S1 **S1. Sample preparation**

S2 Symmetric-type diamond-anvil cells (DAC) were used to perform experiments at high *P*-*T* conditions. A S3 rhenium gasket was indented by two diamond-anvils to a thickness of ∼20 *µ*m and a diameter of 200 or S4 300 *µ*m. Using a laser drilling system in Arizona State University (ASU), a 130–180 *µ*m diameter hole, S5 a sample chamber, was made at the center of the indentation. The drilled rhenium gasket was coated S6 with gold to prevent the H diffusion into the gasket material. MgO (Alfa Aesar, 99.95% purity) powder σ was dried at 388 K for 24 hours in an oven and then mixed with 20 wt% Fe metal (Sigma Aldrich, 99.9%) S8 purity) in a mortar. The mixed powder was compressed to a ∼10-*µ*m thin foil. To prevent contact with s⁹ the diamond anvils and make a space for H_2 , the foil was propped by 2–3 μ m MgO grains on the culets of S10 both diamond anvils. Pure H² gas (Matheson Tri-Gas, Inc.) compressed to 1300 bar was loaded in the S11 sample chamber together with the foil using a gas loading system at ASU (Sanchez Technologies). After s₁₂ closing the DAC, the pressure of the sample chamber measured from ruby chips [\(54\)](#page-1-0) was ∼1 GPa. Before s13 synchrotron measurements, the samples were compressed to target pressures between 5 and 25 GPa.

S14 **S2. Laser-heated diamond-anvil cell experiments**

at beamline 13-IDD of the GeoSoilEnviroCARS (CS). A double-sided laser-heating setup at the beaml
ample. The laser beams were aligned coaxially with
was \sim 20 μ m (18). In previous experimental studie
ound that a conti S15 Experiments were conducted at beamline 13-IDD of the GeoSoilEnviroCARS (GSECARS) sector at the S16 Advanced Photon Source (APS). A double-sided laser-heating setup at the beamline focuses infrared laser s¹⁷ beams on both sides of the sample. The laser beams were aligned coaxially with the X-ray beam. The S18 diameter of the heated area was ∼20 *µ*m (18). In previous experimental studies using pure H² medium \mathfrak{so} (for examples, [\(51\)](#page--1-1)), it was found that a continuous heating for sample in an H_2 medium could shatter s_{20} diamond anvils within a few second. The reason is because H_2 molecules are small enough to diffuse s21 into diamond anvils at high pressure and such a behavior is further promoted by heating. To avoid the S22 diamond embrittlement issue by H diffused into the diamond anvils during continuous heating to the ses melting temperature of MgO, we conducted pulsed laser heating $(19-21)$. This heating method allowed us S24 to reach above 4000 K for the sample in a diamond-anvil cell. Because pulse-heating tends to reduce the S25 mechanical damage of the sample foil by the convection of H fluid at high temperatures, it can preserve a S26 heated sample better. However, the percentage of the samples recovered for chemical analysis is still as low $\frac{\text{S27}}{\text{S28}}$ as 30–50% because of the severe fracturing of the sample foils by intense reaction with an H fluid (Fig. [S2a](#page-3-0)). S28 During heating, thermal radiation spectra were measured for the both side of the sample.

S29 Temperature was estimated by fitting the measured spectra to the gray-body radiation equation for S30 680–840 nm range. Temperature uncertainty is estimated considering 1) the uncertainties from the spectral \mathfrak{so}_1 fitting, 2) temperature differences between the two sides of the sample, and 3) temperature fluctuations S32 during pulsed laser heating (Table [S1\)](#page-15-0). We estimated pressures at high temperatures using the measured s₃₃ volume of MgO combined with its thermal equation of state [\(23\)](#page--1-4). Estimated 1σ uncertainty for pressure is $\pm 10\%$. A monochromatic X-ray beam (a wavelength of 0.3344 Å) was focused to $3\times2 \ \mu m^2$ in size.

S35 A pixel array detector (Pilatus 3X CdTe 1M) was used for measuring X-ray diffraction (XRD) images. S36 Diffraction images were measured at in situ high *P*-*T* conditions as well as after heating at 300 K.

S37 Additional XRD measurements were also conducted at beamline 13-BMD of GSECARS with a monochross matic X-ray beam (a wavelength of 0.3344 Å) focused to $4\times9 \ \mu m^2$. The same type of detector was used. $S39$ XRD data for LaB₆ standard were measured to correct for the detector tilt and to obtain sample-to-detector S40 distance, ∼200 mm. 2D XRD images were converted to one-dimensional XRD patterns using the Dioptas S41 software [\(55\)](#page-1-1). Peak identification and unit-cell fitting were conducted using the PeakPo software package s_{42} (56) .

S43 **S3. Raman spectroscopy**

S44 Raman spectroscopy measurements were conducted at $200-4500$ cm⁻¹ using a system at GSECARS [\(24\)](#page--1-5). S45 A 532-nm laser used as an excitation source. The laser power was 0.05 or 0.1 Watts, and the grating was

S46 1200 grooves/mm in the spectrograph (Princeton Instruments Acton Series 2500; a focal length of 500 mm). $8\times8 \ \mu \text{m}^2$ areas at the heated areas were mapped with a 4- μ m step (i.e., 25 spectra per single map). Each S48 spectrum was measured for 10 seconds. Unheated areas and hydrogen-only areas in the sample chamber S49 were also measured to compare with the spectra from the heated areas. Similar measurements were also S50 conducted using a Raman system at ASU.

S51 **S4. Chemical analysis**

U. [D](#page--1-11)ue to the cracks and fractures in the heated is
inning (e.g., vertical streaks in Fig. 4a). [T](#page--1-6)he thin is
ing to ~100 nm thickness at 2 keV and 44 pA. In a
ree imaged with an acceleration voltage of 80 keV. T
lark field S52 The recovered samples were analyzed in electron microscopes for the chemical composition and the S53 structure in the heated areas. For scanning electron microscopy (SEM) analysis and energy-dispersive X-ray S54 spectroscopy (EDS), we milled the recovered sample using a focused ion beam (FIB) instrument (Helios S55 5 UX, Thermo Fisher Scientific Inc.) at Arizona State University (ASU). For SEM imaging and EDS, ϵ_{ss} we milled the heated areas for cross sections using the focused Ga⁺ ion beam with 30 keV and 9 nA for S57 initial milling, and 30 keV and 2.6 nA for the final milling. SEM images were obtained at 10 keV or 15 keV. S58 Two-dimensional elemental distribution maps were measured at 10 keV using EDS (Xflash 6|30 detector, S59 Bruker co.) for Mg, Fe, and O. For transmission electron microscopy (TEM) analysis and EDS, thin sections \sim 6×6 μ m²) were extracted from the heated areas using the Ga⁺ ion beam with 30 keV and 0.26 nA in the S61 same FIB instrument at ASU. Due to the cracks and fractures in the heated areas, we could not avoid S62 some irregularities during thinning (e.g., vertical streaks in Fig. 4a). The thin specimen was moved on a S63 copper grid for further thinning to ∼100 nm thickness at 2 keV and 44 pA. In a ARM-200F (JEOL co. at S64 ASU) STEM, the samples were imaged with an acceleration voltage of 80 keV. TEM images were obtained S65 under a high-angle annular dark field scanning transmission-electron microscope (HAADF-STEM) mode. S66 Chemical data were acquired using energy dispersive X-ray spectroscopy for 20–30 seconds. Elemental S67 analysis was conducted for Mg, Fe, and O.

S68 **S5. Pressures at the interface between atmosphere and magma in sub-Neptunes**

 $\frac{1}{100}$ For sub-Neptune exoplanets with a 2% H₂ atmosphere and an Earth-like interior composition, the planet's σ mass can be calculated for the cases of $P(\text{AMB}) = 15$, 20, and 25 GPa using equations 9–11 and 21–22 in $_{571}$ [\(8\)](#page--1-7). From the mass and $M = (1.74 \pm 0.38)R^{(1.58 \pm 0.10)}$, suggested by (45) for a mass range of 5–25*M(E)*, s_{72} radius can be calculated. The planet's radii for $P(\text{AMB}) = 15$, 20, and 25 GPa are 2.52, 3.03, and 3.48*R(E)*, s_{73} respectively (Fig. [5\)](#page--1-9).

S74 A slightly different method can be used, but with essentially the same results within the uncertainties. S75 The radius cliff is observed at ∼3*R*(*E*) [\(13,](#page--1-10) 14, 44). To calculate the pressure at the atmosphere-magma s_{76} boundary of sub-Neptunes using equation 22 in [\(8\)](#page--1-7), H₂ percentage and the mass of planet are required. The σ ⁷⁷ mass-radius relation of sub-Neptunes with a 2% H₂ atmosphere (equilibrium temperature (T_E) of 700 K) \mathbb{S}_{78} and Earth-like core exists at 3.0 ± 0.2 $R(E)$ between $5M(E)$ and $15M(E)$ [\(8\)](#page--1-7). We estimated the pressure S79 range for the atmosphere-magma boundary shown in Fig. [1](#page--1-13) from three chosen planets: Kepler-1808b $_{580}$ $(T_E = 601 \text{ K}; 9.41M(E); 3.03R(E)),$ Kepler-1717b $(T_E = 780 \text{ K}; 9.22M(E); 2.99R(E))$ and TOI-220b $S₈₁$ $(T_E = 806 \text{ K}; 13.6M(E); 3.02R(E))$ assuming a 2% H₂ atmosphere.

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Fig. S1. The sample setup in a diamond-anvil cell used in this study.

Fig. S2. Scanning electron microscopy (SEM) images (a and b) and the energy-dispersive spectroscopy (EDS) maps (c) of the heated area in the samples recovered from run #31-005 (12.5 GPa and 4770 K). (a) A top view of the heated area. The hole at the center is from the transport of the melt away from the spot by an H fluid during heating. A white dashed line shows the location of the cross section in b. (b) An SEM image showing small cavities in the heated area (the left side of a dashed line). (c) Elemental distribution maps for Fe and Mg. The two white circles with a weaker Fe signal than the Fe metal in the unheated area and a weaker Mg signal than the surrounding MgO area can be assigned to Mg₂FeH₆, which contains both Mg and Fe, unlike MgO. The circled areas contain small cavities as shown in (b). According to Eq. [1,](#page--1-14) Mg₂FeH₆ and H₂O should crystallize together upon temperature quench at high pressure. When the DAC was opened, H₂O becomes a liquid and is eventually evaporated, leaving the small cavities. Therefore, both the chemical map and the texture can support the interpretation for Mg_2 FeH $_6$.

Fig. S3. X-ray diffraction images (unrolled) and patterns measured (a,b) before heating, and ($\bigoplus c, d$) after heating to 3820 K at 11.5 GPa from run #31-003. The circles above the peaks mark the peak positions for the observed phases calculated from their crystal structures and equations of state. For the peaks which are predicted but not observed, open circles are used. The inset in the diffraction pattern of (ab) shows the raw 1D diffraction pattern for the part affected by the shadow from the beamstop. The yellow arrowsinsets in the unrolled diffraction image (c) and the diffraction pattern (d) of (b) show part of the unrolled diffraction images indicate for the major peaks of Mg_2 FeH₆.

INFORMATION THE CONDUCT THE CONDUCT OF THE CONDUCT OF THE STATE OF THE OPTER OF THE OPTER OF THE STATE OF THE OPTER OF **Fig. S4.** X-ray diffraction images (unrolled) and patterns measured after heating to (a) 3500 K at 8.5 GPa, (b) 4870 K at 26 GPa, and (c) 3570 K at 24.5 GPa. The circles above the peaks mark the peak positions for the observed phases calculated from their crystal structures and equations of state. The open circles indicate predicted positions for the diffraction peaks of the observed phases, but the peaks were not observed at the positions. The white arrows in the 2D images indicate defects in the detector. The diffraction pattern in (a) was in different beamtime and therefore the detector defects were not seen there. (b) and (c) were measured during the same beamtime.

tion of state from (57) at high pressures. The volume of fcc FeH_a, measured in this s

is up to 1.2 based on the methods presented in (21). Fig. S5. (a) The unit-cell volumes of Mg₂FeH₆ measured in this study compared with the equation of state from [\(38\)](#page--1-15) at high pressures. (b) The atomic volume of fcc FeH_x measured in this study compared with the equation of state from (57) at high pressures. The volume of fcc FeH*^x* measured in this study is greater than the known values for $x = 1$ [\(57\)](#page-1-3). The estimated hydrogen content, x , is up to 1.2 based on the methods presented in (21).

We used the $Fm\bar{3}m$ (#225) space group for Mg $_2$ FeH $_6$ (Fig. S6. The Le Bail analysis for the XRD pattern in Fig. 2a. The blue crosses are the measured intensities, and the red line is the calculated diffraction pattern after fitting. The black line at the bottom is the difference between the observed and the calculated diffraction patterns. The colored ticks indicate the diffraction positions of the identified phases. $a = 2.562$ Å and $c = 8.834$ Å for dhcp FeH_x, $a = 3.048$ Å and $c = 4.451$ Å for Mg(OH)₂, and $a = 4.145$ Å for MgO.

Fig. S7. X-ray diffraction data measured at 1 GPa after heating to 4510 K at 28.5 GPa (#26-016). The yellow box with the yellow arrow highlights an unidentified peak. The white arrows in the unrolled 2D image indicate defects in the detector.

Fig. S8. Raman spectra measured at (a) 6.1 GPa and (b) 1.1 GPa during pressure decrease to 1 bar after heating experiments at 28.5 GPa and 4510 K, and 26 GPa and 4870 K, respectively (Table [S1\)](#page-15-0). Pressures were determined by the ruby fluorescence close to the gasket. The insets are for H₂O-VII from [\(31\)](#page--1-18) in (a) and Mg(OH)₂ from [\(30\)](#page--1-19) in (b). In (b), a weak peak at \sim 3560 cm $^{-1}$ can be assigned to H₂O.

Fig. S9. Raman spectrum measured at ∼1.0 GPa for the sample heated at 28.5 GPa and 4510 K (#26-016). The spectrum in this figure and Fig. [S10b](#page-11-0) are from different spots in the same heated area. The peak assignments were made by using previous reports on pure phases: Mg(OH)₂ (30), H₂ [\(52\)](#page--1-20), γ-MgH₂ from [\(33\)](#page--1-21) and the Mg–H stretching modes of MgH_2 from [\(32\)](#page--1-22).

heating to (a) 4870 K at 26 GPa and (b) 4510 K at 88.5 GPa. (c) Raman spectra from
t 1 bar, and γ -MgH₂ at 1 bar (33). In (a) and (b), Raman spectra were obtained at 2
The modes of Br (brucite; (30))), H₂O-VII (loe-V Fig. S10. Raman spectra measured after laser heating to (a) 4870 K at 26 GPa and (b) 4510 K at 28.5 GPa. (c) Raman spectra from previous studies reported for Mg(OH)₂ at 1 bar[\(30\)](#page--1-19), H₂O-VII at 9.6 GPa [\(31\)](#page--1-18), α -MgH₂ at 1 bar, and γ -MgH₂ at 1 bar (33). In (a) and (b), Raman spectra were obtained at 22 GPa, 9 GPa, and 1 GPa for two different heated areas during decompression to 1 bar. The modes of Br (brucite; (30)), H₂O-VII (Ice-VII; (31)), MgH₂ (α-MgH₂ and $γ$ -MgH₂; [\(33\)](#page--1-21)) and Mg-H (the stretching modes of MgH₂; [\(32\)](#page--1-22)) were labeled in a and b. The red arrow denotes the range of the Mg–H stretching ondes of MgH₂. The Mg–H stretching vibrational modes exist between 1000 and
1300 cm^{–1} [\(32,](#page--1-22) [33\)](#page--1-21). The red stars highlight the

scopy (STEM) images of the sample recovered from run #31-003 (11.5 GPa and 382
Ilow box in the left image), elongated platy grains (indicated by the yellow arrows) are also reported for Mg(OH)₂ in (29). Fig. S11. Scanning-transmission electron microscopy (STEM) images of the sample recovered from run #31-003 (11.5 GPa and 3820 K). The zoomed-in area was shown on the right. At the edge of the heated area (the yellow box in the left image), elongated platy grains (indicated by the yellow arrows) are observed, and they are Mg(OH)₂, as found in our XRD patterns. Similar structures were also reported for Mg(OH)₂ in (29).

i mages (a and b) and EDS maps (c) of the heated area in the samples recovered from presented in Fig. 4c. (a) The white dashed line divides the heated area and the surrounding unheated areas. (b) recorrectes in texture bet **Fig. S12.** Scanning electron microscopy (SEM) images (a and b) and EDS maps (c) of the heated area in the samples recovered from run #26-009 (26 GPa and 4870 K). The images here show the opposite side of the cross section presented in Fig. 4c. (a) The white dashed line divides the heated and the unheated areas. A relatively low voltage (10 kV) used for the imaging reveals the differences in texture between the heated area and the surrounding unheated areas. (b) The center of the heated area was melted. The melted region has many bubble-shaped structures. Some of them are highlighted by the white arrows. (c) Elemental maps for Fe, Mg, and O. The melted area has lower Mg than the unheated areas. The O-poor region matches with the region with the bubble-shaped structures. In this study, we observed that MgO is decomposed, and Mg and H become miscible at high temperatures (Eq. [2\)](#page--1-24). The process produces Mg dissolved in H fluid and H₂O. The melted region in b overlaps with the area with weak Mg and O signals in c. As explained in the main text, the bubble-shaped structure can be formed from H_2O , H_2 , and Mg–H fluid exsolved during temperature quench at high pressure. These fluids in the bubble-shaped structures should have left when pressure was decompressed completely to 1 bar. Because θ some materials would be removed in this way, this process can result in the weak Mg and O signals compared with the unmelted regions with MgO on the sides.

Fig. S13. The densities of Mg₂FeH₆ [\(38\)](#page--1-15), MgO [\(53\)](#page--1-25), and MgSiO₃-En (orthoenstatite) [\(42\)](#page--1-26) calculated from their equations of state at high pressures and 300 K. Mg₂FeH₆ is much less denser than planetary silicates and oxides.

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Table S1. Experimental runs in this study. *O–H vibration mode was identified in Raman spectra. **dhcp FeH*^x* **is likely metastably observed because the diffraction peaks are broader than those of other phases. Uncertainties for pressures and temperatures are 10% of the given value and** 1*σ* **(given in parentheses), respectively. XRD: X-ray diffraction, TEM: Transmission electron microscopy, SEM: scanning electron microscopy, fcc: face-centered cubic, dhcp: double hexagonal close-packed.**

Table S2. The unit-cell volume of Mg2**FeH**6 **and the atomic volume of FeH***x***. These data points are used in Fig. [S5.](#page-6-0) Uncertainty for the volume is** 1*σ* **given in parentheses. Uncertainty for pressure is smaller than 10% of the given value.** Z is the number of formula units in the unit cell (for example, $Z = 4$ for fcc FeH_{*x*}).

