## **Appendix S2: Free energy of a membrane in a singlewell potential**

We start from the partition function

$$\mathcal{Z} = \left[\prod_{i} \int_{0}^{\infty} \mathrm{d}l_{i}\right] e^{-\left(\mathcal{H}_{\mathrm{el}}\left\{l\right\} + a^{2}\sum_{i} V(l_{i})\right)/k_{B}T}$$

of a membrane bound by the single-well potential

$$V(l_i) = -U\theta(l_{we}/2 - |l_i - l_o|)$$

The potential well has the depth U, width  $l_{we}$  and is centered at the membrane separation  $l_o$ . The derivative of the free energy  $\mathcal{F} = -k_B T \ln \mathcal{Z}$  with respect to the depth U of the well can be written as

$$\frac{\partial \mathcal{F}}{\partial U} = -AP_b$$

where  $P_b(U, l_{we}) = \langle \theta(l_{we}/2 - |l_i - l_o|) \rangle$  is the average fraction of membrane sites within the potential well, and  $A = \sum_i a^2$  is the total membrane area. The free energy difference between two states with potential depths  $U_1$  and  $U_2$  is therefore

$$\Delta \mathcal{F} = -A \int_{U_1}^{U_2} P_b\left(U, l_{\rm we}\right) dU$$

The scaling analysis of appendix A indicates that  $P_b$  depends primarily on the rescaled depth

$$u = U l_{\rm we}^2 \kappa / \left( k_B T \right)^2$$

A change of variables now leads to

$$\Delta f = -\frac{(k_B T)^2}{\kappa l_{\rm we}^2} \int_{u_1}^{u_2} P_b(u) du$$

with  $\Delta f = \Delta \mathcal{F} / A$ .

The function  $P_b(u)$  can be determined from Monte Carlo simulations of a membrane bound in a single well (see fig. 2). We find that the Monte Carlo data can be well described by

$$P_b(u) \simeq P_b^{(3)}(u) = \frac{u + c_2 u^2 + c_3 u^3}{c_1 + u + c_2 u^2 + c_3 u^3}$$

with the three fit parameters  $c_1 \simeq 0.073$ ,  $c_2 \simeq -0.99$ , and  $c_3 \simeq 6.44$  for the rescaled width  $z_{we} = 0.5$ . For  $u \lesssim 0.2$ , the three-parameter function  $P_b^{(3)}(u)$  for the rescaled width  $z_{we} = 0.5$  coincides with the function  $P_b^{(3)}(u)$  for  $z_{we} = 1$  (see full lines in fig. 2), which illustrates that the rescaled well depth u is the dominant parameter in this range. Dimensional analysis indicates that u is the dominant parameter as long as the lateral correlation length of the membrane fluctuations is significantly larger than the discretization length a (see text S1). For strongly bound membranes with values of  $P_b$  close to 1, the correlation length is comparable to a, and the precise value of  $P_b$  depends also on the rescaled well width  $z_{we}$ .