

Lévy laws and $1/f$ noises:
A unified and universal explanation

Supporting Information

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Throughout the supplementary material the following notation will be used: $\mathbf{E}[\xi]$ is the expectation of the random variable ξ ; $\mathbf{Cov}[\xi_1, \xi_2]$ is the covariance of the random variables ξ_1 and ξ_2 .

1 The stationary superposition model: Proofs

In this Section we set:

$$\mathcal{P} = \{(a_k, \omega_k)\}_k ; \quad (1)$$

$$G_X(\theta) = 1 - \mathbf{E}[\exp(i\theta X(0))] \quad (2)$$

(θ real);

$$R_X(t) = \mathbf{Cov}[X(0), X(t)] \quad (3)$$

(t real);

$$\widehat{R}_X(\theta) = \int_{-\infty}^{\infty} R_X(t) \exp(i\theta t) dt \quad (4)$$

(θ real). Recall that $\phi(x) = \int_0^{\infty} \lambda(x, y) dy$ (x real) and $\psi(y) = \int_{-\infty}^{\infty} x^2 \lambda(x, y) dx$ ($y > 0$), and that the signal pattern X is a zero mean stationary process with short-range correlations.

1.1 Proof of amplitudal universality

Step 1. We compute the conditional Fourier transform of the random variable $Y(t)$ (t being an arbitrary time point), given the amplitude-frequency pairs \mathcal{P} :

$$\mathbf{E}[\exp(i\theta Y(t)) \mid \mathcal{P}] \quad (5)$$

(using the definition of the process Y)

$$= \mathbf{E} \left[\prod_k \exp(i(\theta a_k) X_k(\omega_k t)) \mid \mathcal{P} \right] \quad (6)$$

(using the independence of the transmission sources)

$$= \prod_k \mathbf{E} \left[\exp(i(\theta a_k) X_k(\omega_k t)) \mid \mathcal{P} \right] \quad (7)$$

(using the fact that the signal patterns transmitted by the sources are i.i.d. copies of the generic stationary signal pattern X)

$$= \prod_k \mathbf{E} \left[\exp(i(\theta a_k) X(0)) \right] \quad (8)$$

(using the definition of the function $G_X(\theta)$)

$$= \prod_k \left(1 - G_X(\theta a_k)\right) . \quad (9)$$

Step 2. We compute the Fourier transform of the random variable $Y(t)$ (t being an arbitrary time point):

$$\mathbf{E} [\exp (i\theta Y(t))] \quad (10)$$

(using conditioning)

$$= \mathbf{E} \left[\mathbf{E} [\exp (i\theta Y(t)) \mid \mathcal{P}] \right] \quad (11)$$

(using Step 1)

$$= \mathbf{E} \left[\prod_k \left(1 - G_X(\theta a_k)\right) \right] \quad (12)$$

(using equation (3.35) in [1])

$$= \exp \left(- \int_{-\infty}^{\infty} \int_0^{\infty} G_X(\theta x) \lambda(x, y) dx dy \right) \quad (13)$$

(using the definition of the function $\phi(x)$)

$$= \exp \left(- \int_{-\infty}^{\infty} G_X(\theta x) \phi(x) dx \right) \quad (14)$$

(using the change of variables $u = \theta x$)

$$= \exp \left(- \frac{1}{|\theta|} \int_{-\infty}^{\infty} G_X(u) \phi\left(\frac{u}{\theta}\right) du \right) . \quad (15)$$

Step 3. Equation (15) implies that the random variable $Y(t)$ is independent of the signal pattern X – up to a scale factor – if and only if the function $\phi(x)$ is a power-law. Specifically, if $\phi(x) = c_1|x|^{-1-\alpha}$ then

$$\mathbf{E} \left[\exp (i\theta Y(t)) \right] = \exp \left(-c_2|\theta|^\alpha \right) , \quad (16)$$

where

$$c_2 = c_1 \int_{-\infty}^{\infty} \frac{G_X(u)}{|u|^{1+\alpha}} du . \quad (17)$$

The function $G_X(u)$ satisfies $|G_X(u)| \leq 2$ and $G_X(u) \sim (\frac{1}{2}R_X(0))u^2$ as $u \rightarrow 0$. Hence, the integral appearing on the right hand side of equation (17) is convergent in the exponent range $0 < \alpha < 2$.

1.2 Proof of temporal universality

Step 1. We compute the conditional mean of the random variable $Y(t)$ (t being an arbitrary time point), given the amplitude-frequency pairs \mathcal{P} :

$$\mathbf{E}[Y(t) \mid \mathcal{P}] \quad (18)$$

(using the definition of the process Y)

$$= \mathbf{E} \left[\sum_k a_k X_k(\omega_k t) \mid \mathcal{P} \right] \quad (19)$$

$$= \sum_k a_k \mathbf{E}[X_k(\omega_k t) \mid \mathcal{P}] \quad (20)$$

(using the fact that the signal patterns transmitted by the sources are i.i.d. copies of a generic stationary signal pattern X with zero mean)

$$= \sum_k a_k \mathbf{E}[X(0)] = 0. \quad (21)$$

Step 2. We compute the conditional covariance of the random variables $Y(t)$ and $Y(t + \Delta)$ (t being an arbitrary time point; Δ being an arbitrary lag), given the amplitude-frequency pairs \mathcal{P} :

$$\mathbf{Cov}[Y(t), Y(t + \Delta) \mid \mathcal{P}] \quad (22)$$

(using the definition of the process Y)

$$= \mathbf{Cov} \left[\sum_k a_k X_k(\omega_k t), \sum_j a_j X_j(\omega_j(t + \Delta)) \mid \mathcal{P} \right] \quad (23)$$

$$= \sum_k \sum_j a_k a_j \mathbf{Cov}[X_k(\omega_k t), X_j(\omega_j(t + \Delta)) \mid \mathcal{P}]$$

(using the independence of the transmission sources)

$$= \sum_k a_k^2 \mathbf{Cov}[X_k(\omega_k t), X_k(\omega_k(t + \Delta)) \mid \mathcal{P}] \quad (24)$$

(using the fact that the signal patterns transmitted by the sources are i.i.d. copies of the generic stationary signal pattern X)

$$= \sum_k a_k^2 \mathbf{Cov}[X(0), X(\omega_k \Delta)] \quad (25)$$

(using the definition of the function $R_X(t)$)

$$= \sum_k a_k^2 R_X(\omega_k \Delta). \quad (26)$$

Step 3. We compute the covariance of the of the random variables $Y(t)$ and $Y(t + \Delta)$ (t being an arbitrary time point; Δ being an arbitrary lag):

$$\mathbf{Cov} [Y(t), Y(t + \Delta)] \quad (27)$$

(using conditioning)

$$= \mathbf{Cov} \left[\mathbf{E} [Y(t) \mid \mathcal{P}], \mathbf{E} [Y(t + \Delta) \mid \mathcal{P}] \right] + \mathbf{E} \left[\mathbf{Cov} [Y(t), Y(t + \Delta) \mid \mathcal{P}] \right] \quad (28)$$

(using Steps 1 and 2)

$$= \mathbf{Cov} [0, 0] + \mathbf{E} \left[\sum_k a_k^2 R_X(\omega_k \Delta) \right] \quad (29)$$

(using equation (3.9) in [1])

$$= \int_{-\infty}^{\infty} \int_0^{\infty} x^2 R_X(y\Delta) \lambda(x, y) dx dy \quad (30)$$

(using the definition of the function $\psi(y)$)

$$= \int_0^{\infty} R_X(y\Delta) \psi(y) dy . \quad (31)$$

Step 4. We compute the power spectrum of the output process Y :

$$\int_{-\infty}^{\infty} \mathbf{Cov} [Y(0), Y(\Delta)] \exp(iff\Delta) d\Delta \quad (32)$$

(using Step 3)

$$\begin{aligned} &= \int_{-\infty}^{\infty} \left(\int_0^{\infty} R_X(y\Delta) \psi(y) dy \right) \exp(iff\Delta) d\Delta \\ &= \int_0^{\infty} \left(\int_{-\infty}^{\infty} R_X(y\Delta) \exp(iff\Delta) d\Delta \right) \psi(y) dy \end{aligned} \quad (33)$$

(using the change of variables $t = y\Delta$)

$$= \int_0^{\infty} \left(\frac{1}{y} \int_{-\infty}^{\infty} R_X(t) \exp\left(i\frac{f}{y}t\right) d\Delta \right) \psi(y) dy \quad (34)$$

(using the definition of the function $\widehat{R}_X(\theta)$)

$$= \int_0^{\infty} \left(\frac{1}{y} \widehat{R}_X\left(\frac{f}{y}\right) \right) \psi(y) dy \quad (35)$$

(using the change of variables $u = f/y$)

$$= \int_0^{\infty} \widehat{R}_X(u) \left(\frac{1}{u} \psi\left(\frac{|f|}{u}\right) \right) du . \quad (36)$$

Step 5. Equation (36) implies that the power spectrum of the output process Y is independent of the signal pattern X – up to a scale factor – if and only if the function $\psi(y)$ is a power-law. Specifically, if $\psi(y) = c_3 y^{-\beta}$ then

$$\int_{-\infty}^{\infty} \mathbf{Cov} [Y(0), Y(t)] \exp(ift) dt = \frac{c_4}{|f|^\beta}, \quad (37)$$

where

$$c_4 = c_3 \int_0^{\infty} \frac{\widehat{R}_X(u)}{u^{1-\beta}} du. \quad (38)$$

And, plugging $\psi(y) = c_3 y^{-\beta}$ into equation (31) yields

$$\mathbf{Cov} [Y(t), Y(t + \Delta)] = \frac{c_6}{\Delta^{1-\beta}}, \quad (39)$$

where

$$c_6 = c_3 \int_0^{\infty} \frac{R_X(u)}{u^\beta} du. \quad (40)$$

Since the signal pattern X has short-range correlations, the integral $\int_0^{\infty} R_X(u) du$ is finite, and the function $\widehat{R}_X(u)$ satisfies $\widehat{R}_X(0) = \int_{-\infty}^{\infty} R_X(t) dt$ (finite) and $\int_0^{\infty} \widehat{R}_X(u) du = \pi R_X(0)$ (finite). Hence, the integral appearing on the right hand side of equation (38) is convergent in the exponent range $0 < \beta \leq 1$, and the integral appearing on the right hand side of equation (40) is convergent in the exponent range $0 < \beta < 1$. The admissible exponent range is thus $0 < \beta < 1$.

2 From $1/f$ noise to super diffusion: Proof

In this Section we set:

$$R_Y(t) = \mathbf{Cov} [Y(0), Y(t)] \quad (41)$$

(t real);

$$\widehat{R}_Y(\theta) = \int_{-\infty}^{\infty} R_Y(t) \exp(i\theta t) dt \quad (42)$$

(θ real). We prove that if the stationary process Y has a $1/f$ power spectrum then the integrated process Z – given by $Z(t) = \int_0^t Y(t') dt'$ ($t \geq 0$) – is super-diffusive.

Step 1. We compute the mean square displacement of the integrated process Z , in terms of the power spectrum $\widehat{R}_Y(\theta)$ of the process Y :

$$\mathbf{E} [Z(t)^2] \quad (43)$$

(using the definition of the integrated process Z)

$$\begin{aligned} &= \mathbf{E} \left[\left(\int_0^t Y(t_1) dt_1 \right) \left(\int_0^t Y(t_2) dt_2 \right) \right] \\ &= \int_0^t \int_0^t \mathbf{E} [Y(t_1) Y(t_2)] dt_1 dt_2 \end{aligned} \quad (44)$$

(using the fact that the stationary process Y has zero mean)

$$= \int_0^t \int_0^t \mathbf{Cov}[Y(t_1), Y(t_2)] dt_1 dt_2 \quad (45)$$

(using the definition of the function $R_Y(t)$)

$$= \int_0^t \int_0^t R_Y(t_2 - t_1) dt_1 dt_2 \quad (46)$$

(using Fourier inversion)

$$\begin{aligned} &= \int_0^t \int_0^t \left(\frac{1}{2\pi} \int_{-\infty}^{\infty} \widehat{R}_Y(\theta) \exp(-i(t_2 - t_1)\theta) d\theta \right) dt_1 dt_2 \\ &= \frac{1}{2\pi} \int_{-\infty}^{\infty} \widehat{R}_Y(\theta) \left(\int_0^t \exp(it_1\theta) dt_1 \right) \left(\int_0^t \exp(-it_2\theta) dt_2 \right) d\theta \\ &= \frac{1}{2\pi} \int_{-\infty}^{\infty} \widehat{R}_Y(\theta) \left(\frac{\exp(it\theta) - 1}{i\theta} \right) \left(\frac{\exp(-it\theta) - 1}{-i\theta} \right) d\theta \\ &= \frac{1}{\pi} \int_{-\infty}^{\infty} \widehat{R}_Y(\theta) \frac{1 - \cos(t\theta)}{\theta^2} d\theta \end{aligned} \quad (47)$$

(using the change of variables $u = t\theta$)

$$= \frac{t}{\pi} \int_{-\infty}^{\infty} \widehat{R}_Y\left(\frac{u}{t}\right) \frac{1 - \cos(u)}{u^2} du . \quad (48)$$

Step 2. If the stationary process Y is a $1/f$ noise then

$$\widehat{R}_Y(\theta) = \frac{c_4}{|\theta|^\beta} . \quad (49)$$

Plugging the $1/f$ power spectrum of equation (49) into equation (48) yields

$$\mathbf{E}[Z(t)^2] = c_5 t^{1+\beta} , \quad (50)$$

where

$$c_5 = c_4 \frac{2}{\pi} \int_0^{\infty} \frac{1 - \cos(u)}{u^{2+\beta}} du . \quad (51)$$

The integral appearing on the right hand side of equation (51) is convergent in the exponent range $0 < \beta < 1$.

3 The dissipative superposition model: Proofs

In this Section we set:

$$\mathcal{P} = \{(\tau_k, a_k, \omega_k)\}_k ; \quad (52)$$

$$G_X(\theta) = \int_0^{\infty} \left(1 - \mathbf{E}[\exp(i\theta X(t))] \right) dt \quad (53)$$

(θ real);

$$R_X(t) = \int_0^\infty \mathbf{Cov}[X(u), X(u+|t|)] du \quad (54)$$

(t real);

$$\widehat{R}_X(\theta) = \int_{-\infty}^\infty R_X(t) \exp(i\theta t) dt \quad (55)$$

(θ real). Recall that $\phi(x) = \int_0^\infty [\eta(x,y)/y] dy$ (x real) and $\psi(y) = \int_{-\infty}^\infty x^2 [\eta(x,y)/y] dx$ ($y > 0$), and that signal pattern X is a zero mean dissipative process. We require the following conditions regarding the stochastic decay of the process X to zero:

Condition 1 *The function $G_X(\theta)$ is bounded.*

Condition 2 *The function $R_X(t)$ is finite at the origin ($t = 0$), and is integrable over the real line.*

An important case for which the function $G_X(\theta)$ is bounded is the following: There exists a random time T_X – with finite mean – above which the dissipative signal pattern X vanishes. Indeed, in the aforementioned case we have ($I\{E\}$ denoting the indicator function of the event E):

$$\begin{aligned} |G_X(\theta)| &\leq \int_0^\infty \mathbf{E}[|1 - \exp(i\theta X(t))|] dt \\ &= \int_0^\infty \mathbf{E}[|1 - \exp(i\theta X(t))| \cdot I\{t < T_X\}] dt \\ &\leq \int_0^\infty \mathbf{E}[2 \cdot I\{t < T_X\}] dt = 2 \int_0^\infty \mathbf{P}(T_X > t) dt \\ &= 2\mathbf{E}[T_X] < \infty. \end{aligned} \quad (56)$$

3.1 Proof of amplitudal universality

Step 1. We compute the conditional Fourier transform of the random variable $Y(t)$ (t being an arbitrary time point), given the initiation-amplitude-frequency triplets \mathcal{P} :

$$\mathbf{E}[\exp(i\theta Y(t)) \mid \mathcal{P}] \quad (57)$$

(using the definition of the process Y)

$$= \mathbf{E} \left[\prod_{\tau_k \leq t} \exp(i(\theta a_k) X_k(\omega_k(t - \tau_k))) \mid \mathcal{P} \right] \quad (58)$$

(using the independence of the dissipative signals)

$$= \prod_{\tau_k \leq t} \mathbf{E} \left[\exp(i(\theta a_k) X_k(\omega_k(t - \tau_k))) \mid \mathcal{P} \right] \quad (59)$$

(using the fact that the dissipative signals X_k are i.i.d. copies of the generic signal pattern X)

$$= \prod_{\tau_k \leq t} \mathbf{E} \left[\exp \left(i (\theta a_k) X (\omega_k (t - \tau_k)) \right) \right] . \quad (60)$$

Step 2. We compute the Fourier transform of the random variable $Y(t)$ (t being an arbitrary time point):

$$\mathbf{E} [\exp (i\theta Y(t))] \quad (61)$$

(using conditioning)

$$= \mathbf{E} \left[\mathbf{E} [\exp (i\theta Y(t)) \mid \mathcal{P}] \right] \quad (62)$$

(using Step 1)

$$= \mathbf{E} \left[\prod_{\tau_k \leq t} \mathbf{E} \left[\exp \left(i (\theta a_k) X (\omega_k (t - \tau_k)) \right) \right] \right] \quad (63)$$

(using equation (3.35) in [1])

$$\begin{aligned} &= \exp \left(- \int_{-\infty}^t \int_{-\infty}^{\infty} \int_0^{\infty} \left(1 - \mathbf{E} [\exp (i (\theta x) X (y (t - s)))] \right) \eta (x, y) ds dx dy \right) \\ &= \exp \left(- \int_{-\infty}^{\infty} \int_0^{\infty} \left(\int_{-\infty}^t \left(1 - \mathbf{E} [\exp (i (\theta x) X (y (t - s)))] \right) ds \right) \eta (x, y) dx dy \right) \end{aligned} \quad (64)$$

(using the change of variables $u = y (t - s)$)

$$= \exp \left(- \int_{-\infty}^{\infty} \int_0^{\infty} \left(\frac{1}{y} \int_0^{\infty} \left(1 - \mathbf{E} [\exp (i (\theta x) X (u))] \right) du \right) \eta (x, y) dx dy \right) \quad (65)$$

(using the definition of the function $G_X (\theta)$)

$$= \exp \left(- \int_{-\infty}^{\infty} \int_0^{\infty} \left(\frac{1}{y} G_X (\theta x) \right) \eta (x, y) dx dy \right) \quad (66)$$

(using the definition of the function $\phi (x)$)

$$= \exp \left(- \int_{-\infty}^{\infty} G_X (\theta x) \phi (x) dy dx \right) \quad (67)$$

(using the change of variables $u = \theta x$)

$$= \exp \left(- \frac{1}{|\theta|} \int_{-\infty}^{\infty} G_X (u) \phi \left(\frac{u}{\theta} \right) du \right) . \quad (68)$$

Step 3. Equation (68) implies that the random variable $Y(t)$ is independent of the signal pattern X – up to a scale factor – if and only if the function $\phi (x)$ is a power-law. Specifically, if $\phi (x) = c_1 |x|^{-1-\alpha}$ then

$$\mathbf{E} \left[\exp (i\theta Y(t)) \right] = \exp \left(-c_2 |\theta|^\alpha \right) , \quad (69)$$

where

$$c_2 = c_1 \int_{-\infty}^{\infty} \frac{G_X(u)}{|u|^{1+\alpha}} du . \quad (70)$$

The function $G_X(u)$ is bounded (Condition 1), and satisfies $G_X(u) \sim (\frac{1}{2}R_X(0))u^2$ as $u \rightarrow 0$ (Condition 2 assures that $R_X(0)$ is finite). Hence, the integral appearing on the right hand side of equation (69) is convergent in the exponent range $0 < \alpha < 2$.

3.2 Proof of temporal universality

Step 1. We compute the conditional mean of the random variable $Y(t)$ (t being an arbitrary time point), given the initiation-amplitude-frequency triplets \mathcal{P} :

$$\mathbf{E}[Y(t) \mid \mathcal{P}] \quad (71)$$

(using the definition of the process Y)

$$\begin{aligned} &= \mathbf{E} \left[\sum_{\tau_k \leq t} a_k X_k(\omega_k(t - \tau_k)) \mid \mathcal{P} \right] \\ &= \sum_{\tau_k \leq t} a_k \mathbf{E} [X_k(\omega_k(t - \tau_k)) \mid \mathcal{P}] \end{aligned} \quad (72)$$

(using the fact that the dissipative signals X_k are i.i.d. copies of the generic signal pattern X with zero mean)

$$= \sum_{\tau_k \leq t} a_k \mathbf{E} [X(\omega_k(t - \tau_k))] = 0 . \quad (73)$$

Step 2. We compute the conditional covariance of the random variables $Y(t)$ and $Y(t + \Delta)$ (t being an arbitrary time point; Δ being an arbitrary lag), given the initiation-amplitude-frequency triplets \mathcal{P} :

$$\mathbf{Cov} [Y(t), Y(t + \Delta) \mid \mathcal{P}] \quad (74)$$

(using the definition of the process Y)

$$\begin{aligned} &= \mathbf{Cov} \left[\sum_{\tau_k \leq t} a_k X_k(\omega_k(t - \tau_k)), \sum_{\tau_j \leq t + \Delta} a_j X_j(\omega_j(t + \Delta - \tau_j)) \mid \mathcal{P} \right] \\ &= \sum_{\tau_k \leq t} \sum_{\tau_j \leq t + \Delta} a_k a_j \mathbf{Cov} [X_k(\omega_k(t - \tau_k)), X_j(\omega_j(t + \Delta - \tau_j)) \mid \mathcal{P}] \end{aligned} \quad (75)$$

(using the independence of the dissipative signals)

$$= \sum_{\tau_k \leq t} a_k^2 \mathbf{Cov} [X_k(\omega_k(t - \tau_k)), X_k(\omega_k(t + \Delta - \tau_k)) \mid \mathcal{P}] \quad (76)$$

(using the fact that the dissipative signals are i.i.d. copies of the dissipative signal pattern X)

$$= \sum_{\tau_k \leq t} a_k^2 \mathbf{Cov} [X(\omega_k(t - \tau_k)), X(\omega_k(t + \Delta - \tau_k))] . \quad (77)$$

Step 3. We compute the covariance of the of the random variables $Y(t)$ and $Y(t + \Delta)$ (t being an arbitrary time point; Δ being an arbitrary lag):

$$\mathbf{Cov} [Y(t), Y(t + \Delta)] \quad (78)$$

(using conditioning)

$$= \mathbf{Cov} [\mathbf{E}[Y(t) | \mathcal{P}], \mathbf{E}[Y(t + \Delta) | \mathcal{P}]] + \mathbf{E} [\mathbf{Cov} [Y(t), Y(t + \Delta) | \mathcal{P}]] \quad (79)$$

(using Steps 1 and 2)

$$= \mathbf{Cov} [0, 0] + \mathbf{E} \left[\sum_{\tau_k \leq t} a_k^2 \mathbf{Cov} [X(\omega_k(t - \tau_k)), X(\omega_k(t + \Delta - \tau_k))] \right] \quad (80)$$

(using equation (3.9) in [1])

$$\begin{aligned} &= \int_{-\infty}^t \int_{-\infty}^{\infty} \int_0^{\infty} x^2 \mathbf{Cov} [X(y(t - s)), X(y(t + \Delta - s))] \lambda(x, y) ds dx dy \\ &= \int_{-\infty}^{\infty} \int_0^{\infty} \left(\int_{-\infty}^t \mathbf{Cov} [X(y(t - s)), X(y(t + \Delta - s))] ds \right) x^2 \lambda(x, y) dx dy \end{aligned} \quad (81)$$

(using the change of variables $u = y(t - s)$)

$$= \int_{-\infty}^{\infty} \int_0^{\infty} \left(\frac{1}{y} \int_0^{\infty} \mathbf{Cov} [X(u), X(u + y\Delta)] du \right) x^2 \lambda(x, y) dx dy \quad (82)$$

(using the definition of the function $R_X(t)$)

$$= \int_{-\infty}^{\infty} \int_0^{\infty} \left(\frac{1}{y} R_X(y\Delta) \right) x^2 \lambda(x, y) dx dy \quad (83)$$

(using the definition of the function $\psi(y)$)

$$= \int_0^{\infty} R_X(\Delta y) \psi(y) dy . \quad (84)$$

Step 4. We compute the power spectrum of the output process Y :

$$\int_{-\infty}^{\infty} \mathbf{Cov} [Y(0), Y(\Delta)] \exp(iff\Delta) d\Delta \quad (85)$$

(using Step 3)

$$\begin{aligned} &= \int_{-\infty}^{\infty} \left(\int_0^{\infty} R_X(y\Delta) \psi(y) dy \right) \exp(iff\Delta) d\Delta \\ &= \int_0^{\infty} \left(\int_{-\infty}^{\infty} R_X(y\Delta) \exp(iff\Delta) d\Delta \right) \psi(y) dy \end{aligned} \quad (86)$$

(using the change of variables $t = y\Delta$)

$$= \int_0^\infty \left(\frac{1}{y} \int_{-\infty}^\infty R_X(t) \exp\left(i\frac{f}{y}t\right) d\Delta \right) \psi(y) dy \quad (87)$$

(using the definition of the function $\widehat{R}_X(\theta)$)

$$= \int_0^\infty \left(\frac{1}{y} \widehat{R}_X\left(\frac{f}{y}\right) \right) \psi(y) dy \quad (88)$$

(using the change of variables $u = f/y$)

$$= \int_0^\infty \widehat{R}_X(u) \left(\frac{1}{u} \psi\left(\frac{|f|}{u}\right) \right) du . \quad (89)$$

Step 5. Equation (89) implies that the power spectrum of the output process Y is independent of the signal pattern X – up to a scale factor – if and only if the function $\psi(y)$ is a power-law. Specifically, if $\psi(y) = c_3 y^{-\beta}$ then

$$\int_{-\infty}^\infty \mathbf{Cov}[Y(0), Y(t)] \exp(ift) dt = \frac{c_4}{|f|^\beta} , \quad (90)$$

where

$$c_4 = c_3 \int_0^\infty \frac{\widehat{R}_X(u)}{u^{1-\beta}} du . \quad (91)$$

And, plugging $\psi(y) = c_3 y^{-\beta}$ into equation (84) yields

$$\mathbf{Cov}[Y(t), Y(t + \Delta)] = \frac{c_6}{\Delta^{1-\beta}} , \quad (92)$$

where

$$c_6 = c_3 \int_0^\infty \frac{R_X(u)}{u^\beta} du \quad (93)$$

Condition 2 implies that: The integral $\int_0^\infty R_X(u) du$ is finite; The function $\widehat{R}_X(u)$ satisfies $\widehat{R}_X(0) = \int_{-\infty}^\infty R_X(t) dt$ (finite) and $\int_0^\infty \widehat{R}_X(u) du = \pi R_X(0)$ (finite). Hence, the integral appearing on the right hand side of equation (91) is convergent in the exponent range $0 < \beta \leq 1$, and the integral appearing on the right hand side of equation (93) is convergent in the exponent range $0 < \beta < 1$. The admissible exponent range is thus $0 < \beta < 1$.

References

- [1] Kingman J-F-C (1993) *Poisson processes* (Oxford University Press, Oxford).