

1 **Supplementary Materials**

2 ***Sample Preparation***

3 ***AuFON***

4 We oxygen plasma cleaned glass coverslips (25 mm diameter, #2) to improve surface hy-
5 drophilicity and remove surface contaminants, rinsing the glass coverslips thoroughly with
6 ethanol and deionized water both before and after the plasma etching. We washed a solu-
7 tion of 800 nm diameter silica microspheres (NanoCym, 5% by volume) in Millipore ultrapure
8 water (18 M Ω) and re-diluted to a 5% by volume. We then manually dispersed 25 μ L of this
9 solution across the surface of the glass coverslip. We allowed the silica microsphere solution to
10 evaporate under ambient conditions, forming a close-packed monolayer. (24) We then deposited
11 \sim 200 nm of gold on the silica-coated substrates using a thermal evaporation system at a rate of
12 2 \AA /s, with samples rotated at 16 RPM. (CHA Industries, SEC-600). An image of the AuFON
13 substrates is shown in Figure S1A. We prepared a 1 mM solution of methyl viologen dichloride
14 hydrate (Sigma Aldrich, 98%, 75365-73-0) in ethanol (Decon, 95%) and deposited 40 μ L on
15 the AuFON for ultrafast SERS measurements.

16 We measured the localized surface plasmon resonance (LSPR) extinction spectrum of the
17 AuFON with a Shimadzu UV-2600 equipped with an integrating sphere attachment (Shimadzu,
18 ISR-2600). The data acquisition time for each collection was 4 milliseconds. A representative
19 example substrate and extinction spectrum are shown in S1. We collected extinction spectra
20 for all experimental samples before analyte addition, after analyte addition, and after ultrafast
21 experiments.

22 ***AuFON Electrode***

23 We prepared the AuFON working electrode pictured in Figure 1C as follows. We cut prepared
24 AuFON substrates into \sim 5 x 10 mm pieces. We then adhered these pieces to a similar sized

25 piece of 1 mm thick glass slide (Fisherbrand, 12-549-3) with two part epoxy (Loctite Quick
26 Set Epoxy, HC01-81501) and allowed it to cure for 1 hour. We flattened the end of a piece of
27 14 gauge copper wire and adhered it to one side of the exposed AuFON surface with two part
28 epoxy and allowed to cure for an additional hour. We verified continuity between the copper
29 wire and AuFON surface using a multimeter and encapsulated the remaining exposed copper
30 wire with two part epoxy, allowing the electrode to cure fully in a vacuum desiccator overnight
31 before using in SEC-SERS experiments.

32 ***Spectroelectrochemical SERS Instrumentation***

33 For spectroelectrochemical measurements, we used a simple electrochemical cell coupled to a
34 potentiostat (Pine Research, WaveNow Wireless Potentiostat/Galvanostat). We performed all
35 measurements in 0.1 M methyl viologen dichloride in 0.1 M KCl (H₂O) electrolyte solution.
36 The working electrode was a homemade AuFON electrode (S2), with a Ag/AgCl reference and
37 gold counter electrode. We obtained cyclic voltammograms of this system at sweep rates from
38 50 to 1000 mV/s. We stepped the potential in increments of 25 mV from 0-1200 mV, resting for
39 60 s between steps to allow the system to equilibrate. We collected spontaneous SER spectra
40 using the system described above (probe only), with an average power of 25 mW, and 60 s
41 acquisition time.

42 ***Ring Deformation Peak Ratio Calibration***

43 In order to translate changes in the ultrafast SERS MV⁰ : MV^{•+} ring deformation peak ratio
44 into relative redox population changes and plasmon-generated potentials, we created a set of
45 calibration curves. First, we correlated the potential dependent SEC-SERS MV⁰ : MV^{•+} ring
46 deformation peak ratios to the $E_{r,2}^0$ reduction peak from the AuFON-MV CV. As stated in the
47 main text, we assume that changes in the $E_{r,2}^0$ reduction peak amplitude between -0.85 and -1.06

48 V correspond to a change in the MV^0 population, from ~ 0 to 100%. This assumption ensures
49 that our quantified values are, at worst, underestimated. We approximated the $E_{r_2}^0$ reduction
50 peak as a Gaussian distribution, onto which we mapped the percentages for MV^0 (S2). From
51 here, we were able to correlate the peak ratio to both % MV^0 (S3) and potential (S4). We
52 created the two necessary calibration curves for translating MV^0 : $MV^{\bullet+}$ ring deformation peak
53 ratios into % MV^0 and applied potential by fitting the linear regime of these correlative plots.
54 The equations for these fits are shown on S3 and S4 respectively. Error in the y-axis for S3
55 corresponds to the propagated uncertainty in the Gaussian fit of the $E_{r_2}^0$ reduction peak and
56 error in the x-axis for S3 and S4 corresponds to the propagated uncertainty in the SEC-SERS
57 peak fits.

58 ***Ultrafast SERS Instrumentation***

59 Detailed descriptions of our ultrafast SERS instrument have been published previously. (15–
60 18) Briefly, we used a Yb-doped fiber-based amplified laser source (Clark-MXR Impulse) to
61 produce <250 fs pulses centered at ~ 1040 nm at a 2.04 MHz repetition rate, with an average
62 power of 15 W (S6). We generated the picosecond probe pulse by sending the incident beam
63 through a spectral filter constructed in a retroreflecting geometry, which contains a transmission
64 grating (LightSmyth Technologies LSFSG-1000-3212-94) to disperse the beam and a 100 mm
65 focal length cylindrical lens (Newport CKX100AR.16) to image the beam onto a slit. (56) We
66 determined the spectral FWHM of the filtered output to be 17.9 cm^{-1} , as measured from the
67 1003 cm^{-1} peak of toluene. We used a 100 mm focal length achromatic lens (Edmund Optics
68 49-374) to focus the pump and probe beams onto the sample to radii of approximately $180\text{ }\mu\text{m}$
69 and $160\text{ }\mu\text{m}$, respectively. The probe power was 10 mW and the pump power ranged from 5 to
70 11 mW. After the sample, we blocked the pump transmission with a spatial filter, and collimated
71 the spontaneous SERS signal with a 60 mm focal length lens (Thorlabs, LA1134-C). We used

72 a 100 mm focal length achromatic lens (Thorlabs, AC254-100-C) to focus the SERS signal
73 into a 0.3 m spectrograph (Princeton Instruments SP2300i) equipped with a 300 gr/mm, 1 μ m
74 blaze grating. We placed a 1064 nm Raman filter (Edmund Optics 47-510) directly before the
75 spectrograph to remove any remaining \sim 1040 nm light from the SERS signal. We used a 1024
76 pixel, liquid nitrogen cooled InGaAs photodiode array (Princeton Instruments PyLoN-IR 1.7)
77 for detection of the Stokes-shifted SER scattering, with an integration time of 20 s for each
78 ultrafast SER spectrum. A schematic depiction of the described setup is shown in S5.

79 *Data Collection and Analysis*

80 *Data Collection and Pre-Processing*

81 As described in the main text, we acquired 10 data sets at each incident average pump power,
82 ranging from 5 - 11 mW across 4 substrates. We set the average probe power to \sim 10 mW (flux,
83 13.6 W/cm²) for all experiments. We calculated the pump flux (4.87 - 10.6 W/cm²) using the
84 measured pump beam diameters listed in the previous section. Here, we outline the general
85 data analysis process used in this work. All analysis was done in IGOR Pro (v7.0.8.1), unless
86 otherwise stated.

87 For ultrafast measurements, we used a mechanical delay stage (Newport XMS50) in the
88 pump line to control the temporal offset of the pump and probe. We controlled spectral
89 acquisition using a custom Labview code, which automatically randomizes predetermined
90 stage positions (time points) collected for each experiment. We collected three spectra at
91 each timepoint: pump off– or, ground state (GS), -50 ps offset (n50), and the predetermined
92 time offset, or excited state (ES). Commonly, our acquired spectra exhibit high frequency
93 etaloning, which we remove by applying a Fourier filter. We then scaled each spectrum by the
94 incident probe power and acquisition time. We performed an optical Kerr effect measurement

95 in a 2 mm path-length cuvette of toluene to determine our pump and probe cross-correlation.
96 From this measurement, we calculate the time resolution to be 1.9 ps – the FWHM of the
97 cross-correlation (S7). We note that, on the AuFON, the SERS background differs substantially
98 between the pump-off (GS) and pump-on (n50 and ES) spectra due to scattering from the
99 pump pulse off of the rough and reflective surface. In order to account for the difference in
100 pump-induced SERS background, we describe the ground state in our analysis herein as the
101 n50 spectrum.

102

103 *Data Analysis*

104 In some cases, the 5 minute wait time prior to running experiments was not long enough to
105 allow for full equilibration of the AuFON-MV system, resulting in spectral anomalies at early
106 acquisition times. To account for this issue, we removed the first three acquisitions for all
107 experiments. Because the time points are randomized and many replicates were taken, there
108 was still ample data for analysis at each time point. Using a custom macro, we fit the 1030
109 cm^{-1} and 1070 cm^{-1} ring deformation modes of MV^0 and $\text{MV}^{\bullet+}$ as a pair of Gaussian peaks
110 for every pre-processed ES and corresponding n50 spectrum (S8, A/B). We calculated the peak
111 ratio for the ES and n50 individually from the fit amplitudes, and calculated the change in peak
112 ratio (Δ_{n50} peak ratio) as $\text{peak ratio}_{ES} - \text{peak ratio}_{n50}$. For the kinetic representation in figure
113 3A, we fit the full set of time-dependent mean Δ_{n50} peak ratios for each pump power to a
114 convolution of the Gaussian instrument response and single exponential decay using the Global
115 Fit module in IGOR Pro. For quantification of $\Delta \% \text{MV}^0$ and Δ potential, we fit each individual
116 data set with the same fit function, as shown in S8, C. Error bars here represent the propagated
117 peak amplitude fit error for the $\text{MV}^0 : \text{MV}^{\bullet+}$ peak ratio. We took the amplitude of the Gaussian
118 component of these fits as the maximum Δ_{n50} peak ratio (grey dots, S9). Outliers from each

119 pump power data set were detected with a generalized extreme studentized deviate (ESD) test
120 from the PyAstronomy package in Python 3. We subsequently removed these outliers from
121 each data set. We then performed a one-sided ANOVA analysis to determine significance of
122 the linear trend of Δ_{n50} peak ratio with pump flux (Table 1; $F=8.21$, $P=0.00034$), and a Tukey
123 (HSD) test to analyze the significance of the offset between individual sets (Table 2). Next, we
124 averaged the peak ratios for all 10 experiments at each pump power to calculate a set of mean
125 Δ_{n50} peak ratios, with error propagated from the ultrafast/SEC-SERS peak fits (y-axis) and the
126 measured pump powers/beam area (x-axis) (blue dots, S9). Subsequently, we used these mean
127 values to calculate the $\Delta\%MV^0$ and potential in Figure 3 using the linear fit equations in S3 and
128 S4.

129 For analysis of the long-lived reduced MV^0 population, we averaged the mean of the Δ_{n50}
130 peak ratios for the 30, 40, and 50 ps time points for each data set. We calculated the average
131 of this mean value for all experiments at each pump power to give the average mean Δ_{n50} peak
132 ratio for $t=30-50$ ps, with error propagated from the ultrafast/SEC-SERS peak fits (y-axis) and
133 the measured pump powers/beam diameters (x-axis) (S10). We performed a one-sided ANOVA
134 analysis to determine if the linear relationship between the average mean Δ_{n50} peak ratio ($t=30-$
135 50 ps) and pump flux was statistically significant (Table 3; $F=2.81$, $P=0.058$) – although not
136 significant, there is a perceivable positive trend between the sets. We then performed individual
137 T-tests on each set to analyze the significance of their offset from zero (Table 4). Similarly as
138 described above, we converted these peak ratios to $\Delta\%MV^0$ and potential using the linear fit
139 equations in S3 and S4.

140 *Ground State Addition Spectra for Visualization of Transient Features*

141 In order to prepare a qualitative visualization of the transient changes to the $MV^0 : MV^{*+}$
142 ring deformation peak ratio in our ultrafast SERS experiments, we display transient spectra

143 shown in Figure 2 of the main text in a form that is referred to as ground state addition spectra
144 (GSA). For each time point, we performed a 1:1 subtraction of the ES and n50, creating a
145 depletion spectrum. Using a custom fitting macro, we first fit a characteristic depletion feature
146 to a negative Gaussian line shape. Using the amplitude of that fit, we calculated a percentage
147 of the n50 to add back to the depletion spectrum so that the amplitude at that specific peak
148 frequency would be equal to zero. The new spectrum created in this process is called a GSA
149 spectrum. Doing this at every time point emphasizes positive changes in the transient ultrafast
150 SERS spectra, as shown in Figure 2 of the main text. This preparation is only used to visualize
151 small transient changes in SERS spectra, and is not used for further quantification.

152 *Fitting MV⁰ Excited Vibrational State Feature*

153 In order to approximate the time dependent behavior of the feature red-shifted from the 1030
154 cm⁻¹ MV⁰ ring deformation peak, we fit the shoulder and 1030 cm⁻¹ peak to a pair of Gaussian
155 distributions for the n50 and ES data sets for one experiment in the 10.6 W/cm² pump power
156 series. The frequency of the shoulder feature was held at 1017 cm⁻¹, in accordance with our
157 anharmonic DFT frequency analysis (Table 5). We calculated the 1017 cm⁻¹ and 1030 cm⁻¹
158 peak areas from the fits. We then calculated the percent contribution of the 1017 cm⁻¹ peak
159 to the total area of both peaks for both the ES and n50. Finally, we took the difference of the
160 percent contribution of the 1017 cm⁻¹ peak for the ES and n50 spectra, giving the plot shown
161 in S12. Here, we see an approximate change in contribution for the 1017 cm⁻¹ peak at time
162 zero of 3%.

163 From the same data set, we calculated an approximate increase in MV⁰ of 3% from the
164 1030:1070 cm⁻¹ peak ratio using methods described above. Using this value, we can extrapo-
165 late that there is an additional 0.1% increase in the MV⁰ population owed to MV⁰ molecules in
166 excited vibrational states.

167 *DFT Calculations of MV*

168 In order to confirm the vibrational assignments of the Raman modes observed in the ultrafast
169 and SEC SERS experiments, we referenced literature values (34, 35) and performed DFT-based
170 optimization and vibrational frequency calculations for MV^0 , $MV^{\bullet+}$, and MV^{2+} with Gaussian
171 16 using the B3LYP functional and 6-31G** basis set. In these calculations, we included a Self-
172 Consistent Reaction Field (SCRF) using the default Polarizable Continuum Model and integral
173 equation formalism variant. We attempted to make calculations including the gold surface us-
174 ing higher-level models, but the results significantly deviated from experimental observations.
175 The level of theory used here is acceptable for calculations of radical species, (57–59) and the
176 calculated frequencies match closely to experimental observations. From our calculations, we
177 can assign the 990/1020 cm^{-1} (SEC-SERS) and 1030/1070 cm^{-1} (ultrafast SERS) peak pairs
178 to a C-N stretch/in-plane ring deformation motion. Figure S11 shows the scaled (0.961) cal-
179 culated spectra for MV^0 (yellow), $MV^{\bullet+}$ (red), and MV^{2+} (blue), with peaks for the calculated
180 frequencies broadened to 18 cm^{-1} to match our experimental parameters. Anharmonic fre-
181 quency calculations were performed using the same functional and basis set, with omission of
182 the SCRF.

183 *MV Surface Concentration Estimation*

184 We were able to estimate the number of MV molecules deposited onto the AuFON surface to
185 an upper limit using a molecular packing density of 1×10^{-10} mol cm^2 for pyridine. Assuming
186 MV is adsorbed parallel to the AuFON surface and has an area comparable to two adjacent
187 pyridines, this is a reasonable order of magnitude estimation. (52) In order to account for the
188 nanostructure of the AuFON, we used a geometric model to calculate a scaling factor of 2.6 for

189 the probed area (F_{sc}) – assuming a monolayer of half prolate ellipsoids.

$$F_{sc} = \frac{\pi\sqrt{3}(2r+h)}{6\sqrt{1-\frac{r^2}{(r+h)^2}}} \sin^{-1} \sqrt{1-\frac{r^2}{(r+h)^2}} \quad (1)$$

190 Here, r is the radius of the silica nanospheres (400 nm) and h is the thickness of the gold layer
191 (200 nm). Using this, we estimate the number of MV molecules within the probed region to be
192 $\sim 2 \times 10^{11}$. Assuming half of the starting molecules are $MV^{\bullet+}$, then for the reduction of $MV^{\bullet+}$
193 to MV^0 , we observe a transient conversion on the order 10^9 of molecules at time zero in the
194 ultrafast SERS experiments, and a conversion of approximately one molecule per pulse for the
195 MV^{2+} to $MV^{\bullet+}$ reduction in the steady state. If we have $\sim 3 \times 10^{16}$ photons per pump pulse
196 (10.6 W/cm^2 pump flux), this translates to a quantum yield for reduced species on the order of
197 $10^{-5} \%$ and $10^{-14} \%$ respectively.

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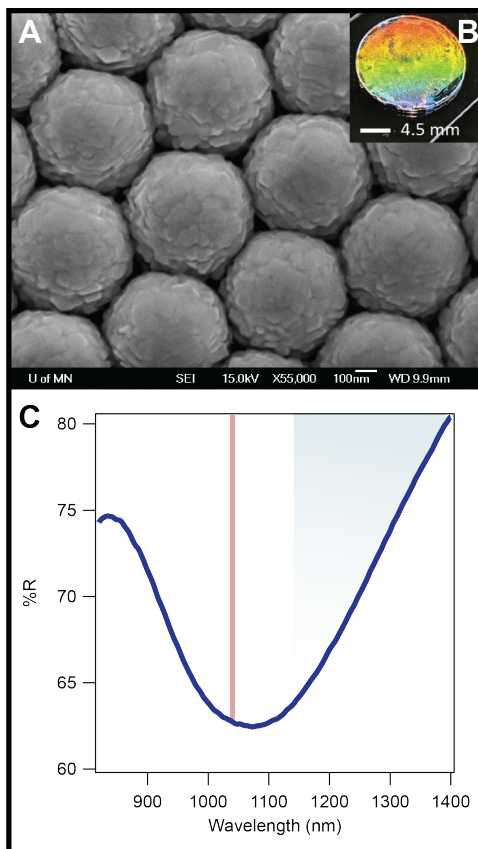


Figure S1: A) SEM Image of AuFON substrate. B) Full-scale image of AuFON substrate. C) Reflection spectrum of AuFON used in these experiments, with maximum extinction given as the minima of the reflectance spectrum at 1075 nm. The red line shows our excitation wavelength at ~ 1040 nm, and the green shaded region shows the region of interest for detecting SER scattered photons.

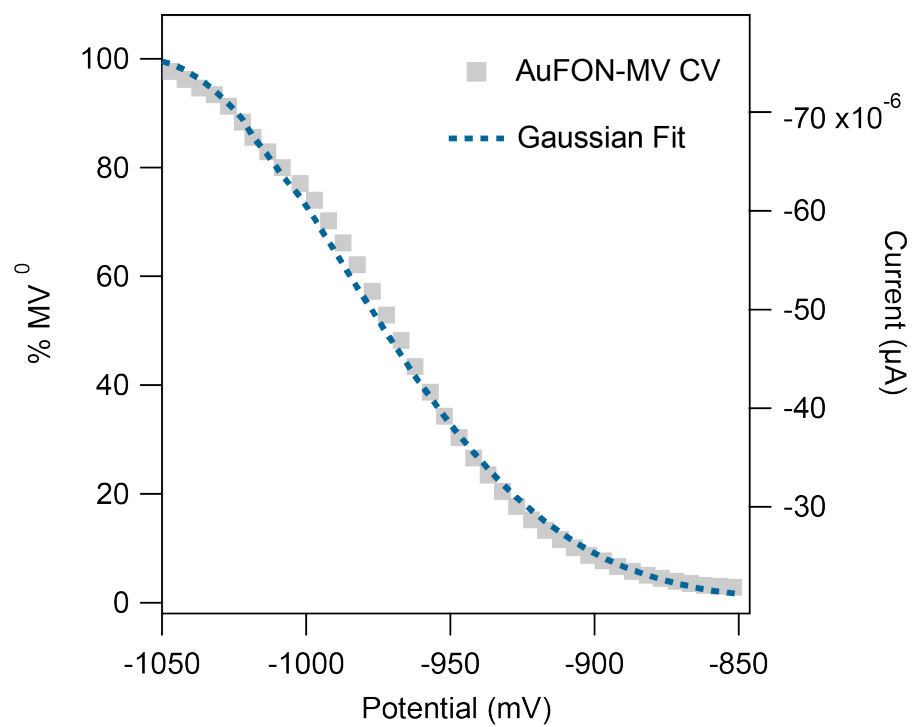


Figure S2: Correlation of [MV⁰] with potential.

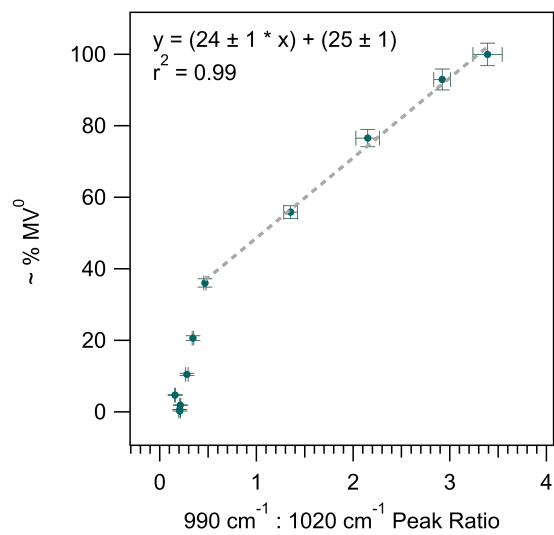


Figure S3: Calibration curve, translating SEC-SERS ring deformation peak ratio to % MV⁰

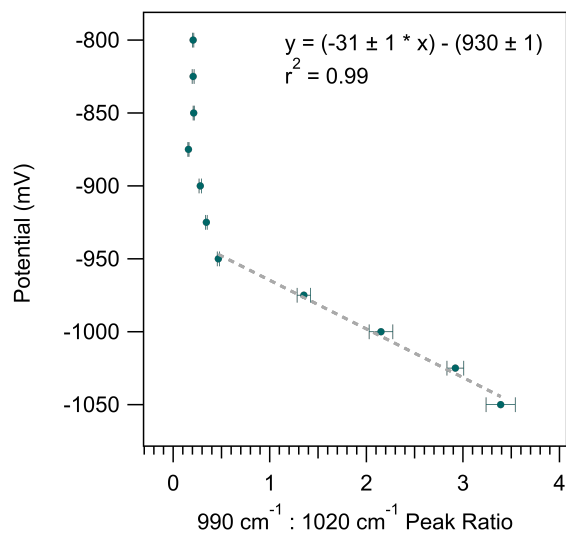


Figure S4: Calibration curve, translating SEC-SERS ring deformation peak ratio to potential

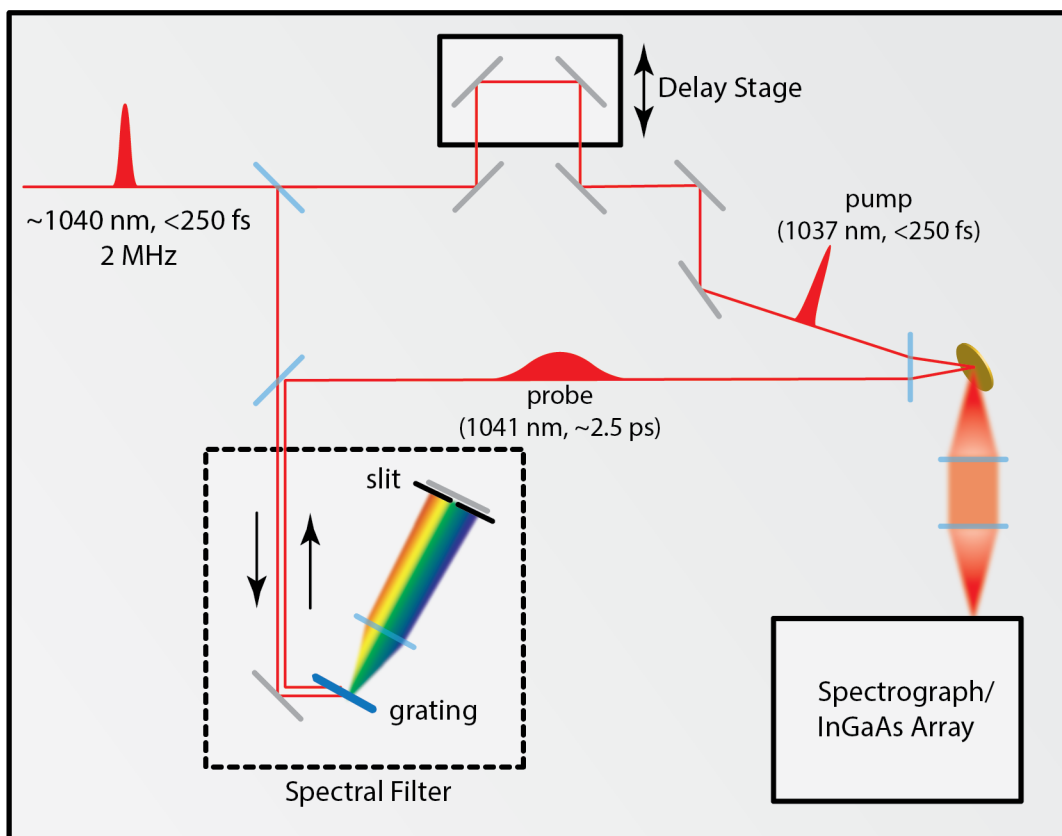


Figure S5: Schematic depiction of ultrafast SERS experimental setup.

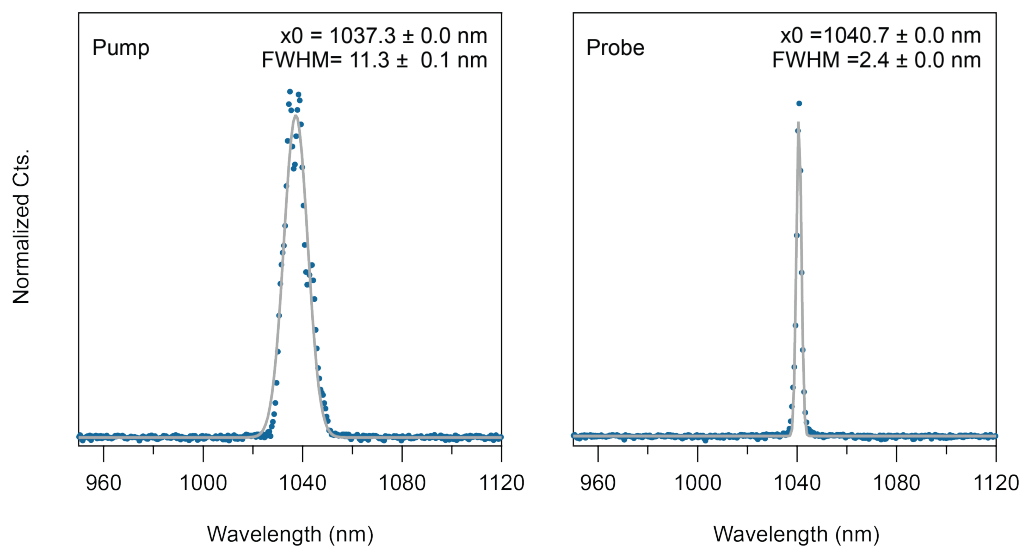


Figure S6: Spectra of pump and probe used in ultrafast SERS and SEC-SERS measurements. The pump is centered at 1037 nm, and the probe at 1041 nm.

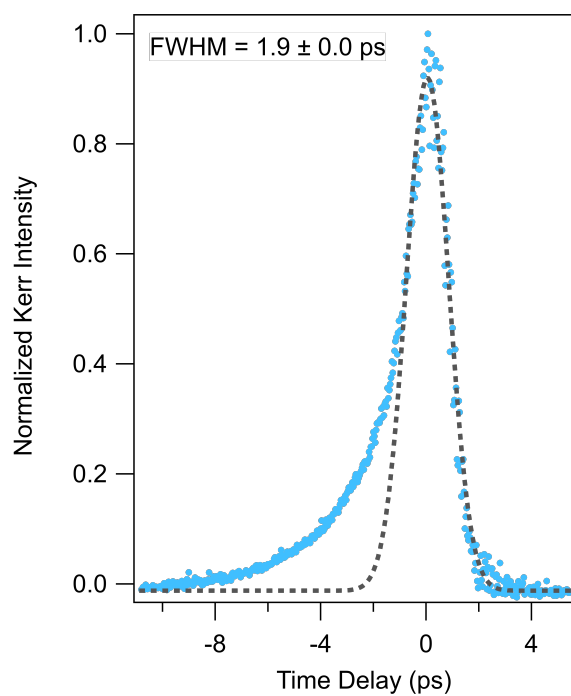


Figure S7: Representative pump/probe cross-correlation measurement using the optical Kerr effect. The peak is fit to a Gaussian lineshape, with a FWHM of 1.9 ps. The tail at negative timepoints is due to spatial chirp of the pump beam.

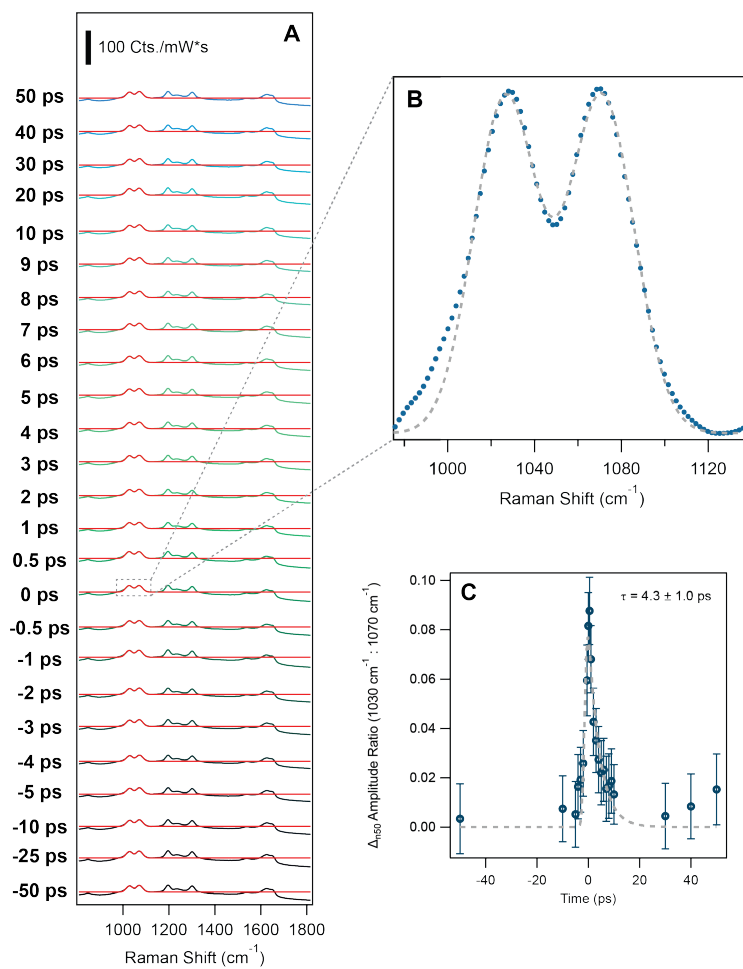


Figure S8: Examples from workup of one data set at 10.6 W/cm^2 pump flux. A) Individual fits of 1030 cm^{-1} and 1070 cm^{-1} peaks from ES spectra. B) Detailed image of fit for 0 ps spectrum. C) Resulting kinetic trace for this data set after subtracting fits from n50 spectra.

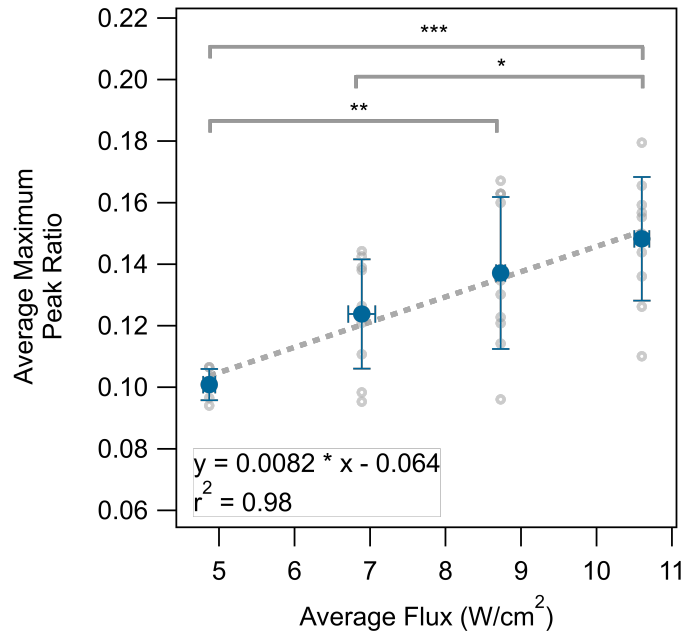


Figure S9: Maximum peak ratio at each pump power. Grey markers indicate individual experiments and blue markers represent the average of those experiments. The linear trend of the average peak ratio is significant (ANOVA [F=8.21, P=0.00034]). Significance of offset between individual experiments is given by the asterisks (* = $P \leq 0.05$, ** = $P \leq 0.01$, and *** = $P \leq 0.001$).

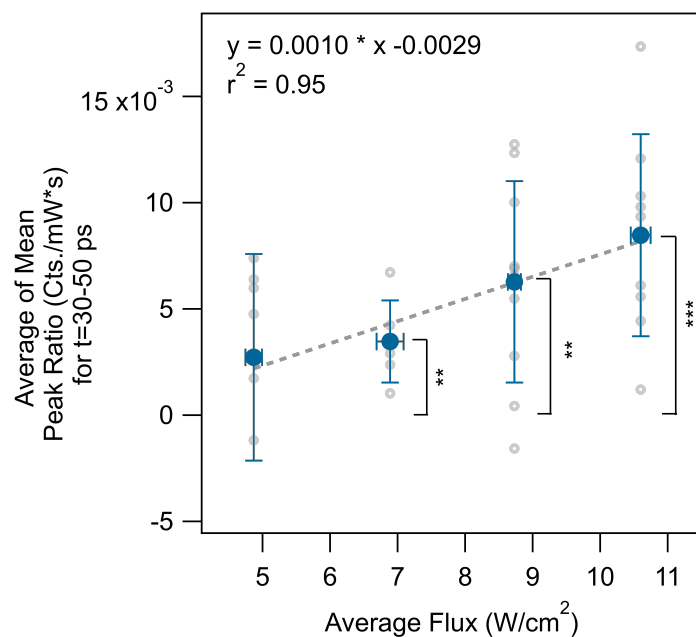


Figure S10: Mean of peak ratio for t=30 - 50 ps at each experimental pump power. Grey markers indicate individual experiments and blue markers represent the average of those experiments. The linear trend of the peak ratio is not significant, but there is a clear positive trend (ANOVA [F=2.81, P=0.058]). Significance of offset between individual experiments is given by the asterisks (* = $P \leq 0.05$, ** = $P \leq 0.01$, and *** = $P \leq 0.001$).

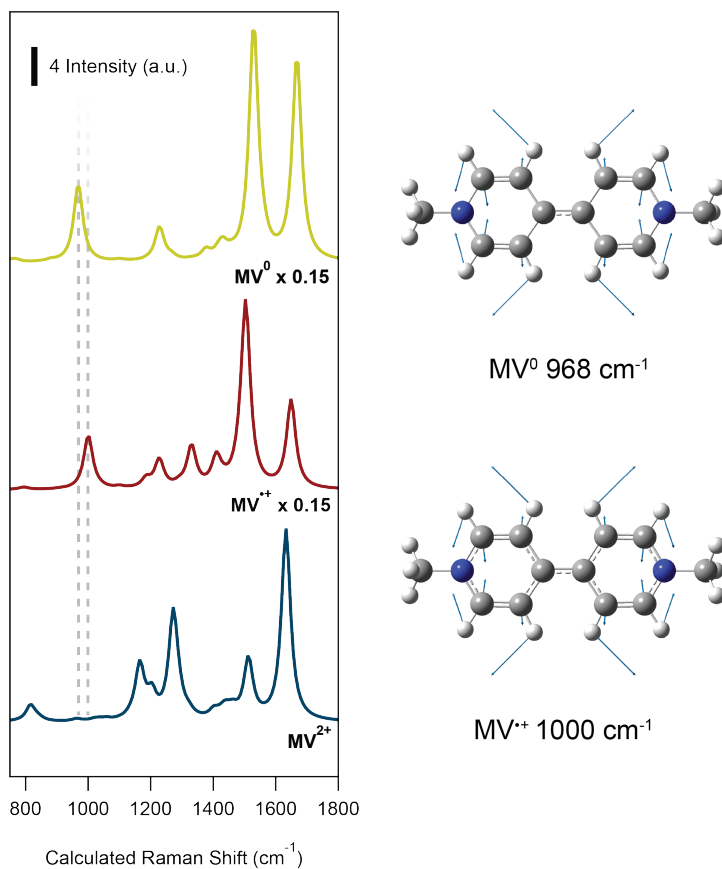


Figure S11: DFT calculated Raman spectra for MV^0 , $MV^{\bullet+}$, and MV^{2+} , showing the same ring deformation modes observed in our SERS experiment at 968 cm⁻¹ and 1000 cm⁻¹. The displacement vectors for the at 968 cm⁻¹ (MV^0) and 1000 cm⁻¹ ($MV^{\bullet+}$) modes are displayed, confirming that this is the same mode in both redox states.

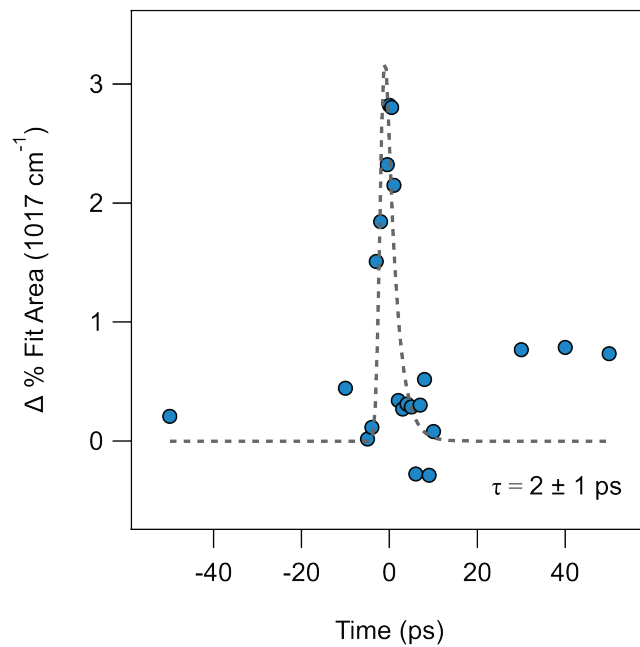


Figure S12: Typical kinetic trace showing the change in percent fit area for the 1017 cm^{-1} shoulder component of the 1032 cm^{-1} and 1017 cm^{-1} peak pair at 10.6 W/cm^2 pump power.

	DF	SS	MS	F	Fc	P
Groups	3	0.000931	0.00310			
Error	32	0.0121	0.000378			
Total	35	0.0214	0.000612	8.21	2.90	0.000341

Table 1: ANOVA analysis of data sets in Figure S9.

Comparison (W/cm ²)	Difference	SE	q	q0.05_32_4	Conclusion	P
10.6 -- 4.87	0.0474	0.00710	6.68	3.83	Significant	0.000251
10.6 -- 6.89	0.0244	0.00615	3.97	3.83	Significant	0.0400
10.6 -- 8.73	0.0111	0.00615	1.80	3.83	Not Significant	0.585
8.73 -- 4.87	0.0363	0.00710	5.11	3.83	Significant	0.00533
8.73 -- 6.89	0.0133	0.00615	2.17	3.83	Not Significant	0.430
6.89 -- 4.87	0.0230	0.00710	3.23	3.83	Not Significant	0.122

Table 2: Tukey (HSD) Test comparing data sets in Figure S9.

	DF	SS	MS	F	Fc	P
Groups	3	0.000163	5.45E-05			
Error	28	0.000543	1.94E-05			
Total	31	0.000707	2.28E-05	2.81	2.95	0.058

Table 3: ANOVA analysis of data sets in Figure S10.

flux (W/cm ²)	n	DF	avg	stdev	sxBar	t statistic	critical	P	Conclusion
4.87	10	6	0.00272	0.00486	0.00184	1.48	2.45	0.189	Not Significant
6.89	10	5	0.00346	0.00193	0.000787	4.40	2.57	0.00702	Significant
8.73	10	9	0.00627	0.00474	0.00150	4.18	2.26	0.00236	Significant
10.6	10	8	0.00846	0.00475	0.00158	5.34	2.31	0.000692	Significant

Table 4: TTests comparing offset from zero for data sets in Figure S10.

ω (cm ⁻¹)	993.66
$\Delta\omega$ (cm ⁻¹)	-16.99
I (km/mol)	0.0005
ΔI (km/mol)	0.0005

Table 5: Output of the anharmonic vibrational analysis of the MV⁰ ring deformation mode.