

# Supporting Information

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## SI Text

**The Chemical Shift Tensor.** The electronic environment surrounding the nucleus gives rise to the CST. In the presence of an external magnetic field,  $B_0$ , the electrons in the orbitals surrounding these nuclei orient themselves either with the field or against the field, shifting the observed frequency of precession. The chemical shift,  $\delta^{\text{CS}}$ , is typically represented as a Cartesian tensor composed of three orthogonal axes.

$$\delta^{\text{CS}} = \begin{pmatrix} \delta_{11} & \delta_{12} & \delta_{13} \\ \delta_{21} & \delta_{22} & \delta_{23} \\ \delta_{31} & \delta_{32} & \delta_{33} \end{pmatrix} \quad [\text{S1}]$$

where  $\delta_{ii}$  are the nine components of the CST. In solution, isotropic molecular tumbling averages this tensor and the familiar isotropic “chemical shift” ( $\delta_{\text{iso}}$ ) is one third the trace of the above tensor:

$$\delta_{\text{iso}} = \frac{1}{3}(\delta_{11} + \delta_{22} + \delta_{33}). \quad [\text{S2}]$$

Here,  $\delta_{11}$  is the most downfield tensor element,  $\delta_{33}$  the furthest upfield, with  $\delta_{22}$  between these two extremes. The orientation of these elements to the molecular frame can be described by a series of angles. Often, these are a set of Euler angles in a common coordinate system. In the convention presented, here these angles are defined as  $\alpha_n$ , the angle between the  $n$ th tensor element and the  $^1\text{H}$ - $^{13}\text{C}$  or  $^{15}\text{N}$ -H vector, and  $\beta_n$ , the orientation of the  $n$ th tensor element and the  $^{13}\text{C}\alpha$ - $^{15}\text{N}$  vector, illustrated in Fig. 1. An alternate description of the CST was popularized by Haerberlen, Merhing, and Waugh. In this convention, very convenient for conversion from a Cartesian to spherical tensor basis set, labels the three axes  $\delta_{xx}$ ,  $\delta_{yy}$ ,  $\delta_{zz}$ . They are then ordered by their deviation from the isotropic chemical shift, with  $\delta_{zz}$  having the greatest deviation, followed by  $\delta_{xx}$  and  $\delta_{yy}$  closest to the  $\delta_{\text{iso}}$ . In this convention tensor magnitude and rhombicity are defined by two parameters  $\delta$  (or  $\delta_{\text{aniso}} = \delta_{zz} - \delta_{\text{iso}}$ ) and  $\eta$  ( $\eta = (\delta_{xx} - \delta_{yy})/\delta$ ).

**Order Parameters.** The experimentally determined  $^1\text{H}$ - $^{15}\text{N}$  and  $^1\text{H}$ - $^{13}\text{C}$  vector orientation report upon both bond distances as well as molecular motion. The assumption is that if the  $^1\text{H}$ - $^{15}\text{N}$  bond length is approximately 1.04 Å and the  $^1\text{H}$ - $^{13}\text{C}$  bond length is approximately 1.12 Å, in the presence of fast rigid motions, the order parameter  $S$  is the ratio of the measured dipolar coupling to the ideal dipolar coupling. This approach is the same as that used in recent studies of both ubiquitin (1) and thioredoxin (2). The order parameters for both NH and CH are presented in Fig. S3. The values of  $S$  measured from the H $\alpha$ -C $\alpha$  dipole (Fig. S3A) reveal a rigid backbone, where  $S \sim 0.95$  for most sites, comparable to values measured in solid thioredoxin (2) but larger than the  $S$  values observed in solid ubiquitin (1). It is known that L12 and G41 in GB1 experience significant motional averaging; however, L12 is not labeled in this sample and the glycine  $^1\text{H}$ - $^{13}\text{C}$  trajectories were not fit because of the high degeneracy of solu-

tions arising from the presence of two directly bound protons; however, the residues adjacent to these sites (T11, K13, D40) do exhibit motional averaging greater than 1 standard deviation from the other sites. Based upon a recent study by Case et al., (3) dipolar averaging of this magnitude would result only a very small motional averaging of the  $^{13}\text{C}$  and  $^{15}\text{N}$  CSTs. For example, this study revealed that a Lipari-Szabo order parameter,  $S^2$ , determined for F52 in GB3 using HN dipolar terms would be 0.897, however when CSA is considered the value rises to 0.975, this is largely attributed to the motion of the lighter proton relative to the heavier peptide backbone. This indicates there is only small motional averaging of CST magnitudes at sites away from the loop containing G41 in GB1 at ambient temperatures. The order parameters measured from the backbone  $^1\text{H}$ - $^{15}\text{N}$  couplings (Fig. S3B) follow a nearly identical pattern to that seen in the  $^1\text{H}$ - $^{13}\text{C}$  results. Here, the G41  $^1\text{H}$ - $^{15}\text{N}$  tensor dipole can be measured, and reveals significant motional averaging, comparable to that reported by Barchi et al. (4). The main differences relative to the  $^1\text{H}$ - $^{13}\text{C}$  scalings are in the turn near A20, and the turn near T49, due most likely to slightly elongated intermolecular and intramolecular hydrogen bonds. However, in the  $^1\text{H}$ - $^{15}\text{N}$  recoupling experiment, A20 is overlapped with N8, and the signal-to-noise ratio at E19 is poor, so this cannot be confirmed. Overall, however, these measurements reveal a relatively rigid, well-ordered backbone.

**Data Analysis and Fitting.** The data that was acquired resulted in a total of six trajectories for each fit site in the  $^{13}\text{C}$  correlation experiments and five trajectories for each set of  $^{15}\text{N}$  correlation experiments. During the fitting first the CST magnitudes were fixed (in the case of  $^{13}\text{C}$  previously reported values were used). Following this the  $R18_1$  trajectory was fit for both effective bond length and relaxation. In the next step angles are fit holding all magnitudes and relaxation fixed. In the fourth step all magnitudes, relaxation parameters, and angles are allowed to vary. It was found during this procedure that modeling of cross polarization helped improve the fit quality, especially in the case of  $^1\text{H}$ - $^{13}\text{C}$  trajectories and [ $^1\text{H}$ - $^{13}\text{C}$ ]: [ $^{13}\text{C}$  CST] correlation spectra.

While powerful, these experiments possess a few complexities that needed to be addressed. Tensor correlation experiments in SSNMR are most accurate when the correlated vectors are oriented at  $30^\circ$  or less. To partially overcome this shortcoming, each CST orientation was constrained relative to two different vectors, resulting in an improvement in the fit quality for these near orthogonal orientations. Unlike the  $^1\text{H}$ - $^{13}\text{C}$  dipolar tensor, the C $\alpha$  CST deviates significantly from axial symmetry, and a single correlation of one tensor to another can be ambiguous in isolated regions of conformational space. To address this, we acquired several different ratios of ROCSA to  $R18_1$  evolution times. This allowed for tighter constraints on the C $\alpha$  tensor, and also alleviated errors arising from  $^1\text{H}$ - $^{13}\text{C}$  distance measurements. Using this approach we found that varying the  $^1\text{H}$ - $^{13}\text{C}$  dipolar coupling by up to 1 kHz resulted in minimal perturbation ( $x$ - $y^\circ$ ) in the fitted orientations.

1. Lorieau JL, McDermott AE (2006) Conformational flexibility of a microcrystalline globular protein: Order parameters by solid-state NMR spectroscopy. *J Am Chem Soc* 128:11505–11512.
2. Yang J, Tasayco ML, Polenova T (2009) Dynamics of reassembled thioredoxin studied by magic angle spinning NMR: Snapshots from different time scales. *J Am Chem Soc* 131:13690–13702.

3. Tang S, Case DA (2007) Vibrational averaging of chemical shift anisotropies in model peptides. *J Biomol NMR* 38:255–266.
4. Barchi JJ, Grasberger B, Gronenborn AM, Clore GM (1994) Investigation of the backbone dynamics of the Igg-binding domain of streptococcal protein-G by heteronuclear 2-dimensional  $^1\text{H}$ - $^{15}\text{N}$  nuclear magnetic resonance spectroscopy. *Protein Sci* 3:15–21.















**Table S3. Structural quality as assessed by Procheck**

Restrains Used Structure	Ramachandran Quality %						X-ray Structure Equivalence					
	T	D	C	O	V	SE	Most Favored	Allowed	Rama	H-bondEnergy	$\chi^1$	$\chi^2$
1				X			77.0	22.2	2.5	2.9	1.8	1.0
2		X	X	X			94.0	6.0	1.0	3.3	1.9	1.0
3	X	X					94.0	6.0	1.0	3.4	3.2	1.7
4	X		X	X			86.0	14.0	1.8	2.7	2.3	1.3
5	X	X	X	X			93.8	6.2	1.0	3.0	2.2	1.4
6	X	X	X	X	X		96.0	4.0	1.0	3.0	2.4	1.2
7	X	X			X	X	100.0	0.0	1.0	2.9	2.0	1.0
8	X		X	X	X	X	95.8	4.2	1.0	2.8	1.9	1.0
9	X		X	X		X	93.2	6.8	1.0	2.6	1.9	1.0
10	X	X	X	X		X	96.0	4.0	1.0	2.8	1.8	1.0
11	X	X	X	X	X	X	96.4	3.6	1.0	2.8	1.8	1.0
2QMT							96.0	4.0	1.0	1.5	1.0	2.1
2GI9							94.0	6.0	1.0	1.2	1.0	1.9
1PGA							92.0	8.0	1.0	1.4	1.4	2.1
1PGB							90.0	10.0	1.0	1.0	1.5	2.1
2JSV		X			X		92.0	8.0	1.0	4.0	3.2	1.6

**Table S4. Agreement of CST magnitudes and orientations with theory for all structures**

Structure	CST						$\alpha_{(1,2,3)}$		$\beta_{(1,2,3)}$			
	T	D	C	O	V	SE	RMSD(ppm)	$R^2$	RMSD (°)	$R^2$	RMSD (°)	$R^2$
1				X			3.6	0.95	4.7	0.99	5.7	0.98
2		X	X	X			2.1	0.98	5.6	0.98	6.6	0.97
3	X	X					3.1	0.96	11.0	0.93	9.6	0.93
4	X		X	X			2.1	0.98	5.5	0.98	6.5	0.97
5	X	X	X	X			2.0	0.98	5.6	0.98	6.6	0.97
6	X	X	X	X	X		2.1	0.98	5.8	0.98	6.7	0.97
7	X	X			X	X	2.9	0.96	9.3	0.96	9.0	0.95
8	X		X	X	X	X	2.1	0.98	5.8	0.98	6.6	0.97
9	X		X	X		X	2.0	0.98	5.7	0.98	6.5	0.97
10	X	X	X	X		X	2.0	0.98	5.6	0.98	6.6	0.97
11	X	X	X	X	X	X	2.1	0.98	5.5	0.98	6.7	0.97
2QMT							2.5	0.97	8.0	0.97	8.5	0.95

**Table S5. Agreement of structures with measured pseudodihedral angles**

Restrains Used Structure	VEAN						rmsd (°)*	dev. > 5°*	dev. > 10°*
	T	D	C	O	V	SE			
1					X		11.5	40	21
2			X	X	X		7.4	27	12
3		X	X				8.1	41	17
4		X		X	X		9.1	31	16
5		X	X	X	X		6.7	26	11
6		X	X	X	X	X	3.5	11	2
7		X	X		X	X	3.0	12	4
8		X		X	X	X	3.0	11	2
9		X		X	X	X	8.6	33	14
10		X	X	X	X	X	8.6	34	14
11		X	X	X	X	X	3.0	13	2
2QMT†							6.1	28	13

\*After accounting for experimental error.

†Assuming canonical <sup>1</sup>H bond lengths, bond angles, and positions.



**Table S6. Agreement of experimental isotropic chemical shifts with shifts predicted by SPARTA for each structure**

Structure	N	T	D	C	O	V	SE	C $\alpha$		C $\beta$		C'	
								RMSD (ppm)	R <sup>2</sup>	RMSD (°)	R <sup>2</sup>	RMSD (°)	R <sup>2</sup>
1	X				X			1.08	0.95	1.35	0.99	1.23	0.8
2	X		X	X	X			1.06	0.95	1.28	0.99	1.15	0.8
3	X	X	X					1.04	0.96	1.25	0.99	1.14	0.8
4	X	X		X	X			1.01	0.96	1.13	0.99	1.13	0.8
5	X	X	X	X	X			1.02	0.96	1.13	0.99	1.16	0.8
6	X	X	X	X	X	X		1.04	0.96	1.15	0.99	1.19	0.8
7	X	X	X			X	X	1.04	0.96	1.27	0.99	1.16	0.8
8	X	X		X	X	X	X	1.01	0.96	1.13	0.99	1.13	0.8
9	X	X		X	X		X	1.01	0.96	1.13	0.99	1.13	0.8
10	X	X	X	X	X		X	1.01	0.96	1.13	0.99	1.13	0.8
11	X	X	X	X	X	X	X	1.03	0.96	1.14	0.99	1.14	0.8
2QMT								0.93	0.97	0.99	0.99	1.12	0.8

**Table S7. Backbone rmsd among crystal structures and SSNMR structure**

Structure	SSNMR	2QMT	2GI9	1PGA	1PGB	Crystal Avg.
SSNMR	—	0.51	0.58	0.54	0.53	0.54
2QMT	0.51	—	0.47	0.39	0.28	0.39
2GI9	0.58	0.47	—	0.24	0.36	0.37
1PGA	0.54	0.39	0.24	—	0.25	0.30
1PGB	0.53	0.28	0.36	0.25	—	0.30
Crystal Avg	0.54	0.39	0.37	0.30	0.30	—

**Table S8. Table of TALOS restraints that violated at some point during annealing and refinement**

Residue	Angle	TALOS Prediction (°)	Doubled TALOS error (°)
Q2	$\phi$	-109	34.0
N8	$\phi$	-100	34.0
T11	$\phi$	-94.0	34.0
E19	$\phi$	-117.0	20.0
A24	$\phi$	-65.0	24.0
E27	$\phi$	-66.0	16.0
N37	$\phi$	-105.0	32.0
V39	$\phi$	-101.0	30.0
W43	$\phi$	-122.0	32.0
T49	$\phi$	-110.0	36.0
T11	$\psi$	-6.0	28.0
T18	$\psi$	137.0	28.0
A26	$\psi$	-41.0	10.0
Y33	$\psi$	-37.0	18.0
D36	$\psi$	-31.0	16.0
W43	$\psi$	160.0	20.0
Y45	$\psi$	126.0	48.0
T49	$\psi$	10.0	26.0
T51	$\psi$	134.0	32.0
T53	$\psi$	144.0	34.0

