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Supporting Information for

- **Elucidating the role of water during collagen assembly by isotopically inhibiting**
- **water-collagen interactions**

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¹⁴ **Supporting Information Text**

 A. Turbidity curve analysis. To identify the rate of increase (*k*) of the turbidity curves, we need to find the time range where the growth is linear. To do so, we used the first-derivative curves, based on a similar approach used in ref.[\(1\)](#page-2-0). We selected the time range between the time value where the first derivative is ∼70% of the maximum derivative, and then we fitted a line ¹⁸ using least squares fitting. The slope of the fitted line is k . The t_{lag} is calculated by taking the time at which the tangent line to the turbidity curve at the maximum slope intersects the horizontal axis, whereas *tplateau* is calculated by taking the time at which the tangent intersects the orizontal line at the maximum turbidity value.

B. HDO measurements. We studied collagen assembly and properties in HDO solutions, containing a H_2O : D_2O at a volume 22 ratio of 1:1. In Fig[.S1,](#page-3-0) we reported the turbidity measurements performed at 1 mg/ml in H₂Oand HDO. For experiments 23 in HDO, we used the stock solution of collagen dissolved in water containing 0.2 wt $\%$ of acetic acid (same stock solution of $_{24}$ the H₂O experiments); and we then neutralized it by using a customized D₂O-buffer to obtain a final collagen concentration 25 of 1 mg/ ml with a 1:1 volume ratio of H₂O:D₂O. The reported continuous curves and shades represented the mean and 26 standard deviation over three different measurements. We noticed that assembly in HDO happens faster than in H₂O. To ²⁷ check whether there is a dependence on the solvent used for the stock solutions, we repeated the same measurements by using ²⁸ the stock solution of collagen dissolved in heavy water containing 0.2 wt % of acetic acid, and neutralizing it with customized $_{29}$ H₂O-buffer to obtain again a final collagen concentration of 1 mg/ml with a 1:1 volume ratio of H₂O:D₂O (dashed grey line in 30 Fig[.S1\)](#page-3-0). We again observed a faster assembly than in H₂O, indicating an independence on the solvent used for the original 31 stock solution. In Fig[.S4A](#page-6-0), we reported the rheological measurements performed at 1.25 mg/ml of collagen in H_2O , HDO and 32 D₂O. In this case for the HDO measurements, we used a stock solution of collagen dissolved in HDO containg 0.2 wt $\%$ of ³³ acetic acid. Similar to the turbidity measurements, we observed a faster fibrillization in HDO; further, the elastic modulus ³⁴ was lower in HDO compared to H₂O, indicating a significant effect on the collagen gel properties already when 50% of water ³⁵ was replaced with heavy water. Additionally, we performed frequency-sweep (10-0.1 Hz) oscillatory rheology measurements ³⁶ of equilibrium collagen network with a concentration of 1.25 mg/mL of collagen concentration at a strain of 0.5% (Fig[.S4B](#page-6-0)). 37 The presence of D₂O does not influence the dynamics of stress relaxation in the collagen network. With a power-law with an ³⁸ exponent of ∼0.1 for all collagen networks, G' is weakly dependent on the frequency, *ω*. This is typical behavior of biopolymer 39 networks that are transiently cross-linked by non-covalent interactions. $(2, 3)$ $(2, 3)$ $(2, 3)$

 C. D-band analysis. Although the D-band is not easy to resolve, we can resolve the sub-banding peaks (Fig[.S5\)](#page-7-0). To investigate whether the collagen alignment is quarter-staggered and whether it is the same in water and heavy water, we plotted the 42 intensity profiles across D_2O and H_2O fibrils, and measured the distances between the same two peaks (red letter d) for 5 43 subsequent repeats. Taking the average, we find D-band periodicities of 67.3 ± 1.3 nm and 67.0 ± 0.8 nm for fibrils assembled in D₂O and H₂O respectively, showing that heavy water has not effect on the quarter staggered alignment of the collagen molecules within the fibrils.

⁴⁶ **D. Fit analysis for IR spectroscopy.** The linear absorption spectra (Fig. 3B of the main text) were fitted using 3 Gaussian-shaped ⁴⁷ peaks absorbing at 1635, 1660, and 1680 cm⁻¹. The fits were performed by leaving the width and center frequencies of the ⁴⁸ peaks as free parameters. As shown in Fig.Fig. [S7B](#page-9-0)-C, the fits reproduce the linear infrared data well.

 E. Central line slope (CLS) for 2D-IR spectroscopy. To calculate the CLS, we first fit a series of cuts through the 2D spectrum that are parallel to the probe frequency axis by using two Lorentzian-shaped peaks that describe the bleach (blue color) and 51 the excited state absorption (red color). We then fit a line through the center positions of the bleach obtained from the fit of the cuts. The line slope is then the CLS value.

 F. Melting temperature calculation. The CD spectra for solutions of collagen at ambient temperatures exhibited a marked negative peak at ∼198 nm and an adjacent positive peak at ∼220 nm (Fig. [S8\)](#page-10-0). With increasing temperature collagen denatures both in H₂O and D₂O, as indicated by the decrease of the amplitude of the negative 200 nm peak and the change of the sign of the 220 nm spectral feature (Fig. [S8A](#page-10-0)-B). These transitions were observed in a very narrow temperature range around 39-40°C for solutions of collagen in H₂O and around 41-43[°]C for solutions of collagen in D₂O (Fig. [S8A](#page-10-0)-B). These different ranges $\frac{1}{28}$ indicated a higher melting temperature of collagen in D₂O, as compared to H₂O. To quantify solvent isotope effects on the melting behavior of collagen, we calculated the center of mass wavelength of the negative 220 nm peak:

$$
CM = \frac{\sum_{\lambda=10 \; nm}^{210 \; nm} \theta_{CD} \bullet \lambda}{\sum_{\lambda=180 \; nm}^{210 \; nm} \theta_{CD}} \tag{1}
$$

 ϵ_1 The center of mass wavelength, CM, was obtained as the mean value of the wavelength(λ), weighted by the absolute value of the ellipticity (θ_{CD}) at 180-210 nm. The determined values for the three individual scans for collagen in H₂O and D₂O are ϵ_{ss} displayed in Fig.S3C, where we observed that the transition temperature of collagen in D₂O shifts to higher temperatures, as compared to collagen in H₂O. To quantify the transition temperature, we averaged the data of the 3 repeat measurements per condition in Figure 2 and fitted a sigmoidal function together with a linear variation of the CM as a function of temperature *T* to the data:

$$
CM(T) = a_1 + \frac{a_2 - a_1}{1 + 10^{(a_3 - T)} a_4} + a_5 T
$$
\n[2]

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where a_j ($j = 1 \cdots 5$) are fit parameters. As displayed in Fig. [S8C](#page-10-0), Eq (2) describes the data very well, and the obtained es parameters are listed in table 1. From these fits we found the melting temperature (a_3) to increase from 40.2 °C in H₂O to

 $70 \quad 42.8 \text{ °C in } D_2O.$

 G. Analysis of coarse-grained simulations. In order to quantitatively analyze our simulations, we recorded the positions of ⁷² all molecules every 10000 timesteps for a total of $6 \cdot 10^6$ timesteps, resulting in a total simulation time of $6000\tau_0$ and 600 simulation snapshots (frames). To measure the assembly rate, in each frame we calculated the mass of all assembled clusters normalized by the total mass, where a cluster was defined as an aggregation of $n \geq 10$ monomers. To calculate the fibril diameter probability distribution function, we identified fibrils in the last frame of a simulation as clusters that contained at τ_6 least $n \geq 10$ monomers and were at least $l \geq 2l_{\text{monomer}}$ long. The diameter *D* of such a fibril was calculated as the average π diameter along its longitudinal axis normalized by the smallest measured diameter D_0 among all identified fibrils. To obtain averages and standard deviations for the assembled mass and to increase the number of fibrils available for the diameter

⁷⁹ distribution, we ran 10 independent simulations for each set of parameters $(\epsilon_{H}, \epsilon_{E})$.

⁸⁰ **References**

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- ⁸³ 2. C Semmrich, et al., Glass transition and rheological redundancy in f-actin solutions. *Proc. Natl. Acad. Sci*. **104**, 20199–20203 84 (2007).
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Fig. S1. Turbidity curves for collagen fibril formation in HDO and H2O at concentration of 1 mg/ml at room temperature (21◦C). Stock solutions of collagen dissolved in water or heavy water (0.2 wt % of acetic acid, pH=3.3=3.4) were neutralized using heavy water buffer or water buffer to obtain a 1:1 ratio of H₂O(which we refer to as HDO₂) and D₂O (which we refer to as HDO₁). Shaded areas represent the standard deviation obtained over 3 different measurements for HDO and H₂O,respectively.

Fig. S2. Turbidity curves at room temperature (21◦C) and 37 ◦C of collagen assembled in heavy water at a concentration of ∼0.1 mg/ml at a pH=7.2-7.5. Temperature was kept at 37 degrees Celsius by a temperature controller (Julabo, TopTech F32-ME).

Fig. S3. A) and B) turbidity curves obtained by adding different salt concentration to water and heavy water PBS solutions containing collagen at a concentration of ∼0.1 mg/ml after buffer neutralization (pH=7.2-7.5). C) Values and errors for growth rate (top row), plateau time (central row) and lag time (bottom row) calculated from the turbidity curves reported in A)-B). The errors represent the standard deviation over at least 3 independent measurements.

Fig. S4. A) Representative evolution of elastic shear modulus G' (continuous lines) and viscous shear modulus G" (dashed lines) during collagen self-assembly in D₂O, HDO and H2O. Measurements were conducted with a collagen concentration of 1.25 mg/mL at a strain amplitude of 0.5%, a frequency oscillation of 0.5 Hz and temperature of 23◦C. B) Frequency sweep of fully formed collagen network in different solvents with a collagen concentration of 1.25 mg/mL at a strain amplitude of 0.5% and temperature of 23◦C. $(G' =$ open symbols, $G'' =$ closed symbols $).$

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Fig. S5. A) and B) cryo-TEM images (top) and intensity profiles (bottom) of collagen fibril assembled in H₂O and D₂O at a concentration of ∼1-1.25 mg/ml at room temperature, respectively.

Fig. S6. IR spectra of collagen monomers dissolved in water and heavy water (0.2 wt% acetic acid, pH=3.2-3.4) recorded between 1200 and 3600 cm⁻¹ (top and bottom respectively). Amide II vibration (O=C-N-H) absorbs at around 1550 cm⁻¹, whereas Amide II' (O=C-N-D) absorbs at 1490 cm⁻¹.

Fig. S7. A. IR spectra of collagen fibrils dissolved in water and heavy water(pH=7.4) at a concentration of 5 mg/ml and 1.25 mg/ml, respectively. B.-C. Fit of the IR spectra of collagen monomers dissolved in water and heavy water (0.2 wt% acetic acid, pH=3.2-3.4) at a concentration of 10 mg/ml and 2 mg/ml, respectively. Blue dotted lines represent the fitted Gaussian-shaped peaks, and the red line the fit to the raw data (black line). D. IR spectra of collagen at different level of N-H/N-D exchange, which slowly increase over storage time. Upon N-H/N-D exchange, the amide II frequency shifts from 1550 cm $^{-1}$ to 1490 cm $^{-1}$, whereas the amide I frequencies red-shift of 1-2 cm $^{-1}$. We observe that the ratio between the peak intensities of the two amide I bands is constant, although the amide II decreases in intensity because of the N-H to N-D exchange. By using collagen samples using the same stock solution after different storage time, we found no evidence of N-H/N-D effect on the macro- and microscopic changes observed when dissolving collagen in heavy water. E-F. IR spectra of collagen monomers dissolved in heavy water and water (0.2 wt% acetic acid, pH=3.2-3.4) at different concentration. The light blue line and shade represent the mean and the standard deviation obtained from 3 different measurements. All measurements reported here were done at 23 °C.

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Fig. S8. A.-B. CD spectra of solutions of collagen monomers dissolved in water (A) and heavy water (B) (0.2 wt% acetic acid, pH=3.2-3.4) at a concentration of 0.1 mg/ml at different temperatures. C. Averaged center of mass wavelength as a function of temperature. Symbols show experimental data and error bars show the standard deviation within three experiments. Solid shaded lines show fits of Eq.(2) to the data, while solid opaque lines are guides-to-eye.

Fig. S9. Radial distribution functions (RDFs) for water around proline, hydroxyproline and glycine. All atoms are considered for the analyses. The RDFs are averaged over the individual value obtained from the five independent simulations.

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Fig. S10. Radial distribution functions (RDFs) for water around the carbonyl groups of proline, hydroxyproline and glycine. Only the O atoms of the carbonyl groups and the water molecules are considered for the analyses to compare against the 2D-IR experiments in Fig. 3B. Hence, the evident shift in hydration shells on the x-axis as compared to Fig. S9 arises from the consideration of only a subset of atoms, *i.e* oxygen atoms. The RDFs are averaged over the individual value obtained from the five independent simulations.