

## Zigzag-HMC explores the energy space more efficiently than BPS

In our experience, BPS tends to generate samples with high auto-correlation between their respective energy function evaluations  $-\log \pi(\boldsymbol{x})$ . In other words, it slowly traverses the target distribution’s energy contours even when the marginal dimensions all appear to demonstrate good mixing. A similar behavior has also been reported by [2], who introduce a velocity refreshment to address the issue. As we demonstrate below, however, even velocity refreshments cannot fully remedy BPS’s slow-mixing on the energy space.

We apply BPS and Zigzag-HMC to a 256-dimensional standard normal truncated to the positive orthant (all  $x_i > 0$ ). We run both samplers for 2000 iterations where per-iteration travel time is one unit time interval and repeat the experiments for 10 times with varying initial values. For BPS we include Poisson velocity refreshments to avoid reducible behavior and set the refreshment rate to an optimal value 1.4 [3]. At every iteration we refresh Zigzag-HMC’s momentum by redrawing it from the marginal Laplace distribution. Both samplers have no problem sampling from the target distribution and the minimal ESS across all dimensions are  $158 \pm 25$  (mean  $\pm$  SD) for BPS and  $207 \pm 21$  for Zigzag-HMC, estimated from the last 1000 samples of the MCMC chains across 10 runs. As a sanity check, the average sample mean and variance are  $(0.800, 0.365)$  for BPS and  $(0.798, 0.363)$  for Zigzag-HMC, close to the analytical values — the univariate marginal distribution of our truncated standard normal is a truncated normal with mean  $2/\sqrt{2\pi} \approx 0.798$  and variance  $1 - 2/\pi \approx 0.363$  [4].

However, Zigzag-HMC returns a clear win over BPS in the mixing of joint density (Figure 1). The sampling inefficiency for  $-\log \pi(\mathbf{x})$  is less of a problem if one only needs to sample from a truncated normal with a fixed covariance matrix, but we are keenly interested in sampling the covariance matrix as a target of scientific interest. In this context, inefficient traversal across energy contours harms the sampling efficiency for all model parameters.

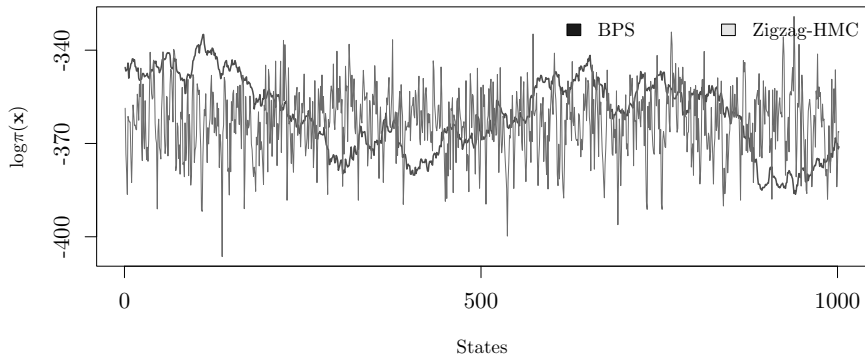


Figure 1: Trace plot of the log density of a 256-dimensional truncated standard normal sampled by BPS and Zigzag-HMC for 1000 MCMC iterations.

We can provide an intuition for BPS’s slow movement in energy space. Assume the  $d$ -dimensional parameter at the  $t$ th MCMC iteration is  $\mathbf{x}(t) = (x_1(t), \dots, x_d(t)) \in \mathbb{R}^d$ ,  $t = 1, \dots, T$ , with  $T$  being the total number of iterations. For a truncated standard normal, its log density  $\log \pi(\mathbf{x}) \propto \sum_i^d x_i^2$ , and a high auto-correlation suggests  $\log \pi(\mathbf{x})$  changes little between successive iterations, that is, the squared jumping distances

$$J_D = \left[ \sum_i^d x_i^2(t+1) - \sum_i^d x_i^2(t) \right]^2, \quad t = 0, \dots, T-1$$

are small. We then decompose  $J_D$  into two components

$$\begin{aligned}
 J_D &= J_1 + J_2, \\
 J_1 &= \sum_i^d [x_i^2(t+1) - x_i^2(t)]^2, \\
 J_2 &= \sum_{j \neq k}^d [x_j^2(t+1) - x_j^2(t)] [x_k^2(t+1) - x_k^2(t)], \quad t = 0, \dots, T-1,
 \end{aligned} \tag{1}$$

where  $J_1$  measures the sum of the marginal travel distances and  $J_2$  the covariance among them. We compare  $J_D$ ,  $J_1$  and  $J_2$  between BPS and Zigzag-HMC in the aforementioned experiments. Clearly seen in Table 1, BPS yields a much lower  $J_D$  than Zigzag-HMC because its  $J_2$  is largely negative, suggesting strong negative correlation among the coordinates.

Table 1: Squared jumping distance ( $J_D$ ) of  $\log \pi(\mathbf{x})$  sampled by the bouncy particle sampler (BPS) and Zigzag Hamiltonian Monte Carlo (Zigzag-HMC). We report the empirical mean of  $J_1$  and  $J_2$  in their means and standard deviations (SD) across ten independent simulations with  $T = 1000$  after burn-in samples. Both samplers have a per-iteration travel time 1.

Quantity	BPS		Zigzag-HMC	
	mean	SD	mean	SD
$J_D$	8.3	0.6	583	21.7
$J_1$	521	16.3	564	2.6
$J_2$	-513	16.0	18.6	20.9

## References

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