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SI Materials and Methods

The two-exponential kinetics of Q_p was reported by Horrigan and Aldrich (1) for BK channels in the absence of Ca^{2+} . They explained this behavior in terms of an allosteric model (Fig. S5A) in which the four voltage sensors have two states, resting, R , and active, A . The $R-A$ equilibrium constant, J is voltage dependent according to the relation $J = J_0 exp(-z_1 eV/kT)$, where J_0 is the zero voltage equilibrium constant and every R to A transition results in the transfer of z_J elementary charges across the membrane. Membrane depolarization displaces the $R-A$ equilibrium toward A , and this process generates the fast gating current that when integrated in time gives Q_{fast} . Activation of the voltage sensors displaces the closed, C , to open, O , equilibrium toward the open state due to allosteric coupling, through the allosteric coupling constant D . The $O-C$ equilibrium constant, L , is also voltage dependent ($L = L_0 \exp(-z_L eV/kT)$ and the closed-toopen transition displaces z_L elementary charges. Because the closed-to-open transition is slow, the charge transported contributes to the Q_{pSlow} component of Q_{p} . The allosteric coupling of the voltage sensors and the channel gate shifts toward A the voltage sensor equilibrium and the additional charge displaced by this rearrangement contributes to Q_{pslow} . Therefore, the slow component of Q_p is dependent on the closed-to-open transition.

1. Horrigan FT, Aldrich RW (1999) Allosteric voltage gating of potassium channels II. Mslo channel gating charge movement in the absence of $Ca(2+)$. J Gen Physiol 114(2): 305–336.

Our results on BK channels associated with the β1 auxiliary subunits show that Q_p also has a slow component but it appears at voltages where the closed-to-open transition is fully displaced toward the closed state. This result is incompatible with the H–C–A allosteric model, and, therefore we split the closed-to-open transition into two steps (Fig. S5B): The slow component of Q_p represents the charge transported due to the transition from the deeper closed state, C_0 , to another closed state, C_1 , and it is independent of the final transition from C_1 to the open state. This latter transition is governed by the equilibrium constant M , which may be voltage-dependent and may displace charge, but the amplitude of the associated gating current is expected not to be measurable because the opening of the channel in the presence of β 1 is very slow. We have augmented the two-tiered 10state H–C–A allosteric model to a three-tiered 15-state model (Fig. $S5B$). Now a new allosteric coupling constant E must be added to complete the scheme.

Steady-state analysis of the gating charge, i.e., determination of Q_{fast} and Q_{pSlow} as a function of voltage is sufficient to make a first approximation of the equilibrium constants J and L as a function of voltage and the allosteric coupling constant D. We augmented the model to account for the kinetics of the processes, by splitting the equilibrium constant J, L , and M into forward and backward rate constant (Fig. S5B).

Fig. S1. I_K and I_g currents measured in the same macropatch. I_K currents were recorded in the cell-attached configuration (A) to ensure that BK channels were formed by the expected α or α/β complex. After patch excision (B), I_0 currents were isolated in conditions of 0 K⁺ by repeated bath exchange of the internal side of the macropatch with a K⁺-free solution (internal solution, see Materials and Methods).

Fig. S2. ON and OFF gating currents. (A and C) Q_{ON} (red circles) and Q_{OFF} (blue circles) charge of ($\alpha\beta\gamma$ BK and $\alpha\beta\gamma$)BK channel are equal. (B and D) Q_{ON} plotted against Q_{OFF} are linearly related.

Fig. S3. I_{gON} is well fitted using a single exponential decay (red line).

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Fig. S4. Kinetics of gating charge. (A–E) Exponential decays of I_{gON} were fitted with an exponential function and the time constants (τ_{g} , mean \pm SEM) were plotted against the voltage. The solid black line is the best fit to a two-state model (Table S1). All graphs include the Boltzmann fit of the Q–V curve (dashed lines in A–E) and the corresponding V_h is indicated by a vertical line for the Q–V curve (solid colored line in A–E). (F) The linear range of each plot was fitted with an exponential decay function and is shown by black lines with the exception of (α/β4)BK, which is shown by a green line.

Fig. S5. Allosteric models explain the complex kinetics of BK channels in the absence of calcium. (A) Horrigan, Cui, and Aldrich allosteric model (1). States within a row correspond to the possible combinations of the four independent voltage sensors. Subscripts indicate the number of voltage sensors in the active position. Resting/active equilibrium is governed by the voltage-dependent constant J. A single channel gates between two possible states, closed and open, with an equilibrium constant L. As voltage sensors reach the active state, channel opening is favored by an allosteric factor D. (B) Allosteric model in which the channel gate opens in two sequential steps. Equilibrium between the first and second row of closed states is governed by L and the allosteric factor D , whereas allosteric factors M and E describe equilibrium between $C_{1,0}$ through $C_{4,0}$ and O_1 through O_4 .

1. Horrigan FT, Cui J, Aldrich RW (1999) Allosteric voltage gating of potassium channels I. Mslo ionic currents in the absence of Ca(2+). J Gen Physiol 114(2):277–304.

Fig. S6. Changes in z_j and $V_{1/2}$ account for calcium sensitivity in (α/β 4)BK channels. Data from ref. 1 (bluish-green circles) are superimposed with simulation of (α)BK (dashed black line) and (α/β4)BK channels (dashed bluish-green line) using the H–C–A model (Fig. 4A) with the following parameters: (α)BK V_h = 140 mV, $z_1 = 0.58$, $z_1 = 0.31$, $D = 19$, $C = 14$, $E = 2.8$, $L_0 = 1e-6$, and $K_d = 8.2$ (data from ref. 2). For (α /β4)BK all parameters were the same except $V_{1/2} = 186$ mV and $z_1 = 0.4$.

1. Wang B, Rothberg BS, Brenner R (2006) Mechanism of beta4 subunit modulation of BK channels. J Gen Physiol 127(4):449–465. 2. Orio P, Latorre R (2005) Differential effects of beta 1 and beta 2 subunits on BK channel activity. J Gen Physiol 125(4):395–411.

Table S1. Fit parameters of gating charge-voltage relationships for (α)BK and (α/βx)BK channels

All of the corresponding $Q(V)$ curves were displaced in the voltage axis by $\Delta V = (V_h - \langle V_h \rangle)$. V_h and z obtained from the fit of the Q(V)data to a Boltzmann function.

z (e_0)	δ	α_0 (s ⁻¹)	β_0 (s ⁻¹)
0.63	0.29	4.5	70
0.58	0.30	7.5	28
0.61	0.30	4.5	37
0.56	0.30	4.5	48
0.48	0.28	7.0	19

Table S2. Fit parameters of the kinetics of gating charge

Best fit to a two-states model with z constrained to the values found for the Q–V relation.