Supplementary Information S1 (Box) Modelling the effects of buffers in realistic coupling regimes.

A lot of insight into the mechanisms of coupling between  $Ca^{2+}$  channels and  $Ca^{2+}$  sensors can be obtained by modelling the diffusion of  $Ca^{2+}$  and its reaction with buffers. As in other fields of Neuroscience, the Hopfield quote "build it, and you understand it" perfectly applies. How can one model the  $Ca^{2+}$  transient?

In a simple scenario, the steady-state solution to the linearized reaction-diffusion problem is obtained analytically<sup>1,2</sup>. In this framework, the  $Ca^{2+}$  concentration ([ $Ca^{2+}$ ]) can be described by a simple equation, which is comprised of a 1 / r term (representing diffusion) and an exponential term (representing buffering):

$$[Ca^{2+}] = \frac{i_{Ca}}{4\pi F D_{Ca}} \, 1 / r \, \exp(-r / \lambda)$$
(Eq. 1)

with  $\lambda = \sqrt{D_{Ca}} / (k_{on} [B]),$ 

where  $i_{Ca}$  is the Ca<sup>2+</sup> current, F is the Faraday constant, D<sub>Ca</sub> is the diffusion coefficient of Ca<sup>2+</sup>, r is radial distance from a source,  $\lambda$  is the length constant,  $k_{on}$  is the Ca<sup>2+</sup>-binding rate of the buffer, and [B] is the concentration of the buffer<sup>1</sup>.

Although the linear approach represents a useful approximation for short distances from the source, it does not account for the time course of the  $Ca^{2+}$  transient, the phenomenon of buffer saturation, and the presence of fixed and mobile buffers<sup>1</sup>.

The limitations can be overcome by obtaining the time-dependent solution to the full reaction-diffusion equations<sup>3-7</sup>. This can be done by numerically solving a set of partial differential equations, containing the Ca<sup>2+</sup> and buffer concentrations as a function of space and time, as well as several partial derivatives.

Everything starts from Fick's first and second law of diffusion<sup>8</sup>. Fick's first law relates the diffusive <u>flux</u> to the concentration field. In the simplest possible form in one spatial dimension, the first law is

$$J = D_{Ca} \frac{\partial [Ca^{2+}]}{\partial x}, \qquad (Eq. 2)$$

where J is the flux in units mol  $s^{-1} m^{-2}$ . From the law of mass conservation and Fick's first law, Fick's second law can be derived<sup>2</sup>.

$$\frac{\partial [Ca^{2+}]}{\partial t} = \frac{\partial J}{\partial x} = \frac{\partial}{\partial x} \left( D_{Ca} \frac{\partial [Ca^{2+}]}{\partial x} \right)$$
(Eq. 3)

Equation 3 gives the partial differential equation that has to be solved. Equation 2 gives the boundary condition near the source. In addition, a second boundary condition has to be implemented remote from the source. This is usually a reflective boundary condition, which is given as  $\partial [Ca^{2+}] / \partial x$ = 0 for  $x \rightarrow x_{max}$ . As there is no gradient at this distance, Ca<sup>2+</sup> cannot escape beyond this point. Furthermore, initial conditions have to be appropriately chosen. For example,  $[Ca^{2+}]$  at t = 0 is set to the resting value. The partial differential equations can be solved numerically, e.g. using NDSolve of Mathematica<sup>3,4,9</sup>.

Finally, the effect of the Ca<sup>2+</sup> transient on transmitter release has to be simulated, using models of transmitter release derived from Ca<sup>2+</sup> uncaging experiments<sup>10-14</sup>. Based on a 6- to 8-state reaction scheme, a set of ordinary differential equations can be formulated, which can be solved numerically.

The cookbook recipe (Eq. 1 - 3) describes the backbone of the simulations, defining the Ca<sup>2+</sup> transients from a point source in the absence of buffers. For a more realistic simulation, several extensions have to be made. In the presence of buffers, the right hand side of equation 3 has to be extended by the sum of reaction terms. To simulate Ca<sup>2+</sup> transients originating from Ca<sup>2+</sup> channel clusters or other distributed sources, the one-dimensional simulations have to be extended into two or three dimensions<sup>5-7</sup>.

Early studies have used several different approximations, such as the steady-state excess buffer approximation (EBA; buffer concentration is so high that it changes little during Ca<sup>2+</sup> inflow) and rapid buffer approximation (RBA; buffers are so fast that they are in chemical equilibrium with Ca<sup>2+</sup> at every point in time and space<sup>7</sup>). As computer power has increased, these approximations have become unnecessary.

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