

**A Semi-Analytical Model to Study the Effect of Cortical Tension on  
Cell Rolling**

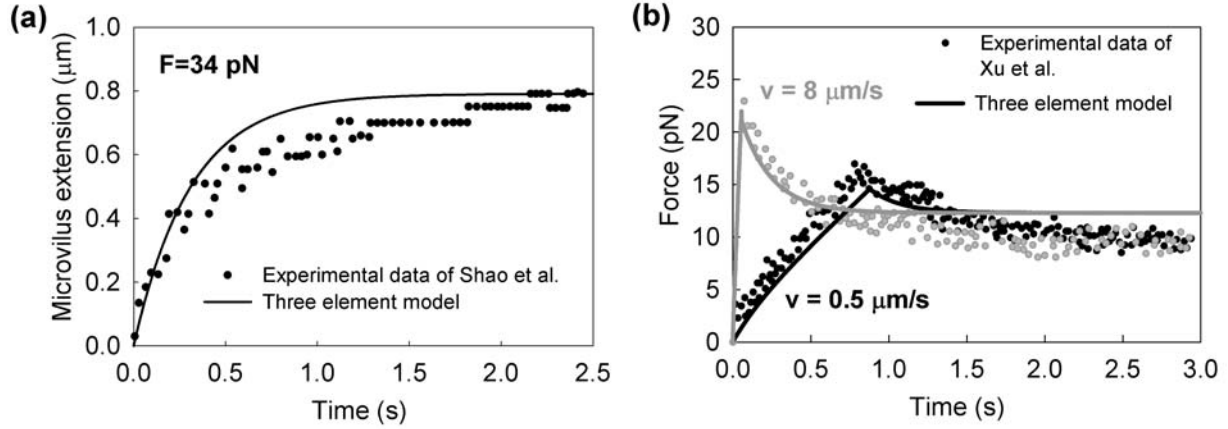
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**A) Three-element model of microvillus**

In order to model the dynamical response of the microvillus, we used a three element model as described in the paper. Briefly, the model consists of a spring (stiffness  $K_m$ ) and damper (viscosity  $\eta_m$ ) in series, which are both in parallel to another spring (stiffness  $K_c$ ). Transition to viscous regime occurs when the force in the spring  $K_c$  exceeds the critical force  $F_0$ , in which case the force in that spring is held constant at  $F_0$ . The governing equations are:

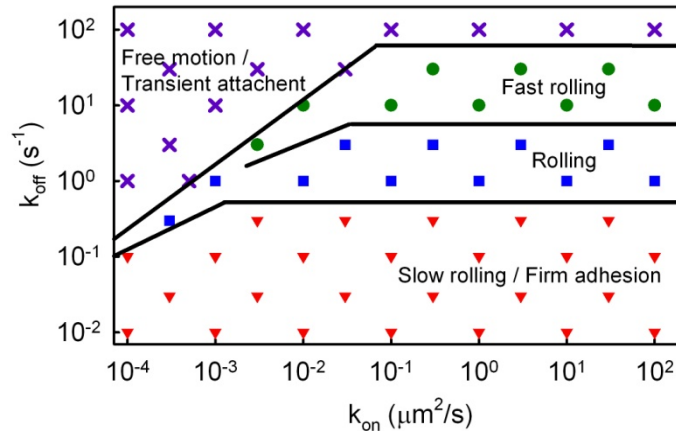
$$\begin{aligned} \frac{dF_m}{dt} + \frac{K_m}{\eta_m} F_m &= \frac{K_c K_m}{\eta_m} l + (K_c + K_m) \frac{dl}{dt} & \text{for } K_c l \leq F_0 \\ \frac{dF_m}{dt} + \frac{K_m}{\eta_m} (F_m - F_0) &= K_m \frac{dl}{dt} & \text{for } K_c l > F_0 \end{aligned} \quad (1)$$

When a constant force is applied such that the extension doesn't cause microvillus transition to the viscous regime, the steady state extension is governed by the spring stiffness  $K_c$ . In another scenario, when a constant force is applied such that it causes transition to the viscous regime, the steady state dynamics are governed purely by the viscous damper ( $\eta_m$ ) as the microvillus would extend at a constant velocity given by  $\mathcal{G} = (F - F_0)/\eta_m$ . In micropipette experiments, Shao et al. (1) applied a constant force on neutrophils and recorded microvillus extensions. From the steady state response they predicted an elastic stiffness of 43 pN/ $\mu\text{m}$  and a viscosity of 11 pN-s/ $\mu\text{m}$ , which for our model are assigned to  $K_c$  and  $\eta_m$  respectively. Recently, Xu et al. (2) used an optical trap to extract microvilli from neutrophils up to a fixed distance and studied their relaxation behavior. We used their data to estimate the value of  $K_m$ . The optical trap was modeled by a Kelvin-Voigt model with a spring (80 pN/ $\mu\text{m}$ ) placed in parallel to a damper (0.32 pN-s/ $\mu\text{m}$ ) as described by Xu (2). The Kelvin-Voigt body was connected in series to the three-element microvillus and the combined system of equations was solved numerically. We found a good agreement between the simulated results and the experimental data for  $K_m=200$  pN/ $\mu\text{m}$ . Figure S1 shows the comparison between the results from the model and the experimental data available.



**Figure S1:** (a) Extension of microvillus under a constant force of 34 pN was simulated using the three-element model and compared to the micropipette aspiration data of Shao et al. (1). The experimental data were corrected for time and natural length of microvillus such that  $t=0$  represented the beginning of the extension process. (b) Comparison between the microvillus relaxation data by Xu et al. (2) and the prediction of the model (for  $K_m=200$  pN/ $\mu$ m). In the experiments, microvilli were pulled to a fixed distance with a certain velocity and then held in place, and the decay in force was noted. Different rates of pulling, shown in different colors (grey and black), gave different transient responses.

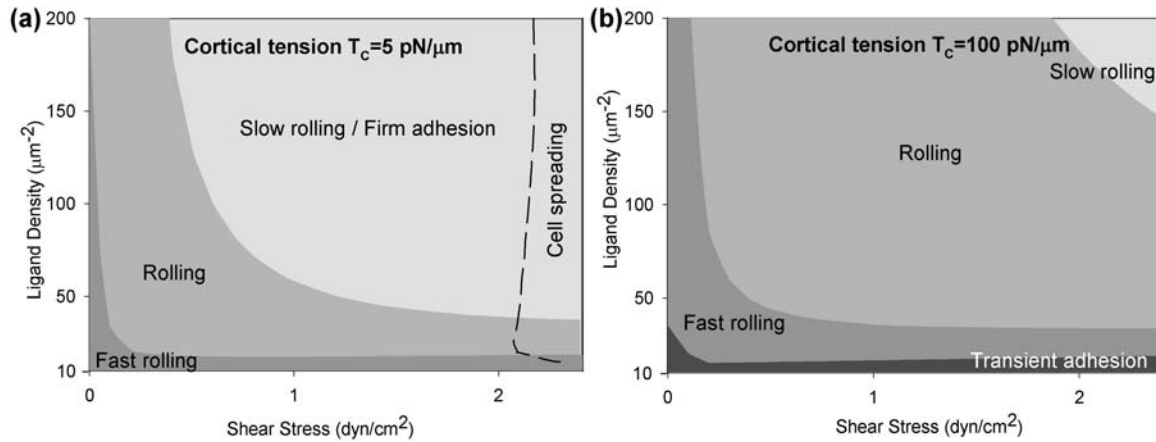
### **B) State diagram of rolling for varying on-rates and off-rates of receptor ligand interaction**



**Figure S2:** The state space diagram for rolling for different on and off rates of the receptor ligand interactions. The shear stress was 1 dyn/cm<sup>2</sup>, ligand density was  $N_L=50$   $\mu$ m<sup>-2</sup> and membrane tension was  $T_c=24$  pN/ $\mu$ m. Three regimes - slow rolling (or firm adhesion), rolling and fast rolling were identified based on the ratio of rolling velocity to hydrodynamic velocity as described in the main text. Free motion or transient tethering denotes the region where the equations did not yield a solution, indicating that steady rolling was not feasible.  $k_{off}$  significantly affected cell rolling, indicating that bond breakage rate is a crucial parameter that dictates the state of rolling. Stronger bonds (lower  $k_{off}$ ) resulted in higher force carried per bond reducing the rolling velocity. On the contrary,  $k_{on}$  had an effect on the rolling state only at low values ( $< 0.1$   $\mu$ m<sup>2</sup>/s) indicating a regime where bond formation is compromised. In this regime the lower formation rate of the bonds had to be compensated by a lower breakage rate to ensure

steady rolling resulting in the dip of the rolling region observed in the on-off diagram. All other parameters are as listed in Table 1 (main text).

### **C) Dependence of the state space diagram on cortical tension**



**Figure S3:** The state space diagram in the case of low ( $5 \text{ pN}/\mu\text{m}$ ) cortical tension (a) exhibits cell spreading at low shear stress in the range of  $2 \text{ dyn}/\text{cm}^2$  and slow rolling/firm adhesion at relatively low ligand densities and shear stresses compared to the case of high ( $100 \text{ pN}/\mu\text{m}$ ) cortical tension (b).

### **References**

1. Shao, J. Y., H. P. Ting-Beall, and R. M. Hochmuth. 1998. Static and dynamic lengths of neutrophil microvilli. *Proceedings of the National Academy of Sciences of the United States of America* 95:6797-6802.
2. Xu, G., and J. Y. Shao. 2008. Human neutrophil surface protrusion under a point load: location independence and viscoelasticity. *American Journal of Physiology-Cell Physiology* 295:C1434-C1444.