## Appendix B. Derivation of Formula 5.2.1

Given assumptions B.4 of the text, Eq. 5.1.3 becomes

$$\delta^{\prime(n)}\Delta\langle V_c\rangle = \frac{a^2}{4\pi^2 \varepsilon_{sc}} \int_{\text{MIR}} d\omega \, Im \, \{\eta(\omega)K_o(\omega)J_n(\omega)\}, \quad [\textbf{B.1}]$$

where

$$J_{n}(\omega) \equiv \int_{o}^{\infty} q^{2} dq \left\{ n^{-1} \sum_{i=1}^{n} \frac{f_{i}(q)}{(1 + qf_{i}(q)K_{o}(\omega/\varepsilon_{sc})^{2})} \right\}.$$
 [B.2]

Let us suppose that the *i*th acoustic eigenvalue (see section 5.1) tends to  $\alpha_i qd$  for  $qd \ll 1$  (i = 1, ..., n - 1). To evaluate expression **B.2**, we invoke assumption *B.7* of the text and assume that *n* is not too large; then, the inequality  $|K_o(\omega)| \gg nd\varepsilon_{sc}$  is fulfilled for most of the relevant range of  $\omega$ . This allows us to split the integral in Eq. **B.2** into two parts,  $J_1$  and  $J_2$ , corresponding to  $0 \leq q \leq q_c$  and  $q_c \leq q \leq \infty$ , where we choose the cutoff  $q_c$  so that  $nq_cd \ll 1$  but  $q_c^2|K_o|d/\varepsilon_{sc} \gg |$ , and to approximate the *i*th term in the integrand of  $J_1$  by  $n^{-1}\alpha_i dq^3/(1 + \alpha_i dq^2 K_o(\omega)/\varepsilon_{sc})^2$  (i = 1 ..., n - 1) and in the integrand of  $J_2$  by  $\varepsilon_{sc}^{-2} K_o^{-2} \{n^{-1}f_i^{-1}(q)\}$  (i = 1 ..., n). The single optical term (i = n) contributes to  $J_1$  only a term of relative order  $q_c d$ , which may consistently be neglected (through compare section 6).

The crucial observation, now, is that<sup>j</sup>

$$n^{-1}\sum_{j=1}^{n}f_{j}^{-1}(q) - 1 = 2\left(1 - \frac{1}{n}\right)(\exp(2qd) - 1)^{-1},$$
 [**B.3**]

from which it follows, taking the limit  $q \to 0$ , that  $\sum_{i=1}^{n-1} \alpha_i^{-1} = n-1$ . Explicit evaluation of  $J_1$  and  $J_2$  using Eq. **B.3** then leads to the result (to within terms of relative order  $\zeta^{-1} \equiv \varepsilon_{sc} d/|K|$ )

$$J_n(\omega) = \left(1 - \frac{1}{n}\right) \cdot \frac{\varepsilon_{sc}^2}{2dK_o^2(\omega)} \cdot \{\ln(\mathbf{K}_o(\omega)/4e^{1/2}d\varepsilon_{sc}) - \beta_n\},$$
[B.4]

where

$$\beta_n = \frac{1}{2} (n-1) \sum_{i=1}^{n-1} \alpha_i^{-1} \ln \alpha_i$$
 [B.5]

The term  $\beta_n$  is numerically small ( $\beta_n = 0, 0.13, 0.18...03, ...$  for  $n = 2, 3, 4...\infty$ ). If we neglect it, we immediately obtain from Eq. **B.4** formula **5.2.1** of the text.

<sup>&</sup>lt;sup>j</sup>It is straightforward to prove Eq. **B.3** case by case for  $n \le 6$  and  $n = \infty$ . I am indebted to Misha Turlakov for a proof in the general case.