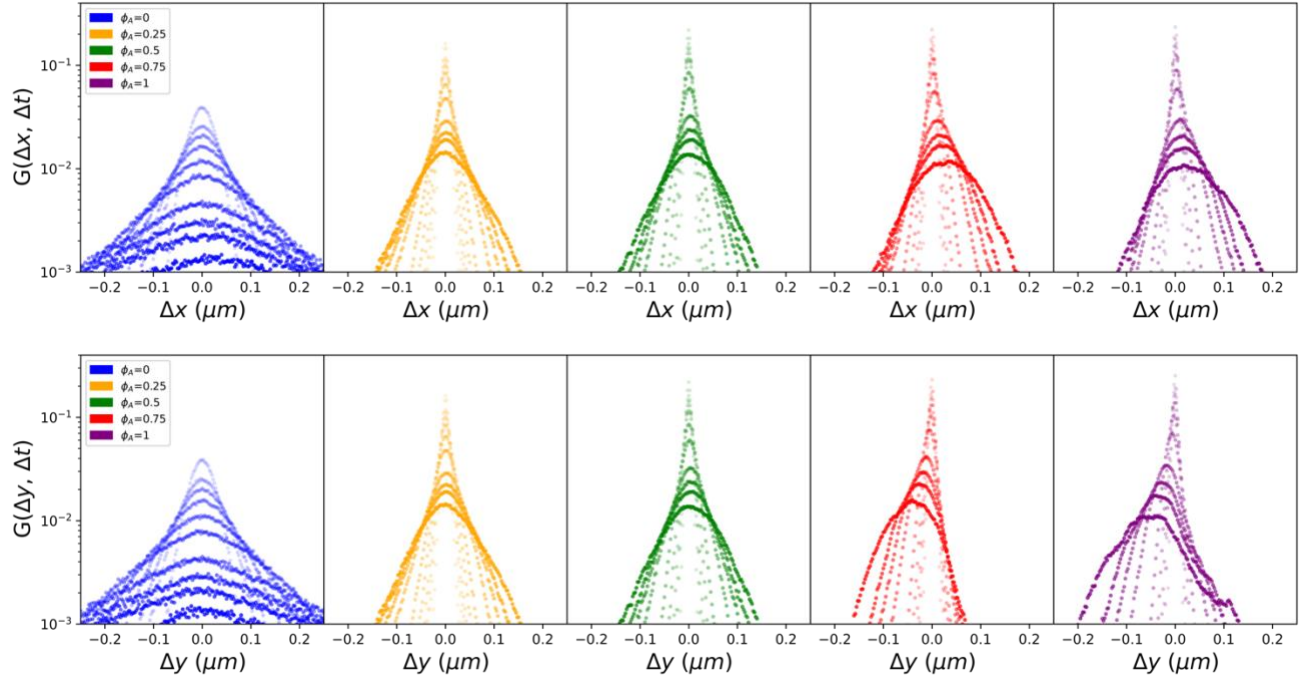
**Figure S4:** van Hove distributions for 11 individual measurements for  $\phi_A = 1$ .



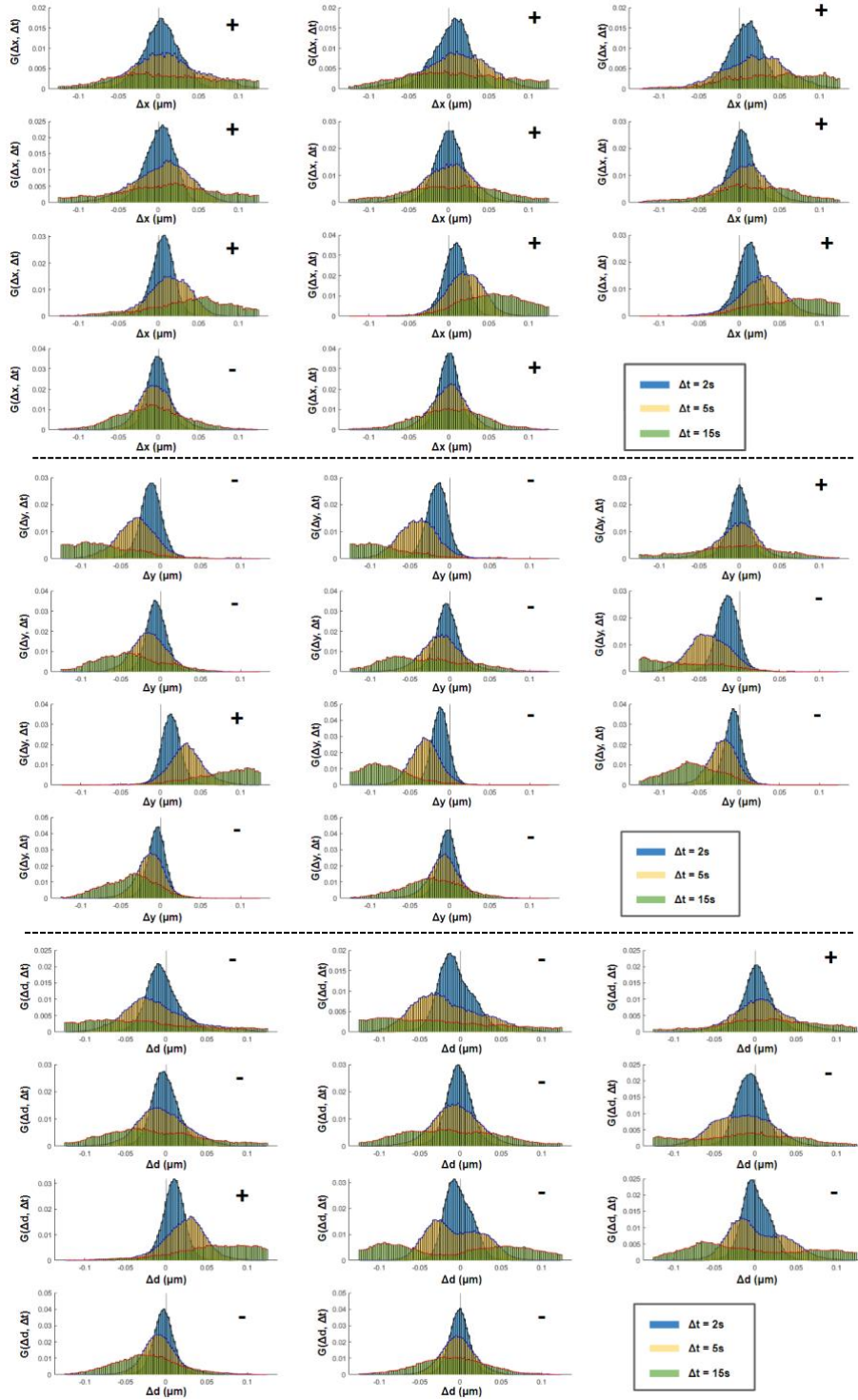
**Figure S1:  $\tau(q)$  versus wavevector  $q$  evaluated over the entire  $q$  range over which corresponding  $D(q, \Delta t)$  curves are fit.** Data shown is the same as in Fig 5B but for an extended range  $q = 0.16 - 16 \mu\text{m}^{-1}$ . Dashed and dotted lines correspond to ballistic and diffusive scaling exponents  $\beta = 1$  and 2, respectively. To determine scaling behavior of composites, we evaluate  $\tau(q)$  for  $q = 1 - 3.9 \mu\text{m}^{-1}$  over which power-law behavior is observed for all composites. The non-physical upticks in  $\tau(q)$  for  $q > 3.9 \mu\text{m}^{-1}$  are due to the optical resolution of our microscope. While the theoretical resolution limit is  $q \approx 10 \mu\text{m}^{-1}$  with an objective of NA=1.0, non-ideal imaging conditions, such as imaging across a capillary tube which has a refractive index not perfectly matched to that of the sample, reduces this limit to  $\sim 4 \mu\text{m}^{-1}$  in our setup. In the low- $q$  limit, the unphysical rollovers and plateaus in some of the data are due to a combination of the image size, the maximum lag time we probe, and noise. We analyze  $256 \times 256$  square-pixel images with a pixel size of  $\sim 0.1 \mu\text{m}$ , setting a minimum of  $q \approx 2 \mu\text{m}^{-1}$ . However, we are further limited in certain cases by the accessible time scales. Namely, density fluctuations at small  $q$  values are expected to slowly decay, and the maximum  $\Delta t$  over which we fit  $D(q, \Delta t)$  is  $\sim 100$  s, above which the data is prohibitively noisy to accurately fit due to low statistics.



**Figure S2: Stretching exponents  $\gamma(q)$  determined from fits to  $D(q, \Delta t)$  and plotted for all  $q$  values used to determine the power-law dependence of  $\tau(q)$ .  $\gamma(q)$  values for all composites are approximately  $q$ -independent, validating our power-law analysis of  $\tau(q)$ . Averaging over  $q$  for each composite yields the data shown in Fig 5G.**



**Figure S3: van Hove distributions of particle displacements in the  $x$ - and  $y$ - directions for all composites.** van Hove distributions  $G(\Delta x, \Delta t)$  (top) and  $G(\Delta y, \Delta t)$  (bottom) of particle displacements  $\Delta x$  and  $\Delta y$ , measured via SPT, for lag times  $\Delta t = 0.1, 0.2, 0.3, 0.5, 1, 2, 3, 5, 10, 15$  s denoted by the color gradient going from light to dark for increasing  $\Delta t$ . Each panel corresponds to a different composite demarked by their  $\phi_A$  value with color-coding as in Fig 3. Data shown is the same as that in Fig 3A separated into  $x$ - and  $y$ - direction distributions. For reference,  $x$ - and  $y$ - directions correspond to the narrow and long dimensions of the capillary sample chamber, respectively.



**Figure S4: van Hove distributions for 11 individual measurements for  $\phi_A = 1$ .** van Hove probability distributions  $G(\Delta x, \Delta t)$ ,  $G(\Delta y, \Delta t)$ , and  $G(\Delta d, \Delta t)$  (from top to bottom) for particle displacements  $\Delta x, \Delta y, \Delta d = \Delta x \cup \Delta y$  for each measurement of the  $\phi_A = 1$  composite. Each plot displays distributions for  $\Delta t = 2, 5, 15$  s with the dashed vertical line demarking zero displacement. The net direction of motion for each trial, positive or negative, is indicated in the upper right as + or -.  $G(\Delta x, \Delta t)$  and  $G(\Delta y, \Delta t)$  distributions are primarily in the positive and negative directions, respectively, with  $G(\Delta y, \Delta t)$  distributions displaying relatively larger deviations from zero.

