# Supplementary information:

Barrett *et al.*, "Extra permeability is required to model dynamic oxygen measurements: evidence for functional recruitment?", JCBFM

# Supplementary Methods

# Nondimensionalization

The equations in the main text are non-dimensionalized to simplify the system and improve numerical stability using the same approach as our previous model (Barrett et al., 2012). The variables were scaled such that

$$c_{O_{2},i}(t) = C_{O_{2},i}(T)/\hat{C}_{O_{2}}; \quad cmr_{O_{2}}(t) = CMR_{O_{2}}(T)/\hat{J}_{O_{2}};$$

$$f_{i}(t) = F_{i}(T)/\hat{F}; \qquad g_{i} = G_{i}/\hat{G};$$

$$j_{O_{2},i}(t) = J_{O_{2},i}(T)/\hat{J}_{O_{2}}; \qquad n_{O_{2}}(t) = N_{O_{2}}(T)/\hat{N}_{O_{2}}; \qquad (S.1)$$

$$p_{O_{2},i}(t) = P_{O_{2},i}(T)/\hat{P}_{O_{2}}; \qquad t = T/\hat{T};$$

$$v_{i}(t) = V_{i}(T)/\hat{V}; \qquad v_{t} = V_{t}/\hat{V}.$$

We defined the scales as

$$\hat{C}_{O_2} = C_{O_2,ref}; \qquad \hat{F} = F^*; \qquad \hat{G} = \hat{N}_{O_2}/\hat{P}_{O_2}\hat{T}; 
\hat{J}_{O_2} = \hat{N}_{O_2}/\hat{T}; \qquad \hat{N}_{O_2} = \hat{C}_{O_2}\hat{V}; \qquad \hat{P}_{O_2} = P_{O_2,ref}; \qquad (S.2) 
\hat{T} = \hat{V}/\hat{F}; \qquad \text{and } \hat{V} = \sum_{i=1}^3 V_i^*$$

where  $C_{O_2,ref}$  and  $P_{O_2,ref}$  are reference values for oxygen concentration and partial pressure,  $F^*$  is the flow at baseline, and  $\sum_{i=1}^{3} V_i^*$  is the total volume of the three compartments at baseline. Derived variables or parameters not specifically mentioned in Equation (S.1) are nondimensionalized using the scale with the correct dimensions.

# **CMRO**<sub>2</sub> Stimulus

The CMRO<sub>2</sub> stimulus,  $s_t(t)$ , is described by the piecewise function

$$s_t(t) = \begin{cases} s_{up}(t), & t < t_0 + \tau_{up} \\ s_{decay}(t), & t_0 + \tau_{up} \le t \le t_0 + t_{stim} \\ s_{down}(t), & t > t_0 + t_{stim} \end{cases}$$
(S.3)

where:

$$s_{up}(t) = \frac{1}{2} s_{up}^* \left[ 1 + \operatorname{erf}\left(\frac{t - [t_0 + \tau_{up}/2]}{32^{-1/2}\tau_{up}}\right) \right];$$
(S.4a)

$$s_{decay}(t) = \left(s_{up}^* - s^*\right) \exp\left(\frac{t_0 + \tau_{up} - t}{\tau_{decay}}\right) + s^*;$$
(S.4b)

$$s_{down}(t) = s_{end} \exp\left(\frac{t_0 + t_{stim} - t}{\tau_{down}}\right);$$
(S.4c)

 $t_0$  is the stimulus onset time;  $\tau_X$  are time constants (where X is one of up, decay, or down);  $t_{stim}$  is the duration of stimulation (where  $t_{stim} \ge \tau_{up}$ );  $s_{up}^*$  is the peak value of  $s_{up}(t)$ ; erf is the Gauss error function:

$$\operatorname{erf}(x) = \frac{2}{\sqrt{\pi}} \int_0^x \exp\left(-t^2\right) dt; \qquad (S.5)$$

 $s^*$  is the steady state value of the stimulus; and  $s_{end}$ , the stimulus value at  $t_{end}$ , is given by  $s_{decay}(t_{end})$ .

To fit model predictions of tissue PO<sub>2</sub> to experimental measurements, the optimization algorithm adjusted the three stimulus time constants ( $\tau_{up}$ ,  $\tau_{down}$ , and  $\tau_{decay}$ ) and two stimulus amplitudes ( $s_{up}^*$ ,  $s^*$ ), as described in the Methods section in the main text (Experimental Design subsection).

#### Calculating Mean Tissue PO<sub>2</sub>

This section describes how the model calculates  $PO_2$  in tissue, and provides the basis for calculating the value of baseline tissue  $PO_2$  (see below). Although using an assumed or estimated  $PO_2$  value would partially remove the need for this additional modelling, we believe that the equations in this section represent a more rigorous approach. As a consequence, the predictions in the main text do not depend on the details of the following model, but do rely on the baseline  $PO_2$  value it produces.

While including these equations adds an additional layer of complexity, there are a number of advantages of doing so. For example, this technique makes it possible to estimate baseline tissue  $PO_2$  using the radial measurements of tissue  $PO_2$  made by Vovenko (1999), whose vascular  $PO_2$  measurements define the model's reference state (see below). In addition, this approach allows us to predict the dynamic  $PO_2$  of tissue with a baseline  $PO_2$  value different to the average, as discussed below. For brevity, the notation used for most of this section is different to that of the main text.

Oxygen transport in an annular slice of tissue of radius  $R_2$  surrounding a blood vessel of radius  $R_1$  can be described by the differential equation

$$\frac{1}{r}\frac{d}{dr}\left(r\frac{dp}{dr}\right) = \frac{m}{D} \tag{S.6}$$

where r is the radial co-ordinate, p(r) is the oxygen partial pressure, and the constants m and D are the tissue consumption rate and diffusivity of oxygen, respectively. Solving this equation subject to the boundary conditions

$$p(R_1) = p_b; \quad \frac{dp}{dr}\Big|_{r=R_2} = 0$$
 (S.7)

leads to

$$p(r) = p_b + \frac{m}{4D} \left( r^2 - R_1^2 \right) - \frac{mR_2^2}{2D} \ln \frac{r}{R_1},$$
(S.8)

where the constant  $p_b$  is the PO<sub>2</sub> in the blood vessel.

Making the substitution  $k = \frac{m}{D}$  and applying the additional constraint  $p(R_2) = p_{t0}$ , where  $p_{t0}$  is the lowest partial pressure in the tissue, Equation (S.8) becomes

$$p(r) = p_b + \frac{1}{4}k\left(r^2 - R_1^2\right) - \frac{1}{2}kR_2^2\ln\frac{r}{R_1},$$
(S.9)

where

$$k = (p_{t0} - p_b)/\alpha \tag{S.10}$$

and

$$\alpha = \frac{1}{4} \left( R_2^2 - R_1^2 \right) - \frac{1}{2} R_2^2 \ln \frac{R_2}{R_1}.$$
 (S.11)

The average  $PO_2$  in the tissue area,  $\bar{p}$ , can be calculated by substituting

Equation (S.9) into the integral

$$\bar{p} = \frac{\int_{R_1}^{R_2} rp(r)dr}{\int_{R_1}^{R_2} rdr},$$
(S.12)

and solving analytically (by parts) to obtain

$$\bar{p} = p_b + \frac{1}{8}k\frac{R_2^4 - R_1^4}{R_2^2 - R_1^2} + \frac{1}{4}k\left(R_2^2 - R_1^2\right) - \frac{1}{2}k\frac{R_2^4}{R_2^2 - R_1^2}\ln\frac{R_2}{R_1}.$$
 (S.13)

Equation (S.13) can also be written as a weighted sum of the two boundary partial pressures, such that

$$\bar{p} = w_b p_b + (1 - w_b) p_{t0} \tag{S.14}$$

where  $w_b$ , the weight applied to the oxygen partial pressure in the blood, is given by

$$w_b = 1 - \frac{R_2^4 - R_1^4}{8\alpha \left(R_2^2 - R_1^2\right)} - \frac{1}{4\alpha} \left(R_2^2 - R_1^2\right) + \frac{R_2^4}{2\alpha \left(R_2^2 - R_1^2\right)} \ln \frac{R_2}{R_1}.$$
 (S.15)

For a tissue compartment made up of three of the annuli described by the preceding equations, the average oxygen partial pressure,  $\bar{p}_t$ , can be written as a weighted sum of the individual compartment averages such that

$$\bar{p}_t = \sum_{i=1}^3 w'_i \bar{p}_i,$$
 (S.16)

where  $w'_i$  is the weight applied to compartment *i* (weights must sum to 1), and  $\bar{p}_i$  is the average partial pressure, which is calculated from Equation (S.14). We used the baseline vascular volume fractions,  $v^*_i$ , as compartment weights  $(w'_i)$ . For simplicity, we assumed constant radii  $(R_1 \text{ and } R_2)$  in the different vascular compartments. On the basis of recent evidence from two-photon microscopy (Devor et al., 2011), we assumed the minimum PO<sub>2</sub> was the same for tissue surrounding each of the vascular compartments. As such, Equation (S.16) can also be written in terms of the vascular PO<sub>2</sub> and the minimum  $PO_2$  in the tissue so

$$\bar{p}_t = w_{t0}p_{t0} + \sum_{i=1}^3 w_i p_{bi},$$
 (S.17)

where  $w_i = w'_i w_b$ ,  $w_{t0} = 1 - \sum_{i=1}^3 w_i$ , and  $p_{bi}$  are the boundary partial pressures in the three vascular (blood) compartments. Using the same notation as the main text, Equation (S.17) would be written as

$$\bar{p}_{O_2,t}(t) = w_{t0} p_{O_2,t0}(t) + \sum_{i=1}^3 w_i \bar{p}_{O_2,i}(t)$$
(S.18)

#### Parameter Selection and Solution Procedure

This section gives more detail on the choice of parameters used in the model, and how the equations were solved to produce the predictions in the results section. Equation numbers not prefaced by 'S' refer to equations in the main text. Supplementary Table 1 summarizes the parameter values, with more detail given below.

In brief, the model was solved in three stages: 1) calculating the full set of parameters in a reference state; 2) adjusting the baseline parameters to account for experimental conditions different from the reference state (discussed in the next subsection); and 3) numerically solving the equations to produce the dynamic results. We chose to use the conditions observed by Vovenko (1999) as a reference state for the model, since these data provided nearly all of the measurements required to parametrize the model.

To determine the dimensionless tissue volume, we calculated the mean vascular volume fraction (2.87%) from data reported by a range of groups with different methodologies (An and Lin, 2002; Ito et al., 2005, 2001; Kim et al., 2007; Lauwers et al., 2008; Reichold et al., 2009; Weber et al., 2008). Since we scale volume terms by the total baseline vascular volume (Barrett et al., 2012), this means the dimensionless tissue volume,  $v_t$ , becomes 34.8.

First, we set  $[R_1, R_2] = [15, 135] \ \mu m$  on the basis of recent measurements of radial PO<sub>2</sub> in tissue (Devor et al., 2011), and included modifications to these parameters in the sensitivity analysis. Next, to determine the baseline value for the minimum tissue PO<sub>2</sub>,  $p_{O_2,t0}^*$ , we fit the model described in Equation (S.9) to radial measurements of tissue PO<sub>2</sub> in all three vascular compartments (Vovenko, 1999). Then, we used measurements of vascular PO<sub>2</sub> (Vovenko, 1999) to calculate baseline O<sub>2</sub> concentrations  $(c_{O_2,i,i+1}^*)$ from Equation (7), baseline O<sub>2</sub> amounts  $(n_{O_2,i,i+1}^*)$  and the leak concentration  $(c_{O_2,l})$  from Equation (3), and average vascular PO<sub>2</sub>  $(\bar{p}_{O_2,i}^*)$  from Equation (6). This makes it possible to calculate the average tissue PO<sub>2</sub>  $(\bar{p}_{O_2,t}^*)$  from Equation (S.18). The baseline oxygen consumption,  $cmr_{O_2}^*$ , was calculated from Equation (2).

We identified a feasible range of values for the shunt conduction coefficient,  $g_s$ , by requiring all conduction coefficients to be greater than zero to ensure that O<sub>2</sub> diffuses down its partial pressure gradient. The steady state forms of equations (1a) and (1c) then yield two inequalities that define the feasible range for  $g_s$ . Initial simulations (data not shown) suggested that the dynamics of PO<sub>2</sub> were not particularly sensitive to the value of  $g_s$ , so we defined  $g_s$  as the mean of the feasible range, but included values at the top and bottom 10% of the range in the sensitivity analysis. The remaining conduction coefficients were determined by solving Equations (1) at steady state. At this point, the model is fully parametrized in the reference state.

# **Adjusting Baseline Conditions**

This section describes the process of adjusting the parameters from the reference state conditions to the conditions under which the other data sets were obtained. All of the parameters adjusted under these different experimental conditions are listed in Supplementary Table 2.

To adjust the model from the reference state to the conditions reported by Masamoto et al. (2008), we changed only the value for femoral artery PO<sub>2</sub>  $(p_{O_2,0})$  as no other vascular PO<sub>2</sub> measurements were reported. As such, we assumed that baseline CMRO<sub>2</sub>  $(cmr_{O_2}^*)$  and the O<sub>2</sub> conduction coefficients  $(g \text{ and } g_s)$  were unchanged from the reference state. This means the steady state version of Equations (1) and (2) can be solved directly for the new baseline conditions.

For the simulations of data from Yaseen et al. (2011) and Vazquez et al. (2010), we adjusted the femoral artery PO<sub>2</sub> ( $p_{O_2,0}$ ), large venous PO<sub>2</sub> ( $p_{O_2,3,4}$ ), and average venous PO<sub>2</sub> ( $\bar{p}_{O_2,3}^*$ ) to match the reported measurements (see Table 2 in the main text). For all simulations, the concentration of oxygen 'leaked' between the femoral artery and large cerebral arteries,  $c_{O_2,l}$ , remained constant. Adding the steady state form of Equations (1) and (2), the adjusted baseline CMRO<sub>2</sub> is obtained such that

$$cmr_{O_2}^* = f^* \left( c_{O_{2,0,1}}^* - c_{O_{2,3,4}}^* \right),$$
 (S.19)

where  $f^*$  is the baseline CBF, and  $c_{O_2,0,1}^*$  and  $c_{O_2,3,4}^*$  are the baseline input arterial and output venous oxygen concentrations calculated from the imposed PO<sub>2</sub> values.

In both cases, the baseline state is underdetermined: the steady state form of Equations (1) and (2) provide 3 equations to determine 6 unknowns  $(g, g_s, p_{O_2,1,2})^*$  and  $\bar{p}_{O_2,t}$ . Therefore, it was necessary to make one or more assumptions about the baseline conditions. We describe the particular assumption used in more detail below, but during development we tested a number of plausible assumptions, which gave mostly similar results since the input (femoral artery PO<sub>2</sub>) and output (large venous PO<sub>2</sub>) values were defined by experimental data.

To determine the  $O_2$  conduction coefficients we assumed that, for the same femoral artery  $PO_2$ , the proportion of CMRO<sub>2</sub> supplied by each compartment at baseline would be constant between different experimental conditions. For example, for the conditions reported by Vazquez et al. (2010), we assumed that the steady state fraction of total oxygen flux supplied by each vascular compartment would be the same as those under the reference state conditions, if the femoral artery  $PO_2$  was reduced to that of the reference state (Vovenko, 1999) while maintaining the newly-calculated CMRO<sub>2</sub>.

Since CMRO<sub>2</sub> remains constant between these two conditions, the output venous O<sub>2</sub> concentration under the reference femoral artery PO<sub>2</sub> conditions,  $\dot{c}_{O_{2,3,4}}$ , can be calculated from a modified form of Equation (S.19)

$$\dot{c}_{O_{2,3,4}}^{*} = \dot{c}_{O_{2,0,1}}^{*} - cmr_{O_{2}}^{*}/f^{*}.$$
 (S.20)

Then,  $\dot{p}_{O_{2,3,4}}^{*}$  can be calculated from  $\dot{c}_{O_{2,3,4}}^{*}$  using Equation (7).

This leaves a system of 6 equations (the previously mentioned 3, plus the steady state form of Equations (1) and (2) under the reference femoral artery PO<sub>2</sub> conditions) and 9 unknowns (the previously mentioned 6, plus  $\dot{p}_{O_2,1,2}, \dot{p}_{O_2,2,3}, \text{ and } \dot{\bar{p}}_{O_2,t}$ ). The final 3 equations come from the assumption of constant fractional O<sub>2</sub> supply from each compartment, such that

$$\frac{g_i\left(\dot{\bar{p}}_{O_2,i} - \dot{\bar{p}}_{O_2,t}\right)}{cmr_{O_2}^*} = k_{O_2,i},\tag{S.21}$$

where the constant  $k_{O_2,i}$  is the proportion of CMRO<sub>2</sub> supplied from the compartment, and is calculated from the reference conditions. This assumption requires that the shunt conduction coefficient  $(g_s)$  be free to deviate from the mean of the feasible range; however, the value was determined to be within ~ 10% of the mean.

Finally, with the adjusted parameters fully specified, the four ordinary differential equations specified in Equations (1) and (2) were solved numerically to produce the dynamic predictions.

# Adjusting Tissue PO<sub>2</sub> Predictions

The methods section in the main text outlines the principle behind adjusting the tissue  $PO_2$  predictions and the rationale for doing so. However, in practice there are five steps required to produce these additional predictions.

First, the mean baseline tissue PO<sub>2</sub>,  $\bar{p}_{O_2,t}^*$ , is calculated by adjusting the reference state for different experimental conditions, as described in the previous section. Second, we choose new weights,  $w_x^*$ , to satisfy the equation

$$p_{O_2,t}^{\ \star} = w_{t0}^{\star} p_{O_2,t0}^{\ \star} + \sum_{i=1}^3 w_i^{\star} \bar{p}_{O_2,i}^{\ \star}, \tag{S.22}$$

where  $p_{O_2,t}$  is the measured tissue PO<sub>2</sub> at baseline. The new weights must

as

sum to one, but since Equation (S.22) is still underconstrained, the new weights are calculated by minimising the change from the original weights (i.e.  $w_x^{\star} - w_x$ ).

Third, the dynamic mean tissue PO<sub>2</sub>,  $\bar{p}_{O_2,t}(t)$ , is calculated by solving the ordinary differential equations in the main text, as discussed in the previous section. Fourth, the minimum tissue PO<sub>2</sub>,  $p_{O_2,t0}(t)$ , is calculated by rearranging Equation (S.18). Finally the additional dynamic tissue PO<sub>2</sub> prediction is calculated from

$$p_{O_2,t}^{*}(t) = w_{t0}^{*} p_{O_2,t0}(t) + \sum_{i=1}^{3} w_i^{*} \bar{p}_{O_2,i}(t).$$
(S.23)

In this way it is possible to estimate the dynamic PO<sub>2</sub> of tissue with any baseline value between the minimum tissue value,  $p_{O_2,t0}^*$ , and the maximum vascular (arterial) value,  $\bar{p}_{O_2,1}^*$ .

Steps three, four, and five are repeated during each iteration of the optimisation in order to determine the CMRO<sub>2</sub> parameters that generate the best fit between Equation (S.23) and the experimental observations. However, the new weights,  $w_x^*$ , are only calculated once to match the model-predicted and measured baseline tissue PO<sub>2</sub>, and are not used as fitting parameters in the dynamic simulations.

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Supplementary Table 1: List of model parameters in the reference state. Parameters with three values represent  $x = [x_1, x_2, x_3]$ . Parameters with lower case symbols are dimensionless (see the Supplementary Methods), while parameters with upper case symbols are not and are given with appropriate units.

Symbol	Description	Value	Reference
$C_{O_2,l}$	$O_2$ conc. leakage	$0.116 \mathrm{~mM}$	See Supplementary Methods
$C_{O_2,max}$	Hill equation max. $O_2$ conc.	$9.26 \mathrm{~mM}$	Cartheuser (1993)
$\operatorname{prop}(g_s)$	Proportion of shunt feasible range	50%	See Supplementary Methods
h	Hill equation exponent	2.6	Cartheuser $(1993)$
$P_{50}$	Hill equation $O_2 P_{50}$	$36 \mathrm{~mmHg}$	Cartheuser $(1993)$
$\Delta P_{SNP}$	SNP-induced pressure drop	50%	Masamoto et al. $(2008)$
$[R_1, R_2]$	Krogh cylinder radii	$[15, 135] \ \mu m$	See Supplementary Methods
$v^*$	Baseline vascular volume fraction	[0.29, 0.44, 0.27]	Barrett et al. (2012)
$v_t$	Tissue volume fraction	34.8	See Supplementary Methods
$w_b$	Vascular $PO_2$ weight	0.133	See Supplementary Methods
$w'_i$	Vascular compartment weight	$v^*$	See Supplementary Methods
$\sigma_{O_2}$	Tissue $O_2$ solubility coefficient	1.46 $\mu {\rm M/mmHg}$	Dash and Bassingthwaighte (2004)

12

2 in the main text). Values give by '-' are the same as the reference state (Vovenko).

 Symbol
 Description
 Unit
 Value for simulations of:

 Vovenko
 (1999)
 Vascent
 (2010)

Supplementary Table 2: List of adjusted model parameters for simulations of experimental conditions (see Table

Symbol	Description	Unit -	Value for simulations of:			
Symbol			Vovenko (1999)	Yaseen $(2011)$	Masamoto (2008)	Vazquez (2010)
$P_{O_2,0,1}^{*}$	Baseline input art. $PO_2$	mmHg	81.2	99.4	103.7	116.7
$P_{O_2,1,2}^{*}$	Baseline art.–cap. $PO_2$	mmHg	59.7	68.0	62.3	55.7
$P_{O_2,2,3}$	Baseline cap.–vei. $PO_2$	$\rm mmHg$	39.6	50.1	41.1	35.3
$P_{O_{2,3,4}}^{*}$	Baseline output vei. $PO_2$	$\rm mmHg$	41.3	54.4	44.6	40.3
$\bar{P}_{O_2,t}^{*}$	Mean baseline tis. $PO_2$	$\rm mmHg$	22.4	38.1	25.3	22.8
$cmr^*_{O_2}$	Baseline $CMRO_2$	none	0.336	0.208	-	0.423
$g_1$	Art. $O_2$ conduction coef.	none	0.075	0.059	-	0.096
$g_2$	Cap. $O_2$ conduction coef.	none	0.790	0.619	-	1.124
$g_3$	Vei. $O_2$ conduction coef.	none	0.201	0.155	-	0.309
$g_s$	Shunt $O_2$ conduction coef.	none	0.207	0.198	-	0.263

Supplementary Table 3: List of parameters modified in the sensitivity analysis and the amount of perturbation imposed. For parameters with multiple values (e.g.  $p_{O_2}^*$ ), each value was perturbed individually. Perturbations to the reference conditions also affect simulations under modified experimental conditions.

Data Set	Symbol	Perturbation
	$p_{O_2}^{*}$	$\pm 10\%$
	$[R_1, R_2]$	$\pm 10\%$
Reference	$\operatorname{prop}(g_s)$	$\pm 40\%$
	$v^*$	$\pm 10\%$
	$v_t$	$\pm 10\%$
Magamoto et al. (2008)	$\Delta P_{SNP}$	$\pm 10\%$
Wasamoto et al. (2000)	$p_{O_2,0}^{*}$	$\pm 10\%$
	$p_{O_2,0}^{*}$	$\pm 10\%$
Vazquez et al. $(2010)$	$p_{O_{2,3,4}}^{*}$	$\pm 5\%$
	$\bar{p}_{O_{2,3}}^{*}$	$\pm 5\%$



Supplementary Figure 1: Model predictions (with additional mechanisms) of data from Masamoto et al. (2008) in response to 10s electrical forepaw stimulation. Format as per Figure 3 in the main text.



Supplementary Figure 2: Model predictions (with additional mechanisms) of data from Vazquez et al. (2010) in response to 20s electrical forepaw stimulation. Format as per Figure 4 in the main text.