

Supplementary Information for:

Microgravity effects on nonequilibrium melt processing of neodymium titanate: thermophysical properties, atomic structure, glass formation and crystallization

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SUPPLEMENTARY DISCUSSION

Pyrometry and temperature corrections

Temperature measurements in ELF used an optical pyrometer sensitive to 1.45-1.8 μm radiation, or nominally $\lambda = 1.55 \mu\text{m}$. The apparent temperature measured by the pyrometer, T_{app} , must be corrected to obtain the absolute temperature, T_{abs} , according to Wien's law:

$$\frac{1}{T_{\text{abs}}} - \frac{1}{T_{\text{app}}} = \frac{\lambda \ln(\epsilon_{\text{eff}})}{C_2} \quad (1)$$

In Supplementary Equation 1, λ is the pyrometer wavelength, $C_2 = 1.4388 \text{ cm K}$ is the second radiation constant, and ϵ_{eff} is the effective emissivity, which contains contributions from the sample and any windows (reflective surfaces) between the sample and pyrometer. In ELF, two sapphire windows separate the pyrometer's detector from the sample¹. The index of refraction, n , and the reflection per window surface, R_0 , are related by the Fresnel equation:

$$R_0 = \left| \frac{n_1 - n_2}{n_1 + n_2} \right|^2 \quad (2)$$

where $n_1 = 1.77$ for sapphire and $n_2 = 1.0$ for air, so $R_0 = 0.0773$. The sample emissivity is estimated as 0.86, based on the previously measured refractive index of 2.2 for lanthanum titanate glass². This yields an effective emissivity of $\epsilon_{\text{eff}} = 0.86 \times (1 - 0.0773)^4 = 0.63$.

The effective emissivity was also estimated via a second approach, using the recalescence temperatures of materials with known melting points. During crystallization from a supercooled melt, the sample is expected to self-heat back up to its equilibrium melting point (T_m) as the latent heat of fusion is released. Using this assumption, the peak temperature of recalescence would correspond to T_m , from which ϵ_{eff} could be calculated using Supplementary Equation 1 and the measured T_{app} . However, this assumption can often fail for at least two reasons: (i) if the liquid has supercooled sufficiently, its latent heat is not sufficient to heat back to T_m (also called hypercooling); (ii) the solid's emissivity is often lower than that of the melt, due to loss of a specular, spherical surface. In both of these scenarios, ϵ_{eff} would be underestimated by this method based on the recalescence temperature. With that acknowledgement, several recalescence events were analyzed for materials other than the NT samples reported in this paper: Al_2O_3 , Gd_2O_3 , Tm_2O_3 , lanthanum titanate, and barium titanates. Of this set, the ϵ_{eff} estimates typically ranged 0.6 to 0.7.

Based on these two approaches to estimating ϵ_{eff} , a value of 0.63 was used for all analyses in this study. The temperature uncertainty arising from $\epsilon_{\text{eff}} = 0.63 \pm 0.05$ is $\pm 30 \text{ }^\circ\text{C}$ at $1800 \text{ }^\circ\text{C}$ and is smaller at lower temperatures.

When liquids were held isothermally in ELF, pyrometer noise was typically $\pm 15 \text{ }^\circ\text{C}$ (standard deviation). Thus, the overall temperature uncertainty for the ELF measurements in this study is estimated as $\pm 30 \text{ }^\circ\text{C}$.

In the terrestrial aerodynamic levitator, the temperature of liquid samples varies spatially by $\sim 40 \text{ }^\circ\text{C}$ from top to bottom, since the single heating laser only heats the top of the sample.³ Using the sample emissivity of 0.86 and appropriate window corrections, $\epsilon_{\text{eff}} = 0.79$. The temperature measurement uncertainty arising from $\epsilon_{\text{eff}} = 0.79 \pm 0.05$ is $\pm 27 \text{ }^\circ\text{C}$ at $1800 \text{ }^\circ\text{C}$ and is smaller at lower temperatures.

Density analysis and possible measurement artifacts

Glass density at room temperature

A literature estimate for NT glass density is available from pycnometry measurements by Arai *et al.*⁴ on lanthanum titanate (LT) glasses. In their study, LT with 18.2 mol. % La_2O_3 had a density of 4.89 g cm^{-3} , which is 98.6% of the density for the compositionally identical crystal phase $\text{La}_4\text{Ti}_9\text{O}_{24}$ (4.96 g cm^{-3}). Assuming the same ratio of glass-to-crystal density, NT glass with 18.2 mol. % Nd_2O_3 is expected to have a density of 5.10 g cm^{-3} (using $\rho = 5.18 \text{ g cm}^{-3}$ for crystalline $\text{Nd}_4\text{Ti}_9\text{O}_{24}$ ⁵). Arai *et al.* also observed a density increase of 2.3% in LT glasses as La_2O_3 content increased from 15.4 to 18.2 mol. %. Based on our atomic structural characterizations, the NT microgravity glasses likely contained ~ 19 mol. % Nd_2O_3 , so a linear extrapolation of Arai's compositional trend would yield a final estimate of 5.13 g cm^{-3} for the MG1 and MG2 glasses here. Thus, the room temperature glass density according to the ELF measurements, 5.28 g cm^{-3} , is 2.9% larger than the literature-based estimate.

For comparison, the MG1 and MG2 glass densities at room temperature have been assessed using two other techniques in this study, X-ray tomography and light microscopy. The results are summarized in Supplementary Table 2 alongside the ELF measurement and literature-based estimate. The volume calculations from tomography and microscopy are assumed to have uncertainties of ca. 5%. In X-ray tomography, uncertainty arises from the image calibration factor and the thresholding algorithm used for image segmentation. For light microscopy, uncertainty arises from focal depth and assuming the sample is a perfect spherical shape. For ELF, an uncertainty of 2.5% is listed in Supplementary Table 2 based on prior reports⁶. For the estimate based on Arai *et al.*⁴, no uncertainty is given in Supplementary Table 2 since no information on replicate measurements were reported. All four values for glass density in Supplementary Table 2 are almost within their collective uncertainties.

Anomalous thermal expansion near T_g

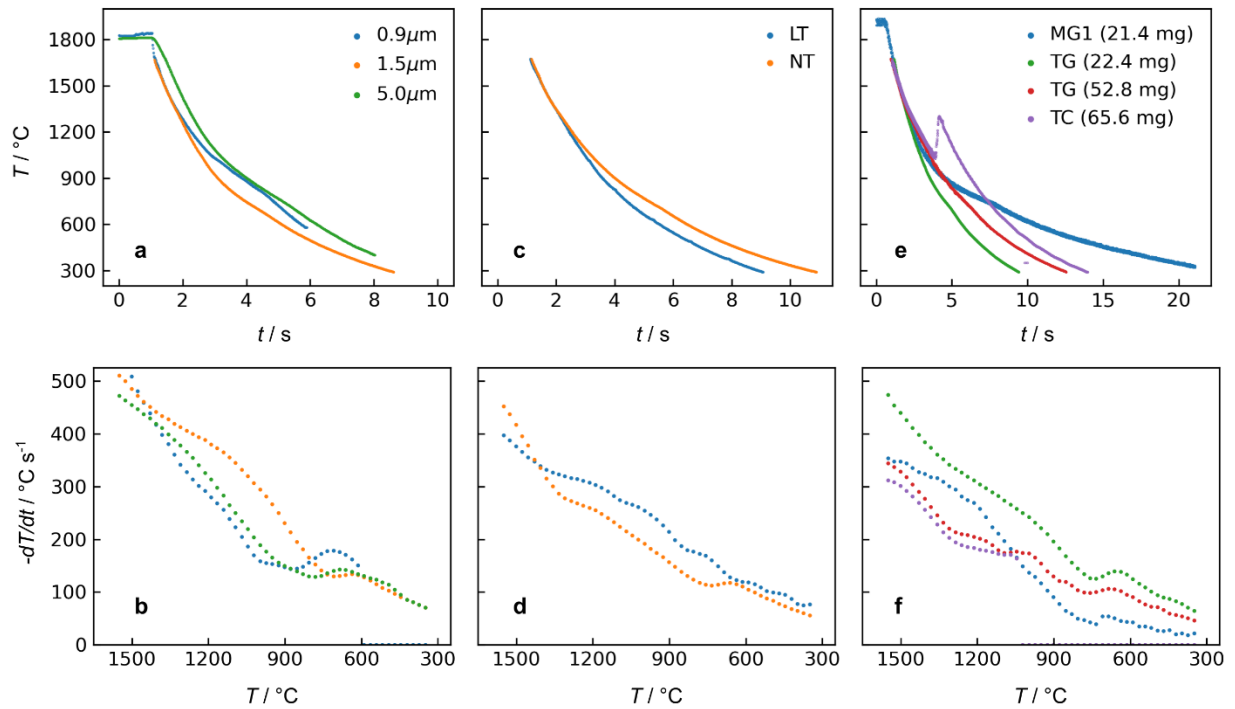
As discussed in the main text, the density-temperature relationship in Fig. 2 exhibits a steeper slope between 720-930 °C than either the glass or liquid regions. To explain this unexpected observation, several possible scenarios involving measurement artifacts were explored, including: (i) sample transparency to the silhouette backlight, (ii) changes in bubble volume inside the sample, and (iii) sample transparency at the pyrometer wavelength. These are discussed in detail below.

ELF measurements of liquid densities have been shown to be in good agreement with prior (terrestrial) studies, for example with Al_2O_3 ¹, Y_2O_3 ⁷, lanthanoid sesquioxides⁸, Ga_2O_3 ⁹, Zr ¹⁰, and Au ¹¹. These examples provide validation of the ELF instrument for measuring liquid density. However, the measurements here for NT are the first to be reported from ELF for glass density upon cooling, so we explored whether transparency of the glassy state may be leading to an underestimation of the sample volume. Specifically, if the sample becomes partially transparent near the wavelength of the ultraviolet backlight, some of the backlight may pass through the edges of the sample, making its silhouette appear smaller in the camera image than the true size. Sample volume is calculated based on edge detection of this silhouette image (see Methods section). The typical ELF sample has a radius of ~ 120 pixels in the camera image, so if sample transparency introduces an error of ~ 1 pixel to the edge detection algorithm, a volume error of $\sim 2.5\%$ would result. This error may onset suddenly during cooling if the sample transparency changes at a particular temperature, which matches what is observed in the anomalous density increase (i.e., volume decrease) ca. 720-930 °C near $T_g = 786$ °C. For illustrative purposes, Supplementary Fig. 6a compares the measured data against the hypothetical “true” density suggested by this explanation (green lines). To test this hypothesis, precision spheres of optically-transparent ruby, sapphire, and BK7 glass were levitated at room temperature under vacuum in a ground-based electrostatic levitator that uses the same silhouette imaging technique. The measured volumes of all three standards matched the expected values within 0.5%. This provides strong evidence against the hypothesis of edge transparency causing an underestimation of sample volume.

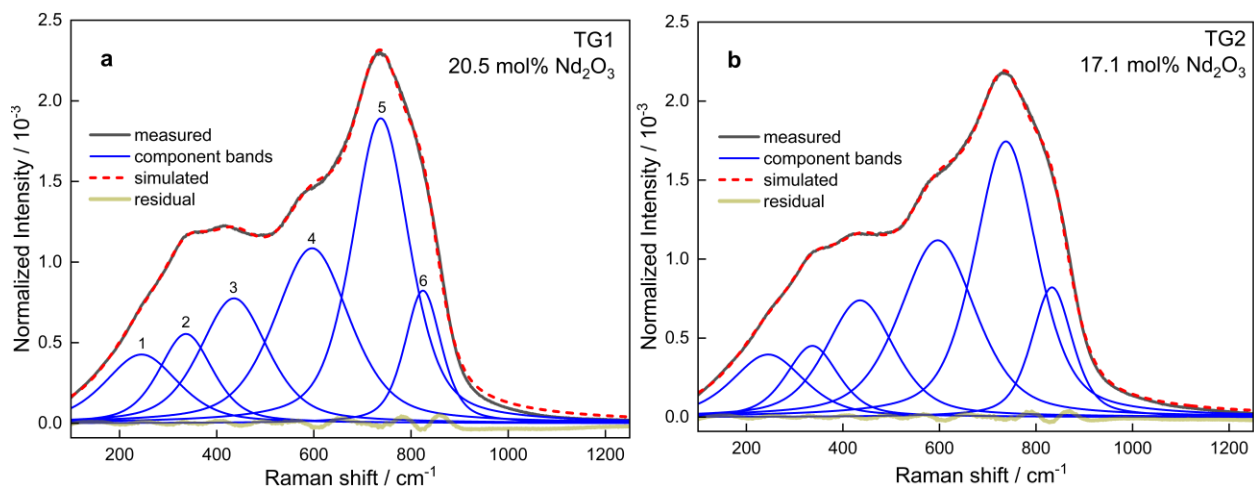
A second hypothesis for explaining the anomalous 720-930 °C region is that gas bubbles inside the samples shrank suddenly near 930 °C and then became fixed in size as the sample became viscous near T_g . The shrinking bubbles would result in the apparent sample volume decreasing, and the calculated density would increase. Supplementary Fig. 6b provides a comparison of the measured data and the hypothetical density suggested by this explanation. Since the internal voids of the glass samples were measured to be $\leq 0.25\%$ of the total volume, the gas bubbles would have to be $10\times$ larger in the melt to account for the magnitude of the density discrepancy. This seems unlikely. Furthermore, the two replicate samples exhibited reproducible density-temperature relationships (Fig. 2), and it is unlikely that such a bubble process would occur identically in both samples.

A third hypothesis involves the sample becoming partially transparent at the pyrometer wavelength during cooling. If the transition to partial transparency began near 930 °C, then the pyrometer would start seeing some of the radiation emitted by the sample's hotter interior. This would result in a higher reading than the surface temperature, until the sample interior had also cooled enough to become partially transparent. Supplementary Fig. 6c shows a comparison of the measured data and the hypothetical density suggested by this explanation. This scenario seems the most convincing of the considered measurement artifacts, so it was presented in the main text.

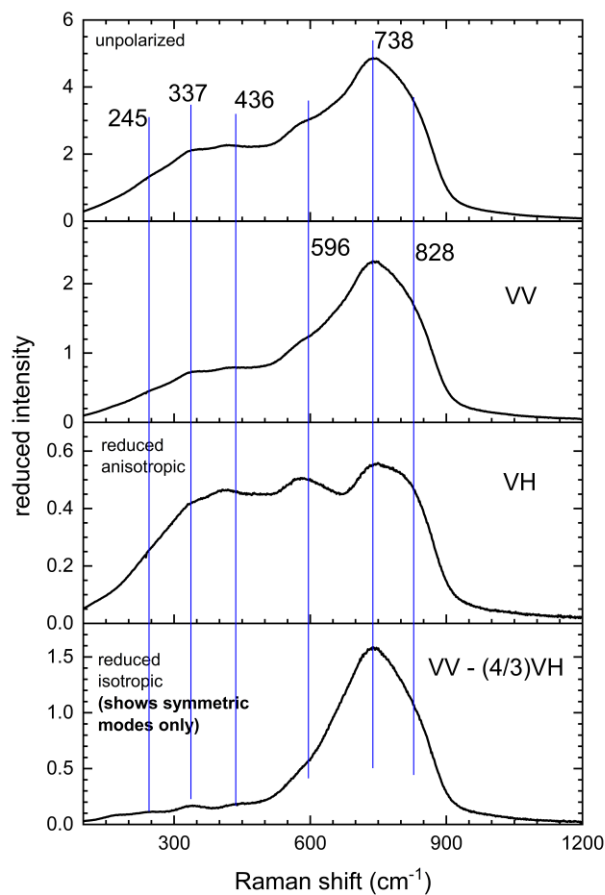
SUPPLEMENTARY FIGURES



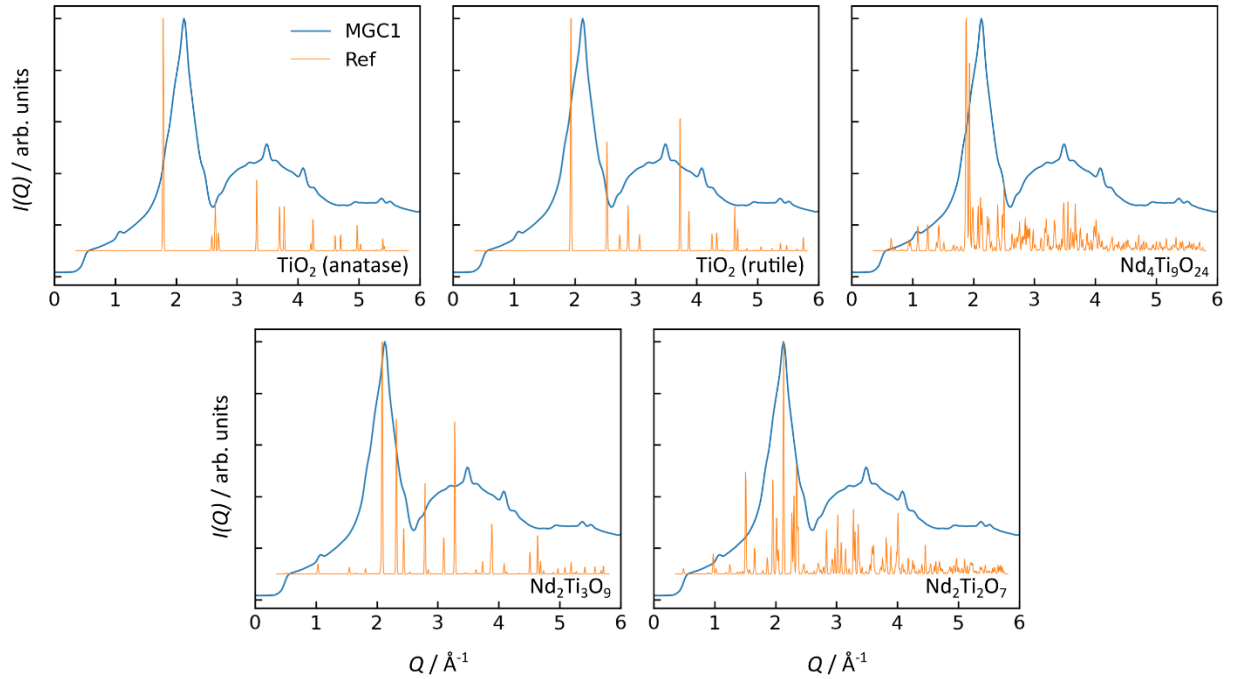
Supplementary Fig. 1. Pyrometry and cooling rates of molten titanates processed using the terrestrial aerodynamic levitator. **a**, Cooling curves and **b**, cooling rates for a single NT sample (21.8 mg) in pure O_2 gas measured with three different pyrometers: $\lambda = 0.9, 1.5,$ or $5.0 \mu\text{m}$. **c**, **d**, Comparison of rare earth titanates (nominally $83\text{TiO}_2\text{-}17\text{RE}_2\text{O}_3$) containing either RE = La ("LT", 33.8 mg total mass) or Nd ("NT", 32.6 mg). **e**, **f**, Comparison of one sample vitrified in microgravity (MG1) with three terrestrial samples of different masses. Terrestrial samples of 22.4 or 52.8 mg mass formed glass (TG), while a 65.6 mg sample crystallized (TC) upon supercooling to $\sim 1040^\circ\text{C}$. All data in **c-f** are for processing in air and collected at $\lambda = 1.5 \mu\text{m}$.



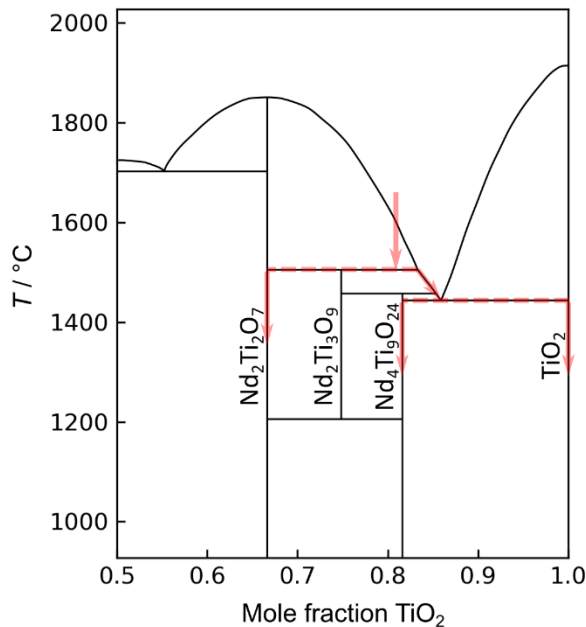
Supplementary Fig. 2. Fitting of component bands in Raman spectra. a, TG1 and b, TG2 terrestrial glasses. Optimized fitting parameters are listed in Supplementary Table 1.



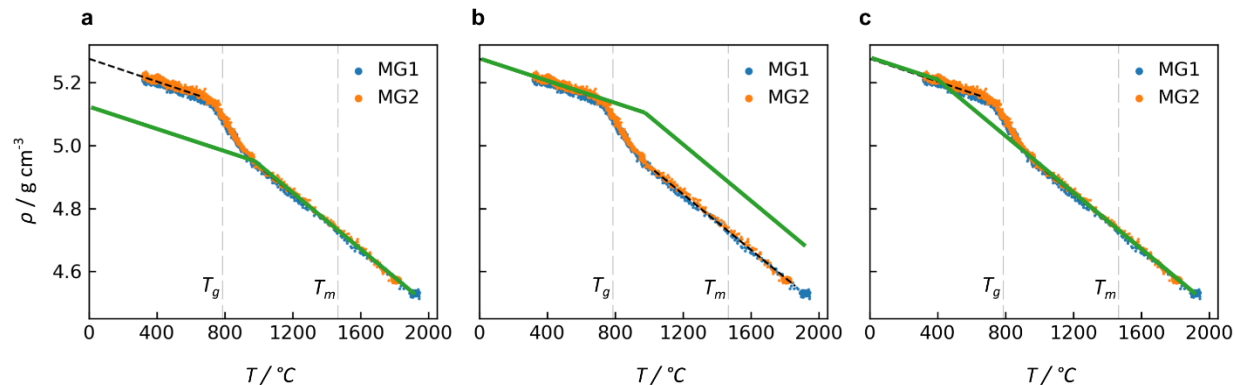
Supplementary Fig. 3. Comparison of unpolarized and polarized Raman spectra. For microgravity glass MG1. From top to bottom: unpolarized, parallel (VV) and perpendicular (VH) polarization, and the reduced isotropic spectrum, which shows only symmetric modes that are found in the high frequency envelope. Annotated bands are summarized in Supplementary Table 1 and discussed in the main text.



Supplementary Fig. 4. X-ray diffraction intensity for NT glass-ceramic processed in microgravity (MGC1). The small Bragg peaks in the MGC1 diffraction are compared with reference patterns for possible matching crystalline phases: TiO_2 (anatase)¹², TiO_2 (rutile)¹³, $\text{Nd}_4\text{Ti}_9\text{O}_{24}$ ⁵, $\text{Nd}_2\text{Ti}_3\text{O}_9$ ¹⁴, and $\text{Nd}_2\text{Ti}_2\text{O}_7$ ¹⁵.



Supplementary Fig. 5. A portion of the Nd_2O_3 - TiO_2 equilibrium phase diagram. Adapted from Gong and Zhang¹⁶, showing likely solidification pathway (in light red shading) for the NT sample that crystallized during melt processing in microgravity. Annotations mark the line phases discussed in the main text.



Supplementary Fig. 6. Hypothetical measurement artifacts to explain the anomalous density-temperature region ca. 720-930 $^{\circ}\text{C}$. **a**, Sample transparency to the silhouette backlight, leading to underestimation of sample volume below 930 $^{\circ}\text{C}$; **b**, large bubbles in the melt that shrink as the sample approaches T_g ; **c**, sample transparency at the pyrometer wavelength. The green lines illustrate the hypothetical “true” density that is suggested by each of these scenarios.

SUPPLEMENTARY TABLES

Supplementary Table 1. Raman peak assignments for NT glasses. Peak fitting results are shown in Supplementary Fig. 2.

Position (cm ⁻¹)	Assignment	Reference	TG1 Height (×10 ²)	Width	TG2 Height (×10 ²)	Width
245	Network deformation and/or RE-O stretching	17,18	0.04264	179	0.03963	176
337	RE-O stretching and/or bending associated with corner-sharing TiO ₆ units	17,18	0.0554	131	0.04511	128
436	Ti-O-Ti bending	19,20	0.07748	165	0.07392	167
596	Symmetric stretching of edge-sharing TiO ₆	19,20	0.10854	184	0.11191	187
738	Symmetric stretching of corner-sharing TiO ₆	19,20	0.18913	140	0.17447	153
~828	TiO ₅	19–24	0.08227	89	0.08213	100

Supplementary Table 2. Estimates of NT glass density.

Method	Density (g cm ⁻³)
ELF	5.28(13)
Tomography	4.86(24)
Light microscopy	5.20(26)
Pycnometry, Arai <i>et al.</i> ⁴	5.13

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