## SI Appendix

## **Derivation of the model**

The current-balance equation took the form

C dV/dt = 
$$-I_{Leak}(V) - I_{Na}(V,m,h) - I_{K}(V,n) - I_{CAN}(V,Ca) - I_{syn}(V,s_1 \dots s_N) - I_{pump}(Na) - I_{noise}$$

where

$$dx/dt = (x_{\infty}(V) - x) / T_{X}(V)$$

$$ds/dt = ((1 - s) s_{\infty}(V) - k_{s} s) / T_{s}$$

$$dCa/dt = \epsilon (s k_{IP3} - k_{Ca} (Ca - Ca_{\infty}))$$

$$dNa/dt = \alpha (-I_{CAN}(V,Ca) - I_{pump}(Na))$$

describe the evolution of the state variables, for each x in  $\{m, h, n\}$ .

Membrane currents were described with chord-conductance equations, in some cases modified for  $Ca^{2+}$  or  $Na^{+}$  gating ( $I_{CAN}$  and  $I_{pump}$ ), with electrogenic pumps ( $I_{pump}$ ) adapted from Li et al. (1):

$$\begin{split} I_{\text{Leak}}(V) &= g_{\text{Leak}} \; (V - E_{\text{L}}) \\ I_{\text{Na}}(V, m, h) &= g_{\text{Na}} \; m^3 \; h \; (V - E_{\text{Na}}) \\ I_{\text{K}}(V, n) &= g_{\text{K}} \; n^4 \; (V - E_{\text{K}}) \\ I_{\text{CAN}}(V, \text{Ca}) &= g_{\text{CAN}} \; (V - E_{\text{CAN}}) \; / \; (1 + \exp((\text{Ca} - k_{\text{CAN}}) / \sigma_{\text{CAN}})) \\ I_{\text{syn}}(V, s_1 \ldots s_N) &= g_{\text{syn}} \sum_{i=1}^N s_i (V - E_{\text{syn}}) \end{split}$$

where N is the number of presynaptic neurons and  $\{s_1 \dots s_N\}$  is the set of presynaptic s variables, and

$$I_{\text{pump}}(\text{Na}) = r_{\text{pump}} (\Phi(\text{Na}) - \Phi(\text{Na}_{\infty}))$$
  
 $I_{\text{noise}} = \xi(t)$ 

where  $\xi(t)$  is a Gaussian noise term with an average of zero and gamma ( $\gamma$ ) as an adjustable parameter controlling the magnitude of the fluctuations,  $\langle \xi(t) \rangle = 0$  pA,  $\langle \xi(t) \xi(t') \rangle = \gamma \cdot \delta(t-t')$ .

The remaining functions in the model, including those representing the voltagedependence of channel kinetics, were

$$x_{\infty}(V) = 1 / (1 + \exp((V - \theta_{x})/\sigma_{x}))$$
  
 $\tau_{x}(V) = \tau_{x-max} / \cosh((V - \theta_{x}) / (2 \sigma_{x}))$   
 $\Phi(Na) = Na^{3} / (Na^{3} + k_{Na}^{3}).$ 

Model parameters were set to the following values, unless otherwise specified: C = 45 pF,  $g_{leak}$  = 3 nS,  $E_L$  = -60 mV,  $g_{Na}$  = 150 nS,  $E_{Na}$  = 85 mV,  $g_K$  = 30 nS,  $E_K$  = -75 mV,  $g_{CAN}$  = 4 nS,  $E_{CAN}$  = 0 mV,  $g_{syn}$  = 2.5 nS,  $E_{syn}$  = 0 mV,  $\theta_m$  = -36 mV,  $\sigma_m$  = -8.5 mV,  $\tau_{m-max}$  = 1 ms,  $\theta_h$  = -30 mV,  $\sigma_h$  = 5 mV,  $\tau_{h-max}$  = 15 ms,  $\theta_n$  = -30 mV,  $\sigma_n$  = -5 mV,  $\tau_{n-max}$  = 30 ms,  $\theta_s$  = 15 mV,  $\sigma_s$  = -3 mV,  $\tau_s$  = 15 ms,  $\tau_s$  = 15 ms,  $\tau_s$  = 15 ms,  $\tau_s$  = 15 ms,  $\tau_s$  = 10 mM,  $\tau$ 

Calcium currents and synaptic inputs. Voltage-dependent  $Ca^{2+}$  currents are present in the preBötC and contribute to burst generation (2, 3) but were omitted from the current-balance equation for the following reasons.  $Ca^{2+}$  currents are predominantly activated by AMPA receptor-medicated depolarization during inspiratory burst generation (4), so instead of deriving an explicit  $Ca^{2+}$  current that would respond specifically to synaptic depolarization, we coupled the synaptic variable s to the  $Ca^{2+}$  equation directly. The synaptic variable s therefore represents both ionotropic and metabotropic glutamatergic receptor (mGluR) activation: AMPA receptors cause  $Ca^{2+}$  changes via depolarization, whereas group I mGluRs elevate intracellular  $Ca^{2+}$  by evoking intracellular release from stores, which both contribute to  $Ca^{2+}$  increases that couple to and activate  $I_{CAN}$  during inspiratory burst generation (4). The variable s also appears in  $I_{syn}$  in the standard way, and the parameters  $k_{ip3}$  and  $g_{syn}$  scale its relative contributions to  $Ca^{2+}$  dynamics and to  $I_{syn}$ , respectively.

**Parameter values**. Synaptic coupling affects  $Ca^{2+}$  accumulation and  $I_{CAN}$  activation, which collectively underlie burst generation. Consequently, burst duration and the interburst interval depend on the parameter values for  $g_{CAN}$ ,  $k_{IP3}$ ,  $g_{syn}$ ,  $r_{pump}$ , and  $k_{Ca}$  when  $I_{pump}$  is present. Rhythmic solutions in the model persisted over substantial variation in these parameters from nominal baseline values. Regulating these key parameters supported a wide range of interburst intervals and burst durations (Fig. S2). This robustness and capacity to generate bursts with varying characteristics support the relevance of this model to respiratory rhythmogenesis, where modulation and regulation of the basic rhythm is essential for the physiology of respiration (5-8).

**Synaptic delay**. In simulations of the self-coupled cell, we introduced a latency  $T_{\text{delay}}$ , with  $s_{\infty} = 1 / (1 + \exp((V(t - T_{\text{delay}}) - \theta_s)/\sigma_s))$ , to mimic the delay in synaptic coupling inherently present in the multi-cell simulations due to spike asynchrony (9). We found qualitatively similar dynamics for  $T_{\text{delay}}$  from 0 up to 19 ms. In general, we used  $T_{\text{delay}} = 6$  ms to generate time courses in the self-coupled case (Figs. 3,4,S2,S3).

<u>Activity-dependent outward currents</u>. In some cases,  $I_{pump}$  was replaced by alternative net outward currents, as follows (see Figs. 4,S3,S4).

1) Slowly-activating M-like K<sup>+</sup> current (I<sub>M</sub>):

$$I_{M}(V, n_{M}) = g_{M} n_{M} (V - E_{K})$$

$$dn_{M}/dt = (n_{M^{\infty}}(V) - n_{M}) / \tau_{nM}(V)$$

$$\tau_{nM}(V) = \tau_{nM-max} / \cosh((V - \theta_{n}) / (2 \sigma_{n}))$$

where  $E_L = -57.7$  mV,  $\gamma = 600$  pA<sup>2</sup>,  $g_M = 2$  nS,  $\tau_{nM-max} = 1000$  ms and  $\theta_n = -30$  mV,  $\sigma_n = -5$  mV, as above, was based on a delayed-rectifier current in a preBötC model by Butera and colleagues (10, 11) but with slower kinetics.

2) Ca<sup>2+</sup>-dependent K<sup>+</sup> current (I<sub>K-Ca</sub>):

$$I_{\text{K-Ca}}(V, z_{\text{K-Ca}}) = g_{\text{K-Ca}} z^2_{\text{K-Ca}} (V - E_{\text{K}})$$
  
 $dz_{\text{K-Ca}}/dt = (z_{\infty}(Ca) - z_{\text{K-Ca}}) / \tau_{\text{K-Ca}}$   
 $z_{\infty}(Ca) = 1 / (1 + (k_{\text{K-Ca}}/Ca)^4)$ 

where  $E_L = -60.15$  mV,  $\gamma = \frac{500}{100}$  pA<sup>2</sup>,  $g_{K-Ca} = 3$  nS,  $k_{K-Ca} = 0.9$   $\mu$ M,  $\tau_{K-Ca} = 100$  ms, was based on biophysical modeling and general properties of  $I_{K-Ca}$  (12, 13).

3) Persistent Na<sup>+</sup> current ( $I_{Na-P}$ ) with slow inactivation:

$$\begin{split} I_{\text{Na-P}}(V, \, h_{\text{Na-P}}) &= g_{\text{NaP}} \, m_{\text{NaP}}(V) \, h_{\,\text{Na-P}} \, (V - E_{\text{Na}}) \\ dh_{\text{Na-P}}/dt &= (h_{\text{Na-P-}\infty}(V) - h_{\text{Na-P}}) \, / \, \tau_{\text{Na-P}} \\ m_{\text{Na-P-}\infty}(V) &= 1 \, / \, (1 + \exp((V - \theta_{\text{m-Na-P}}) / \sigma_{\text{m-Na-P}})) \\ h_{\text{Na-P-}\infty}(V) &= 1 \, / \, (1 + \exp((V - \theta_{\text{h-Na-P}}) / \sigma_{\text{h-Na-P}})) \end{split}$$

where  $E_L = -61.25$  mV,  $\gamma = \frac{500}{N}$  pA<sup>2</sup>,  $g_{Na-P} = 1$  nS,  $\theta_{m-Na-P} = -40$  mV,  $\sigma_{mNaP} = -6$  mV,  $\theta_{hNaP} = -48$  mV,  $\sigma_{h-Na-P} = 6$  mV, and  $\tau_{Na-P} = 1000$  ms, was based on the preBötC model by Butera et al. (10, 11).

When any of these outward currents was included, rhythmic activity in the model was qualitatively similar to that generated with  $I_{pump}$  (Figs. 3,4,S3). In particular, the self-coupled single-cell model and the two-cell model produced repetitive activity patterns featuring spontaneous spiking prior to the burst, followed immediately by a higher-frequency spike rate during the burst, spike attenuation within each burst (i.e., depolarization block), and a transient quiescent period before low-rate pre-burst spiking resumed. In all cases except the model with  $I_{\text{Na-P}}$ , both  $\text{Ca}^{2+}$  and a slow activation variable associated with the additional outward current accumulated during the burst. In the case of  $I_{\text{Na-P}}$ ,  $\text{Ca}^{2+}$  accumulation was accompanied by a decrease in the gating variable  $h_{\text{NaP}}$ , which caused  $I_{\text{Na-P}}$  inactivation and functioned analogously to activation of an outward current (Fig. 4C,S3D). No matter which current was used, the rhythmic activity depended crucially on synaptic interactions and  $I_{CAN}$ . The only significant difference observed among the currents considered was that with  $I_{pump}$ , bursting could occur even without including synaptic transmission latency, while the other currents required a small latency ( $T_{delay}$  above a small positive lower bound) to cause tonic spiking to ramp up the synaptic response, thereby promoting sufficient recurrent excitation to elicit bursting (Fig. S3B-D).

We also simulated 200-neuron networks (Fig. S4). For the  $I_{pump}$  simulations,  $E_L$  was drawn from a Gaussian (normal) distribution with a mean of -60.0 and standard deviation of 1 mV (i.e.,  $-60.0\pm1$  mV). For the other net outward currents, the  $E_L$  distribution had the following means and standard deviations:  $-57.3\pm1$  mV (Fig. S4B),  $-60.1\pm1$  mV (Fig. S4C), and  $-61.\pm1$  mV (Fig. S4D). The differences in means were selected to compensate for the different tonic components of the outward currents ( $I_M$ ,  $I_{K-Ca}$ , and  $I_{NaP}$  inactivation) that affected baseline membrane potentials in each case.

**Caveats and limitations**. Our preBötC-like model neuron lacks transient A-type K<sup>+</sup> current ( $I_A$ ) and hyperpolarization-activated mixed cationic current ( $I_h$ ).  $I_A$  influences the recruitment pattern of preBötC neurons during recurrent excitation in the 300-1000 ms prior to the inspiratory burst (14).  $I_h$  modulates respiratory frequency *in vitro* and influences the baseline membrane potential trajectory (15, 16). Our model predicts that the omitted currents may not be essential for burst generation, but rather influence inspiratory burst characteristics and regulate respiratory frequency (2-5).

The most problematic simplification may be that the model is limited to a single isopotential compartment in which synaptic integration and spike generation occur explicitly in the same location. Recent reports suggest that synaptic recruitment of  $I_{CAN}$  may occur principally in the dendrites (4, 17, 18), which is likely the primary site of Na/K pump activity, whereas most sodium influx via  $I_{Na}$  occurs in the axon hillock. These observations are reflected in our Na<sup>+</sup> equation, which includes dendritic Na<sup>+</sup> currents yet neglects  $I_{Na}$ , since diffusion of sodium from the axon hillock to the dendrites would be minimal; our model also generates qualitatively similar bursting when up to 10% of the Na<sup>+</sup> influx via  $I_{Na}$  is included in the dNa/dt equation. Despite this complication, maintaining a single-compartment model has a number of important advantages: i) the single-compartment construction facilitates bifurcation analyses; ii) because its morphology is not specific to preBötC neurons, the model provides a more general template to begin exploring group-pacemaker activity by modifying the model to match other neuronal

phenotypes found in CPGs in other parts of the brain stem and spinal cord; and iii) the simplicity of the present model elucidates the general principles underlying a nevel group pacemaker mechanism for rhythmogenesis applicable to CPGs, yet retains its connection to inspiratory rhythm generation in the preBötC, which motivated our search for new rhythmogenic mechanisms.

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