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Supplementary Materials for

Mesozoic intraoceanic subduction shaped the lower mantle beneath the East Pacific Rise

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Supplementary Text

S1. Plate reconstruction of SE Pacific in the Paleozoic

Prior to the Mesozoic, two major tectonic events occurred in the present-day SE Pacific: the subduction-induced closure of the Rheic Ocean and the subduction of proto-Pacific (Panthalassa) beneath western Gondwana (*30*, *31*). Paleomagnetic and geochronological data reveal that the Rheic Ocean closure occurred during the Devonian-Carboniferous period (420–370 Ma) as the Laurussia supercontinent drifted southward and collided with western Gondwana (*77*, *78*) (fig. S9, A and B). Panthalassa was subsequently subducted along an intraoceanic boundary due to the extension of the Laurussia Plate, starting approximately 390 Ma (*30*) ("PP" in fig. S9, C and D). Following the divergence of Laurussia and southern Gondwana, a second subduction zone was formed along the western margin of Gondwana ("WG" in fig. S9, D-G). The closure of the marginal basin between the two subduction zones was coincident with the accretion of South Patagonia onto western Gondwana (*79*, *80*). The western Gondwana trench continued its eastward drift during the Triassic, initiating a retreat at around 200 Ma (fig. S9, G and H). However, the relatively short interval of the subduction of the Rheic Ocean beneath Laurussia ("EL" in fig. S9A) precludes a prolonged slab-mantle interaction that still impacts the mantle transition zone (MTZ). The motion of Laurussia relative to Gondwana is also subject to considerable uncertainty. In addition, the inferred subduction polarity from seismic tomography is opposite to that of the Patagonian Plate or the western Gondwana subduction zone, thus ruling out its association with the proposed subduction.

S2. Thermal history modeling

In this study, we associate the observed anomaly beneath the Nazca Plate with an ancient oceanic slab. To assess the long-term behavior of a subducted slab, we model temperature evolution within a slab over time (t) using the 1-D unsteady heat conduction equation (81)

$$
T(z,t) = \frac{1}{2\sqrt{\pi\kappa t}} \int_{-\infty}^{\infty} T_0(z') e^{-\frac{(z-z')^2}{4\kappa t}} dz',
$$
 (S1)

where $T_0(z)$ represents the 1-D temperature distribution at $t = 0$. We adopt a typical mantle thermal diffusivity (κ) value of 1×10^{-6} m² · s⁻¹. Our model assumes a slab with a finite thickness of 100 km and an internal temperature of 1000 K. The mantle temperatures above and below the slab are set to 1600 K and 1800 K, respectively. We solve Eq. (S1) using numerical integration. As the slab equilibrates with the mantle, the temperature anomaly gradually widens, while the temperature within the slab remains lower than the surrounding mantle by approximately 150 K over a period of 200-250 Myr (fig. S11). The modeling results indicate that the thermal feature (manifested as cold temperatures) of the slab persists even after 200 Myr. Assuming that seismic velocity is exclusively temperature-dependent and using a temperature derivative of S velocity of $-0.7\%/100$ K at the base of the MTZ (82), we compute a velocity perturbation of $+1\%$, which is in general agreement with the tomographic images.

S3. Slab volume and sinking rate

Published global plate reconstruction models do not include an intraoceanic subduction episode since the Mesozoic in our study region, with the exception of the model proposed by van der Meer et al. (*21*) which was based on mantle tomography. Intracoeanic subduction offers perhaps the simplest explanation for our observations, including a thickened MTZ and fast velocity anomalies in tomographic images. To support our interpretation in terms of plate tectonic history, we further investigate the radial distribution of the tomographic anomalies. We use the cross-sectional area

method of Hafkenscheid et al. (*83*) to quantify the volume of the "slab" anomaly (fig. S12). The slab edge is defined based on a dV_s cutoff value of $+0.4\%$ (see Fig. 3C). The estimated slab volume, as quantified by the cross-sectional area of tomographic anomalies, shows a large peak at \sim 2100 km depth (fig. S12B). The area-depth curve is also characterized by a slight increase from 500 km to 800 km depth, while the slab volume remains relatively small in the lower mantle across the depth range of 1000–1800 km. The curve also shows that approximately 14% of the total slab volume is still confined within the MTZ.

Global correlations between subduction time and depth (*35, 36, 84*) would suggest that the SPS slab started subduction in the early Triassic and ended in the Cenozoic, a much later termination than that we propose. We compare the depth-dependent trend of slab volume with that of the slab volume flux (total volume consumed at subduction zones per unit time) according to East et al. (*85*) (fig. S12A), which is derived from the global plate reconstruction model of Müller et al. (*86*). We convert the slab flux from time to depth, assuming a constant slab sinking rate. The comparison shows that the globally averaged sinking rate of 1.2–1.3 cm/yr (*35, 36*) does not fully account for the observed depth trend of the slab volume, in particular the excess in the MTZ and the lowermost mantle, as highlighted by a moderate anticorrelation (with a correlation coefficient of −0.42) between the two curves (fig. S12B). This discrepancy indicates depth-varying slab sinking rates. One plausible scenario suggests that the slab encountered resistance to sinking at the base of the MTZ. A time span of around 250 Myr for the subducted material to have reached and piled up at depths of 2000–2500 km corresponds to a lower mantle sinking rate of 1 cm/yr. The presence of slab remnants within the MTZ since the mid-Cretaceous translates to an upper mantle sinking velocity of approximately 0.5 cm/yr. This decelerated upper mantle sinking is consistent with the scenario of long-term slab stagnation.

Fig. S1. Map of earthquakes, stations, and bounce points. (**A**) Theoretical raypaths of SS and its precursors for a source-receiver distance of 150°. (**B**) Geographic distribution of the earthquakes (purple circles) and seismic stations (green triangles) used in this study and locations of SS bounce points (black dots). Plate boundaries (*73*) are marked as orange curves.

Fig. S2. MTZ topography obtained using the conventional common midpoint stacking method. (**A** to **C**) Topography of the 410 (A), the 660 (B), and MTZ thickness (C). The depth estimates are shown relative to the regional average values at the bottom right corner of each panel. The orange lines in (A) denote the plate boundaries (*73*).

Fig. S3. Maps of standard errors. (**A** to **C**) Standard errors of the 410 depth (A), the 660 depth (B), and the MTZ thickness (C) obtained from the conventional common midpoint stacking. (**D** to **F**) Standard errors of the 410 depth (D), the 660 depth (E), and the MTZ thickness (F) after processing using partial stacking and FMSSA reconstruction.

Fig. S4. MTZ topography of the study region using different *k* **(data rank threshold) values in the FMSSA reconstruction**. (**A** to **C**) Lateral variations in the 410 depth (A), the 660 depth (B), and the MTZ thickness (C) using $k = 85$. (**D** to **F**) Lateral variations in the 410 depth (D), the 660 depth (E), and the MTZ thickness (F) using $k = 120$.

Fig. S5. MTZ topography obtained using different tomographic models for mantle velocity correction. The discontinuity depths and MTZ thickness are shown relative to their regional average values. Maps are shown for models (**A**) S40RTS (*65*), (**B**) TX2007 (*87*), (**C**) SEMUCB-WM1 (*15*), (**D**) S20RTS (*88*), and (**E**) savani (*69*). Panel (A) corresponds to Panels (A) to (C) in Fig. 2. (**F**) The standard deviation of the estimates for the above five models. The orange lines denote the plate boundaries (*73*).

Fig. S6. Comparison of MTZ topography. MTZ thickness models from (**A**) this study, (**B**) Flanagan and Shearer (*26*), (**C**) Houser (*27*), and (**D**) Waszek et al. (*28*). Note that the other three models are global-scale. The average thickness estimates are shown at the bottom right corner of each panel.

Fig. S7. Precursor data coverage. (**A** to **D**) Logarithmic hit count patterns (number of traces per bin) from precursor data recorded before 2000 (A), before 2010 (B), before 2020 (C), and up to now (D). (**E**) Histogram of the precursor measurements used in this study by year. The blue curve shows the cumulative sum. The end times used by Flanagan and Shearer (*26*) (FS98), Houser (*27*) (H16), and Waszek et al. (*28*) (W21) are marked. The numbers in parentheses indicate the count of measurements available by the completion of each major study, based on the dataset used in this study. Our dataset spans 34 years, but most of our data come from the last decade.

Fig. S8. Comparison of MTZ topography errors. Maps of the errors in MTZ thickness from (**A**) this study, (**B**) Flanagan and Shearer (*26*), (**C**) Houser (*27*), and (**D**) Waszek et al. (*28*).

Fig. S9. Modeled subduction zones in the study region since 400 Ma from Matthews et al. (*31***).** The gray shades are the reconstructed plates, and the orange shades indicate the present-day continental plates. The thick brown lines indicate the present-day plate boundaries (*73*). (**A** and **B**) During the Devonian, the Rheic Ocean subducted along East Laurussia (EL) and western Gondwana, and the Laurussia supercontinent drifted southward and collided with western Gondwana. (**C** to **H**) Plate reconstructions of the Patagonian Plate (PP) and the western Gondwana (WG) subduction zones at (C) 340 Ma, (D) 310 Ma, (E) 290 Ma, (F) 250 Ma, (G) 200 Ma, and (H) 100 Ma (*29*). The retreat of the western Gondwana/South America (WG/SA) trench started at around 200 Ma.

Fig. S10. Plate reconstructions of southern Panthalassa between 240 and 150 Ma. Reconstructions of the Izanagi-Farallon-Phoenix triple junction are obtained from Torsvik et al. (*89*). The colored lines with triangles denote interpreted subduction zones in plate reconstruction model of van der Meer et al. (*21*). The dashed lines represent presumed transform zones. The profile D-D′ show the location of the cross-sectional view of model interpretation in Fig. 5.

Fig. S11. Modeling of slab temperature growth over a period of 250 Myr. The modeling assumes a 100-km-thick slab with an initial internal temperature of 1000 K (black line). The ambient mantle thermal conditions are 1600 K above the slab and 1800 K below the slab. A thermal diffusivity of 1×10^6 $m^2 \cdot s^{-1}$ is used. Over a period of 200 Myr, the temperature anomaly gradually widens, while the temperature within the slab remains lower than the ambient condition with a deficit of approximately 150 K.

Fig. S12. Comparison of the depth-dependent trend of the SPS slab volume with that of the reported slab flux. (**A**) The global estimate of the slab volume flux (the total slab volume consumed globally at subduction zones per unit time) from East et al. (*85*). (**B**) The solid black curve represents the depth variations in the SPS slab volume [cross-sectional area of tomographic anomalies in the LLNL-G3D-JPS-model (*20*)] (see Fig. 3C). The thin blue curve is the depthconverted slab flux (panel A) assuming a global average slab sinking rate of 1.2 cm/yr. The anticorrelation between the two curves indicates that a constant sinking velocity of 1.2 cm/yr fails to explain the observed slab sinking behavior. The dashed black lines indicate the depths of 410 km, 660 km, and 1000 km.

Fig. S13. Example of a common midpoint gather of the observed dataset. (**A**) The raw common midpoint gather. (B) The common midpoint gather after reconstruction and denoising using the FMSSA algorithm. The phases at around 150 s and 230 s prior to SS correspond to S410S and S660S precursors.

Fig. S14. Synthetic test based on the PREM model. (**A**) Decimated synthetic data with 10% traces missing and corrupted traces. (**B**) Reconstructed and denoised data using the FMSSA algorithm.

Fig. S15. Spike test. (**A**) Input topography model. (**B**) Output model from the common midpoint stacking of noise-free traces using the actual SS midpoint distribution (i.e., source-receiver geometry). (**C**) Same as (B) but with 10% random Gaussian noise, and 10% of the total traces are also contaminated with non-Gaussian erratic noise. (**D**) Resulting model after FMSSA reconstruction of the noisy data in (C).

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