

1 **Supplementary Information for: Extensive hydrogen incorporation is not**  
2 **necessary for superconductivity in topotactically reduced nickelates**

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28 **Supplementary Note 1: Sample Summary**

29 Supplementary Table 1 lists details of all samples referenced in this paper, along with their fabrication  
30 details and which figures they correspond to.

31 **Supplementary Note 2: Hydrogen Quantification using SIMS**

32 Although ToF-SIMS is very sensitive to mass and depth, extracting quantitative information (e.g., chem-  
33 ical stoichiometries) can be challenging due to so-called matrix effects where the chemical and electronic  
34 environment of the material (matrix) affects the ionization of the sputtered species. First, the choice of  
35 the sputtering beam can change the matrix. For example, the use of a Cs<sup>+</sup> sputter beam results in Cs  
36 implantation into the film, lowering the work function and dramatically increasing the yield of negative  
37 ions [1]. Because different ions can interact differently with the implanted Cs, the ion yield ratios can vary  
38 drastically. To minimize potential chemical interactions with the sample, we chose to use an Ar<sup>+</sup> noble gas  
39 cluster source for sputtering, which is also volatile and escapes to vacuum after bombardment. The type of  
40 analysis beam (Bi<sub>3</sub><sup>+</sup>, O<sub>2</sub><sup>+</sup>, Ga<sup>+</sup>, etc.) also has an effect on ionization, and this needs to be specified when  
41 comparing quantitative results.

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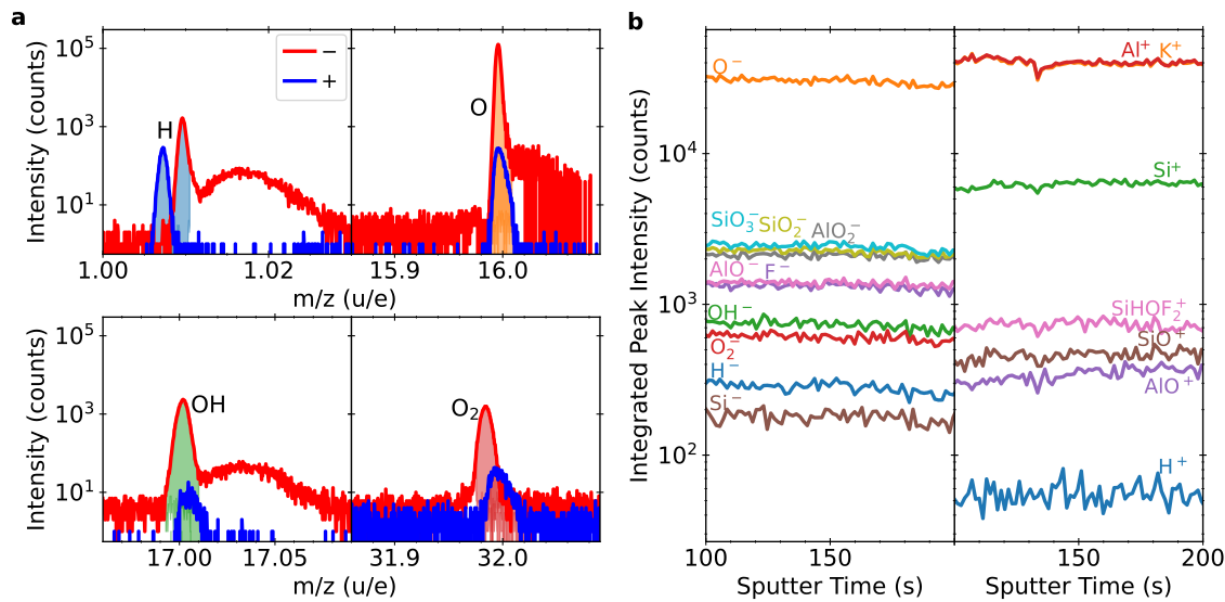
Label	Material	Synthesis	Figures
As-grown (La,Ca)NiO <sub>3</sub> Superconducting (La,Ca)NiO <sub>2</sub>	5.4 nm La <sub>0.78</sub> Ca <sub>0.22</sub> NiO <sub>3</sub> / SrTiO <sub>3</sub> 29 nm SrTiO <sub>3</sub> / 4 nm La <sub>0.78</sub> Ca <sub>0.22</sub> NiO <sub>2</sub> / SrTiO <sub>3</sub>	1	2b,c,d, Supplementary Figure 6 2a,b,c,d,e,f, Supplementary Figure 6
As-grown (La,Sr)NiO <sub>3</sub> Superconducting (La,Sr)NiO <sub>2</sub>	5.3 nm La <sub>0.8</sub> Sr <sub>0.2</sub> NiO <sub>3</sub> / SrTiO <sub>3</sub> 6 nm SrTiO <sub>3</sub> / 4.2 nm La <sub>0.8</sub> Sr <sub>0.2</sub> NiO <sub>2</sub> / SrTiO <sub>3</sub>		2b, Supplementary Figure 5 2a,b,g,h, Supplementary Figure 3, Supplementary Figure 5
Superconducting Nd <sub>6</sub> Ni <sub>5</sub> O <sub>12</sub>	100 nm Pt / 10 nm Ti / 23 nm Nd <sub>6</sub> Ni <sub>5</sub> O <sub>12</sub> / NdGaO <sub>3</sub>	2	1c, 3a,b,c, Supplementary Figure 7
As-grown NdNiO <sub>3</sub>  Reduced NdNiO <sub>2</sub>	17.3 nm NdNiO <sub>3</sub> / SrTiO <sub>3</sub>  17.3 nm NdNiO <sub>2</sub> / SrTiO <sub>3</sub>	2	4a,c,d, Supplementary Figure 8, Supplementary Figure 10 4a,e,f, Supplementary Figure 2, Supplementary Figure 8, Supplementary Figure 10
Reduced NdNiO <sub>2</sub>	16 nm NdNiO <sub>2</sub> / SrTiO <sub>3</sub>		4b
As-grown (Nd,Sr)NiO <sub>3</sub>	10 nm Nd <sub>0.8</sub> Sr <sub>0.2</sub> NiO <sub>3</sub> / SrTiO <sub>3</sub>	3	5a,c,d, Supplementary Figure 11a, Supplementary Figure 9, Supplementary Figure 12
As-grown capped (Nd,Sr)NiO <sub>3</sub>	10 nm SrTiO <sub>3</sub> / 10 nm Nd <sub>0.8</sub> Sr <sub>0.2</sub> NiO <sub>3</sub> / SrTiO <sub>3</sub>		5a, Supplementary Figure 11b, Supplementary Figure 12
Reduced (Nd,Sr)NiO <sub>2</sub>	10 nm Nd <sub>0.8</sub> Sr <sub>0.2</sub> NiO <sub>2</sub> / SrTiO <sub>3</sub>		5a, Supplementary Figure 12
Reduced capped (Nd,Sr)NiO <sub>2</sub>	10 nm SrTiO <sub>3</sub> / 10 nm Nd <sub>0.8</sub> Sr <sub>0.2</sub> NiO <sub>2</sub> / SrTiO <sub>3</sub>		5b,e,f, Supplementary Figure 9, Supplementary Figure 12, Supplementary Figure 13, Sup- plementary Figure 14
Mica	(KF) <sub>2</sub> (Al <sub>2</sub> O <sub>3</sub> ) <sub>3</sub> (SiO <sub>2</sub> ) <sub>6</sub> (H <sub>2</sub> O)	MTI Corp.	Supplementary Figure 1

Supplementary Table 1. **Table of samples and where they are referenced in this paper.** Grouped samples are pieces split from the same deposited film.

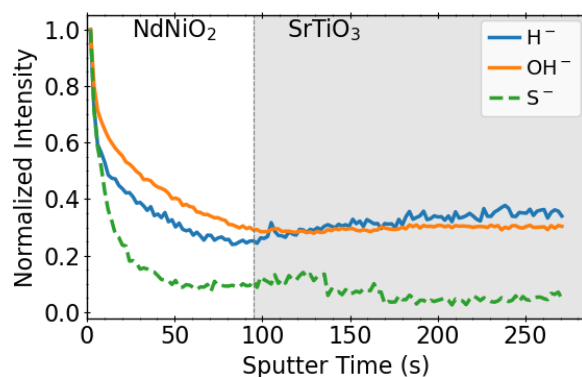
42 Second, quantification requires matrix-matched standards. Since even small impurities can affect ioniza-  
43 tion, standards need to be prepared with known ratios of chemical species and exist in a similar chemical  
44 environment to that of the films of interest [2]. Relative sensitivity factors can then be calculated to scale  
45 the integrated peak areas to atomic concentrations. We first attempt to quantify hydrogen content in our  
46 nickelate films by comparing to a mica standard, as used in Ding *et al.* [3]. The chemical formula for the mica  
47 (muscovite) crystals we used is (KF)<sub>2</sub>(Al<sub>2</sub>O<sub>3</sub>)<sub>3</sub>(SiO<sub>2</sub>)<sub>6</sub>(H<sub>2</sub>O). Peak shapes as well as intensity measurements  
48 are shown in Supplementary Figure 1.

49 To illustrate the importance of matrix-matched standards, we compare the levels of oxygen (O<sup>-</sup> and O<sub>2</sub><sup>-</sup>)  
50 measured in the mica and in the SrTiO<sub>3</sub> and NdGaO<sub>3</sub> substrates prior to any film reduction. Although  
51 they have comparable oxygen number densities, we found that dissimilarities in the ionization efficiencies  
52 produced drastically different results (i.e., estimated stoichiometric ratios). A quick calculation revealed  
53 that the apparent oxygen concentration was significantly lower in the mica than in the SrTiO<sub>3</sub> and NdGaO<sub>3</sub>  
54 substrates, and this is due to the presence of different elements in the matrix. Similar matrix-related effects  
55 are usually observed for all elements of the periodic table, including H, which highlights the difficulty of  
56 quantification using SIMS. In addition, hydrogen is present in adventitious carbon contamination at surfaces  
57 and interfaces, and requires careful interpretation of the profile data.

58 Another complication is the effect of the reduction process on the chemical and physical properties of the  
59 film, which directly affects the sensitivity factor and the sputtering rate, respectively. A lower sputter rate  
60 means that fewer number of elements and molecules are being ejected at any given time, and even though  
61 the ionization efficiency may not have changed, the intensity appears to be lower. In our study, this apparent  
62 intensity of molecular species within a single nickelate sample varied dramatically as a function of reduction  
63 time and conditions, as discussed in further detail in Supplementary Note 5. This renders even samples  
64 cleaved from a single growth and reduced differently to not be quantitatively comparable.



Supplementary Figure 1. **ToF SIMS measurements of mica** a. SIMS showing the spectral peak shapes of  $\text{H}^\pm$ ,  $\text{O}^\pm$ ,  $\text{OH}^\pm$ , and  $\text{O}_2^\pm$  ions. The broad peaks centered around  $m/z = 1.02$  and  $m/z = 17.03$  are attributed to charging effects, and were not included for the construction of the depth profile. b. SIMS depth profiles showing the integrated peak intensities as a function of sputtered depth for both elemental and molecular ions from a mica substrate.

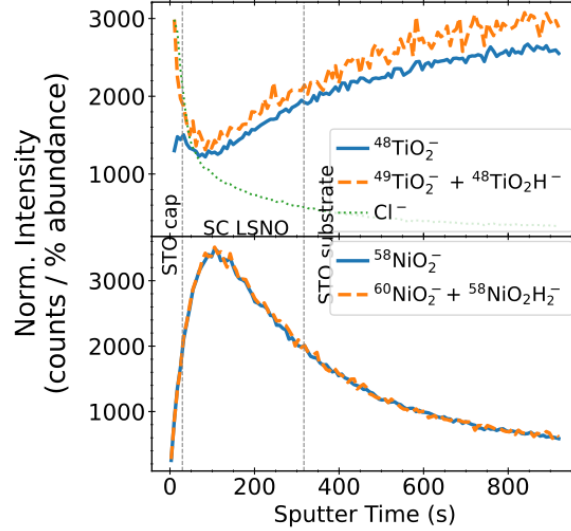


Supplementary Figure 2. **Comparison of ToF-SIMS hydrogen signals in reduced non-superconducting  $\text{NdNiO}_2$  with surface adsorbate signals.** Intensities have been normalized to the surface value. The slower decay of hydrogen intensities indicate that some of the hydrogen signal is intrinsic to the sample.

65

### Supplementary Note 3: Disentangling Extrinsic and Intrinsic Hydrogen

66 As mentioned earlier, SIMS can also detect surface adsorbates, and requires a way to distinguish hydrogen  
 67 intrinsic to the sample from surface contamination. Surface adsorbates are often hydrocarbons, so we can  
 68 compare carbon peaks such as  $\text{C}_2^-$ , as well as other contaminant peaks, and see if hydrogen trends with  
 69 these compounds. The rapid decay in signal at the surface is usually due to these hydrocarbons, and  
 70 their contribution to the H and OH signals can be subtracted from those inside the film. We can further  
 71 confidently attribute certain signals to measured hydrogen by isotope analysis. Metal oxides such as Ti will  
 72 have the same ionization efficiency regardless of the isotope (e.g.  $^{48}\text{TiO}_2$  vs.  $^{49}\text{TiO}_2$ ). Species with isobaric  
 73 interference, or those that overlap in mass, such as  $^{49}\text{TiO}_2$  and  $^{48}\text{TiO}_2\text{H}$ , can be distinguished by analyzing  
 74

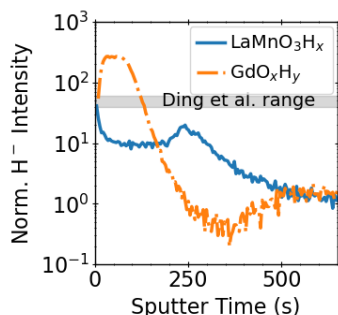


Supplementary Figure 3. **Comparison of the isotopic distributions of Ti and Ni oxides in superconducting  $\text{La}_{0.8}\text{Sr}_{0.2}\text{NiO}_2$ .** The less abundant isotopic oxides have an overlapping mass with the more abundant isotopic hydroxides. By scaling the depth profiles by the natural isotopic abundance, we can separate out these contributions. In this case, there is a contribution from the Ti hydroxide, but no evidence of Ni hydroxide. The sharp decay in peak intensity for  $^{48}\text{TiO}_2\text{H}^-$  at the surface indicates that the Ti oxide of the  $\text{SrTiO}_3$  cap has adsorbed some hydrogen, but that has not affected the nickelate film.

75 the natural isotopic distribution in the mass spectrum or profile shape, as shown in Supplementary Figure  
 76 3. In this figure, the less abundant  $^{49}\text{TiO}_2^-$  has an overlapping mass with  $^{48}\text{TiO}_2\text{H}^-$  at  $m/z = 64.945$   
 77 that cannot be resolved. However, by constructing the depth profile and seeing their profile shapes, it is  
 78 possible to identify which molecule the peak corresponds to. In this case, the main contribution to the peak  
 79 is from the Ti hydroxide. Similarly for Ni oxides, comparison of the profiles show that  $^{60}\text{NiO}_2^-$  is the likely  
 80 assignment for this peak instead of the hydroxide. The sharp decay in peak intensity for  $^{48}\text{TiO}_2\text{H}^-$  at the  
 81 surface indicates that the Ti oxide of the  $\text{SrTiO}_3$  cap has adsorbed some hydrogen, but that has not affected  
 82 the nickelate film.

#### 83 **Supplementary Note 4: Previous detection of hydrogen using SIMS**

84 In order to place our observations into context, it is helpful to consider previous studies which used the ToF-  
 85 SIMS instrument employed in this study to examine oxide films with large degrees of hydrogen incorporation,  
 86 such as humidity-treated  $\text{GdO}_x\text{H}_y$ ,  $\text{LaMnO}_3$  membranes exposed to water, and  $\text{CaH}_2$ -reduced  $\text{BaZrO}_3$  [4-6].  
 87 In Supplementary Figure 4, we show reproduced  $\text{H}^-$  intensity curves from a  $\text{LaMnO}_3$  freestanding membrane  
 88 with incorporated water due to the dissolution of  $\text{Sr}_3\text{Al}_2\text{O}_6$  during processing [5]. Also shown is a film of  
 89  $\text{GdO}_x\text{H}_y$  which is very near the  $\text{Gd}(\text{OH})_3$  composition [4]. In these systems, we found enhancements of  
 90 the  $\text{H}^+$ ,  $\text{H}^-$ , and  $\text{OH}^-$  intensities which were 1 to 3 orders of magnitude above the substrate background,  
 91 and were able to establish the high hydrogen concentrations based on the effects on the material properties.  
 92 This prior experience provides reasonable upper and lower boundaries for expected effects and demonstrates  
 93 our sensitivity to intrinsic hydrogen. Interestingly, the results from Ding *et al.* fell between these limits.  
 94 We conclude therefore that hydrogen would readily have been detected in our measurements of layered  
 95 square-planar nickelate samples if it were present at previously reported levels.



Supplementary Figure 4. **Incorporated H-sensitivity baseline from previous ToF-SIMS measurements.** Examples of previously measured and published SIMS depth profiles for the  $\text{H}^-$  ion in the hydrogen containing oxides  $\text{LaMnO}_3\text{H}_x$  and  $\text{GdO}_x\text{H}_y$  [4, 5], normalized to 1 in the substrate.

96

### Supplementary Note 5: Extended SIMS Characterization

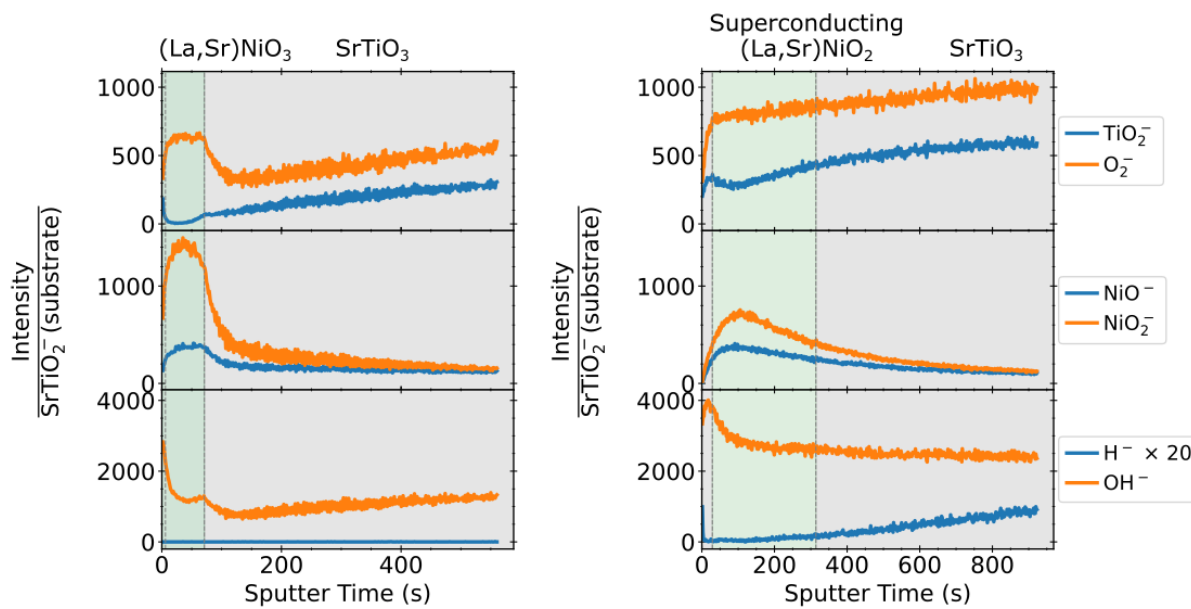
97 For each sample, we have recalibrated the SIMS depth profiles by scaling all intensities to an appropriate  
 98 steady-state ion in the substrate to compare between sample types. This also ensures that small fluctuations  
 99 in primary ion current between replicate runs can be accounted for. Hydrogen compounds, having much  
 100 lower intensity, have been scaled as noted for visualization clarity. These are shown in Supplementary Figure  
 101 5, Supplementary Figure 6, Supplementary Figure 7, Supplementary Figure 8, and Supplementary Figure 9.

102 The interface position of the layers is defined as the sputter time when the characteristic signal reaches  
 103 50% of its maximum. The characteristic peaks used were  $\text{TiO}_2^-$  in the  $\text{SrTiO}_3$ ,  $\text{NiO}^-$  in the nickelate film,  
 104 and  $\text{GaO}_2^-$  in the  $\text{NdGaO}_3$ . The interface width is defined as the time or depth required for the intensity to  
 105 change from 84% to 16% of the maximum (or vice versa), and its sharpness is dependent on several factors  
 106 such as surface roughness, difference in density between the two layers, and redeposition of the sputtered  
 107 neutrals, in decreasing order of influence. Normally, the sharpest interface is obtained for lower-mass species,  
 108 and when transitioning from a lower to a higher density film. On the other hand, we observed long  $\text{NiO}_2^-$  tails  
 109 in the  $\text{SrTiO}_3$  substrates, which may likely be attributed to cascade mixing effects induced by the analysis  
 110 beam knocking atoms deeper into the film[7]. These effects may also account from some of the  $\text{TiO}_2^-$  signal  
 111 observed in films capped by  $\text{SrTiO}_3$ , although the nickelate films are denser than the  $\text{SrTiO}_3$  caps and these  
 112 films are thin enough that the  $\text{TiO}_2^-$  is not expected to decay to zero.

113 We observed that the sputter rate is faster for lower crystallinity layers, including amorphous  $\text{SrTiO}_3$  caps,  
 114 as well as reduced films which are lower in crystallinity. In fact, the aggressively reduced  $\text{Nd}_{0.8}\text{Sr}_{0.2}\text{NiO}_2$   
 115 without protection of the  $\text{SrTiO}_3$  cap, which exhibited significant crystalline mosaicity with multiple crys-  
 116 talline and amorphous phases, was sputtered through so quickly that no depth profile could be obtained. In  
 117 general, the  $\text{SrTiO}_3$  substrate is expected to sputter more quickly than the higher-density films, although  
 118 film quality also plays a role.

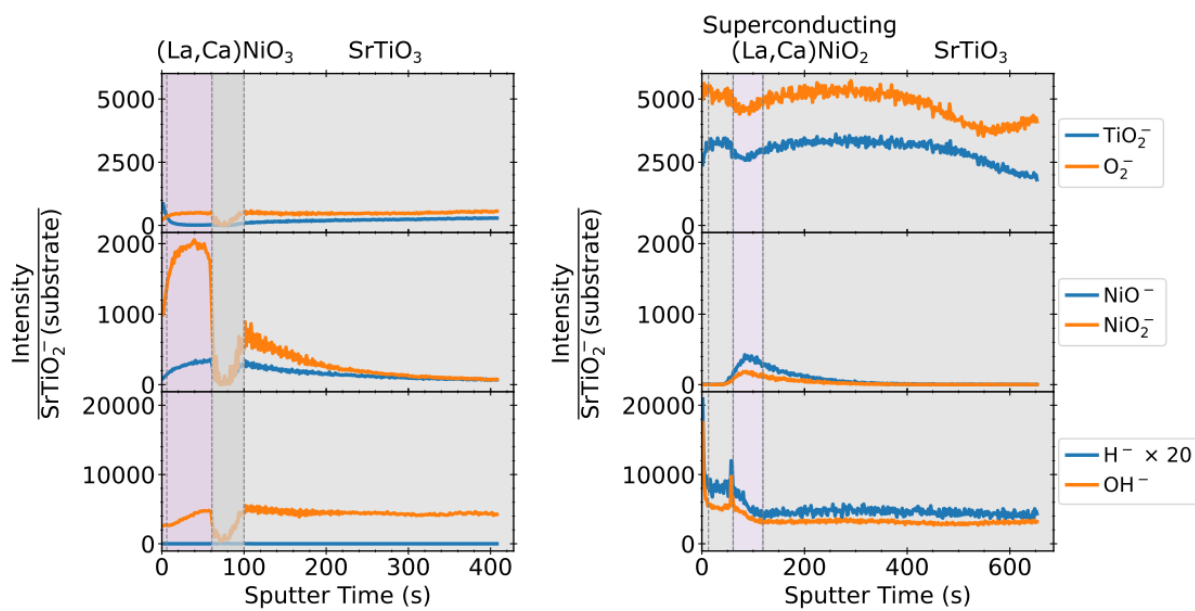
119 When discussing SIMS “resolution”, the most important question is not the interface width but rather our  
 120 ability to resolve the individual layers and identify their contributions to the signal. This is clearly achieved  
 121 even in the thinnest  $\approx 5$  nm thick  $\text{La}_{0.78}\text{Ca}_{0.22}\text{NiO}_2$  and  $\text{La}_{0.8}\text{Sr}_{0.2}\text{NiO}_2$  superconducting layers, with clear  
 122 onset and decay of the  $\text{NiO}_2^-$  signals which yield estimates of the  $\pm 1\sigma$  interface widths of approximately 5  
 123 nm - 6 nm. We may therefore be confident that the vast majority of the signal in the regions designed “film”  
 124 originates in the nickelate film, allowing the hydrogen content to be probed.

125 Thus, SIMS was able to reveal important differences in chemical composition between the as-grown and  
 126 superconducting samples. As can be seen most clearly in Supplementary Figure 5,  $\text{CaH}_2$  reduction preferen-  
 127 tially removes oxygen from the nickelate layers. In the as-grown  $\text{La}_{0.8}\text{Sr}_{0.2}\text{NiO}_3$  film, the  $\text{O}_2^-$  signal decreases  
 128 dramatically at the nickelate/substrate interface, but this decrease is not present after reduction, indicating  
 129 a lower oxygen content in the nickelate film. Comparison of the as-grown and reduced films (Supplementary  
 130 Figure 5, Supplementary Figure 6, Supplementary Figure 8, and Supplementary Figure 9) reveals a dra-  
 131 matic decrease in the measured  $\text{NiO}_2^-$  intensity and  $\text{NiO}_2^-:\text{NiO}^-$  ratio. As the measured intensities of ions  
 132 is dependent on their oxidation state [8, 9], this change is likely due to the change in the valence state of Ni  
 133 upon reduction.

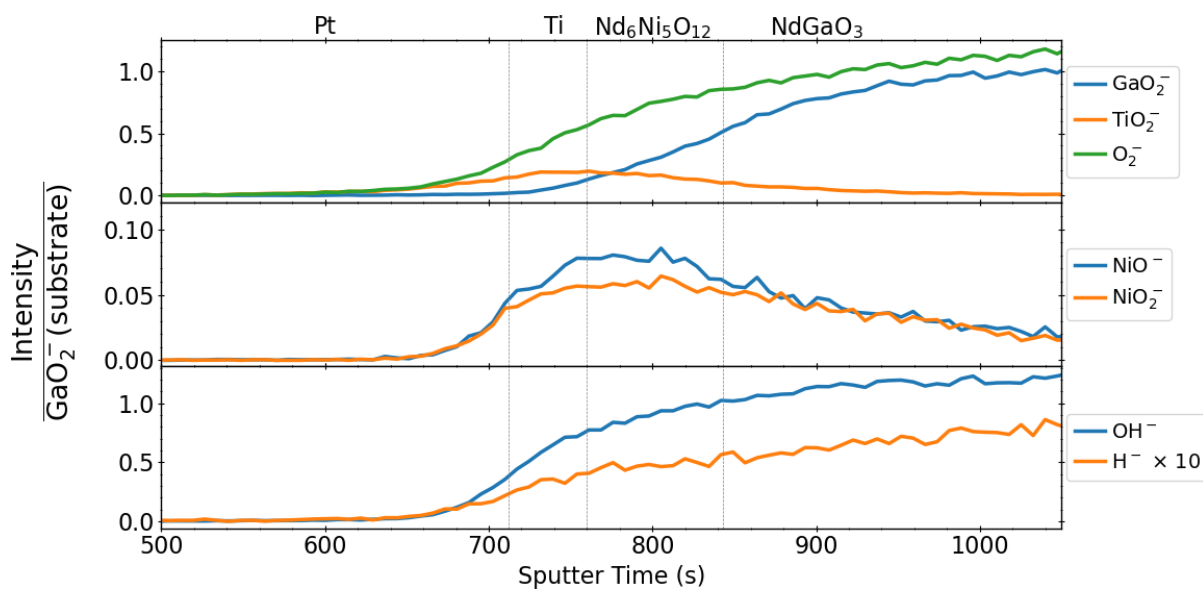


Supplementary Figure 5. **ToF-SIMS depth profiles for the as-grown and superconducting  $\text{La}_{0.8}\text{Sr}_{0.2}\text{NiO}_{3\rightarrow 2}$  films**, including those shown in Figure 2. Intensities are normalized to the steady-state level of  $\text{SrTiO}_2^-$  in the substrate.

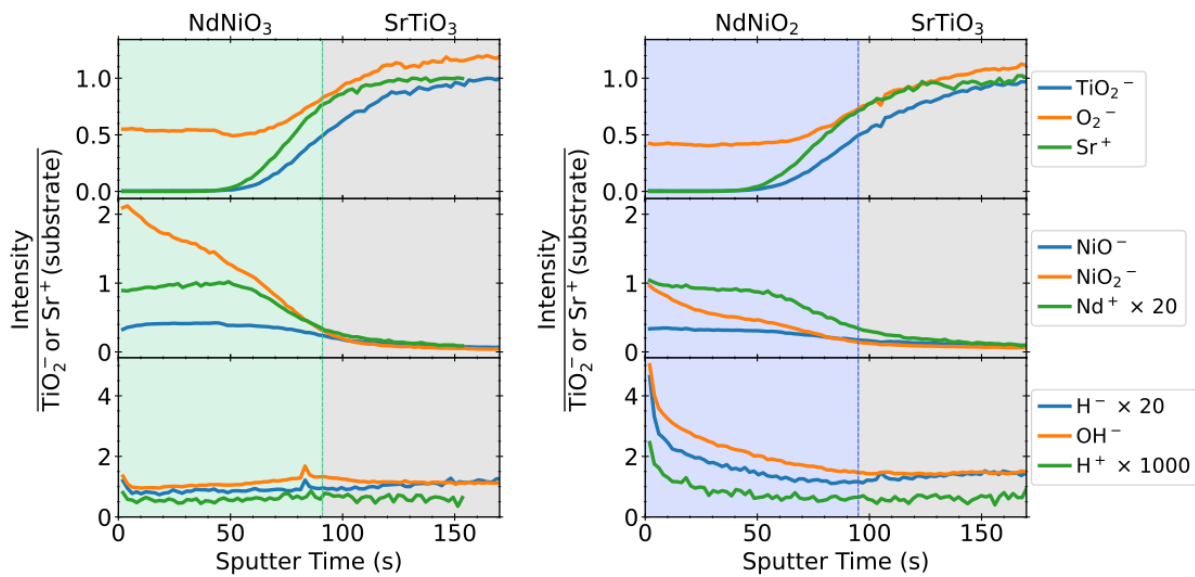
<sup>134</sup> Finally, though we see no evidence for extensive hydrogen incorporation in any of these samples, the  
<sup>135</sup> hydrogen signals which are measured are primarily present as  $\text{OH}^-$ , and secondarily  $\text{H}^-$ . The  $\text{H}^+$  peak is  
<sup>136</sup> orders of magnitude smaller, and in many cases undetectable.



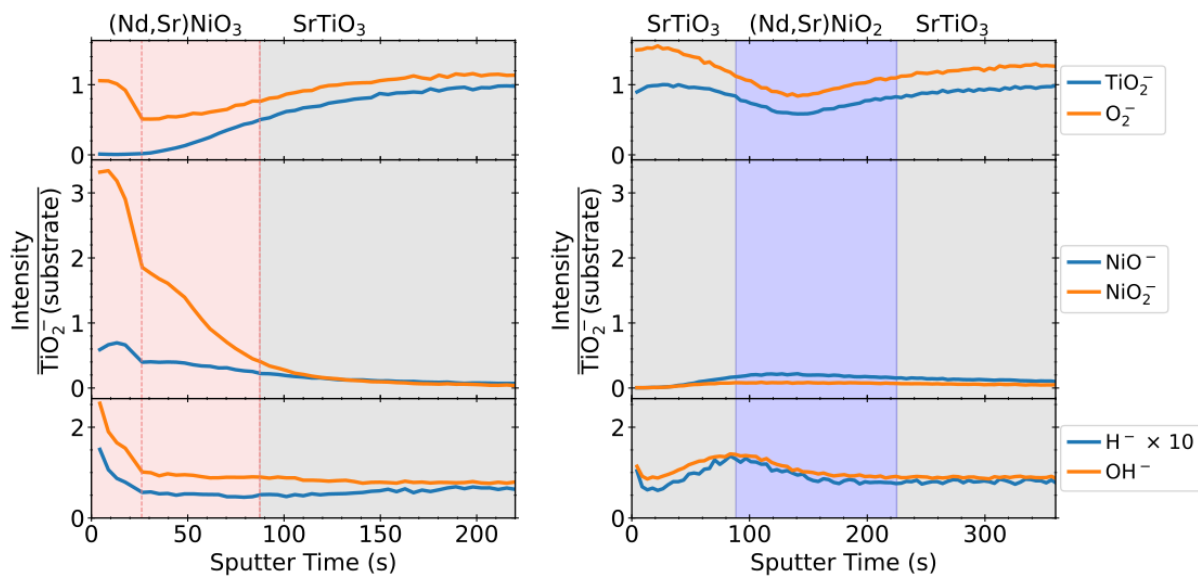
Supplementary Figure 6. **ToF-SIMS depth profiles for the as-grown and superconducting  $\text{La}_{0.78}\text{Ca}_{0.22}\text{NiO}_{3\rightarrow 2}$  films**, including those shown in Figure 2. The greyed out region at the surface of the  $\text{SrTiO}_3$  substrate of the as-grown film is a measurement artifact due to charging, which decreased the total measured signal.



Supplementary Figure 7. **ToF-SIMS depth profile for the  $\text{Nd}_6\text{Ni}_5\text{O}_{12}$  film** shown in Figure 3. Intensities are normalized to the steady-state  $\text{GaO}_2^-$  in the  $\text{NdGaO}_3$  substrate.



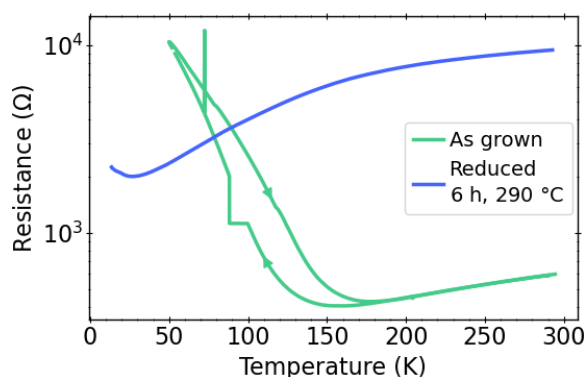
Supplementary Figure 8. ToF-SIMS depth profiles for the  $\text{NdNiO}_3$  and non-superconducting  $\text{NdNiO}_2$  films shown in Figure 4. Intensities are normalized to steady-state  $\text{TiO}_2^-$  or  $\text{Sr}^+$  in the  $\text{SrTiO}_3$  substrate.



Supplementary Figure 9. ToF-SIMS depth profiles for the  $\text{Nd}_{0.8}\text{Sr}_{0.2}\text{NiO}_3$  and  $\text{SrTiO}_3/\text{Nd}_{0.8}\text{Sr}_{0.2}\text{NiO}_2$  shown in Figure 5. Intensities are normalized to steady-state  $\text{TiO}_2^-$  in the substrate. The surface of the uncapped, as-grown film is likely different due to aging.



### Supplementary Note 6: Electrical characterization of non-superconducting films



Supplementary Figure 10. **Transport measurements of  $\text{NdNiO}_{3 \rightarrow 2}$ .** Two-point resistance measurements of as-grown and reduced  $\text{NdNiO}_{3 \rightarrow 2}$  shown in Figure 4.

Supplementary Figure 10 shows resistance vs. temperature for the non-superconducting, undoped  $\text{NdNiO}_3/\text{SrTiO}_3$  shown in Figure 4. These electrical transport measurements were taken using a home-built electrical dipstick probe compatible with a helium dewar. Indium contacts were soldered on the four corners of each sample in a van der Pauw (VDP) configuration. AC transport measurements were taken at 17.777 Hz using an SR-830 lock-in amplifier. The voltage and current were measured simultaneously to determine the resistance. Note that values shown are the measured resistance of a single VDP channel of these  $5 \text{ mm} \times 10 \text{ mm}$  samples, not sheet resistance or resistivity.

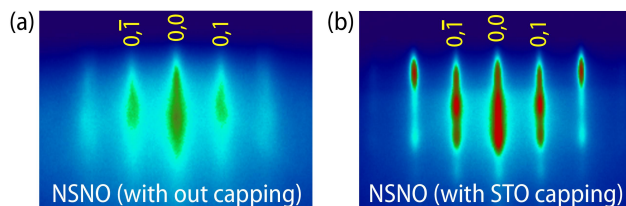
The original as-grown film shows the hysteretic metal-insulator transition expected for a stoichiometric  $\text{NdNiO}_3$  film, while the reduced infinite-layer film exhibits the metallic behavior expected for undoped  $\text{NdNiO}_2$ .

$\text{Nd}_{0.8}\text{Sr}_{0.2}\text{NiO}_2$  films were insulating enough that two-point resistances were measured only at room temperature. As-grown, without significant growth optimization, uncapped films started with resistances between  $500 \text{ k}\Omega$  to  $700 \text{ k}\Omega$ , while capped films started with resistances between  $6 \text{ k}\Omega$  to  $8 \text{ k}\Omega$ . After reduction, the resistance of uncapped films actually increased and was not measurable using a standard multimeter, while capped films had a lower resistance.

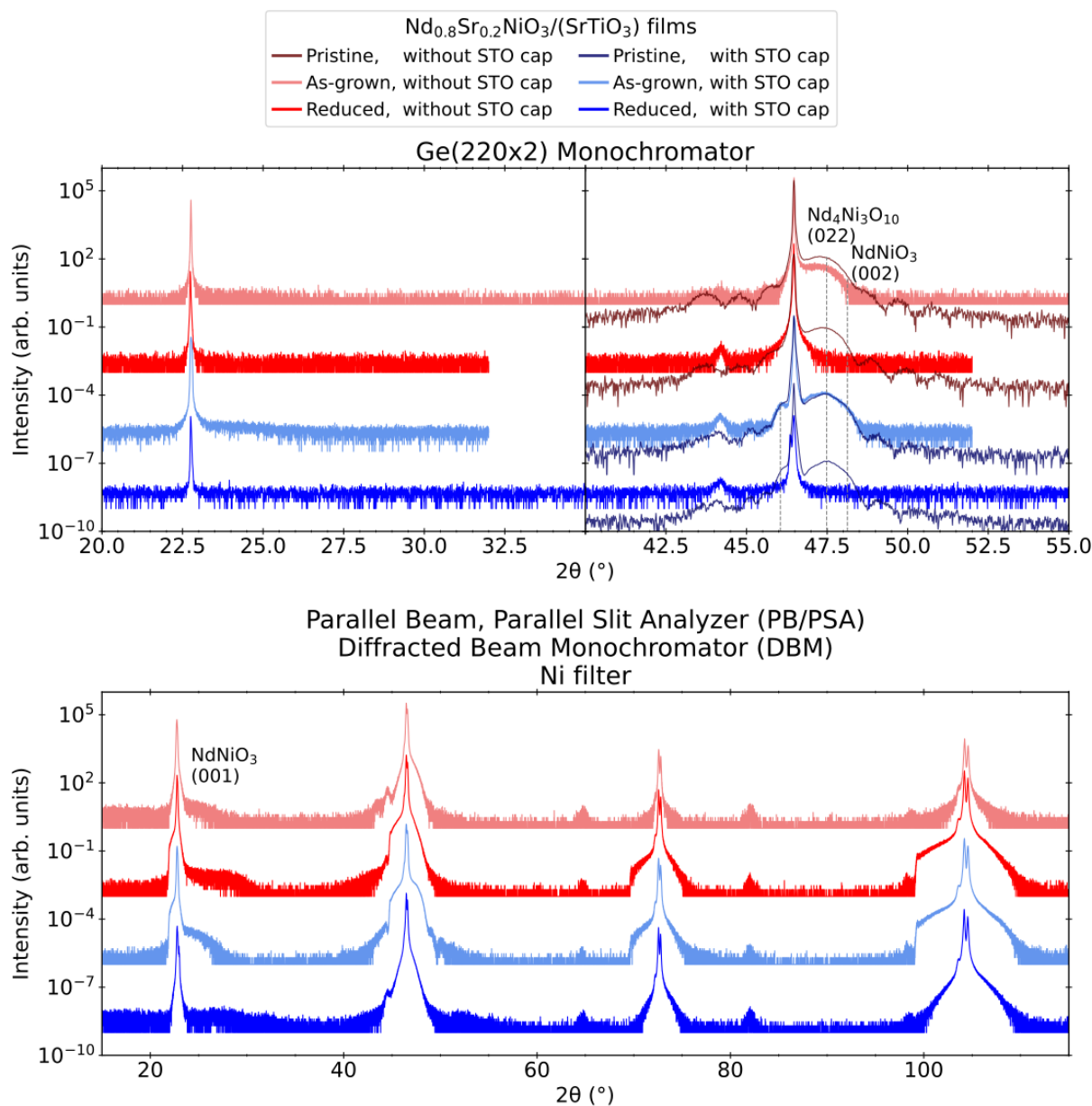
### Supplementary Note 7: Further structural characterization

Supplementary Figure 11 shows the RHEED pattern of specular (0 0) and off-specular (0 1), (0 -1) spots from PLD-deposited  $\text{Nd}_{0.8}\text{Sr}_{0.2}\text{NiO}_3$  films, recorded after cooling the samples to room temperature. The clear streaky pattern confirms the excellent surface morphology of the films immediately after the deposition.

Supplementary Figure 12 shows full XRD characterization of the capped and uncapped  $\text{Nd}_{0.8}\text{Sr}_{0.2}\text{NiO}_{3 \rightarrow 2}$  films. Complementary XRD measurements were taken using two different XRD configurations. The top scans

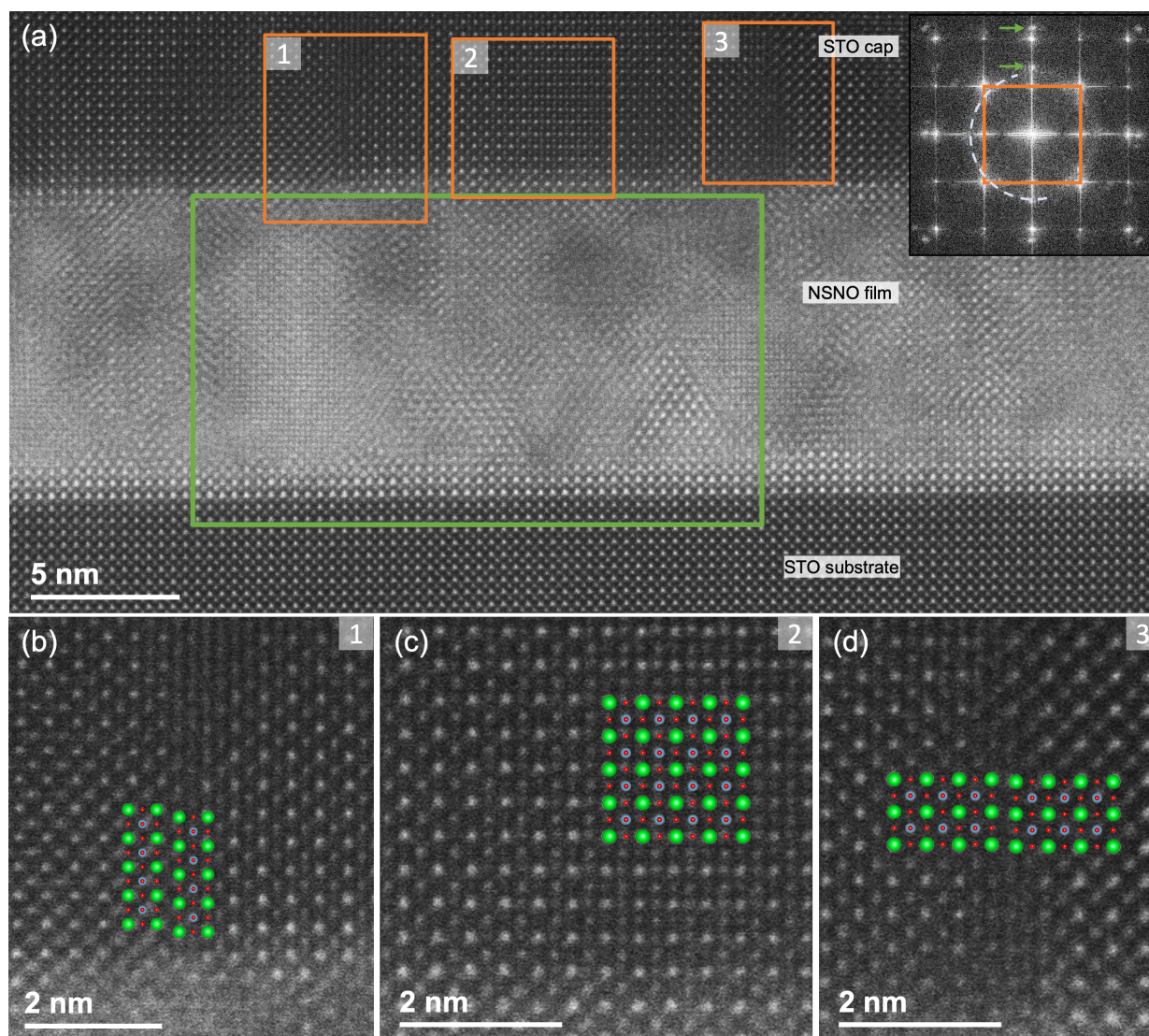


Supplementary Figure 11. **RHEED image of the  $\text{Nd}_{0.8}\text{Sr}_{0.2}\text{NiO}_3$  (NSNO) films** (a) with out and (b) with  $\text{SrTiO}_3$  (STO) capping layer after cooling the samples to room temperature.



Supplementary Figure 12. **Additional X-ray characterization.** Full XRD scans of samples in Figure 5a taken using two instrument configurations. On top are measurements taken using a lower-intensity, high-resolution mode, while on bottom are the same scans taken in a high-intensity, low-resolution mode to survey for impurity phases.

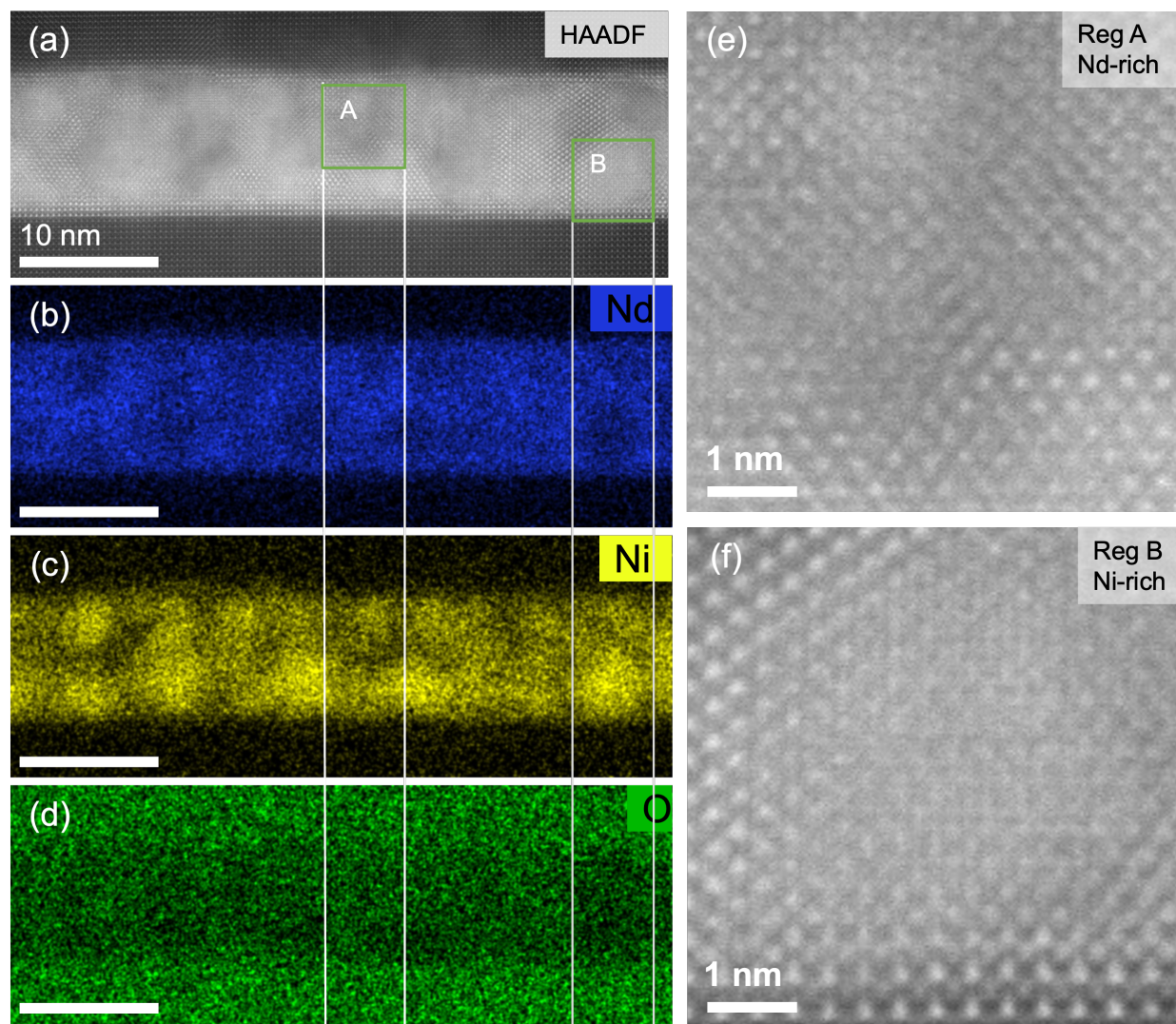
159 were measured in a high-resolution mode, using a Ge (220)×2 monochromator to select Cu  $K\alpha_1$  radiation  
 160 on the incident side. The Diffracted Beam Monochromator assembly was installed on the receiving side  
 161 containing a Ni filter, Ge (400)×2 analyzer, and 0.5° PSA. The dark lines represent measurements performed  
 162 immediately after film deposition (pristine state), while measurements performed on the same films prior to  
 163 ToF-SIMS are colored as labeled, either as-grown or reduced. The high-resolution measurements allow us to  
 164 measure the as-grown (002) film peak at 47.485°, and confirm the much lower intensity of the corresponding  
 165 (001) peak. We note that some slight degradation of the uncapped film occurs over time, consistent with the  
 166 slightly different surface region seen in ToF-SIMS (Supplementary Figure 9). However, in this measurement  
 167 mode, the peaks of both reduced films are so low in intensity that they cannot be seen at all. The bottom  
 168 scans were measured in a high-intensity mode, using a parallel beam (PB) selection slit, and a parallel slit



Supplementary Figure 13. **Extended electron microscopy characterization.** (a) The larger field of view of the image shown in Figure 5b, showing the presence of defects in the STO capping layer as well. The inset shows the FFT pattern of the image. The spots corresponding to the cubic STO substrate (and cap) is denoted by the square, the extra spots and the diffused ring indicate the presence of additional crystalline and amorphous regions in the films, as respectively shown by the arrows and dashed semi-circle. (b-d) Magnified regions of the SrTiO<sub>3</sub> capping layer showing the presence of various forms of defects, with a schematic structural model overlaid (Sr = green, Ti = blue, O = red spheres). (b) Region 1 shows a Ruddlesden-Popper-like phase at the nickelate-cap interface. (c) Region 2 shows strong signal from the oxygen columns, which cannot be only due to O as low Z elements cannot to scatter to the high collection angle used in this experiment. (d) Region 3 shows the planes in the SrTiO<sub>3</sub> cap are highly distorted.

169 analyzer (PSA). Cu K $\alpha$  lines were selected using a Ni filter. These measurements allow observation of a  
 170 (001)-type film peak which shifts to larger angle with reduction.

171 The crystalline quality and nature of defects in the reduced SrTiO<sub>3</sub>-capped Nd<sub>0.8</sub>Sr<sub>0.2</sub>NiO<sub>2</sub> film is studied  
 172 by cross-sectional HAADF-STEM studies and is observed to be highly defective as shown in Figure 5b.  
 173 Mosaicity is observed in the (Nd,Sr)NiO<sub>2</sub> layer, appearing as a mixture of amorphous-like and crystalline  
 174 domains, marked by the circles in Figure 5b. Particularly, two different atomic arrangements deviating from  
 175 the projected [100] perovskite structure is observed in the crystalline regions marked as 2 and 3 in Figure  
 176 5b. Similar defects are present uniformly throughout the observed thin region ( $\sim 5 \mu\text{m}$ ) of the FIB foil.



Supplementary Figure 14. **Extended electron microscopy characterization.** (a-d) HAADF-STEM image and the corresponding EDS elemental maps obtained from the reduced SrTiO<sub>3</sub>-capped (Nd,Sr)NiO<sub>2</sub> film. A clear segregation of Ni and Nd is observed in the nickelate film. (e, f) The amorphous-like and crystalline domains correspond to the Nd-rich and Ni-rich region, respectively.

177 The Fast Fourier Transform (FFT) of the atomic resolution HAADF-STEM image is shown in the inset of  
 178 Supplementary Figure 13a. The sharp spots with the 4-fold symmetry corresponds to the SrTiO<sub>3</sub> perovskite  
 179 structure along the [100] orientation. The FFT also shows the presence of additional spots corresponding  
 180 to periodicities deviating from the STO structure (marked by the green arrows in inset of Supplementary  
 181 Figure 13a, indicating the occurrence of a phase transformation, most likely during the reduction process.  
 182 The presence of weak diffused rings corresponds to the presence of the amorphous regions (marked by the  
 183 dashed circle in inset of Supplementary Figure 13a. The SrTiO<sub>3</sub> capping layer is also observed to consist of  
 184 defects as shown in Supplementary Figure 13b-d.

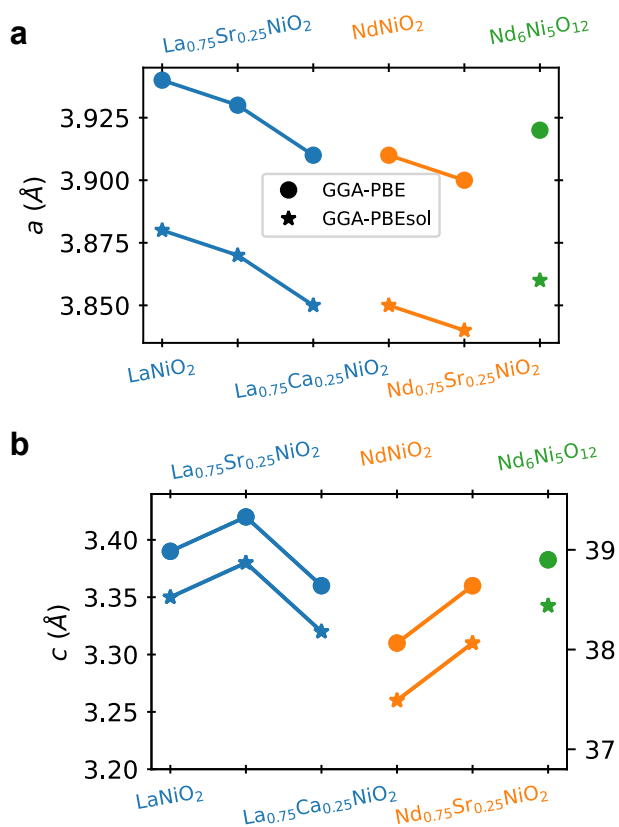
185 Further, the chemical nature of the film was studied using STEM-EDS mapping. The elemental maps  
 186 show chemical segregation in the reduced SrTiO<sub>3</sub>/Nd<sub>0.8</sub>Sr<sub>0.2</sub>NiO<sub>2</sub> film in the form of Ni-rich and Nd-rich  
 187 regions and shown in Supplementary Figure 14a-d. By correlating the chemical maps with the high-resolution  
 188 HAADF image, it is observed that the Nd-rich regions correspond to the amorphous-like regions, while the  
 189 Ni-rich region correspond to the crystalline regions. Finally, Sr was not observed to be concentrated in  
 190 this film. The STEM results thus provide a direct proof for the crystalline and chemical disorder in the

191  $\text{Nd}_{0.8}\text{Sr}_{0.2}\text{NiO}_2$  film after the aggressive reduction process.

192

### Supplementary Note 8: Additional DFT data

193 We provide here the structural data obtained from DFT-based structural relaxations for two choices  
 194 of exchange-correlation functional: GGA-PBE and GGA-PBEsol. The optimized lattice constants for all  
 195 the nickelates considered in the main text within these two functionals are summarized in Supplementary  
 196 Figure 15. We find good agreement with the experimental lattice constants for the undoped infinite-layer  
 197 compounds ( $R = \text{La}, \text{Nd}$ ) is obtained with the GGA-PBE functional, which serves as a benchmark for our  
 198 DFT-optimized lattice constants ( $a_{\text{exp}} = 3.92 \text{ \AA}$ ,  $c_{\text{exp}} = 3.28 \text{ \AA}$  for  $\text{NdNiO}_2$  and  $a_{\text{exp}} = 3.96 \text{ \AA}$ ,  $c_{\text{exp}} = 3.37$   
 199  $\text{\AA}$  for  $\text{LaNiO}_2$  [10, 11]). Previous works have used PBEsol to obtain H-binding energies in infinite-layer  
 200 nickelates [12, 13], but we find that the GGA-PBEsol functional gives rise to optimized lattice constants  
 201 that are quite far from the experimental values.



Supplementary Figure 15. **Structural parameters from the DFT-optimized structures.** (a) In-plane lattice constants  $a$ . (b) Out-of-plane lattice constants  $c$  (right axis corresponds to  $\text{Nd}_6\text{Ni}_5\text{O}_{12}$ ).

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