Protein structure determination using a database of interatomic distance probabilities

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Abstract

The accelerated pace of genomic sequencing has increased the demand for structural models of gene products. Improved quantitative methods are needed to study the many systems (e.g., macromolecular assemblies) for which data are scarce. Here, we describe a new molecular dynamics method for protein structure determination and molecular modeling. An energy function, or database potential, is derived from distributions of interatomic distances obtained from a database of known structures. X-ray crystal structures are refined by molecular dynamics with the new energy function replacing the Van der Waals potential. Compared to standard methods, this method improved the atomic positions, interatomic distances, and side-chain dihedral angles of structures randomized to mimic the early stages of refinement. The greatest enhancement in side-chain placement was observed for groups that are characteristically buried. More accurate calculated model phases will follow from improved interatomic distances. Details usually seen only in high-resolution refinements were improved, as is shown by an *R*-factor analysis. The improvements were greatest when refinements were carried out using X-ray data truncated at 3.5 Å. The database potential should therefore be a valuable tool for determining X-ray structures, especially when only low-resolution data are available.

Keywords: knowledge-based modeling; low-resolution; molecular dynamics; structure refinement; X-ray crystallography

Genomic sequencing efforts are far outpacing our understanding how gene products (i.e., RNA and proteins) give rise to the characteristics of living things. The more than 8,500 currently known protein structures ought to be a rich resource for solving this problem. For example, bond distances and angles obtained from known small-molecule structures commonly are used to generate geometrical restraints in solving new X-ray and NMR structures (Hendrickson & Konnert, 1980; Engh & Huber, 1991). Much of what protein databases "know" about protein structure, however, remains hidden when the data are viewed in this limited way.

Methods for determining X-ray and NMR structures are well understood, and procedures have become largely automated. Many problems still exist, however, in obtaining structures using limited data. This is especially true when the atomic coordinates are underdetermined by the data alone, such as happens for poorly diffracting crystals of molecular complexes. In these cases, prior information about molecular structure must be used to obtain a useful model.

In the limit where no experimental data are available, there is an exclusive reliance on prior information, and the problem becomes an even more difficult one of molecular modeling. Homology modeling has brought significant developments (Browne et al., 1969; Blundell et al., 1987; Jones & Thornton, 1996; Sanchez & Sali, 1997), but use of this method requires that the structure of a molecule with high sequence homology be determined (this fact motivates the structural genomics initiatives reviewed in Terwilliger et al., 1998). Another limitation to homology modeling is that the same sequence is capable of adopting different structures depending on tertiary context (Minor & Kim, 1996). There are also numerous examples of proteins without significant sequence homology that, nevertheless, have high structural homology.

To enhance structure determination and molecular modeling, we have implemented a knowledge-based energy function that makes greater use of the information in structure databases. The energy function is derived from distributions of interatomic distances estimated from an ensemble of reference structures (see Methods). These distributions are referred to as "probability density functions" (PDFs). Molecular dynamics (MD) using a PDF energy adjusts the structures so that the interatomic distances of the model are similar to those found in the ensemble.

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Standard MD refinement techniques (Brünger et al., 1987) use Van der Waals (VDW) interactions to help determine the distance between nonbonded atoms. Figure 1 compares the VDW interaction between two methyl carbons with both the PDF for methyl carbons on valine residues separated by three peptide bonds (m = 3) and the m = 3 PDF for leucine–leucine methyl carbons. The PDF energy is much more richly structured and depends on the context of the atom pairs.

We reason that by replacing VDW with the PDF energy, much more information will be available for nonbonded interactions between atoms, providing a powerful method for optimizing protein structures. Here we present the results of a test of this idea: A PDF energy function was implemented in crystallographic refinement, and the method was tested by running test refinements on randomized X-ray structures. The results show that using the PDF energy function significantly improved the structures obtained at the end of refinement. Analysis of the refined structures shows where the improvements are and why this is a useful method in determining X-ray structures.

Results

A 1.3 Å resolution crystal structure of wild-type myoglobin (Romo, 1998) and a 2.0 Å resolution structure of calmodulin complexed with brain calmodulin-dependent protein kinase II-alpha (Wall et al., 1997; Protein Data Bank (PDB) entry 1cm1) were randomized (see Methods). These structures were refined using real X-ray data truncated at 2.0 and 3.5 Å to simulate low-resolution crystallog-raphy. Three refinement methods were used: Cartesian molecular dynamics using standard potentials (referred to as "NORM" in tables and figures), Cartesian molecular dynamics with the VDW potential replaced by the PDF potential ("PDF" in tables and figures), and torsion-angle refinement (Rice & Brünger, 1994) with standard potentials ("TORS" in tables and figures).

The following sections describe analyses of the atomic positions, interatomic distances, stereochemistry, side-chain dihedral



Fig. 1. Comparison of probability densities for methyl carbon pairs. PDF probability densities vs. distance for V:C γ –V:C γ and L:C δ –L:C δ atom pairs (m = 3) are contrasted with a methyl carbon–methyl carbon VDW potential. The PDF depends upon the context of the methyl carbon and has more features than the probability due to VDW. Note that the PDF densities are shifted to higher distances than in VDW, showing the effect of the backbone connectivity on the distance.

angles, *R*-factors, and PDF scores of the refined structures. The X-ray data and known structures were used to analyze the structures and evaluate the refinements. The analyses show that refinement was improved by using the database potential.

Atomic positions

The root-mean-square deviation (RMSD) of backbone and sidechain atomic positions in refined structures was calculated with respect to the X-ray structures (Table 1). In all but one case, both backbone and side-chain positions were more accurate in PDF refinement. The exception is for the backbone of myoglobin refined at 2.0 Å resolution, where no advantage is observed. Including higher resolution data in refinement led to improved atomic positions in all cases.

Interatomic distances

The RMSDs of interatomic distances between refined structures (NORM and PDF) and the X-ray structures were calculated as a function of the interatomic distance. Results for torsion-angle refinement are almost indistinguishable from those for normal refinement. Use of the PDF potential improved overall interatomic distances in all refined structures (Table 1).

In the 3.5 Å resolution PDF refinement of myoglobin, there was a relatively uniform improvement in the RMSD of interatomic distances for distances between 5 and 25 Å (Fig. 2). The difference decreases slightly from 25 to 30 Å, and then is uniform until 40 Å. Above 40 Å, where the number of atom pairs is small, the RMSD becomes highly variable. Atom pairs involving surface residues likely begin to dominate the statistics at 25 Å. Results from the calmodulin complex (not shown) are similar, except that the RMSD becomes small at long distances, perhaps due to end-to-end pack-

Table 1. RMSDs of backbone and side-chain atomic positions,and RMSD of interatomic distances, all calculated betweenrefined structures and available high-resolution structures a

| | Meth | Back | Side | Dist | BOND | ANGL | DIHE | IMPR |
|---------|------|------|------|------|------|------|------|------|
| 3.5A Mb | norm | 0.52 | 1.95 | 1.25 | 0.47 | 1.07 | 25 | 0.60 |
| | pdf | 0.33 | 1.47 | 0.88 | 0.70 | 1.12 | 23 | 0.94 |
| | tors | 0.52 | 1.94 | 1.24 | 0.35 | 0.84 | 26 | 0.52 |
| CaM | norm | 0.73 | 1.96 | 1.31 | 0.39 | 0.78 | 29 | 0.43 |
| | pdf | 0.45 | 1.60 | 1.05 | 0.76 | 1.10 | 24 | 0.68 |
| | tors | 0.69 | 1.84 | 1.21 | 0.35 | 0.70 | 29 | 0.40 |
| 2.0A Mb | norm | 0.23 | 1.63 | 1.03 | 0.94 | 1.31 | 22 | 0.83 |
| | pdf | 0.24 | 1.41 | 0.87 | 0.83 | 1.12 | 20 | 0.84 |
| | tors | 0.23 | 1.68 | 1.03 | 0.63 | 1.08 | 23 | 0.68 |
| CaM | norm | 0.48 | 1.90 | 1.19 | 0.96 | 1.49 | 27 | 0.74 |
| | pdf | 0.25 | 1.58 | 1.00 | 1.40 | 1.59 | 26 | 1.09 |
| | tors | 0.48 | 1.74 | 1.07 | 0.60 | 0.86 | 28 | 0.49 |
| | | | | | | | | |

^aAlso shown are RMSDs from ideal bond distances (BOND), bond angles (ANGL), multimodal dihedral angles (e.g., backbone) (DIHE), and other dihedral angles (IMPR). Back, Side, and Dist values are in angstroms. BOND values are in 10^{-2} Å. ANGL, DIHE, and IMPR values are in degrees. Structures were chosen on the basis of the lowest $R_{\rm free}$ (see Table 3).



Fig. 2. The difference (as calculated by RMSD) between interatomic distances in the X-ray structure and the 3.5 Å resolution refined structure of myoglobin plotted vs. interatomic distance. RMSD is defined for structures (1) and (2) at distance *s* as $[\langle (d_{ij}^{(1)} - d_{ij}^{(2)})^2 \rangle_{|d_{ij}^{(1)} - s| < b}]^{1/2}$, where $\{i\}$ and $\{j\}$ span all atoms in the structure, $d_{ij} =$ distance between atoms *i* and *j*, and *b* is half the bin size. Values are averaged in 0.25 Å bins.

ing in the crystal lattice. Plots are generally similar for 2.0 Å resolution refinements (not shown).

Stereochemistry

The RMSDs of standard geometry parameters from the mean were calculated for each of the refined structures (Table 1). No VDW clashes were found in any structure, and the values indicate that all of the structures have normal stereochemistry. The PDF potential generally lowered the RMSD for multimodal dihedral angles DIHE (phi, psi, and chi, as defined by IUPAC) while raising the RMSD for the other parameters.

Side-chain conformations

The fraction of correct side-chain rotational isomers (Volkenstein, 1963) or "rotomers" (as defined by IUPAC) was calculated for each of the refined structures (Table 2). With one exception, the fraction is higher in every instance for PDF-refined structures. The exception is the 2.0 Å myoglobin refinement, where more NORM refinement yielded better placement of χ^1 and χ^3 rotomers, and a higher overall fraction of correct rotomers. PDF refinement in this instance produced both a higher fraction of correct χ^2 rotomers.

An analysis of the fraction of correct combined χ^1 and χ^2 rotomers vs. residue type was done for 3.5 and 2.0 Å resolution refinements (Fig. 3). In both 3.5 Å resolution refinement (Fig. 3A) and 2.0 Å resolution refinement (Fig. 3B), PDF improved the side-chain conformations for all but the polar category.

Nonpolar and beta-branched categories showed the biggest improvement in fraction of correct combined χ^1 and χ^2 rotomers. Remarkably, a separate calculation for the 3.5-Å resolution calmodulin refinement showed that all of the (χ^1, χ^2) combinations are correct for this case. The PDF refinement method thus positioned interior residues particularly well. This is understandable,

Table 2. Fraction of correct side-chain rotomers^a

| | | Meth | χ^1 | χ^2 | χ^3 | χ^1 and χ^2 | All |
|---------|-----|------|----------|----------|----------|-----------------------|------|
| 3.5 Å | Mb | norm | 0.66 | 0.72 | 0.49 | 0.53 | 0.65 |
| | | pdf | 0.79 | 0.81 | 0.58 | 0.69 | 0.76 |
| | | tors | 0.71 | 0.75 | 0.29 | 0.58 | 0.66 |
| | CaM | norm | 0.62 | 0.48 | 0.36 | 0.28 | 0.52 |
| | | pdf | 0.75 | 0.62 | 0.40 | 0.51 | 0.64 |
| | | tors | 0.61 | 0.52 | 0.26 | 0.31 | 0.52 |
| 2.0 Å M | Mb | norm | 0.87 | 0.77 | 0.67 | 0.69 | 0.80 |
| | | pdf | 0.83 | 0.87 | 0.51 | 0.78 | 0.79 |
| | | tors | 0.73 | 0.77 | 0.49 | 0.61 | 0.70 |
| | CaM | norm | 0.70 | 0.57 | 0.36 | 0.45 | 0.60 |
| | | pdf | 0.76 | 0.66 | 0.45 | 0.51 | 0.67 |
| | | tors | 0.72 | 0.57 | 0.49 | 0.48 | 0.62 |

^aStructures were chosen on the basis of the lowest R_{free} (see Table 3).

due the increased number of distance distributions (and thus information) available for buried atoms.

Crystallographic R-factors

R-factors were calculated at multiple resolutions using both real and simulated diffraction data (Table 3). R_{free} (Brünger, 1992)



Fig. 3. Analysis of the fraction of correct combined $\chi 1$ and $\chi 2$ rotomers vs. residue type. Myoglobin and calmodulin results are combined to calculate the fractions for (**A**) 3.5 Å refinement and (**B**) 2.0 Å refinement. Residue type definitions: charged = EDHKR; polar = NQSTY; nonpolar = CFILMVW; beta-branched = ITV; ring = FHWY; 1 dihedral = CSTV; 2 dihedrals = FHDILNWY; 3 dihedrals = EMQ; >3 dihedrals = KR; all = CDEFHIKLMNQRSTVWY.

Table 3. R, R_{free} , and R_{calc} evaluated to the resolution indicated at the top of the column^a

| | | | R | | | R_{free} | | | R_{calc} | | |
|----------|------|--------------|-------|-------|-------|------------|-------|-------|------------|-------|-------|
| | Meth | Wght | 3.5 Å | 2.0 Å | 1.3 Å | 3.5 Å | 2.0 Å | 1.3 Å | 3.5 Å | 2.0 Å | 1.3 Å |
| 3.5 Å Mb | norm | $1.0 \times$ | 0.19 | 0.43 | 0.48 | 0.38 | 0.47 | 0.50 | 0.26 | 0.43 | 0.48 |
| | pdf | $1.0 \times$ | 0.22 | 0.37 | 0.42 | 0.35 | 0.42 | 0.45 | 0.25 | 0.36 | 0.41 |
| | tasa | 0.5 	imes | 0.23 | 0.43 | 0.48 | 0.39 | 0.47 | 0.50 | 0.28 | 0.43 | 0.47 |
| CaM | norm | $0.5 \times$ | 0.28 | 0.45 | N/A | 0.43 | 0.52 | N/A | 0.35 | 0.46 | N/A |
| | pdf | $1.0 \times$ | 0.28 | 0.42 | N/A | 0.39 | 0.46 | N/A | 0.31 | 0.40 | N/A |
| | tasa | 0.5 	imes | 0.26 | 0.45 | N/A | 0.42 | 0.52 | N/A | 0.33 | 0.45 | N/A |
| 2.0 Å Mb | norm | $2.0 \times$ | 0.26 | 0.30 | 0.34 | 0.31 | 0.36 | 0.38 | 0.23 | 0.26 | 0.30 |
| | pdf | $1.0 \times$ | 0.27 | 0.30 | 0.34 | 0.30 | 0.34 | 0.36 | 0.22 | 0.25 | 0.29 |
| | tasa | 1.0 	imes | 0.28 | 0.32 | 0.37 | 0.33 | 0.38 | 0.40 | 0.25 | 0.30 | 0.34 |
| CaM | norm | $2.0 \times$ | 0.25 | 0.35 | N/A | 0.38 | 0.44 | N/A | 0.31 | 0.37 | N/A |
| | pdf | $2.0 \times$ | 0.28 | 0.35 | N/A | 0.38 | 0.43 | N/A | 0.30 | 0.34 | N/A |
| | tasa | $1.0 \times$ | 0.27 | 0.36 | N/A | 0.38 | 0.46 | N/A | 0.32 | 0.37 | N/A |

^aValues are for refinements of myoglobin (Mb) and calmodulin (CaM) using X-ray data truncated at 3.5A and 2.0 Å. The relative weight of the X-ray data is indicated in the Wght column. Only results from structures with the lowest value of R_{free} are listed.

(using real data) and R_{calc} (using simulated data) are smaller for structures obtained by PDF refinement. The resolution-dependent analysis shows that the PDF potential improved the high resolution features of the structural models, especially when only low resolution data were used for refinement. Note that R_{free} is not the best measure of the quality of a structure for our tests, as it does not make use of the available high-resolution structure model.

PDF scores by residue

The average PDF score per residue was calculated for the original X-ray structures of myoglobin and the calmodulin complex (Fig. 4). In both structures, there is a good correspondence between regions with relatively uncommon interatomic distances (high PDF score) and nonhelical domains.

The same plots were compared with ones obtained from refined structures by calculating correlation coefficients. For myoglobin refined at 3.5 Å resolution using normal methods, the correlation calculated between the profiles is 0.63 (Fig. 4A). For the calmodulin complex refined at 3.5 Å resolution, the correlation is 0.53 (Fig. 4C). For both myoglobin and calmodulin refined at 3.5 Å resolution using the PDF method, the correlation coefficient calculated between the profiles is 0.95 (Fig. 4B,D). The correlations calculated for structures refined at 2.0 Å resolution were: 0.86 for Mb, normal refinement; 0.96 for Mb, PDF refinement; 0.78 for CaM, normal refinement; and 0.94 for CaM, PDF refinement. Using the PDF energy in low-resolution refinements, therefore, has the striking result of producing PDF profiles that usually come only from high-resolution structure determinations.

Discussion

We have demonstrated that a novel energy function, the PDF database potential, enhances protein structure determination. Compared to standard methods, use of the PDF energy improved atomic positions, interatomic distances, side-chain conformations and R-factors of refined randomized structures of myoglobin and a calmodulin complex. The results indicate that the PDF potential should be a valuable tool for determining X-ray structures, especially in the early stages of a structure determination or when only low-resolution data are available. Future studies will address the value of the method in molecular modeling and NMR refinement.

Plots of the PDF score by residue (Fig. 4) suggest two rules for evaluating structures. The first is that any *alpha-helical* region with a PDF score >0 is likely to have incorrectly modeled side chains. The second is that *any* region with a PDF score >0.2 (using normalization W = 0.58) is likely to be incorrectly modeled. These two rules together can be used as a heuristic method for validating protein structures using plots of the PDF score by residue.

PDF refinements accurately reproduced both the minima *and* maxima in plots of PDF score by residue that were obtained from the high-resolution X-ray structures (Fig. 4). The PDF method thus may improve structural models even when it is applied to proteins where the real interatomic distances are relatively uncommon. Future studies, however, will be needed to demonstrate the method's effectiveness when applied to a broad range of protein classes.

It has been suggested that multimodal information (such as the PDF potential), while useful for validation methods, is inappropriate for use in refinement (Sheldrick & Schneider, 1997). One reason is that structures tend to become trapped in local minima, making it impossible to locate the global minimum without a prohibitively time-consuming search of the space of model parameters. Another is because the negative curvatures in multimodal distributions render useless many optimization methods that rely on curvature information.

Use of multimodal information is certainly not appropriate in the latter stages of refinement of a high-resolution structure, where full-matrix least-squares methods are applied. It is certainly useful, however, in the early stages of refinement, when the solution is far from the global minimum in the X-ray target, or in low-resolution refinements, when the full-atom solution is underdetermined by the X-ray data. In the early stages of refinement, use of the PDF potential can improve interatomic distances, improving model phases and electron-density maps (observed in an application to a

Myoglobin, 3.5 Angstrom Normal Refinement Myoglobin, 3.5 Angstrom PDF Refinement 0.2 0.2 Crystal structure Crystal structure Correlation = 0.63 В Correlation = 0.95 Refined structure Refined structure 0.15 0.15 0.1 0.1 0.05 0.05 0 0 -0.05 -0.05 -0.1 -0.1 -0.15 -0.15 -0.2 -0.2 60 80 Residue Number 0 20 120 160 0 20 60 80 10 Residue Number 40 100 140 40 100 120 140 160 Calmodulin Complex, 3.5 Angstrom Normal Refinement Calmodulin Complex, 3.5 Angstrom PDF Refinement 0.2 0.2 Crystal structure Refined structure Crystal structure Refined structure D Correlation = 0.95 relation = 0.530.15 0.15 0.1 0.1 0.05 0.05 o 0 -0.05 -0.05 -0.1 -0.1 -0.15 -0.15 -0.2 -0.2 100 150 Residue Number 0 50 200 250 100 150 Residue Number ٥ 50 250 150 200

PDF Scores by Residue

Fig. 4. PDF scores calculated by residue for X-ray structures from (**A**) normal refinement of myoglobin, (**B**) PDF refinement of myoglobin, (**C**) normal refinement of the calmodulin complex, and (**D**) PDF refinement of the calmodulin complex, all refined at 3.5 Å resolution. On the calmodulin complex plots, linker residues 74–83 were deleted, and residues beginning at 200 correspond to the peptide. The correlation coefficient calculated between the plots is indicated on each graph. Nonhelical domains are indicated by boxes at the bottom. PDF scores above 0 indicate relatively uncommon interatomic distances.

troponin-C structure problem by Soman et al., 1999). The improvements thus can potentially speed the process of obtaining a good initial model for the latter stages of refinement.

The argument for using the PDF potential in refinement is even stronger when only low-resolution data are available. In this case, many atomic models will fit the data equally well, so that to choose the most likely structure one must rely on prior information about protein structure. In Bayesian terms, the PDF potential can be seen as an estimate of the prior probability for the interatomic distances in the protein (given the assumption of independent atom pairs). When high resolution data are available, the PDF potential is less useful, because (1) the data are sufficient to determine the structure, and (2) high-resolution structures already have PDF profiles that are similar to those obtained by minimization of the PDF energy, leaving little room for improvement.

Our tests definitively show that the PDF refinement method can be useful in increasing the quality of a structural model. One remarkable supporting result is that the high-resolution features (i.e., *R*-factors) of both myoglobin and the calmodulin complex were improved by use of the PDF potential in low-resolution refinement. Another is that refinement at low resolution accurately reproduced the entire profile of the PDF score by residue for the high-resolution X-ray structures (correlation = 0.95), whereas normal refinement generated models with a much poorer agreement (correlation <0.65).

There are two ways to explain the enhancement in refinement. One is that the PDF potential provides extra information that is not available in refinement with ordinary potentials. The PDF potential is derived only from compact structures, whereas the distance distributions that one would derive from ordinary geometry restraints and contact potentials would allow both compact and extended structures. The difference is most likely due to entropic effects: protein configurational entropy and the hydrophobic "interaction." Further studies will be required to determine the relative contributions of entropic effects and geometrical restraints in giving rise to the features in the PDFs. The PDF potential could also enhance refinement by its rerepresentation of local geometrical restraints in terms of central forces between unbonded atom pairs. The many pairwise potentials could result in a more efficient search of configuration space than is possible when the restraints are represented using only local geometry parameters and a VDW nonbonded interaction. The many multimodal distance distributions that determine the forces on a single atom could potentially add up to an energy surface that is relatively free of local traps, allowing a complex downhill trajectory through configuration space. The repulsive component of the PDF potential is also relatively soft, so that conformations that would be frustrated by VDW contacts would be more plastic when PDF is used.

Finally, we note that there are many alternative ways of choosing structures to use in generating the PDFs (e.g., the 100 highresolution structures without VDW clashes described in Word et al., 1999). Many of these may provide additional information valuable for structure determination. Further studies will be required to determine which factors are most important in enhancing structure refinement.

In summary, our results support an "Aufbau principle" of macromolecular structure (Schutt, 1987). In this principle, the properties of larger structures are induced from those of smaller molecules that have been carefully described. A caveat is that new structures that depend on this information should not be used in generating future database potentials. By making fuller use of the information in structures that have already been solved, however, methods that use the new database potentials will likely aid in the structure solution and modeling of molecular complexes for which data are scarce.

Methods

The PDF energy function is derived from a previously described knowledge-based interaction potential for proteins (Sippl, 1990; Subramaniam et al., 1996). A probability is constructed by assuming that each atom pair contributes independently according to spatial separation. The total likelihood of a structure is equal to the product of all of the pairwise probabilities.

To implement this, one wants the precise probability density functions (PDFs) that describe the distances between atoms in proteins (i.e., statistical mechanical pair distribution functions). Although these PDFs are not known, they can be estimated using an ensemble of known protein structures.

Classification of atom pairs for the distributions is critical. For our PDF database, an ordered pair of atom types plus the number of peptide bonds *m* along the bonding path connecting the atoms is used. A value m = 0 corresponds to the atoms on the same residue, and any value m > 4 (4 = alpha-helical repeat) is of one type, termed "tertiary." Each nonhydrogen atom on 21 residue types (20 plus disulfide-bonded cysteine) counts as 1 of 173 different atom types. This leads to distance probability distributions between roughly 150,000 types of atom pairs. Scoring systems based on these and similar distributions already have shown promise for validating protein structures (Sippl, 1993; Rojnuckarin & Subramaniam, 1999).

To obtain the PDFs, structures of 2.5 Å or higher resolution in the Brookhaven Protein Data Bank (PDB) (Bernstein et al., 1977) were sampled for interatomic distances, and the results were smoothed. Only a single structure was selected among structures with more than 25% sequence identity. Selection criteria for this structure were (1) high resolution, (2) absence of prosthetic groups, and (3) low *R*-factors in order of importance. Four hundred thirty-two polypeptide chains from 392 structures were analyzed to generate the distributions.

Selection of the kernel width for smoothing is critical, as too small a width will introduce artifacts from noise, while too large a width will smear important features. Here, the optimal kernel width was determined from the data obtained for each atom pair (Silverman, 1982; Rojnuckarin & Subramaniam, 1999). Distributions were stored as arrays of 140 values of the probability density at discrete interatomic distances, along with a start distance and bin size to generate the ordinate values.

A stand-alone program called SOESA (standing for <u>Structure</u> <u>Optimization and Evaluation using Separations of Atoms; distribution through the internet at www.bioc.rice.edu/soesa) was written to calculate the total PDF energy and its derivatives with respect to atomic positions. This program runs in the UNIX background as a "server" and processes requests for calculations from "client" processes, leaving the answers in a file. Support was added for using the information in the output file for refinement in X-PLOR (Brünger, 1993) (see below), CNS (Brünger et al., 1998) (still to be completed), and TNT (Tronrud, 1997).</u>

The program SOESA calculates the PDF energy by summing contributions $E_{ij}(s)$ from each unique atom pair (i,j), where *s* is the interatomic distance. The values $E_{ij}(s)$ are calculated by spline interpolation of the transformed arrays $E_{ij,n(s)} = -W \ln P_{ij,n(s)}$, where $P_{ij,n(s)}$ is the probability distribution between atoms *i* and *j*, and n(s) is the bin number into which the distance *s* falls in the distribution. Thus, the total energy bears resemblance to an expression of maximum likelihood, where the product of the values $P_{ij}(s)$ is analogous to a prior probability of the interatomic distances. By analogy with Boltzmann statistics, the energy scale *W* was defined as 0.58 kcal mol⁻¹ (= *RT* at room temperature), although in practice *W* is an arbitrary weight whose value should be optimized.

The gradient of the energy with respect to atomic positions was calculated using the analytic derivative of the spline expression used to interpolate the energy values. The gradient is used by, e.g., X-PLOR to calculate molecular dynamics forces \mathbf{F}_i on each atom *i* as $\mathbf{F}_i = -\sum_j dE_{ij}(s)/ds(\mathbf{x}_i - \mathbf{x}_j)/|\mathbf{x}_i - \mathbf{x}_j|$, where \mathbf{x}_i is the vector position of atom *i*, and the sum \sum_j is carried out over all *j* not equal to *i*.

X-PLOR version 3.851 was modified to read energies and derivatives from a file and use them for refinement. Also added was a C routine to (1) write the current atomic coordinates to a PDB file; (2) send a signal to a UNIX process to initiate calculations; and (3) wait for a signal indicating completion. Control was implemented through the USER potential, allowing full use of CONStraints INTERactions statements for selecting atom pairs participating in the database potential, and setting the relative weights of their contributions.

The fact that the PDF energies and gradients are calculated directly from the distance distributions makes this strictly the first true example of a database potential used in structure determination. This is so because the information in the distributions is not parameterized or reduced in any way once estimated. By comparison, for instance, standard geometry restraints reduce the information to mean distances and angles and their standard deviations.

We found that methods such as steepest descent and molecular dynamics, especially when combined with simulated annealing algorithms, efficiently decreased the PDF energy. Methods that use curvature information, however, such as Powell minimization, did not work well. This is due to the multimodal nature of the distance distributions and the amplification of noise caused by differentiation, both of which give rise to sign changes in the curvature.

Test conditions

We used real data from experiments on solved structures for our tests. Artificial X-ray data nearly always agree poorly with low-resolution experimental data, and the tests of low-resolution refinement were critical to our study. The structure and X-ray data from crystallographic experiments on a 1.3 Å resolution wild-type myoglobin (Romo, 1998) and a 2.0 Å resolution structure of calmodulin complexed with the calmodulin-binding domain of brain calmodulin-dependent protein kinase II-alpha (Wall et al., 1997) (PDB entry 1cm1) were used. Linker residues 74–83, which lie in a region of no connected electron density, were deleted from the calmodulin complex.

Homologous myoglobin (PDB entry 1mbd) and calmodulin (PDB entry 1osa) coordinates are in the ensemble of reference structures. These structures account for only a small fraction of the total data used to compile the PDFs, however, and therefore would not significantly bias refinement. In addition, the calmodulin entry is in the open, unbound form, which has major differences with the closed, bound form used in this study.

Test conditions were chosen to mimic early refinement, where the backbone is fairly well defined but the side chains are not yet determined. Starting structures for the tests were generated by first using the program CHAIN (Sack, 1988) to select alternate possible side-chain conformations for one-third of the residues, attempting to minimize changes to all but the χ^1 dihedral angle. Changes preserved good local geometry, but were poor in the context of the rest of the protein. The temperature factors for all atoms in the myoglobin structure were set to 10 Å, and those for the calmodulin complex were set to 20 Å. One refinement cycle (described below) by standard methods using X-ray data truncated at 3.5 Å resolution was then run to generate a starting structure with scrambled side chains and displaced backbone atoms.

Parallel refinements in X-PLOR were run using three methods to try to improve the starting structure. In each method, geometry restraints (BOND, ANGL, DIHE, and IMPR flags) were used. For "NORM" refinement, the VDW potential was used. For "PDF" refinement, the VDW potential was turned off and replaced by the USER (i.e., PDF) potential. To avoid overweighting local geometry restraints, the m = 0 PDFs were not used. Torsion-angle refinement ("TORS") was also carried out using the same potentials as with standard refinement.

To simulate limited available X-ray data, refinement was done using reduced data sets: a low-resolution data set truncated at 3.5 Å and a high-resolution data set truncated at 2.0 Å.

Standard methods were used [see X-PLOR manual (Brünger, 1993)] to scale the contribution of the X-ray target function to the total "energy," dividing the recommended Wa by 3 for all refinements. For each case, two additional refinements were carried out with the recommended weight both doubled $(2.0\times)$ and halved $(0.5\times)$.

For myoglobin, all heme atoms were kept fixed during refinement. For calmodulin, all calcium ions were kept fixed.

Each full round of NORM and PDF refinement consisted of a segment of simulated annealing (300 K final temperature, 25 K step, 25 cycles per step) bracketed by Powell minimizations (120

steps each). Four sequential rounds of refinement were carried out in each case, and the best R_{free} structure was chosen for comparison.

For NORM refinement, starting temperatures of both 2,000 and 4,000 K were used. At 3.5 Å resolution, better results were obtained using a starting temperature of 2,000 K, while at 2.0 Å resolution, better results were obtained using a starting temperature of 4,000 K. For PDF refinement, only a starting temperature of 4,000 K was used.

For the TORS method, four parallel refinements were done, using starting temperatures of both 2,000 and 4,000 K. For 3.5 Å refinement, a starting temperature of 2,000 K gave the best results, and for 2.0 Å resolution refinement, a starting temperature of 4,000 K gave the best results; the sample X-PLOR script was used with minimal editing.

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