# **Science Advances NAAAS**

# Supplementary Materials for

# **Highly sensitive single-molecule detection of macromolecule ion beams**

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## **This PDF file includes:**

Supplementary Text Figs. S1 to S13 Table S1 References

#### 1. Mass spectrometry

#### Sample preparation and beam source

The following dyes, vitamins and proteins were purchased from Merck and prepared with concentrations, purity and solvents as listed in Table S1:



#### Table S1: Analytes and chemical parameters used in the detector experiments.

Vitamin B12 (Vit B12), rhodamine 6G (R6G), insulin bovine pancreas (INS), cytochrome C from bovine heart (CytC), myoglobin from equine skeletal muscle (Mb) and concanavalin A (Con A).



#### Figure S1: Gallery of molecules used in the experiments.

The shapes are meant to provide a comparison of size and complexity, but do not indicate the molecular conformation in transit through the mass spectrometer or during impact on the detector.

All molecules were electro-sprayed using an 80 µm needle operated at 3 to 4 kV. The solution is driven by a Harvard syringe pump at 9 µl/min. Optionally, the molecular charge is reduced in a corona discharge for which nitrogen 5.0 is passed through pure ethanol at elevated temperature and sent into a chamber with about 1 cm<sup>3</sup> volume. A sharp stainless steel needled at -3 kV induces the discharge which is actively stabilized at a current of 15 µA. The resulting bipolar air is guided into an aluminum channel of 3 mm inner diameter and 80 mm length where it mixes with the electrospray before entering the vacuum system of the mass spectrometer.

#### Data acquisition and mass calibration

An QMS-SSPD mass spectrum is recorded. A programmable pulse generator (PPG) triggers a scan of the quadrupole mass spectrometer, while the time to digital converter (TDC) collects the SSPD ion counts. Since the QMS amplitude scan is linear, the SSPD spectrum can be referenced to a calibrated spectrum, here recorded in situ using TOF-MS.

For calibration purposes we can alternatively steer the ion beam in the forward direction, where it is either detected using a conversion dynode at -10 kV, a phosphor screen and a photomultiplier

(QMS-PhS) or using a time-of-flight mass spectrometer with MCP detector at -13 kV (TOF-MCP). The TOF-MS spectrum is first calibrated using cesium iodide clusters and then used to calibrate the QMS-SSPD system, as shown in [Figure S2,](#page-2-0) for a mixture of cytochrome C, insulin, and myoglobin.



<span id="page-2-0"></span>

The mass spectrum of a protein mix of myoglobin (A), insulin (B) and cytochrome C (C) is compared with the mass spectra of the individual proteins using a calibrated TOF-MS. The peaks are used to calibrate the QMS-SSPD measurements, Figure 3 in the main text.

# 2. QMS-SSPD mass spectrum recorded using detector  $D_1$

I[n Figure S3](#page-3-0) we show an insulin mass spectrum recorded using detector  $D_1$ . While the system fulfills its function, the integrated surface of even an 8-pixel detector of  $D_1$  is still an order of magnitude smaller than a single pixel of detector  $D_2$ . This explains why all figures in the main text were recorded using the large area detectors. A new generation of experiments shall combine more than a hundred small or large detectors into versatile systems whose shape, resolution and area can be tailored to the needs of the experiment.



#### <span id="page-3-0"></span>Figure S3: QMS-SSPD spectrum measured with detector D<sub>1</sub>.

Mass spectrum for Insulin<sup>1+</sup> measured with 2 pixels of chipset  $D_1$ . The working principle is clearly corroborated. For mass spectrometry and molecular analysis, integration of hundreds of these small pixels will become important to improve detector area and signal-to-noise. The additional side peaks are due to adducts generated in the ESI source region. The data was smoothed using a 25-point moving average.

#### 3. Superconducting detector

The detector assembly is bolted onto the cold head of a pulse tube cooler (Sumitomo SHI R65p) which can be cooled to 3.7 K (nominal 900 mW at 4.2 K) with all conditioning and signal cables connected. The SSPD detector is shielded by two gold coated OFHC copper cylinders with a 10 mm entrance hole for molecules and light. Visible photons can be shielded by covering all windows and their detection can be suppressed by choosing the bias current appropriately.

Note, however, that the small detectors  $D_{1/(b)}$  are sensitive to single visible and near-infrared photons, and will therefore always see some background, while the larger wires of  $D_2$  require slightly higher energy to trigger a hot-spot and can be operated with dark count rates as low as 0.02 cps.

The superconducting nanowire detectors are connected to a constant current source (Single Quantum), which drives each detector pixel separately with 6 to 20  $\mu$ A for detector  $D_{1(b)}$  or 10 to 65  $\mu$ A for detector  $D_2$ . The signal after ion impact is guided via a bias-T to two subsequent 20 dB voltage amplifiers (Mini circuits), followed by the TDC. The cryogenic assembly is mounted on a motorized translation stage to either mechanically scan the molecular beam profile or to lift the Faraday detector into the beam.

# 4. Calibration of the ion beam profiler

The ion beam profile of Figure 6 in the main text was recorded using detector  $D_2$ . In the vertical direction, the profile can be scanned either by mechanically shifting the detector or by tuning the voltage of the ion deflector. Since the horizontal and vertical deflectors have identical geometries, the vertical calibration of the electrodes is also used for the horizontal axis. During the calibration process the deflector voltage was increased in steps of 10 V or 20 V, while the nanowire detector was mechanically moved in steps of 200 µm, with 10 µm accuracy over a distance of 10 mm, collecting a slice of the ion beam profile.



#### <span id="page-4-0"></span>Figure S4: Calibration of the SSPD ion beam deflector (1000 eV insulin).

The deflector calibration curve represents the barycenter and the uncertainties of a fit to the ion beam profile as a function of voltage applied to the electrode. The high linearity allows for an accurate representation of the ion beam profile in Figure 6 of the main text.

The mean and standard deviation of the center position of these slices are then determined and plotted against the deflector voltage, along with an upper bound for the error due to inaccuracies in the mechanical position of the SSPD (10  $\mu$ m). The result plotted in **Figure S4** shows a highly linear position-voltage dependence. The deflection voltage steps are chosen to produce a shift equivalent to the pixel-length of the detector (200 µm).

#### 5. Systematic uncertainties in the total detection yield

The total detection yield of our system was experimentally determined to be  $\eta = 0.62$ . If we assume a geometrical filling factor of 50%, based on an evaluation of the electron micrographs, we deduce an area-normalized detection yield of even 124%. We here address how to understand and further analyze this observation:

#### Detector geometry

• If the meander filling factor is larger than the specified 50% this will not modify the experimentally determined total quantum yield, but it will increase the apparent areanormalized quantum yield. Our measurements indicate that the effective surface would be 20% larger than deduced from scanning electron micrographs, which nominally define the meander with an accuracy of 3% in one dimension.

- The macroscopic wires that are leading to the detector are thinned out from several micrometers down to the nanowire width of 500 nm in  $D_2$ . At 100 eV of impact energy even wider wire areas may still contribute to the signal. Systematic test using a series of different chip geometries will be required to elucidate this uncertainty.
- There can be contribution to the area-normalized detection yield from ions striking the  $SiO<sub>2</sub>$ substrate in proximity to the nanowires. This effect may be relevant considering that even at the lowest impact energy of 50 eV the hot-spot diameter is estimated to be 35 nm. If energy is released to the  $SiO<sub>2</sub>$  substrate and if that couples to the superconducting film, the effective nanowire width may be increased. The relevance of this effect requires a new series of experiments with nanowires of different "dead zones".
- We can finally convolute the finite diameter of a protein with the nanowire geometry. This results in an increased effective detector width. However, for insulin this effect is minimal, contributing less than 1% when the line width is 500 nm.

# Source Stability

• If the protein flux varies during the normalization, this can change the total and area-normalized quantum yield. The flux was measured to be stable to within 3% over one hour.

# Signal ringing

• For fast electronics, as typical in SSPD research, it is conceivable that cable reflections generate ringing, which may make that signals are counted more than once. The signals in our setting are optimized for mass spectrometry, using large area detectors, and we have not found indication for ringing in these measurements.

#### Ion beam profile

• To exclude the potential deformation of the ion beam during deflection, a 2d beam profile was measured at 0 V and 80 V deflection voltage by mechanically moving the detector along the same axis. This excludes any distortion of the ion beam by the deflection [\(Figure S5\)](#page-5-0).



<span id="page-5-0"></span>Figure S5: Beam profile at different deflection voltages.

The beam profile stays the same after deflection. However, the count rate changed a bit over time.

# 6. Cryogenic low-noise amplifier for scalable devices

With the goal to scale up the detector area and to improve the detector resolution, it is interesting to explore multi-pixel arrays and new ways to control their readout electronics in a cryogenic environment.

A large-area detector such as  $D_2$  is usually driven by a bias current of 30-50  $\mu$ A which generates already a sizeable voltage pulse across a 50 Ohm impedance when the superconductor transitions to normal conductivity. However, the better resolving detector  $D_1$  is typically driven by  $\lt 6$   $\mu$ A and therefore generates only a 300 µV signal. Low-noise cryogenic amplifiers (LNA) are thus important to boost the signal-to-noise ratio and to reduce the rise time of the output pulse for future onboard digitization and time-tagging.

For that purpose, we have designed and realized a novel amplifier with high gain, high bandwidth, low noise, and low power consumption, using a heterojunction-bipolar transistor (HBT) in a BiCMOS process.



#### <span id="page-6-0"></span>Figure S6: Cryogenic on-board low-noise amplifier.

Integration of a SiGe low-noise amplifier (LNA) with the superconducting nanowire detector on a shared printed circuit board at T = 3.8 K. These detectors with onboard electronics shall be used to integrate larger pixel numbers and larger detection areas.

The amplifier shown in [Figure S6](#page-6-0) exhibits good noise and speed at room temperature and improves even its current gain by about 300 % upon cooling to T = 3.7 K (17,46,47) The device measures only  $480 \times 280 \ \mu m^2$  including the pads and it combines an average noise temperature of 5 K in the frequency range of 0.1 – 8.8 GHz. It has a good weak-signal linearity with well-matched input impedance  $(S11 < -10$  dB) and a gain  $> 33$  dB.

# Integration of the LNAs into detector  $D_{1b}$

To demonstrate this integrated concept, we have joined an array of 8 SSPD pixels of type  $D_1$  with eight of our new LNAs to realize eight detectors  $D_{1b}$  and embedded them into the cryogenic detector setup as illustrated in [Figure S7.](#page-7-0)



#### <span id="page-7-0"></span>Figure S7: Cryogenic low-noise amplifiers (LNAs) integrated with the SSPD chip.

Each SNWD is individually biased by a current drive at room temperature. (A) shows the eight-pixel readout printed circuit board (PCB) mounted onto the 3.7 K stage, as well as the biasing lines that consist of two bias lines per pixel and two supply wires for the LNAs, totaling 18 leads. (B) Detector D1 array, mounted onto the prototype PCB. (C) Wire bonded LNA unit.

## 7. Photocleavage QMS-SSPD spectra

Photochemistry in the gas phase has attracted recent interest as it promises to change the charge of proteins optically in free flight in high vacuum (45). It may even eventually allow the neutralization and reionization of proteins. In a proof-of-concept experiment, we here demonstrate the use of QMS-SSPD mass spectrometry to analyze the photochemistry of a tailored photocleavable Insulin compound that we had modified using a Ru-complex. The tagged compound and the cleavage mechanism are shown in [Figure S8.](#page-8-0)



<span id="page-8-0"></span>Figure S8: Mechanism for photo-cleavage of Ru-complex modified Insulin.

The functionalized insulin molecules were electrosprayed, guided into the quadrupole mass filter and detected by the SSPD.

Different to prior experiments (Figure 2), we have added a UV laser beam (355 nm, 500 kHz, 80  $\mu$ ) pulse energy, 10 ps pulse duration) to induce photo-cleavage of the tailored Ru-tag (Figure 1). Figure S9 shows the effect of that laser in the QMS-SSPD spectrum as the appearance of a new peak at Ins<sup>5+</sup> (emerging from [InsRu]<sup>6+</sup>) as well as a small peak at Ins<sup>6+</sup>.





An QMS-SSPD spectrum of Ru-complex modified Insulin, recorded (A) without and (B) with exposure to a laser beam at a pulse energy of 80  $\mu$  at 355 nm is shown. The continuous nature of the detector fits the quasi-continuous laser interaction at 500 kHz. The appearance of the charge reduced protein Ins<sup>5+</sup> and Ins<sup>6+</sup>. The data was smoothed using a 10-point moving average.

## 8. Synthesis and verification of the molecular compound

#### Synthesis of Ru-complex modified insulin

Ru-complex (6.0 mg, 6.0 μmol), TSTU (2.2 mg, 7.2 μmol) and DIPEA (4.2 µl, 24 μmol) were dissolved in 0.5 ml DMF under N2 with stirring for 1 h at room temperature. Then the reaction mixture was added into insulin (35 mg, 6.0 μmol) and DIPEA (4.2 μl, 24 μmol) in 8 ml DMF, keep it stirring overnight at room temperature. After purification by preparative HPLC, a purple solid was obtained (4.2 mg).





#### Verification of the Ru-complex modified insulin by mass spectrometry and NMR

Mass-to charge ratio m/z: calculated [M]: 6554.37, found [M]: 6554.53.



#### Figure S11: High resolution mass spectrum of the electrosprayed insulin compound.

Ru-complex: 1H NMR (500 MHz, CD3CN) δ 8.54 (d, J = 6.8 Hz, 2H), 8.52 – 8.49 (m, 2H), 8.44 – 8.30 (m, 4H), 8.04 (d, J = 8.4 Hz, 2H), 8.00 – 7.94 (m, 2H), 7.82 (t, J = 8.0 Hz, 1H), 7.78 (d, J = 6.7 Hz, 2H), 7.74 (d, J = 6.8 Hz, 2H), 7.59 (ddd, J = 7.6, 5.5, 1.3 Hz, 2H), 7.45 (d, J = 8.7 Hz, 2H), 6.83 (d, J = 6.8 Hz, 2H), 5.67 (s, 2H), 5.44 (s, 1H), 2.85 – 2.75 (m, 2H), 2.54 – 2.46 (m, 2H), 2.43 (s, 3H), 1.84 – 1.71 (m, 2H), 1.35 (s, 3H). 13C NMR (126 MHz, CD3CN) δ 188.63, 188.11, 167.10, 164.47, 161.85, 160.19, 152.54, 152.32, 151.58, 144.81, 138.71, 137.97, 132.52, 131.55, 131.39, 129.86, 129.21, 128.44, 125.62, 124.13, 123.05, 99.98, 63.96, 35.46, 34.21, 30.10, 28.62, 27.14. HRMS‐ESI m/z: calcd. [M-H]+: 766.1973, found [M-H]+: 766.1976.



Figure S12: <sup>1</sup>H NMR (500 MHz, CD<sub>3</sub>CN) of Ru-complex.



Figure S13: <sup>13</sup>C NMR (126 MHz, CD<sub>3</sub>CN) of Ru-complex.

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