

Supplementary information

Flexible power generators by Ag₂Se thin films with record-high thermoelectric performance

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1. Supplementary Experimental details

Thin film and device fabrication. Flexible Ag₂Se thin films were deposited using the vacuum thermal co-evaporation method at various temperatures. High-purity Ag powders (99.99% purity, Aladdin) and Se-Te mixed powders (99.99% purity, Aladdin) were loaded into a vacuum deposition chamber using tantalum evaporator boats. A polyimide (PI) substrate measuring 20 mm × 20 mm × 2 mm was cleaned ultrasonically for 15 minutes consecutively in acetone, ethanol, and deionized water. The background pressure in the vacuum deposition chamber was maintained at 5.4×10^{-4} Pa. The silver source had an evaporating current of 130 A, while the selenium source stabilized at 40 A, with a concurrent deposition time of 15 minutes. The chemical composition of the samples was controlled by adjusting the mass ratio. All Ag and Se powders were evaporated after the preparation process. To address the issue of insufficient atomic energy during room temperature deposition, the substrate was heated to 353 K during deposition. The average film thickness was 350 nm. A single Ag₂Se-Sb₂Te₃ thermoelectric device was fabricated with a Cu electrode on a PI substrate. The legs had fixed dimensions of 20 mm × 4 mm. An organic coating material, polyvinyl laurate, was employed to enhance the flexibility of the as-fabricated Ag₂Se thin films. First, the Ag₂Se thin films were securely laid flat on glass. Second, an appropriate amount of polyvinyl-laurate-based polymers composed of 90% poly(vinyl laurate) and 10% N-methylpyrrolidone was applied to the thin films under vacuum conditions. The thin films were then placed in a centrifuge for 10 seconds with a rotation speed of 2000 r/s to ensure uniform coverage of the organic layer. Subsequently, the samples were heated to 95 °C for 2 minutes and exposed to ultraviolet light for 15 seconds.

Characterizations and thermoelectric performance evaluation. X-ray diffraction (XRD) analysis was performed using a Rigaku Corporation D/max 2500 instrument with CuK α radiation over an angular

range of 20° to 60° in 0.02° increments to determine the crystal orientation of the Ag₂Se samples. The preferred orientation information of the Ag₂Se thin films was obtained through pole figure scans using a four-circle X-ray diffractometer (X'PERT, Philips). Morphological analysis, mapping, and composition assessment of the samples were conducted using a Zeiss Supra55 scanning electron microscope (SEM). An energy dispersive X-ray (EDS) detector from Bruker, known as the EDS QUANTAX, was used for EDS analysis. The electrical conductivity σ and Seebeck coefficient S were measured using a Seebeck coefficient and electrical conductivity apparatus (SBA458, Netzsch). The in-plane thermal conductivity κ of the thin film in this work, deposited on a commercial flat chip, was determined using the 3 ω method through a thin film comprehensive physical property analyzer (TFA-LINSES). Carrier concentration n and mobility μ were investigated using a Van der Pauw Hall measuring instrument (HL5500 PC, Nanometrics). Contact resistivity ρ_c was measured using a 2400 series digital source meter.

Preferred orientation factor calculation: The orientation factor F was calculated by the equations

$F = \frac{P-P_0}{1-P_0}$, $P = \frac{I(00l)}{\Sigma I(hkl)}$, and $P_0 = \frac{I_0(00l)}{\Sigma I_0(hkl)}$, where $I(00l)$ is the sum of the diffraction intensity of (00l) planes, $\Sigma I(hkl)$ is the total intensity of all (hkl) diffraction peaks and P is the ratio of the intensity of (00l) plane in the measured data. The $I_0(00l)$, $\Sigma I_0(hkl)$, and P_0 represent the corresponding ones from the standard pattern.

Single parabolic (SPB) model calculation: We employed an SPB model to perform a simulation to understand our measured properties by using the following equations¹⁻⁴:

$$S(\eta) = \frac{k_B}{e} \cdot \left[\frac{\left(r+\frac{5}{2}\right) \cdot F_{r+\frac{3}{2}}(\eta)}{\left(r+\frac{3}{2}\right) \cdot F_{r+\frac{1}{2}}(\eta)} - \eta \right] \quad (1)$$

$$n_H = \frac{1}{e \cdot R_H} = \frac{(2m^* \cdot k_B T)^{\frac{3}{2}}}{3\pi^2 \hbar^3} \cdot \frac{\left(r+\frac{3}{2}\right)^2 \cdot F_{r+\frac{1}{2}}^2(\eta)}{\left(2r+\frac{3}{2}\right) \cdot F_{2r+\frac{1}{2}}(\eta)} \quad (2)$$

$$\mu_H = \left[\frac{e\pi\hbar^4}{\sqrt{2}(k_B T)^{\frac{3}{2}}} \frac{C_1}{E_{def}^2 (m^*)^{\frac{5}{2}}} \right] \frac{(2r+\frac{3}{2}) \cdot F_{2r+\frac{1}{2}}(\eta)}{(r+\frac{3}{2})^2 \cdot F_{r+\frac{1}{2}}(\eta)} \quad (3)$$

$$F_j(\eta_F) = \int_0^\infty \frac{x^j}{1+e^{x-\eta_F}} dx \quad (4)$$

where j is the Fermi integral of order, η_F and \hbar are reduced Fermi level ($E_f/k_B T$) and Plank constant. By using the experimentally measured S and n in the above formula and assuming acoustic phonon scattering is prominent (*i.e.*, $r = -1/2$), the theoretical transport properties can be calculated.

Computational calculation: First-principles calculations were performed based on density-functional theory (DFT) with all electron projected augmented wave (PAW) method, as implemented in the Vienna Ab initio Simulation Package (VASP) ⁵⁻¹⁰. Semi-local generalized gradient approximation (GGA) with the fully relativistic Perdew-Burke-Ernzerhof (PBE) exchange correlation functional was employed ¹¹. The valence wave functions were expanded in a plan-wave basis with a cut-off energy of 400 eV. To simulate doping of Te, a $1 \times 1 \times 3$ supercell containing 36 atoms was created, where one of the 12 Se atoms was substituted by one Te atom. All atoms were allowed to relax in their geometric optimizations, with the Brillouin zone sampled by a Monkhorst-Pack \mathbf{k} -mesh of $2 \times 5 \times 4$ ($4 \times 7 \times 6$ for self-consistent calculations). The convergence criterion was set to be 1×10^{-7} eV per electron and 1×10^{-3} eV·Å⁻¹ per unit cell. The Hubbard U model was included in non-self-consistent calculations, with the on-site coulombic (U) and the exchange (J) terms combined in a single effective U parameter of 5.8 eV for Ag_4d orbitals, determined based on linear response theory ¹². Because Te is a heavy element, the spin-orbital coupling (SOC) effect was considered for the calculation of band structures, which is projected on the \mathbf{k} -path indicated by the AFLOW framework ^{13,14}.

For the calculations of (001) surface energy variation before and after Te-doping, the spin-polarized density functional theories (DFT) were performed by using the Vienna Ab initio Simulation Package (VASP) ⁷. The Perdew-Burke-Ernzerhof generalized-gradient approximation functional was

used to describe the interaction between electrons ¹¹. The energy cutoff was set to 400 eV. The Monkhorst-Pack k-points were set to be $6 \times 4 \times 1$ during the calculations. The vacuum region was set to be 15 Å in the z-direction to prevent the interaction between two adjacent surfaces. The energy convergence and atomic forces were set to 10^{-5} eV and 0.02 eV/Å. The surface energy γ_s , is the energy required to cleave a surface from the bulk phase, which describes the stability of the surface. The γ_s is calculated by the below equation ¹⁵:

$$\gamma_s = \frac{1}{2A}(E_{slab} - NE_{atom}) \quad (5)$$

Here A is the area of the surface. E_{slab} is the energy of the slab. N is the number of metal atoms. E_{atom} is the energy for one atom (Ag, Se, and Te) referred to the corresponding bulk structures.

For the calculations of (002) surface energy variation before and after Te-doping, the commercial software package Material Studio, including the Cambridge Serial Total Energy Package (CASTEP) module was used to conduct this geometric optimization to determine the initial lattice energy. In the process, the generalized-gradient approximation (GGA-PBE) along with OTFG ultrasoft pseudo potentials was used to perform the geometric optimization. The (002) surface was created in Material Studio software, the cleave rule is default and the vacuum thickness is 15Å.

Debye Callaway model calculations: Debye Callaway model is investigated to clarify the source of reduction in κ_l ¹⁶:

$$\kappa_l = \frac{4\pi k_B^4 T^3}{\nu h^3} \int_0^{\theta_D} \tau_T \frac{\chi^4 \exp(\chi)}{[\exp(\chi)-1]^2} d\chi \quad (6)$$

here ν , θ_D , γ , and τ_T stand for sound velocity, Debye temperature, reduced phonon frequency ($\chi = \hbar\omega/k_B T$), and total phonon relaxation time.

The phonon-scattering phenomenon involves different scattering progressions from electron-phonon interaction (τ_{EP}^{-1}), dislocations (τ_{DS}^{-1}), impurity/point defects (τ_{PD}^{-1}), grains (τ_{GB}^{-1}), and Umklapp-

scattering (τ_U^{-1}), in the doped system. Hence, τ_T for doped Ag₂Se is demonstrated as ¹⁷:

$$\frac{1}{\tau_T} = \frac{1}{\tau_U} + \frac{1}{\tau_N} + \frac{1}{\tau_{PD}} + \frac{1}{\tau_{DS}} + \frac{1}{\tau_{GB}} \quad (7)$$

where τ_U , τ_N , τ_{PD} , τ_{IF} , τ_{NP} , τ_{DS} , and τ_{SF} are relaxation time corresponding to the scattering from phonon-phonon U- and N-process (U+N), vacancies/alloy elements (point defects, PD), grain boundaries (GB), and dislocations (DS). As a result, τ_T (the total phonon relaxation time) for the doped system shows a significant reduction as compared to τ_T for the undoped Ag₂Se system due to the existence of intense scattering effects under alloying or doping.

Further, Debye-Callaway model-based spectral lattice thermal conductivity κ_s was derived for doped samples to reveal the influence of all possible phonon scattering mechanisms, as follows ¹⁸:

$$\kappa_S = \frac{4\pi k_B^4 T^3}{V_s h^3} \tau_T \frac{z^4 \exp(z)}{[\exp(z)-1]^2} \quad (8)$$

In the present doped system, the main τ_T reduction is related to the phonon-scattering phenomenon from scattering due to impurity/point defects (τ_{PD}^{-1}), and grain boundaries (τ_{GB}^{-1}),

Hereafter, the phonon scattering from prominent scattering centers can be simplified as ¹⁹:

$$\tau_{PD}^{-1} = \frac{V \Gamma}{4\pi v^3} \omega^4 \propto A \omega^4 \quad (9)$$

$$\tau_U^{-1} = \frac{\hbar \gamma^2 \omega^2 T}{M v^2 \theta_D} \exp\left(-\frac{\theta_D}{3T}\right) \propto B \omega^2 \quad (10)$$

$$\tau_{EP}^{-1} = \beta \tau_U^{-1} = \beta \frac{\hbar \gamma^2 \omega^2 T}{M v^2 \theta_D} \exp\left(-\frac{\theta_D}{3T}\right) \propto C \omega^2 \quad (11)$$

$$\tau_{DS}^{-1} = 0.6 B_D^2 N_D \omega \gamma^2 \left\{ \frac{1}{2} + \frac{1}{24} \left(\frac{1-2r}{1-r} \right)^2 \left[1 + \sqrt{2} \left(\frac{v_L}{v_T} \right) \right] \right\} + N_D \left(\frac{v^{-4/3}}{v_s^2} \right) \omega^3 \quad (12)$$

The simplified presentation of scattering phenomena via introducing pre-factors is described as follows:

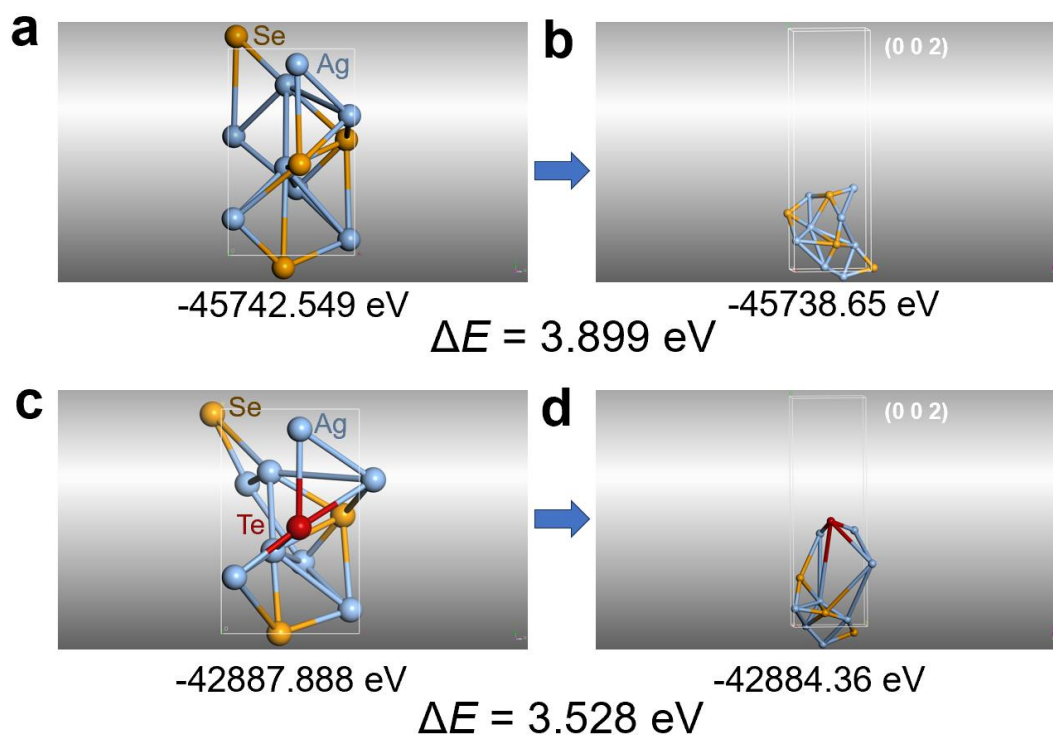
$$\tau_U^{-1} + \tau_N^{-1} = Z_1 \omega^2 T \exp\left(-\frac{\theta_D}{3T}\right) \quad (13)$$

$$\tau_{PD}^{-1} = Z_2 \omega^4 \quad (14)$$

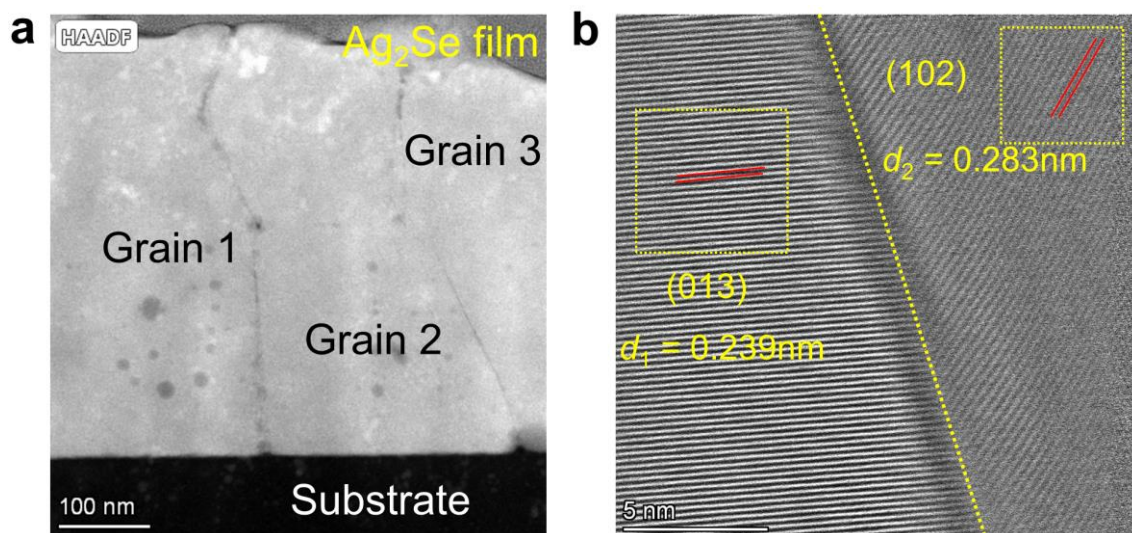
$$\tau_{DS}^{-1} = Z_3\omega^1 + Z_4\omega^3 \quad (15)$$

The proposed pre-factors Z_1 , Z_2 , Z_3 , and Z_4 correspond to the scattering from phonon-phonon U- and N-process, vacancies/anti-sites (point defects), and dislocation defects, respectively. The values of the average v is 1462 ms^{-1} , and the θ_D is 63.72 K , both taken from the literature ¹⁶.

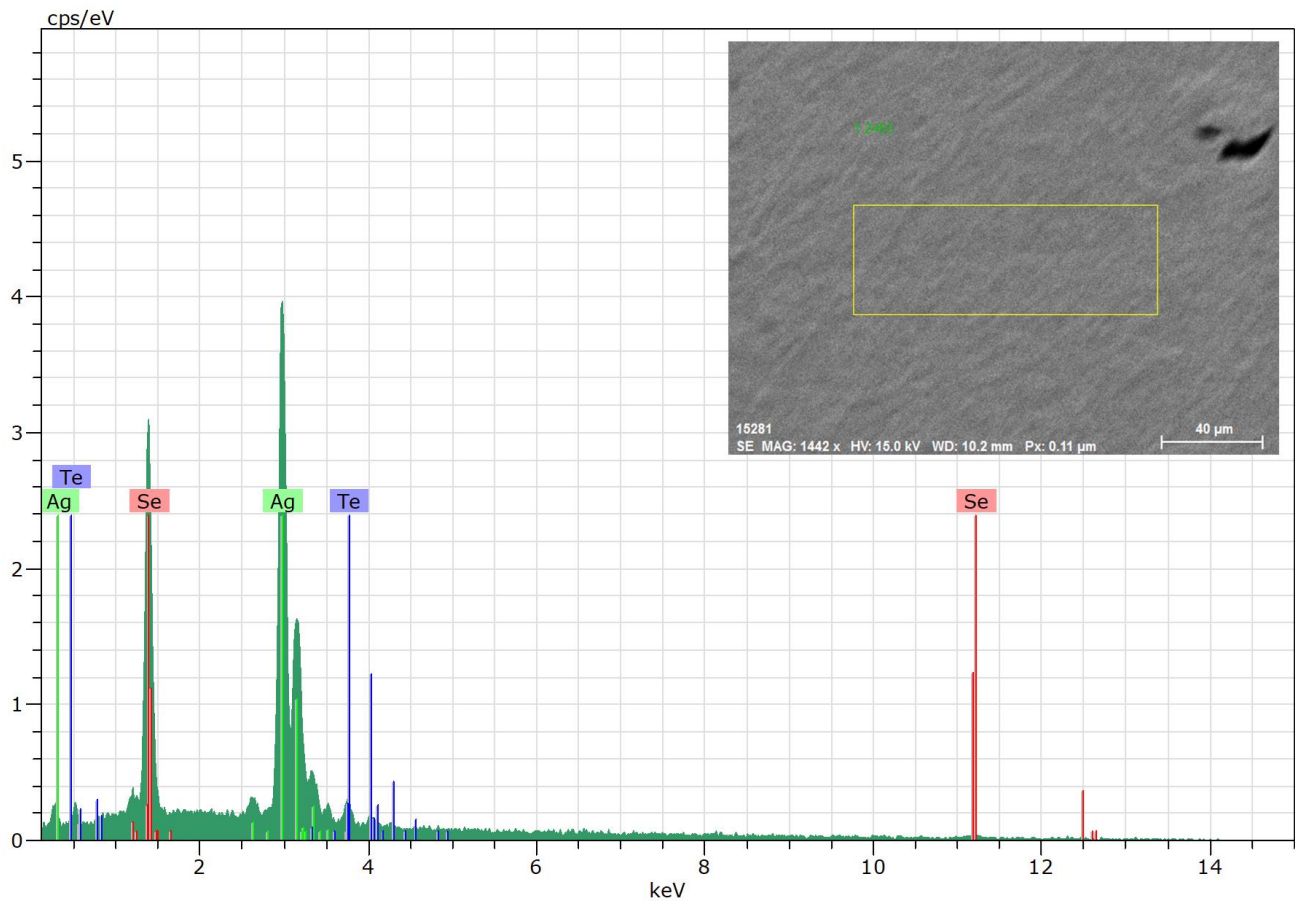
2. Supplementary Figures



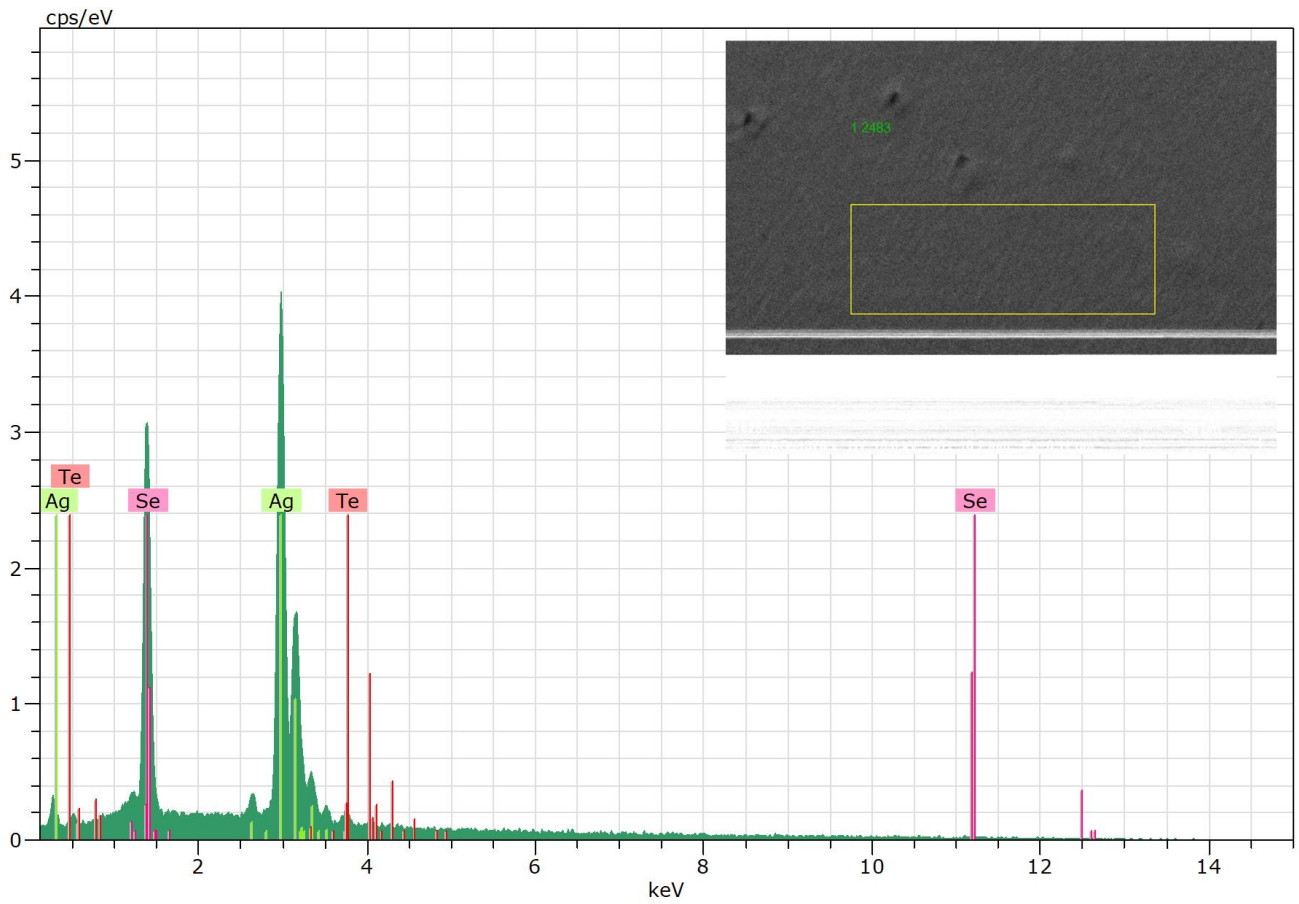
Supplementary Fig. 1. (a) Unit cell of pristine Ag_2Se . (b) Unit cell of pristine Ag_2Se cut along the (002) surface. The vacuum layer thickness is 15 Å. (c) Unit cell of Te-doped Ag_2Se . (d) Unit cell of Te-doped Ag_2Se cut along the (002) surface. The vacuum layer thickness is 15 Å. Here, the undoped Ag_2Se unit cell energy is -45742.549 eV, and the undoped Ag_2Se unit cell energy after cutting along the (002) crystal plane is -45738.65 eV, therefore the energy difference $\Delta E = 3.899$ eV. Oppositely, the Te-doped Ag_2Se unit cell energy is -42887.888 eV, while the Te-doped Ag_2Se unit cell energy after cutting along the (002) crystal plane is -42884.36 eV, therefore the $\Delta E = 3.528$ eV, which is smaller than 3.899 eV, indicating that the introduction of Te on Se sites is beneficial for the formation of the (002) surface of Ag_2Se .



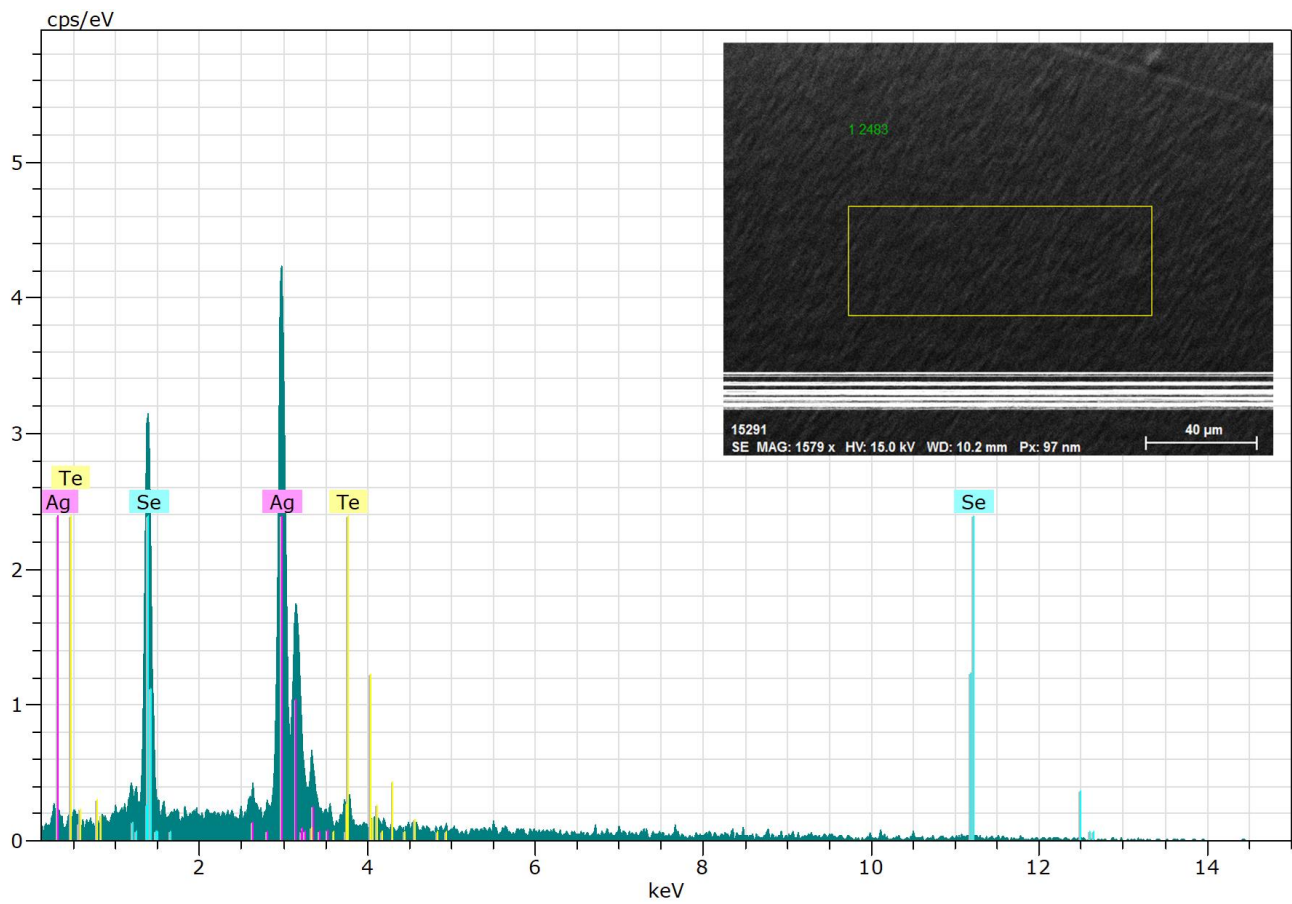
Supplementary Fig. 2. (a) Low-magnification transmission electron microscopy (TEM) high-angle annular dark field (HAADF) image of pristine Ag₂Se thin film with anisotropic grains. (b) High-resolution TEM (HRTEM) image to show a grain boundary between two grains.



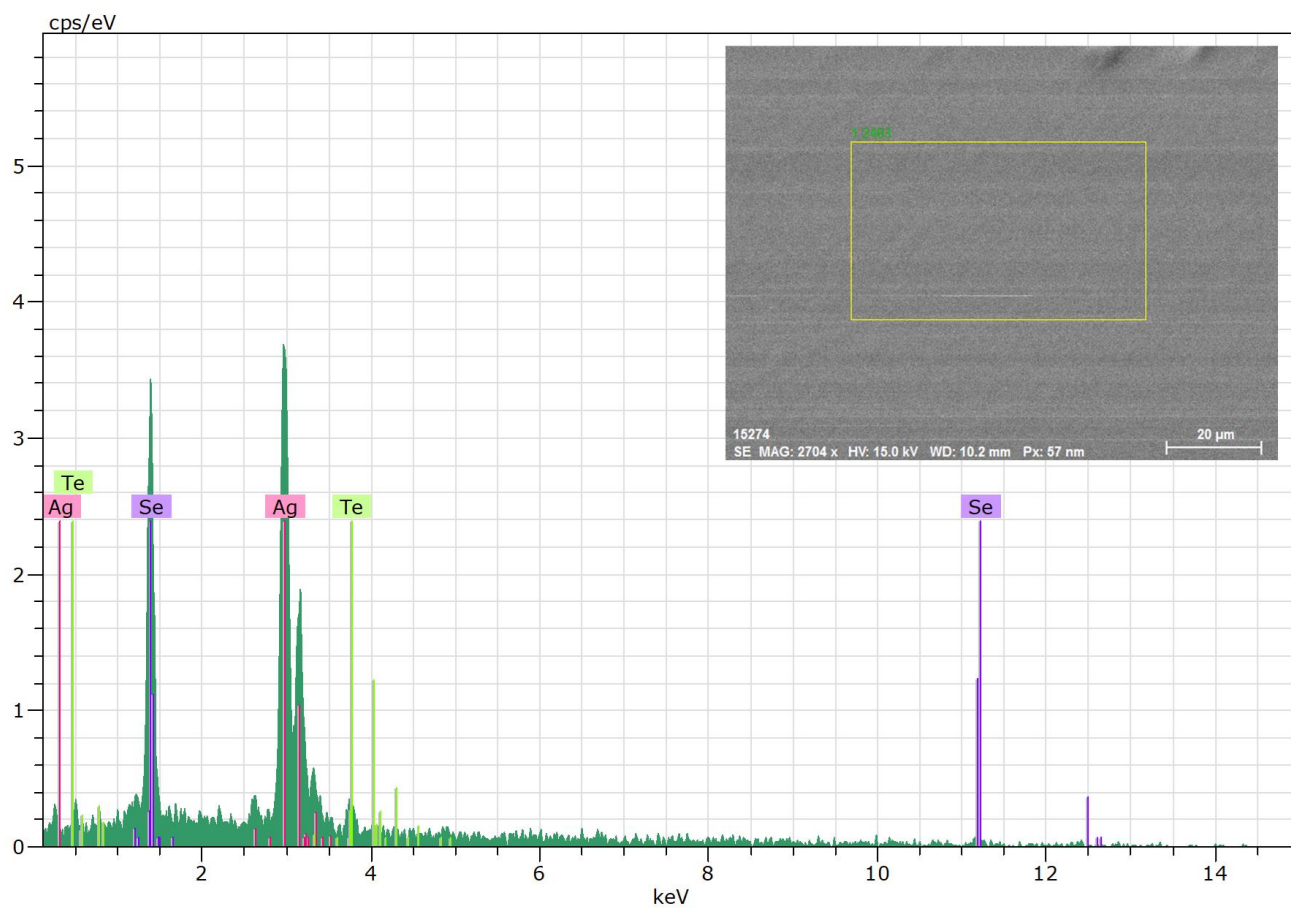
Supplementary Fig. 3. Energy dispersive spectroscopy (EDS) results with inset of scanning electron microscopy (SEM) image for Ag_2Se thin films with 0.7 at.% Te.



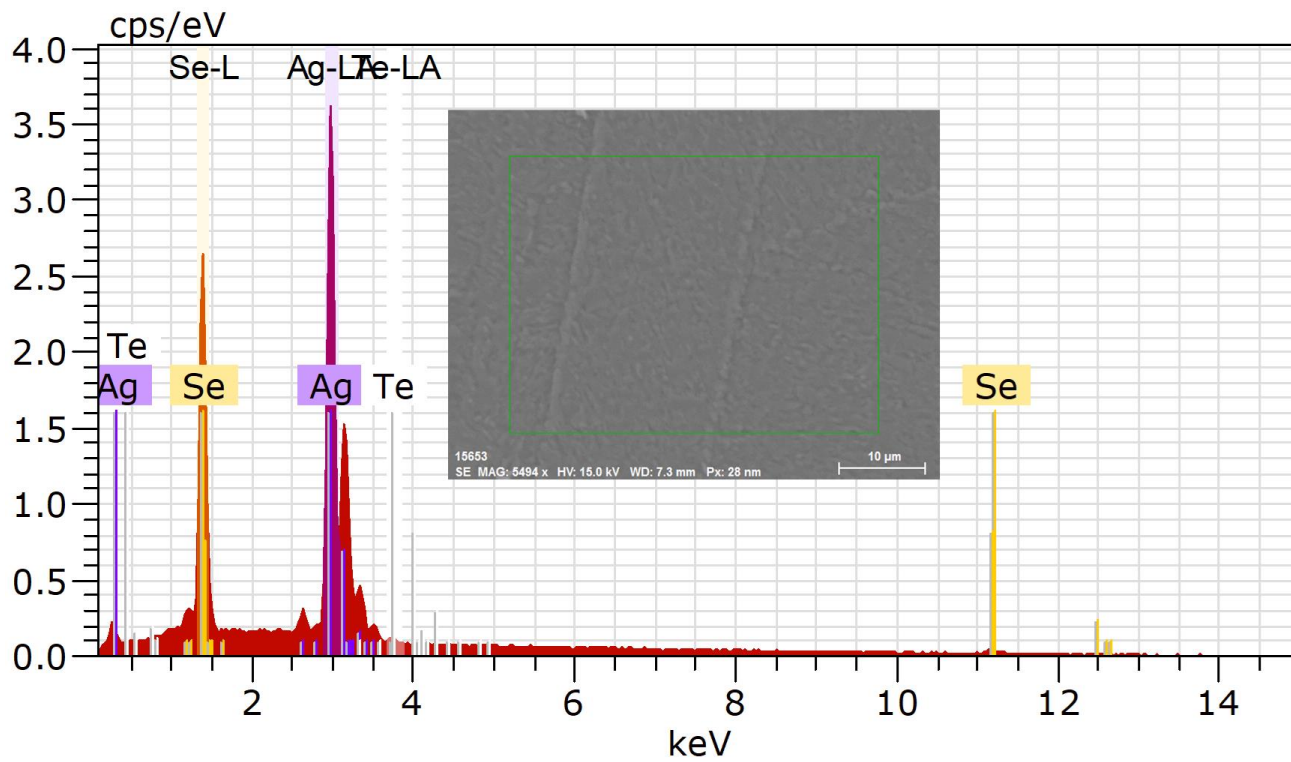
Supplementary Fig.4. EDS results with inset of SEM image for Ag_2Se thin films with 1.3 at.% Te.



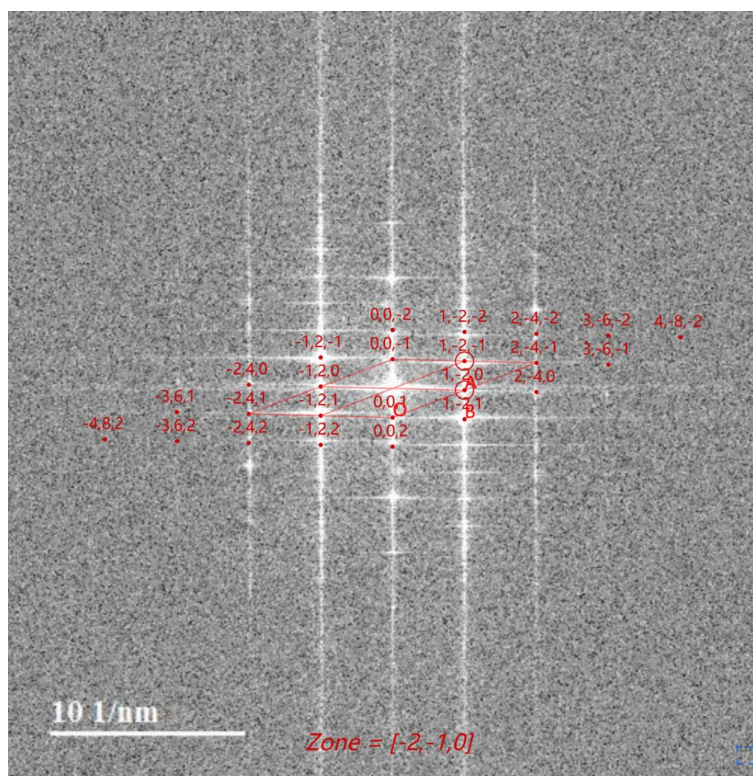
Supplementary Fig. 5. EDS results with inset of SEM image for Ag_2Se thin films with 1.9 at.% Te.



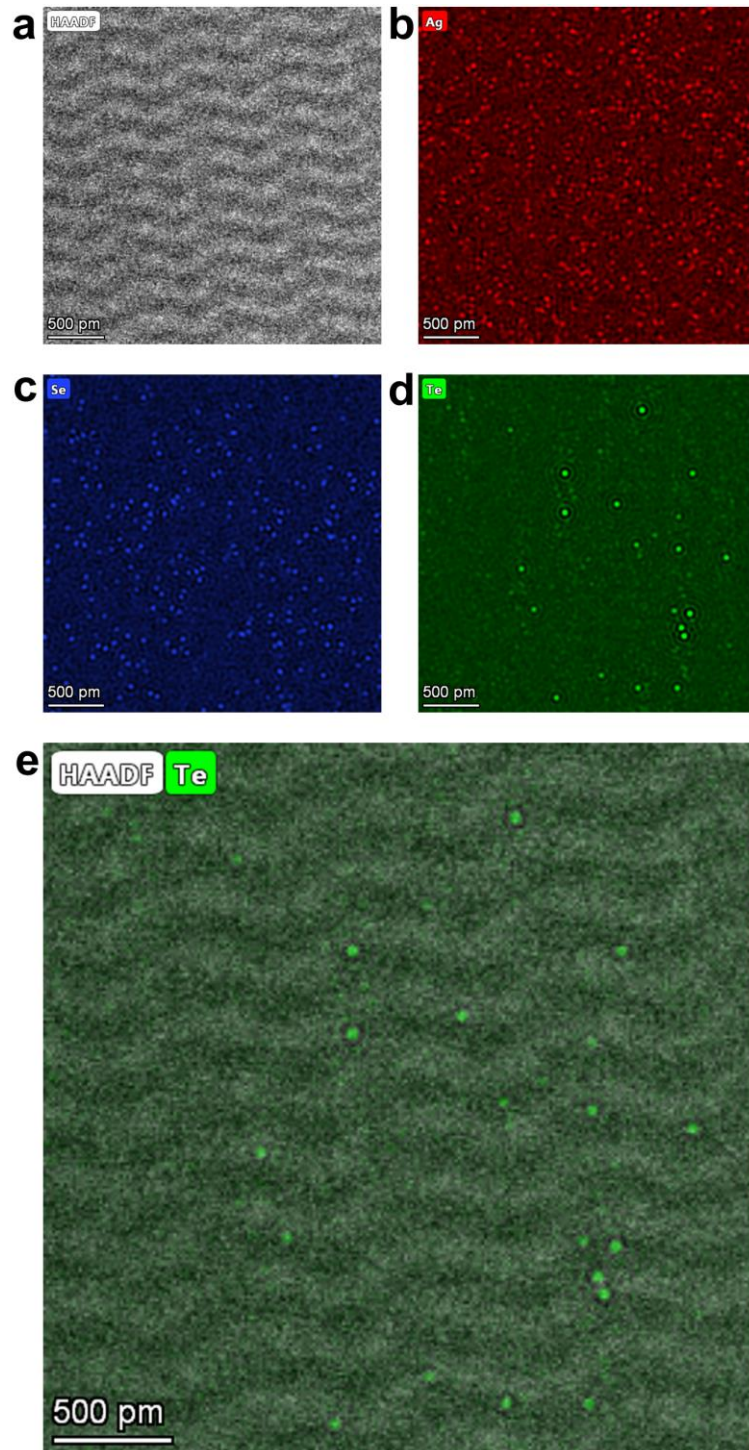
Supplementary Fig. 6. EDS results with inset of SEM image for Ag_2Se thin films with 2.6 at.% Te.



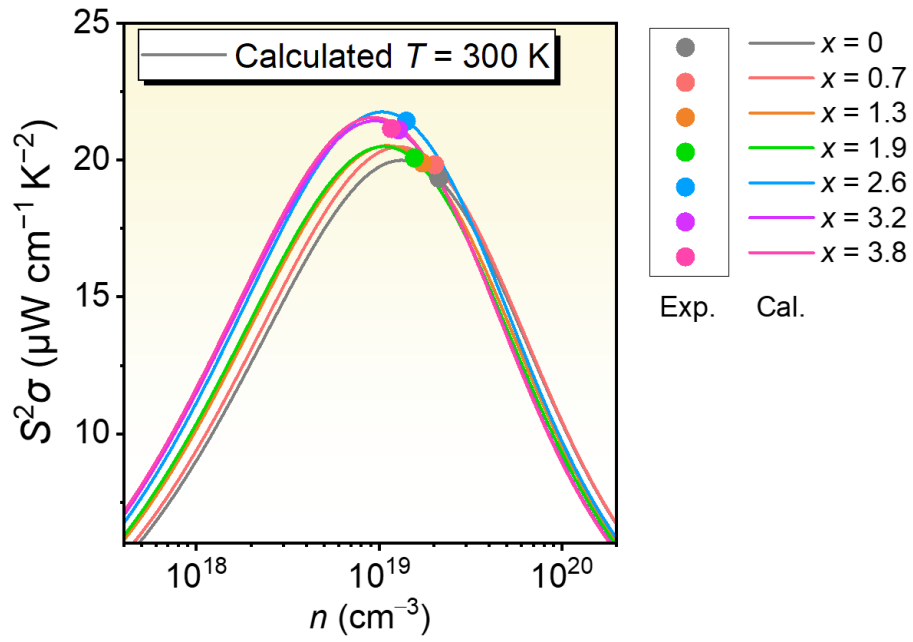
Supplementary Fig. 7. EDS results with inset of SEM image for Ag_2Se thin films with 3.8 at.% Te.



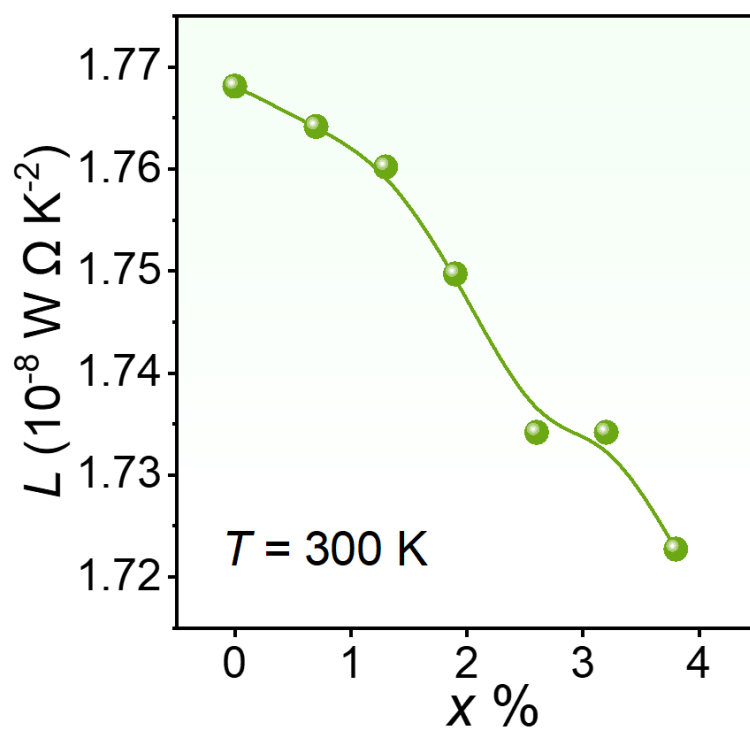
Supplementary Fig. 8. Magnified indexed fast Fourier transform (FFT) pattern taken from **Figure 3e** in the manuscript.



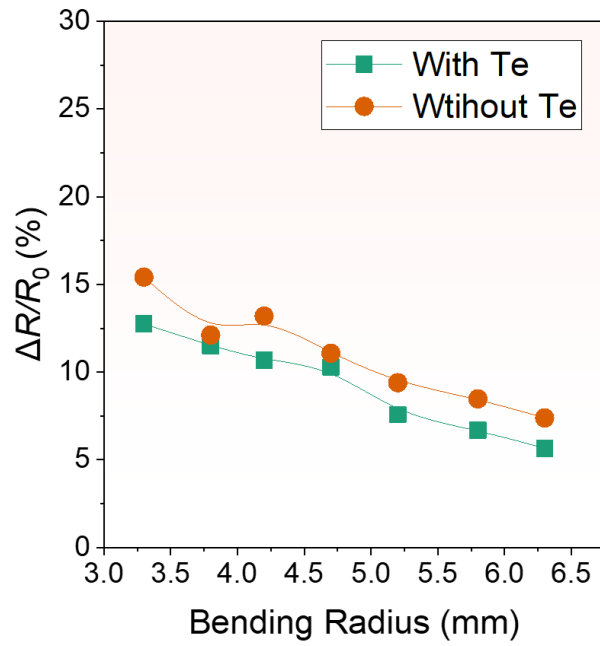
Supplementary Fig. 9. (a) Spherical aberration-corrected scanning TEM (Cs-STEM) HAADF image of Ag₂Se thin film with 3.2 at.% Te, and corresponding EDS maps of (b) Ag, (c) Se, and (d) Te. (e) Overlap of the HAADF image and the EDS map of Te.



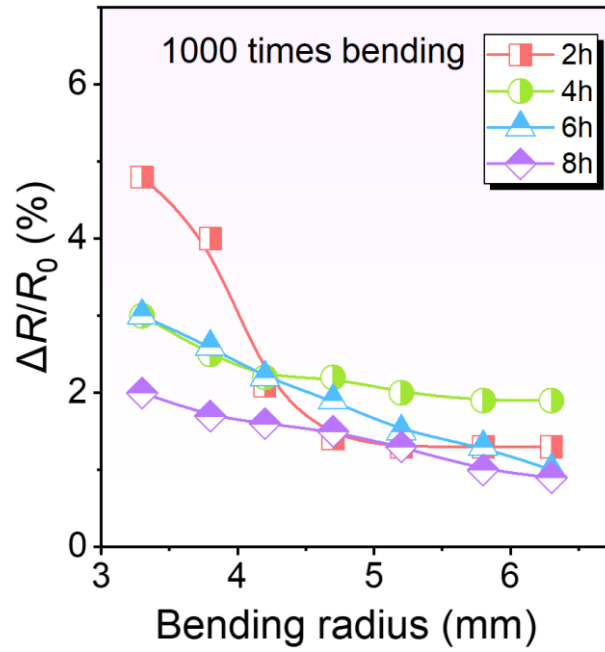
Supplementary Fig. 10. Comparison of calculated and experimental power factor $S^2\sigma$ as a function of carrier concentration n . The calculated data are from the single parabolic band (SPB) model.



Supplementary Fig. 11. Te concentration x -dependent Lorenz number L at room temperature.



Supplementary Fig. 12. Measured increased normalized resistance $\Delta R/R_0$ of the 3.2 at.% Te-doped and undoped Ag_2Se thin films at different bending radii.

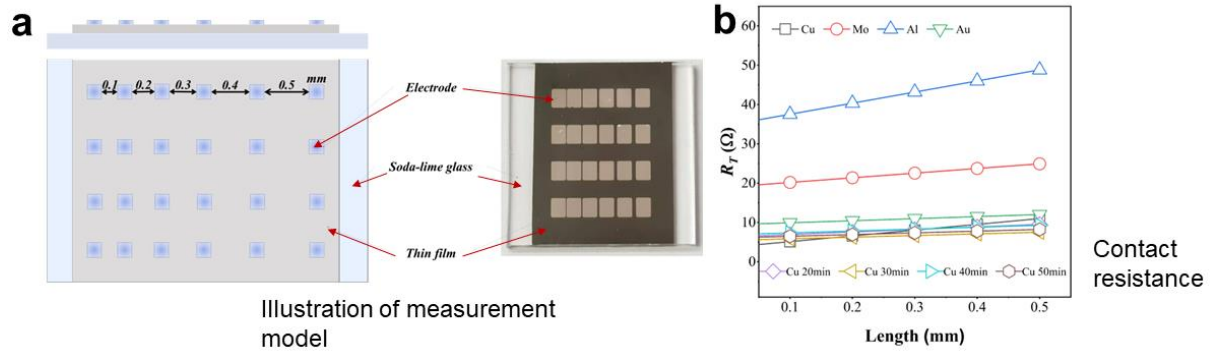


Supplementary Fig. 13. Measured increased normalized resistance $\Delta R/R_0$ of the coated Ag_2Se thin films with 3.2 at.% Te at different bending radii and heating at 95 °C by different periods for 1000 times bending.

Measurement method: Transmission line measurement (TLM)

Measurement tool: 2400 Series Digital Source-meter

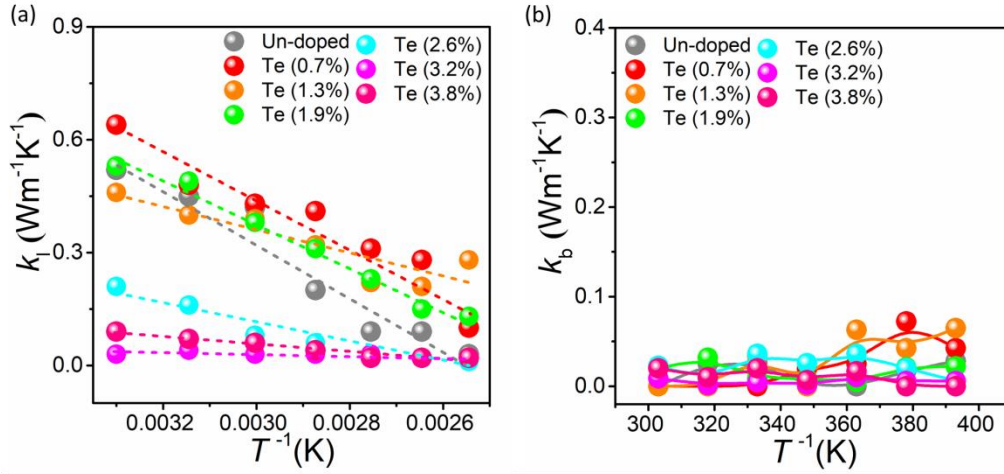
Electrode: 3 mm x 1.5 mm



Supplementary Fig. 14. (a) Illustration of measurement model for evaluating the contact resistance.

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function of length for different electrode materials.



Supplementary Fig. 15. (a) The dotted line is linearly fitting to the lattice thermal conductivity (κ_l) at low temperature. (b) Bipolar thermal conductivity (κ_b) for Ag₂Se thin films with different Te concentrations ($x = 0, 0.7, 1.3, 1.9, 2.6, 3.2,$ and 3.8 at.%)

3. Supplementary Table

Supplementary Table 1. A summary of the thermoelectric performance of Ag₂Se-based films. Here, poly(3,4-ethylenedioxythiophene) is abbreviated as PEDOT; polyvinylidene fluoride is abbreviated as PVDF; polyvinylpyrrolidone is abbreviated as PVP; single wall carbon nanotubes is abbreviated as SWCNTs; polypyrrole is abbreviated as PPy; polypropylene is abbreviated as PI; polyethylene naphthalate two formic acid glycol ester is abbreviated as PEN; and polyethylene terephthalate is abbreviated as PET.

Materials	Substrate	ZT	T (K)	S (μV K ⁻¹)	σ (S cm ⁻¹)	S ² σ (μW cm ⁻¹ K ⁻²)	κ (W m ⁻¹ K ⁻¹)	n (cm ⁻³)	μ (cm ² V ⁻¹ s ⁻¹)	Year	Ref.
Ag ₂ Se thin film with 3.2 at.% Te	PI	1.27	363	–132	1425	24.8	0.71	1.28×10 ¹⁹	444	2023	This work
Ag ₂ Se	PI	~1.2	300	–135	~1440	~25.9	~0.66	6.5×10 ¹⁸	1250	2022	²¹
Ag ₂ Se+PVP	Nylon	1.1	300	–144	925	19.1	/	7×10 ¹⁸	1100	2020	²²
Ag ₂ Se	PI	1.1	300	–161.7	840	22.05	0.61	7.3×10 ¹⁸	721.3	2022	²³
Printed-Ag ₂ Se	PEN	1.03	300	–185	460	17	0.5	/	/	2020	²⁴
Ag ₂ Se	Scaffold	~1	300	–183	472	15.8	~0.47	/	/	2020	²⁵
Ag ₂ Se+Se+PPy	Nylon	0.94	300	–144	1064	22.4	/	9.5×10 ¹⁸	762	2021	²⁶
Ag ₂ Se+bacteria cellulose	/	0.7	400	–167	230	6.24	0.36	/	/	2022	²⁷
Ag ₂ Se+nylon	Substrate-free	0.68	300	–137.	958.9	18.25	/	~1.2×10 ¹⁹	~480	2022	²⁸

				9							
Ag ₂ Se	Nylon	0.6	30 0	– 141	497	9.87	/	3.8×10^1 8	850	201 9	²⁹
Ag _{2.02} Se	PI	~0. 6	34 8	~ 110	~178 5	21.6	~1.2	/	/	202 2	³⁰
Printed film	Ag-Se Glass	0.5 5	30 0	– 220	~108	~5	/	/	/	202 0	³¹
Ag ₂ Se	Substrat e-free	0.5	30 0	– 120	1060	~15. 3	~0.8 9	4.79×10 18	139 0	202 3	³²
Ag ₂ Se	Glass	/	37 6	– 216	974	~46. 6	/	1×10^{19}	634	201 7	³³
Ag _{2.2} Se	PI	/	30 0	– 107	3010	~34. 2	~0.5 6	2.5×10^1 9	750	202 1	³⁴
Ag _{2.3} Se	Paper	/	30 3	– 122	~166 0	~24. 5	/	2×10^{19}	~41 5	202 0	³⁵
Ag ₂ Se+Ag+CuA gSe	Nylon	/	30 0	–45	1050	22.3	/	2.7×10^2 0	250	201 9	³⁶
Ag ₂ Se	PI	/	32 0	– 142	1091	21.9	/	/	/	202 3	³⁷
Cu-doped Ag ₂ Se	PI	/	30 0	– 125	~135 0	20.8	/	8.6×10^1 8	~90 0	202 2	³⁸
S-doped Ag ₂ Se	PET	/	30 0	– 150	~900	20.5 8	/	$\sim 1 \times 10^{19}$	~57 0	202 3	³⁹
Ag ₂ Se	Nylon	/	30 0	– 150	~908	20.4 3	/	5.82×10 18	127 7	202 2	⁴⁰
Ag ₂ Se +SWCNTs	Nylon	/	30 0	– 108. 1	~165 7	~19. 4	/	1.31×10 19	/	202 3	⁴¹
Ag _{1.8} Se	PI	/	38 0	– 138	1000	19	/	7.5×10^1 8	950	202 1	⁴²
Ag ₂ Se	Nylon	/	30 0	– 143	~920	~18. 8	/	7.9×10^1 8	102 4	202 0	¹⁶
Ag ₂ Se+Ag	Nylon	/	30 0	– 67.5	3958	~18. 6	/	1.6×10^2 0	~83	202 1	¹⁷
Ag ₂ Se	Glass	/	30 0	– 120	~120 0	17.3	/	1×10^{19}	714	201 6	¹⁸

Ag ₂ Se+CuAgSe + PEDOT	Nylon	/	300	–121	~1100	16	~0.46	7.4×10^1 ⁸	940	2020	¹⁹
<i>In-situ</i> Ag _{2.09} Se	/	/	423	–110	~1300	15.3	/	1.5×10^1 ⁹	550	2022	⁴³
Ag ₂ Se+Ag+PEDOT	Nylon	/	300	–49.5	~5957	14.4	/	4.4×10^1 ⁹	78	2021	⁴⁴
Ag ₂ Se+SWCNTs	Nylon	/	397	–118	~880	~12.4	/	7.8×10^1 ⁸	~700	2021	⁴⁵
Ag _{1.98} Ga _{0.02} Se	Nylon	/	300	–115	~880	~11.6	/	/	/	2021	⁴⁶
Ag ₂ Se	Glass	/	300	–123	~730	11	/	4.8×10^1 ⁹	94	2021	⁴⁷
S-doped Ag ₂ Se	Nylon	/	300	–106	849	~9.5	/	6×10^{18}	~900	2022	⁴⁸
Ag ₂ Se+PEDOT	/	/	300	–91	520	~4.3	/	/	/	2022	⁴⁹
Ag ₂ Se+PVDF	Glass	/	300	–92	~370	~2.8	/	/	/	2020	⁵⁰

Supplementary Table 2. Debye-Callaway model-based pre-factors for all Te-doped Ag₂Se samples.

Parameters	Samples						
x	0	0.1	0.2	0.3	0.4	0.5	0.5
A_1 (10^{-41} s ³)	0.005	0.61	0.56	0.77	18.32	41.17	55.51
A_2 (10^{-18} s/K)	4.71	0.21	0.29	0.47	17.60	50.14	44.50
A_3 (10^{-15} s)	6.44	6.44	6.44	6.44	6.44	6.44	6.44
A_4 (10^{-25} s ²)	0.0	0.002	0.002	0.004	0.021	0.029	0.032
d (nm)	----	300	270	220	180	160	178

Supplementary Table 3. A summary and comparison of the power densities of reported thermoelectric flexible devices.

Device materials	Output power density (mW cm⁻²)
Cu ₃ Se ₁ single leg ⁵¹	~0.09 ($\Delta T = 30$ °C)
N-Bi ₂ Te ₃ +P-Sb ₂ Te ₃ ⁵²	~0.1 ($\Delta T = 20$ °C)
N-Bi ₂ Te ₃ (PEDOT:PSS)+P-Sb ₂ Te ₃ (PEDOT:PSS) 53	~0.2 ($\Delta T = 20$ °C)
Bi ₂ Te ₃ with 1 wt% Se single leg ⁵⁴	~0.14 ($\Delta T = 20$ °C)
P-Sb + N-Ni ⁵⁵	~0.9 ($\Delta T = 20$ °C)
N-Bi ₂ Te ₃ +P-Sb ₂ Te ₃ ⁵⁶	~1.42 ($\Delta T = 60$ °C)
Ag ₂ Se single leg ³⁶	~0.16 ($\Delta T = 20$ °C)
Ag ₂ Se single leg ²⁹	~0.23 ($\Delta T = 30$ °C)
Ag ₂ Se single leg ²¹	~1.0 ($\Delta T = 20$ °C)
Ag ₂ Se single leg ³⁸	~1.4 ($\Delta T = 20$ °C)
This work	1.5 ($\Delta T = 20$ °C)

Supplementary Table 4. A summary of reported surface energies of other materials.

Materials	Surface Energy (J m^{-2})
Mg_2Si (100) ⁵⁷	0.8-2.0
Al_3Ti (010) ⁵⁸	1.798
Al_3Ti (001) ⁵⁸	2.374/1.167
Boron Carbide Polytype $\text{B}_{11}\text{C}_p(\text{CBC})$ ($10\bar{1}1$) ⁵⁹	3.21
Co (0001) ⁶⁰	2.110
Co_3Cr (0001) ⁶⁰	2.170
Co_3Mn (0001) ⁶⁰	2.766
Co_3Ni (0001) ⁶⁰	2.012
CoSb_3 (100) ⁶¹	0.330
CoSb_3 (110) ⁶¹	0.138
CoSb_3 (111) ⁶¹	0.797
Ni_3Nb (100) ⁶²	2.51
Ti_2AlC (0001) ⁶³	0.078
This work Ag_2Se (001)	0.032-0.064 (0.002 eV/Å ² -0.004 eV/Å ²)

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