

# ADVANCED MATERIALS

## Supporting Information

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Facet-Dependent Surface Charge and Hydration of  
Semiconducting Nanoparticles at Variable pH

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## Supplementary Information

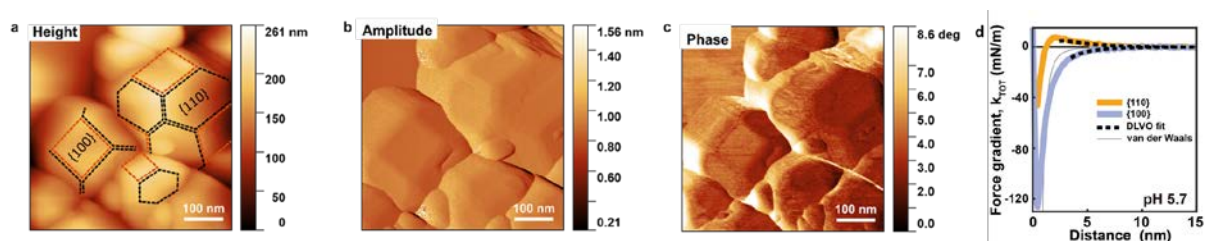
### Facet-dependent surface charge and hydration of semiconducting nanoparticles at variable pH

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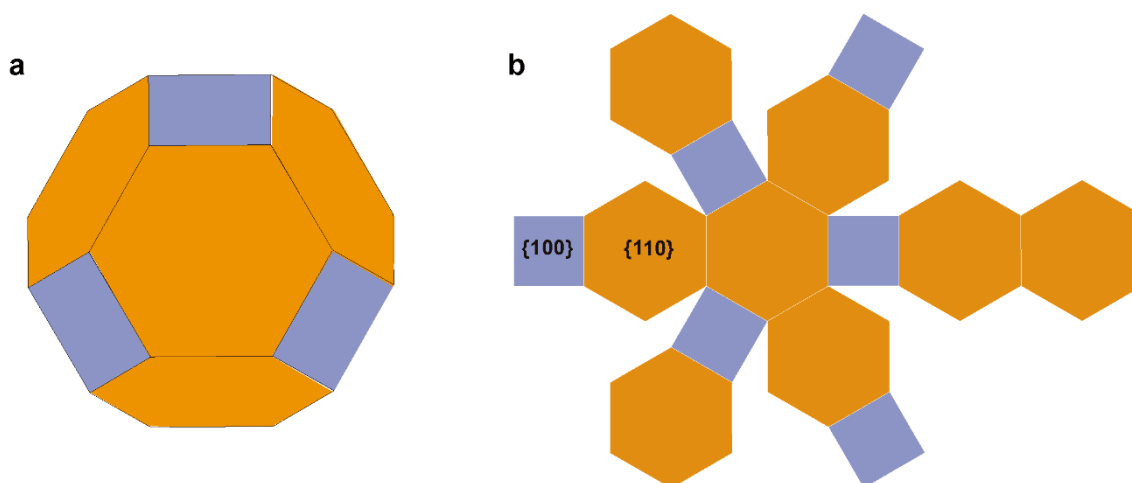
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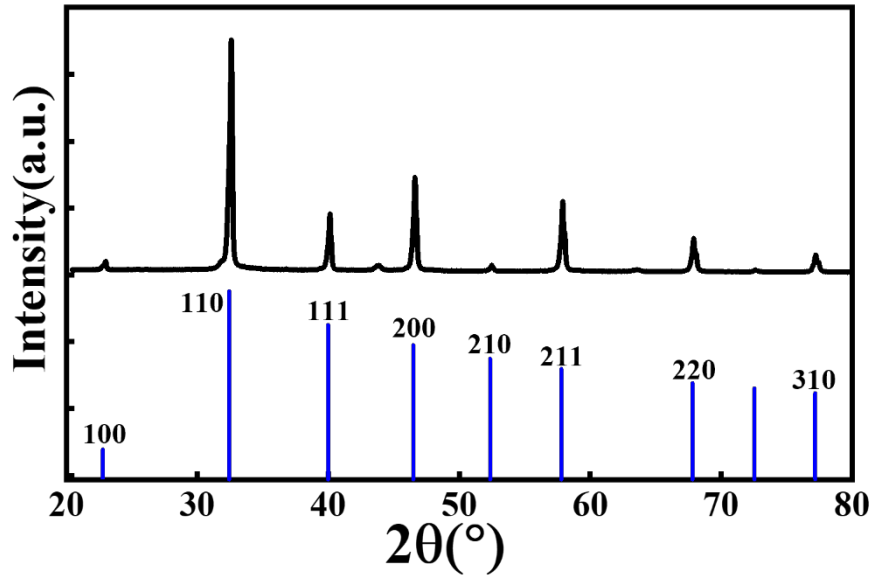
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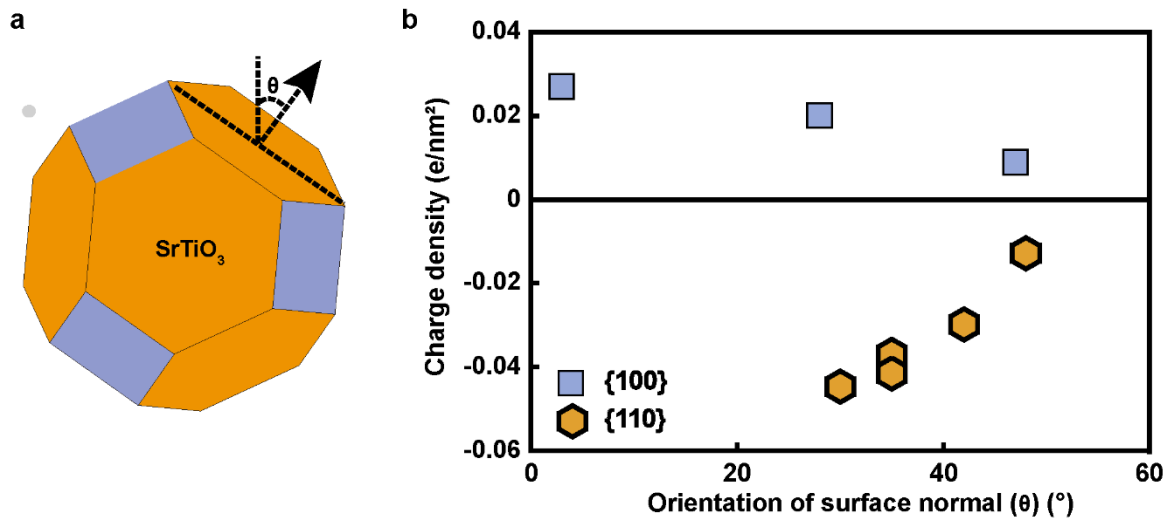
**Figure S1.** AFM images of SrTiO<sub>3</sub> nanoparticles in 10 mM NaCl, pH 5.7. a) Height; b) Amplitude; c) Phase. These images correspond to the location where two dimension (2D) force map measurement shown in Fig. 1c is performed. d) Average total force gradient vs. distance curves across a flat region at the centre of {100} and {110} facets of SrTiO<sub>3</sub> particles (marked with black and red circles on the 2D maps in Figure 1b).



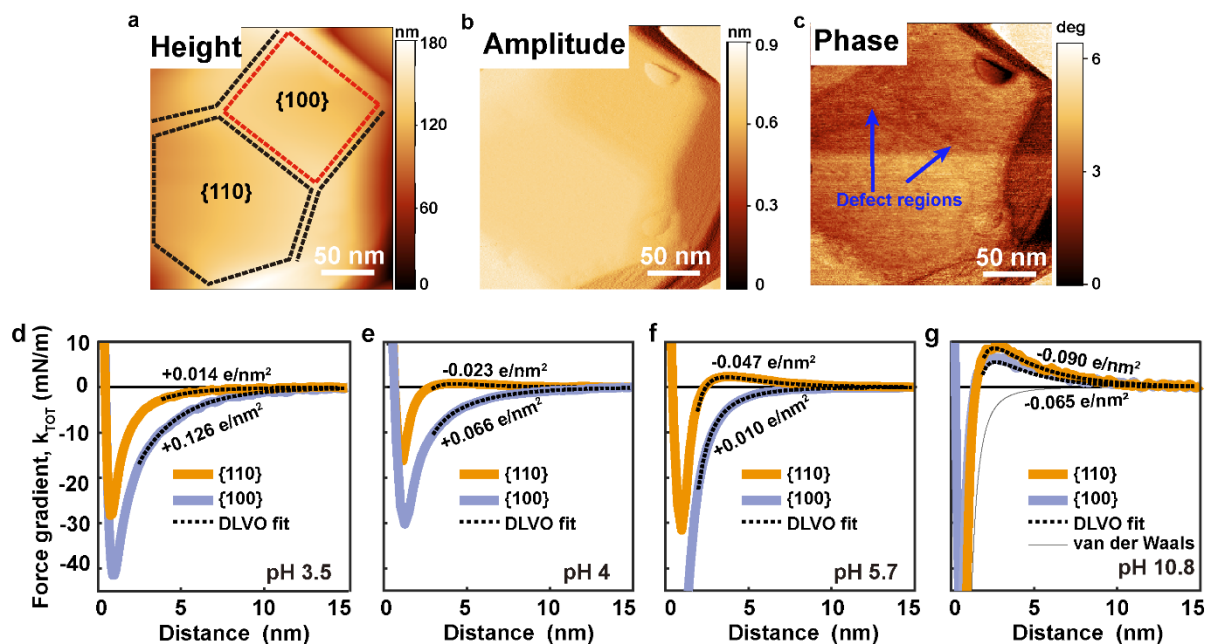
**Figure S2.** Folded and unfolded sketch of truncated rhombic dodecahedral shape. The facet fraction ratio of {100} shown in blue and {110} shown in orange region is 100:110=0.224:0.776.



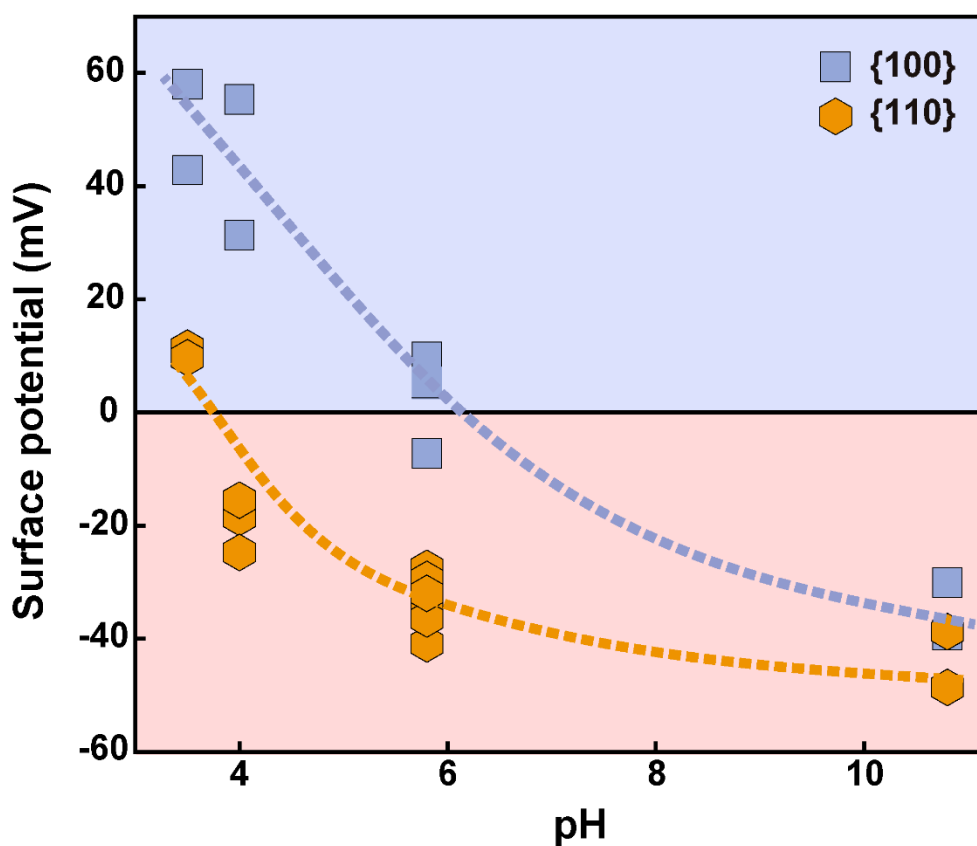
**Figure S3.** X-ray diffraction (XRD) patterns of hydrothermal synthesized  $\text{SrTiO}_3$  nanoparticles. Black curve is the XRD patterns of synthesized  $\text{SrTiO}_3$  nanoparticles. The measurements are performed using Bruker D2 Powder XRD. Blue lines correspond to standard PDF card of  $\text{SrTiO}_3$  (JCPDS No.35-0734).



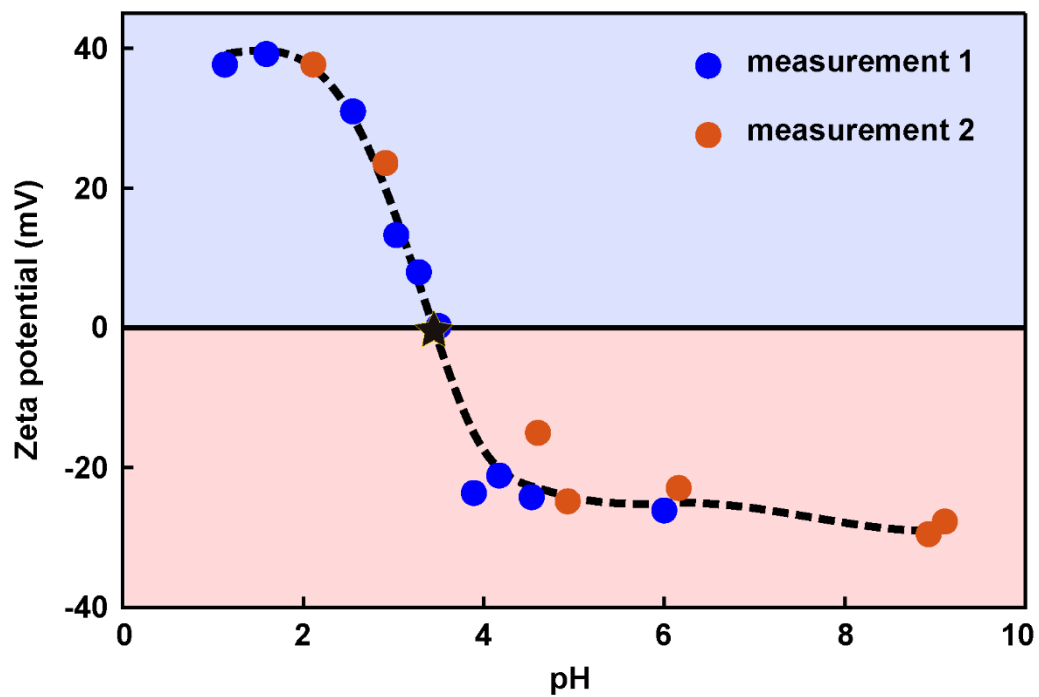
**Figure S4.** Effect of facet orientation on the surface charge density. a) A sketch representing the orientation of surface normal of  $\{110\}$  facet of  $\text{SrTiO}_3$ . b) The facet orientation effect on the surface charge densities. The surface charge densities were calculated from the local average forces on different facets in Figure 1c.



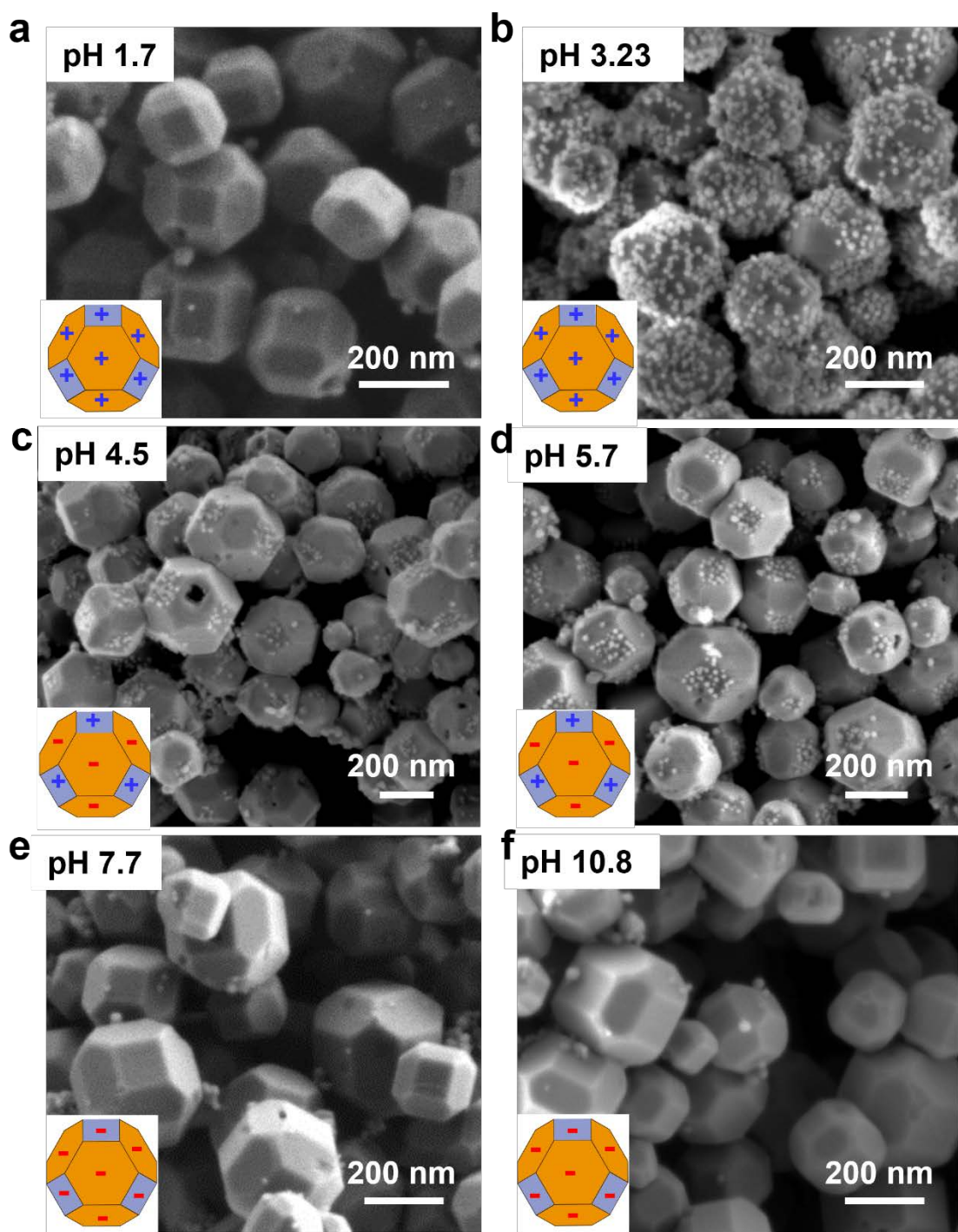
**Figure S5.** AFM images of SrTiO<sub>3</sub> nanoparticle a) Height; b) Amplitude; c) Phase images corresponding to the to the force maps in Figure 2. d-g) Total force gradient vs distance per facet averaged across a flat region at the centre of {100} and {110} facets of SrTiO<sub>3</sub> particles (marked with white circles on the 2D maps in Figure 2 a-d).



**Figure S6.** Electrical (potential) properties of {100} and {110} facets SrTiO<sub>3</sub> nanoparticles. Surface potential of {100} and {110} facets of SrTiO<sub>3</sub> nanoparticle as a function of pH for 10 mM NaCl. The surface charge densities shown in Fig. 3a are converted into surface potentials using the Graham equation (dash lines are guide the eyes).

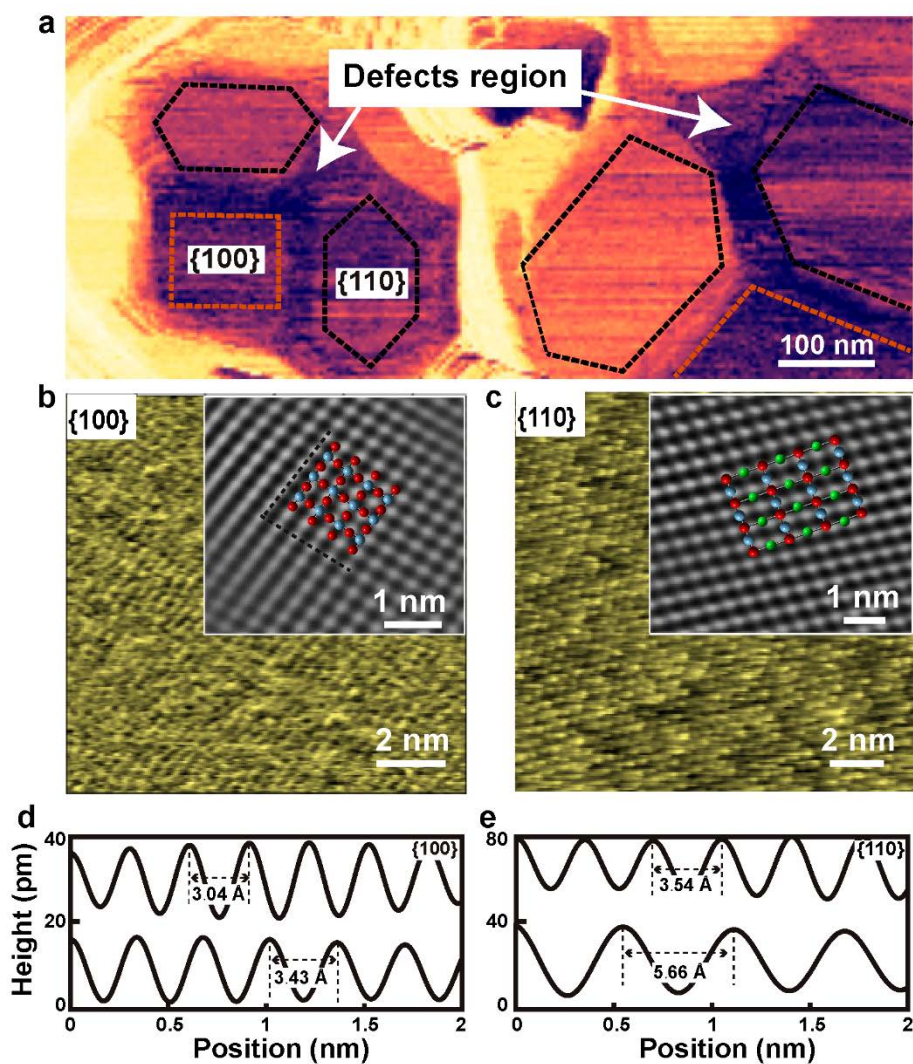


**Figure S7.** Measured zeta potential of SrTiO<sub>3</sub> nanoparticle suspension as function of pH in 10 mM NaCl. Blue and red dots represent two separate measurements (black dash line is to guide the eyes). The isoelectric point of SrTiO<sub>3</sub> nanoparticle suspension is 3.5.

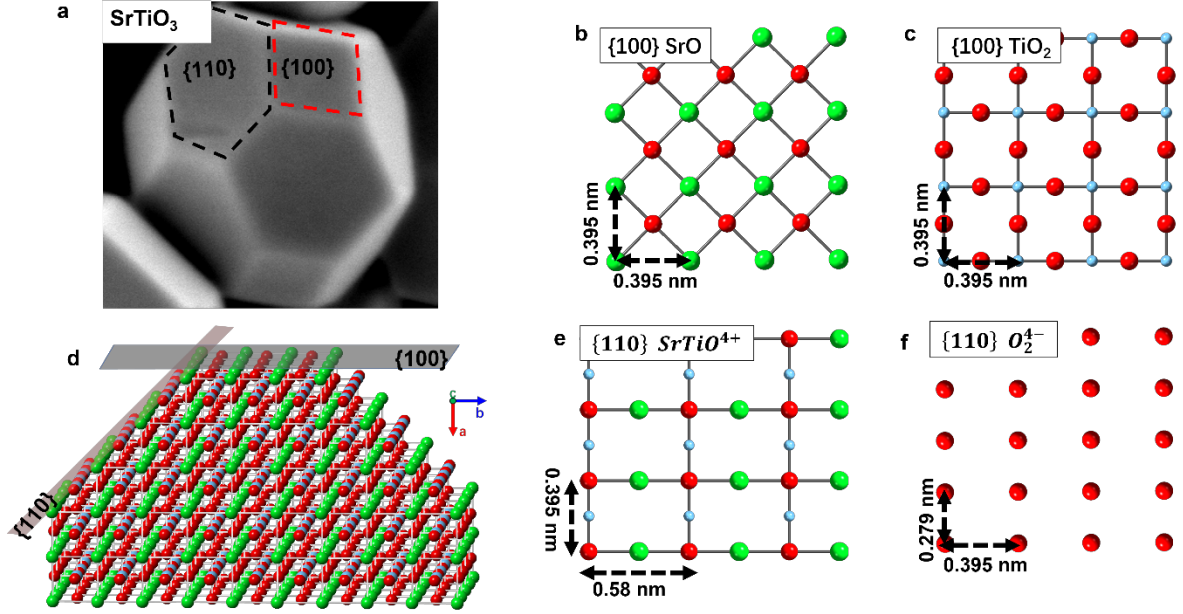


**Figure S8.** SEM images of SrTiO<sub>3</sub> nanoparticle after silica nanoparticles (12 nm) adsorption experiment at different pH. At pH 4 and 5.7 negatively charged SiO<sub>2</sub> nanoparticles are adsorbed only on positively charged squared {100} facets. At pH 3.23 both {100} and {110} facets of SrTiO<sub>3</sub> are positively charged and negatively charged SiO<sub>2</sub> are deposited on both facets. At pH 7.7 and 10.8 negatively charged facets and silica nanoparticles repel each other thus no adsorption of SiO<sub>2</sub> particles is observed. At pH 1.7, {100} and {110} facets of SrTiO<sub>3</sub> and silica NP are positively charged and hence adsorption is again suppressed for both facets. The insets show the charge sign of individual facets of SrTiO<sub>3</sub>.





**Figure S9.** High resolution images on SrTiO<sub>3</sub> nanoparticles measured using AM-AFM. a) AFM phase image of SrTiO<sub>3</sub> nanoparticles displaying hexagonal {110} and squared {100} facets and also a transition region between several neighbouring facets that display steps and disordered structure. Phase images contrast in regions with defects display a slightly different contrast with respect to the more homogeneous/flat regions, presumably related to variations of local surface chemistry such as a broken bonds. b) Atomic resolution topography images on {100} of SrTiO<sub>3</sub> measured in 10 mM NaCl at pH 6. It displays square lattice with lattice parameters  $a = 0.303$  nm and  $b = 0.343$  nm. c) Atomic resolution topography images on {110} facet of SrTiO<sub>3</sub> that show a rectangular structure with lattice parameters  $a = 0.354$  nm and  $b = 0.566$  nm. Insets represent Fourier-filtered view of atomic scale images superimposed with X-ray resolved crystallographic structure of {100} and {110}.



**Figure S10.** Crystallographic structure of different surface terminations on {100} and {110} facets of SrTiO<sub>3</sub>. O, Sr and Ti atoms are shown in red, green and blue, respectively. a) SEM image of faceted SrTiO<sub>3</sub> nanoparticles with square {100} and hexagonal {110} facets. b, c) The X-ray structure of SrO and TiO<sub>2</sub> terminations of {100} facets displaying a square symmetry with the lattice dimension of 0.395 nm. d) 3D crystallographic structure of SrTiO<sub>3</sub>, emphasizing {100} and {110} planes. e, f) The X-ray structure of SrTiO<sub>4</sub><sup>+</sup> and (2O)<sub>2</sub><sup>-</sup> terminations of {110} facets displaying rectangular structure. The SrTiO<sub>4</sub><sup>+</sup> termination has lattice parameters of 0.395 nm and 0.58 nm between oxygen atoms in two perpendicular directions. For (2O)<sub>2</sub><sup>-</sup> termination, the periodicities are 0.279 nm and 0.395 nm along the [001] and the [010] directions.

## Supplementary Information Note 1

**Surface charge determination from force-distance curves:** The DLVO-theory with charge regulation is used to extract the surface charge and potential of the SrTiO<sub>3</sub> nanoparticles from the measured force distance curves.<sup>[1]</sup> To do so, we calculate hypothetical force-distance curves for given surface charge and regulation parameter and compare these curves with the measured curves using the surface charge and regulation parameter as fitting parameters.<sup>[2][3][4]</sup> To determine the force between the tip and particle surface, we first calculate the disjoining pressure  $\Pi(h)$  in the gap with height  $h$  between them. This pressure can be split in a contribution  $\Pi_{vdW}$  due to van der Waals interactions and an electrostatic double layer contribution  $\Pi_{EDL}$ . The force on the tip is calculated by integrating  $\Pi$  over the spherical tip with radius  $R_{tip}$ :

$$F_{int}(h) = \int_h^\infty k_{int}(h') dh' = 2\pi R_{tip} \int_h^\infty [\Pi_{EDL}(h') + \Pi_{vdW}(h')] dh' \quad (1)$$

Ignoring retardation effects, the Van der Waals contribution is calculated using :

$$\Pi_{vdW}(h) = \frac{-A_H}{6\pi h^3} \quad (2)$$

where  $A_H$  is the Hamaker constant, and  $h$  is the tip to surface distance. The Hamaker constants  $A_H$  for the (silica-water-SrTiO<sub>3</sub>) system was fixed to  $2.259 \cdot 10^{-20}$  J. The electrostatic double layer contribution contains the required information on the surface charge and potential. For a 1-1 electrolyte it is given by:

$$\Pi_{EDL}(h) = c_\infty k_B T \left( 4 \sinh^2 \left( \frac{e\psi(z)}{2k_B T} \right) - \kappa^2 \left( \frac{d\psi}{dz} \right)^2 \right) \quad (3)$$



where  $\psi(z)$  is the electric potential for  $0 < z < h$ ,  $k_B$  being the Boltzmann constant,  $T$  the temperature,  $c_\infty$  the bulk ion concentration and  $e$  the elementary charge;  $\kappa$  is the inverse Debye length:

$$\kappa = \sqrt{\frac{2e^2 c_\infty}{\epsilon \epsilon_0 k_B T}} \quad (4)$$

Calculation of the electric double layer contribution requires knowledge of the potential  $\psi(h)$  in the electrolyte. This potential is governed by the Poisson-Boltzmann (PB) equation:

$$\frac{d^2}{dz^2} \psi(z) = \frac{2e}{\epsilon \epsilon_0} c_\infty \sinh\left(\frac{e\psi(z)}{k_B T}\right) \quad (5)$$

and the boundary conditions at both surfaces. Here we employ the charge regulation (CR) approximation assuming a linear relation between surface charge and surface potential<sup>[31]</sup>:

$$\sigma_s = \sigma_s^\infty + c_B(\psi_s - \psi_s^\infty) \quad (6)$$

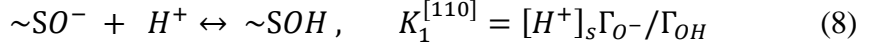
where  $\sigma_s^\infty$  and  $\psi_s^\infty$  are the surface charge density and surface potential of the isolated interface, and  $c_B$  is an effective capacitance. Here, we refer to surface charge density  $\sigma$  as the charge which is compensating the diffuse layer charge. The surface charge density and surface potential of the isolated interface are related by the Grahame relation:

$$\sigma_s^\infty = 2(2c_\infty \epsilon \epsilon_0 k_B T)^{1/2} \sinh\left(\frac{e\psi_s^\infty}{2k_B T}\right) \quad (7)$$

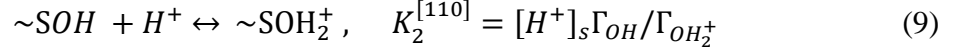
When the electrical potential profile  $\psi(z)$  is known, the disjoining pressure  $\Pi_{EDL}(h, c_B, \psi_s^\infty)$  and the resulting force (or force gradient) are obtained with Equation 3 and 1, respectively. By fitting  $\Pi_{EDL}(h, c_B, \psi_s^\infty)$  to the experimentally obtained  $(h, \Pi_{EDL})$  data, we obtain the desired values for the surface potential  $\psi_s^\infty$  and, with Equation 7, the surface charge  $\sigma_s^\infty$ . The lower limit for the fitting range is set to 1.5-2 nm, in order to minimize the influence of short range forces, since they are not included in the model. The upper boundary was set to 15 nm, above which the tip-sample interaction force is negligible. Changing the upper limit of the fitting boundary to 20 nm leads to a 1% change in the resulting charge. Variation of the lower limit from 1 to 3.5 nm results in a maximum of 10% change in the surface charge.<sup>[2][4]</sup>

## Supplementary Information Note 2

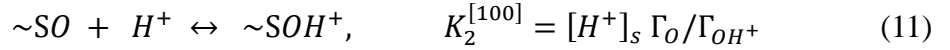
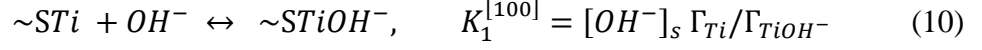
**Surface complexation modeling:** The surface charge on solid-water interfaces is modelled using standard surface complexation models involving the adsorption/desorption of ions  $X_i \in \{H^+, OH^-, Na^+, Cl^-\}$  to surface sites  $S$  according to  $\sim SX_i \leftrightarrow \sim S + X_i$ . Each reaction is characterized by an equilibrium constant  $K_i$  with a corresponding  $pK_i = -\log_{10} K_i$ . According to the law of mass action, the ion concentration  $[X_i]_s$  at the surface depends on the surface concentrations  $\Gamma_{SX_i}$  and  $\Gamma_S$  as:  $[X_i]_s = K_i \Gamma_{SX_i} / \Gamma_S$ . On the other hand,  $[X_i]_s$  is related to the corresponding bulk concentration  $c_i^\infty$  by the Boltzmann factor  $[X_i]_s = c_i^\infty \exp(-Z_i e \psi_s / k_B T)$ , where  $Z_i = \pm 1$  is the valency of the ion and  $\psi_s$  is the potential at the surface. Moreover, the Grahame equation relates the surface potential to the surface charge:  $\sigma_s = 2(2c_\infty \epsilon \epsilon_0 k_B T)^{1/2} \sinh(e\psi_s / 2k_B T)$ , where  $c_\infty$  is the total ionic strength. Using both relations for  $[X_i]_s$  the surface charge can be calculated for given  $I_\infty$  and  $pH$ .<sup>[5][6]</sup> The experiments have been performed in 10 mM  $NaCl$  solutions with different  $pH$  values. Since  $Na^+$  and  $Cl^-$  weakly absorb on  $SrTiO_3$ , the surface charge is merely controlled by adsorption or desorption of hydrogen and/or hydroxyl ions. For the  $O^-$  terminated [110] facet we assume that the charge of the surface oxygen is compensated by the uptake of a proton:



Additional protonation of the surface hydroxyl  $\sim SOH$  group yields the formation of  $\sim SOH_2^+$  according to:



while  $\Gamma_{[110]} = \Gamma_{OH} + \Gamma_{O^-} + \Gamma_{OH_2^+}$  is constant. For the  $TiO_2$  terminated [100] facet we assume that the (slightly positive) titanium atom in the interface is able to adsorb a hydroxyl ion while the (slightly negative) oxygen atom can adsorb a proton:



while  $\Gamma_{[100]} = \Gamma_{Ti} + \Gamma_{TiOH^-} = \frac{1}{2}(\Gamma_O + \Gamma_{OH^+})$  is constant.

The resulting surface charges are given by:

$$\sigma_s^{[110]} = e\Gamma_{[110]} \frac{[H^+]_s^2 - K_1^{[110]} K_2^{[110]}}{K_2^{[110]} [H^+]_s + [H^+]_s^2 - K_1^{[110]} K_2^{[110]}}, \quad (12)$$

$$\sigma_s^{[100]} = e\Gamma_{[100]} \left( \frac{2[H^+]_s}{K_2^{[100]} + [H^+]_s} - \frac{[OH^-]_s}{K_1^{[100]} + [OH^-]_s} \right) \quad (13)$$

with  $[H^+]_s = 10^{-pH} e^{-e\psi_s/k_B T}$  and  $[OH^-]_s = 10^{-14} / [H^+]_s$ . Substitution of these two expressions in Equation 12 and 13 results in a single implicit equation for  $\sigma_s$  that depends only on the two equilibrium constants  $K_1$  and  $K_2$ . Optimal values for  $K_1$  and  $K_2$  (or equivalently  $pK_1$  and  $pK_2$ ) for both facets can be found by fitting the obtained relation  $\sigma_s(pH, K_1, K_2)$  to the experimental ( $pH, \sigma_s$ ) data using a least square fit procedure. Parameter values are:  $pK_1^{\{100\}} = 5.6$ ,  $pK_2^{\{100\}} = 8.4$ ,  $pK_1^{\{110\}} = 4.3$  and  $pK_2^{\{110\}} = 3.1$ .

## Supplementary References

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