Quantum and All-Atom Molecular Dynamics Simulations of Protonation and Divalent Ion Binding to Phosphatidylinositol 4,5-Bisphosphate (PIP₂)

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Supporting Information

Using the Gaussian '09 program, the molecular geometry of PIP_2 was obtained using the Hartree-Fock scheme in conjunction with the 6-31+G(d) basis set. The initial structure contained a proton placed equidistant between oxygen atoms from the 4-phosphate and 5-phosphate groups. The final Z-matrix is as follows:

С			
С	1 cc2		
С	2 cc3	1 ccc3	
С	3 cc4	2 ccc4	
С	4 cc5	3 ccc5	
С	5 cc6	4 ccc6	
0	2 oc7	3 occ7	
р	7 po8	2 poc8	
0	8 op9	7 opo9	
С	9 co10	8 cop1	
С	10 cc11	9 cco1	
С	11 cc12	10 ccc1	
0	12 oc13	11 occ1	
С	13 co14	12 coc1	
0	14 oc15	13 oco1	
0	3 oc16	2 occ1	
0	4 oc17	3 occ1	
0	5 oc18	4 occ1	
р	18 po19	5 poc1	
0	19 op20	18 opo2	
0	6 oc21	5 occ2	
р	21 po22 22 op23	6 poc2 21 opo2	
0	22 op23 1 oc24	21 opo2 2 occ2	
0 0	8 op25	2 0002 7 opo2	
0	8 op26	7 opo2 7 opo2	
0	19 op20	18 opo2	
0	19 op27 19 op28	18 opo2 18 opo2	
0	22 op29	21 opo2	
0	22 op20 22 op30	21 op02 21 op03	
0	11 oc31	10 occ3	
c	31 co32	11 coc3	
С	32 cc33	31 cco3	
0	32 oc34	31 oco3	
С	14 cc35	13 cco3	
h	16 ho36	3 hoc3	
h	17 ho37	4 hoc3	
h	24 ho38	1 hoc3	
h	2 hc39	7 hco3	
h	3 hc40	2 hcc4	
h	4 hc41	3 hcc4	1 2 dih41
h	5 hc42	18 hco4	2 19 dih42
h	6 hc43	21 hco4	3 22 dih43
h	1 hc44	2 hcc4	
h	33 hc45	32 hcc4	5 31 dih45
h	35 hc46	14 hcc4	6 13 dih46
h	29 ho47	22 hop4	7 21 dih47
h	11 hc48	31 hco4	
h	33 hc49	32 hcc4	9 31 dih49
h	33 hc50	32 hcc5	0 31 dih50
h	35 hc51	14 hcc5	1 13 dih51

h		hc52		hcc52		dih52
h		hc53 hc54	11 11	hcc53	31 31	
h h		hc55		hcc54 hcc55	31	
h		hc56	11	hcc56	31	
		1 5060				
cc2 cc3		1.52634				
ccc3		110.464				
cc4		1.53442				
ccc4		107.97				
dih4 cc5		-63.220				
ccc5		109.990				
dih5		65.784	1			
ссб		1.55042				
ccc6 dih6		112.130 -54.29				
oc7		1.4304				
occ7		114.252				
dih7		178.78				
po8 poc8		1.6098 129.570				
dih8		-76.014				
op9		1.65455				
opo9		100.649				
dih9 co10		89.510 1.39049				
cop1	C	118.33				
dih1		56.942	2			
cc11	1	1.5251				
ccol: dih1:		108.169				
cc12	-	1.51660				
ccc1		112.07				
dih12 oc13	2	53.050 1.44462				
occ1	3	109.903				
dih1	3	169.423				
col4 cocl	л	1.30780				
dih1		88.633				
oc15		1.19459				
ocol!		126.101				
dih1: oc16	C	1.91(1.4158)				
occ1	6	111.903				
dih1		-54.700				
oc17		1.39634				
occ1 dih1		108.994 -171.724				
oc18	,	1.3975				
occ1		116.11				
dih18	8	178.192				
po19 poc1	9	1.72429				
dih1		30.224				
op20	~	1.5089				
opo2) dih2)		104.741 50.791				
oc21	0	1.42458				
occ2		112.14				
dih2: po22	L	-64.455				
pozz poc22	2	126.24				
dih22		-43.03	6			
op23	2	1.49699				
opo2: dih2:		109.001 -30.862				
oc24		1.41810	64			
occ2	4	109.01	5			

dih24	-63.339
op25	1.484818
opo25	106.531
dih25	-159.329
	1 474500
op26	1.474582
opo26	112.354
dih26	-22.804
op27	1.504387
opo27	102.301
-	172.140
dih27	173.149
op28	1.529475
opo28	102.611
dih28	-68.029
op29	1.621747
opo29	100.814
dih29	84.774
op30	1.485743
opo30	108.720
dih30	-162.053
oc31	1.445320
	104.664
occ31	
dih31	170.387
co32	1.300480
coc32	123.522
dih32	142.520
cc33	1.510033
cco33	110.400
dih33	176.609
oc34	1.200221
oco34	126.449
dih34	-3.727
cc35	1.510545
cco35	110.753
dih35	-177.884
ho36	0.955083
hoc36	111.918
dih36	-68.963
ho37	0.977482
hoc37	105.034
dih37	170.673
ho38	0.948667
hoc38	105.953
dih38	52.346
hc39	1.082596
hco39	108.137
dih39	45.153
hc40	1.082127
hcc40	109.459
dih40	61.497
hc41	1.086439
hcc41	107.608
dih41	-53.191
hc42	1.080006
hco42	104.771
dih42	145.460
hc43	1.080634
hco43	105.331
dih43	-159.036
hc44	1.080878
hcc44	107.794
dih44	54.135
hc45	1.083874
hcc45	109.720
	±0,7.720
dih45	56.197
hc46	1.083719
hcc46	110.223
dih46	-49.020
ho47	0.961524
hop47	112.864
dih47	-26.260
hc48	1.079439

hc49 1 hcc49 1 dih49 - hc50 1 hcc50 1 hcc51 1 hcc51 1 hcc52 1 hcc52 1 hcc52 1 hcc52 1 hcc53 1 hcc53 1 hcc53 1 hcc53 1 hcc54 1 hcc55 1 hcc55 1 hcc55 1 hcc55 1 hcc55 1 hcc55 1 hcc56 1 hcc56 1	22.809 .084019 09.203 61.054 .082332 10.002 77.802 .082120 09.751 70.839 .084700 08.552 69.339 .086291 09.176 68.761 .081245 08.631 49.990 .075921 10.080 75.558 .077606 11.112 62.016
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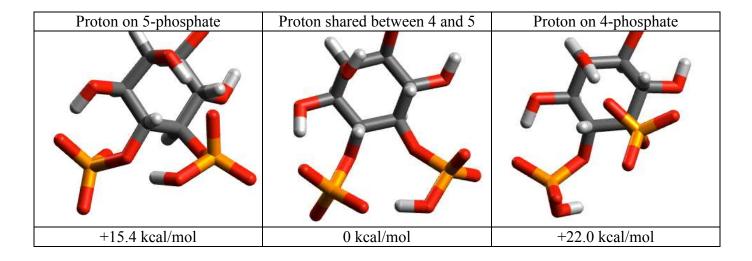


Figure S1. Relative energies of protonation isomers. The relative stabilities of protonated PIP₂ isomers were computed with Gaussian 09 at the HF/3-21G level of theory with implicit water solvent using a Polarizable Continuum Model¹⁻³. Relaxed dihedral potential energy scans were used to identify energetic minima among isomers. In the lowest energy structure (middle panel), the proton is shared between the 4 and 5-phosphates. The distances from the proton to the 5 and 4-phosphate oxygens are 1.11 Å and 1.30 Å, respectively. This shared state is 15.4 kcal/mol and 22.0 kcal/mol more stable than localization to the 5-phosphate (left panel) or the 4-phosphate (right panel), respectively.

Our results are obtained from all-atom molecular dynamics simulations of membrane bilayers of 800 total lipids containing PIP2 in the presence of phosphatidylcholine (PC), phosphatidylserine (PS), phosphatidylethanolamine (PE), and cholesterol with 150 mM NaCl in a solvent box containing approximately 80,000 waters. The outer leaflet contains 300 POPC [3-palmitoyl-2-oleoyl-D-glycero-1-phosphatidylcholine] molecules and 100 cholesterol molecules. The inner leaflet contains 200 DOPE [2,3-dioleyl-D-glycero-1phosphatidylethanolamine] molecules, 100 cholesterol molecules, 60 DOPS [2,3-dioleyl-D-glycero-1-phosphatidylserine] molecules, 40 SA-PIP2 [1-stearoyl-2-arachadonyl-snglycero-1-phosphatidylinositol-4,5-bisphosphate] molecules. The lipids were packed at approximately 65 $Å^2$ per lipid and randomly seeded with PIP₂ typically about 10% of the total lipids on one leaflet. The system was equilibrated for 10 ns before the production run, for which the analysis is presented below. We used the CHARMM36 lipid forcefield in the GROMACS simulation package⁴⁻⁷ to prepare membrane bilayers in an NPT ensemble and have simulated for 31 ns so far, using a 2 fs timestep. Pressure is maintained using the Parrinello-Rahman barostat^{8,9}, coupled independently (semiisotropic) in directions normal and perpendicular to the bilayer with a time constant of 2.0 ps and a pressure of 1 bar. Electrostatics are treated with the particle mesh Ewald (PME)¹⁰ method with cubic periodic boundaries, using a real space cutoff of 1.2 nm and a Fourier grid spacing of 0.16 nm with cubic interpolation. Temperature coupling is accomplished using the thermostat due to Bussi¹¹, in which velocity is rescaled with a stochastic term. These simulations employ the CHARMM-specific TIP3P (i.e., TIPS3P) water model (with van der Waals interactions on the water hydrogen atoms) and van der Waals interactions which are smoothly switched off between 0.8 and 1.2 nm^{12} .

Histograms were computed (Figure S2) for the QM/MM simulations presented in the main text (a single PIP2 molecule in a water sphere) and for each lipid every 20 frames of the classical all-atom trajectory. The head-tail angle was measured as described in the main text. Our QM/MM and classical bilayer simulation results quantitatively agree with each other, but disagree if we employ the parameters of Lupyan, et al.

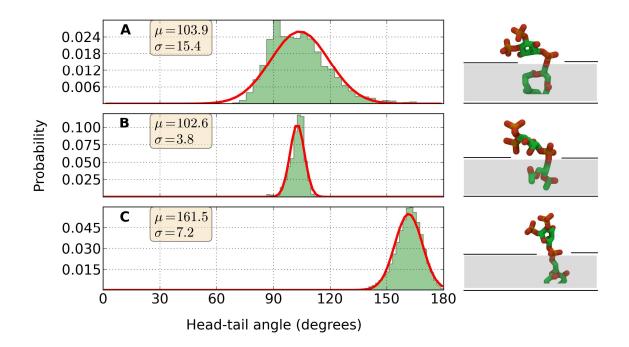


Figure S2. **A**. The probability distribution for 80 PIP₂ lipids in one leaflet of a bilayer also containing phosphatidylserine, phosphatidylethanolamine, and cholesterol. The histogram was averaged over each lipid and every 20 frames of the all-atom simulation for 15 ns using CHARMM36 lipid parameters. The mean head-tail angle is 104 degrees. **B**. The probability distribution of the angle between the PIP₂ headgroup and the PIP₂ acyl chains as measured from an isolated PIP₂ molecule in a water sphere over 10 ps of a QM/MM simulation. Sodium ions were present for charge neutrality but fixed to be ~10 Å away from the PIP₂. The mean head-tail angle is 103 degrees. **C**. The probability distribution from a simulation of a bilayer containing pure PIP₂ using the same system parameters as panel A, but with the parameters of Lupyan, et al. used instead of the CHARMM36 lipid parameters¹³. The average head-tail angle is 162 degrees.

We computed the average PIP₂ molecular area in membrane bilayer systems to compare with the QM/MM results presented in the main text. Since there was only a single molecule in the QM/MM simulations, we used the separation of the phosphate groups as a proxy for molecular area. To calculate the molecular area in the bilayer simulations, we projected the center of gravity of each lipid onto a plane and created a Delaunay triangulation. We summed the area created by forming a polygon from the bisector of each line meeting a vertex corresponding to PIP₂. There are subtleties in defining the average molecular area of lipids in a heterogeneous bilayer simulation and these methods are still

evolving.

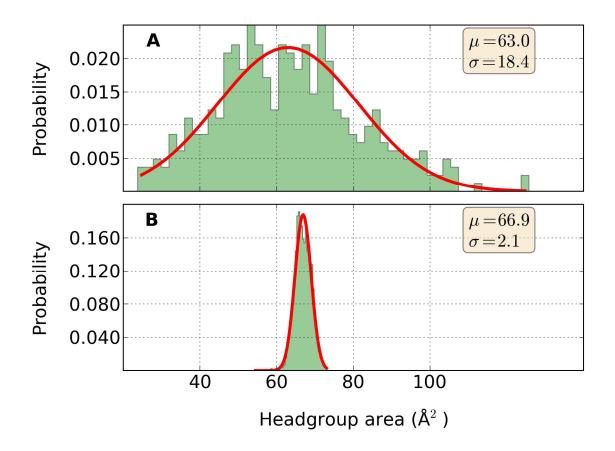


Figure S3. **A**. The probability distribution for 80 PIP₂ lipids in a bilayer also containing phosphatidylcholine, phosphatidylserine, and phosphatidylethanolamine. **B**. The probability distribution of the approximate molecular area of PIP₂ as measured from an isolated PIP₂ molecule in a water sphere over 10 ps of a QM/MM simulation. Sodium were present for charge neutrality but fixed to be ~10 Å away from the PIP₂.

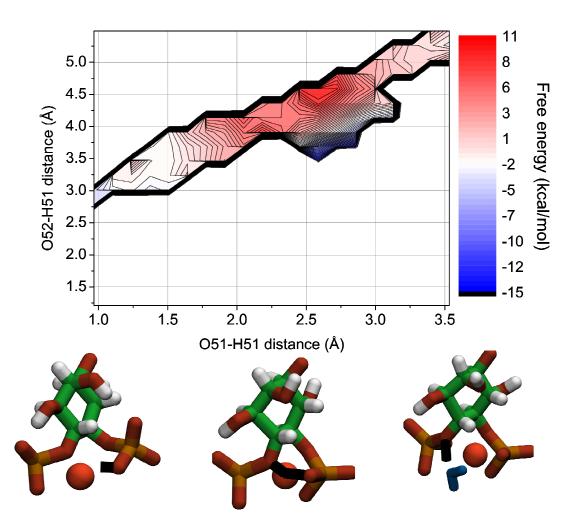


Figure S4. The free energy landscape for dissociating a proton from the 5-phosphate group of PIP₂. Full deprotonation in the presence of magnesium has not been achieved. Insets show magnesium (in red) between the two phosphomonoester groups of PIP₂ and the proton to be dissociated in black. The initial state of the QM/MM 2D umbrella sampling is on the left, with proton H51 bound to oxygen O51. The inset in the middle corresponds to the energy well on the contour plot (blue), where proton H51 is shared between oxygen O51 and oxygen O4P. The inset on the right corresponds to the final window of the trajectory, where proton H51 is shared between oxygen O4P of PIP₂ and a QM water (in blue).

It is possible to calculate pKa values from the free energy plots presented in the main text. To do this, we consider the canonical definition of the association constant, K_a , between a ligand, L, and its binding partner, R, in terms of the ratio of the probabilities of being not bound or bound. In this case, we take the ligand to be a proton and its binding partner to be a PIP₂ molecule. So we write,

$$K_a = \frac{1}{[L]} \times \frac{\mathcal{P}_{\text{scand}}}{\mathcal{P}_{\text{unbound}}}$$

where \mathcal{P}_{hound} and $\mathcal{P}_{unbnurd}$ can be determined by integrating over the degrees of freedom and the potentials of mean force. pKa is related to K using the following relationship: pKa = $\log_{10} \left| K_a \times \frac{1}{1661 \, \text{Å}^3} \right|$ using the standard concentration of 1 mol/liter = 1/1661 Å³.

We can perform the pKa calculation in two ways. In the first method, we integrate the bound state until $r = r^*$ which is the point where the proton leaves the domain of influence of the oxygen to which it was originally bound. This is judged by viewing the simulation trajectories and looking at the minima of the potentials of mean force. The unbound state is integrated from $r = r^*$ until the final point in the energy landscape.

In the second method, we integrate the bound state until $r = r^*$ as before. To evaluate the unbound state, we integrate from $r = r^*$ until the next metastable minimum at $r = r^{**}$. In this technique, events that we categorize as dissociation, such as proton transfer to other atoms on PIP₂, may not correspond to dissociation events judged from experiments. Values for r^* and r^{**} and the results of these calculations are shown in Figure S5. Error

bars in the pKa calculation represent the outcome of changing the value of r^* by 0.25 Å or approximately the width of a well in the potential of mean force.

We stress that the absolute value of these numbers should not be compared to the results of experiments due to the approximations necessary in obtaining absolute free energy curves from the potentials of mean force calculated along a single reaction coordinate. Kooijman et al. calculate the pKa of PIP₂ in lipid vesicles, which is different from our case in several respects (e.g., the effect of a negative surface charge density of the vesicle membrane) so a direct comparison is not strictly valid¹⁴. Moreover, the pKa values reported by Kooijman et al. correspond to $PI(3,5)P_2$ and $PI(3,4,5)P_2$ which could have different ionization properties due to the positioning of the phosphate groups. In fact, for $PI(4,5)P_2$, Kooijman et al. note that a biphasic pH-dependent ionization behavior that cannot be explained by a Henderson-Hasselbach equation was obtained. Our calculations are consistent with this; our free energy landscape does not simply represent two states (bound or unbound) but displays intermediate metastable states. These states correspond to sharing of the proton between the vicinal phosphomonoester groups. Hence, our pKa should be regarded as an apparent pKa for proton transfer from the host oxygen (O51 in the main text) to the next energetically favorable state (rather than a free unbound state).

The subtleties caused by a complex (rugged) free energy landscape and nonindependence of proton binding to vicinal sites cause problems in defining an absolute pKa; moreover, the apparent pKa is subject to the definitions of vicinal binding sites as metastable states (or our choice of $r^{\star\star}$). Another complexity is a computational one. In the QM/MM potentials of mean force, the energy landscape is only accurate for short

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excursions from the initial bound state along the reaction coordinate. For longer excursions, the wave functions from the QM region are spread over larger volumes causing an imbalance in basis set superposition. Also for these large distances along the reaction coordinate, close interaction between classical charges and the QM wave function become significant.

In general, we find that the pKa for protonation of PIP_2 in the presence of calcium is the smallest, implying strong dissociation of the proton. The pKa in the presence of magnesium is similar to that of a pure water environment, where we find protonation of PIP_2 on the 5-phosphate group to be the most stable conformation.

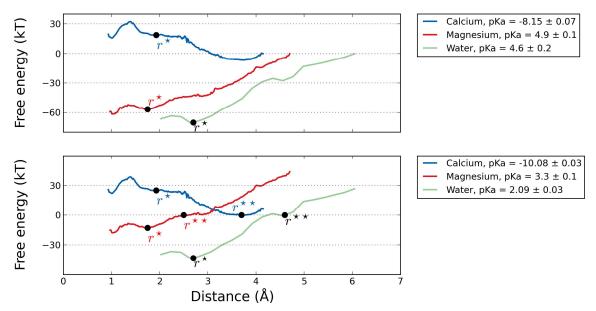


Figure S5. The potentials of mean force for proton dissociation from PIP_2 in a water sphere with either calcium, magnesium, or pure water. In the latter case, sodium ions are present for charge neutrality but confined ~10 Å away from the PIP_2 .

- (1) Miertuš, S.; Scrocco, E.; Tomasi, J. Electrostatic Interaction of a Solute with a Continuum. A Direct Utilization of Ab Initio Molecular Potentials for the Prevision of Solvent Effects *Chem. Phys.* **1981**, *55*, 117–129.
- (2) Pascual-ahuir, J. L.; Silla, E.; Tu on, I. GEPOL: an Improved Description of Molecular Surfaces. III. a New Algorithm for the Computation of a Solvent-Excluding Surface *J. Comp. Chem.* **1994**, *15*, 1127–1138.
- (3) Tomasi, J.; Mennucci, B.; Cammi, R. Quantum Mechanical Continuum Solvation Models. *Chem. Rev.* **2005**, *105*, 2999–3093.
- (4) Van Der Spoel, D.; Lindahl, E.; Hess, B.; Groenhof, G.; Mark, A. E.; Berendsen, H. J. C. GROMACS: Fast, Flexible, and Free J. Comp. Chem. 2005, 26, 1701–1718.
- (5) Bjelkmar, P.; Larsson, P.; Cuendet, M. A.; Hess, B.; Lindahl, E. Implementation of the CHARMM Force Field in GROMACS: Analysis of Protein Stability Effects From Correction Maps, Virtual Interaction Sites, and Water Models J. *Chem. Theo. Comput.* 2010, *6*, 459–466.
- (6) Klauda, J. B.; Monje, V.; Kim, T.; Im, W. Improving the CHARMM Force Field for Polyunsaturated Fatty Acid Chains. *J. Phys. Chem. B* **2012**, *116*, 9424–9431.
- Klauda, J. B.; Venable, R. M.; Freites, J. A.; O'Connor, J. W.; Tobias, D. J.;
 Mondragon-Ramirez, C.; Vorobyov, I.; Mackerell, A. D., Jr; Pastor, R. W.
 Update of the CHARMM All-Atom Additive Force Field for Lipids: Validation on Six Lipid Types. *J. Phys. Chem. B* 2010, *114*, 7830–7843.
- (8) Parrinello, M.; Rahman, A. Polymorphic Transitions in Single Crystals: a New Molecular Dynamics Method *Journal of Applied physics* **1981**.
- (9) Nose, S.; Klein, M. L. Constant Pressure Molecular Dynamics for Molecular Systems *Mol. Phys.* **1983**.
- (10) Darden, T.; York, D.; Pedersen, L. Particle Mesh Ewald: an NT Log(N) Method for Ewald Sums in Large Systems *J. Chem. Phys.* **1993**, *98*, 10089.
- (11) Bussi, G.; Donadio, D.; Parrinello, M. Canonical Sampling Through Velocity Rescaling *J. Chem. Phys.* **2007**, *126*, 014101.
- (12) Piggot, T. J.; Piñeiro, Á.; Khalid, S. Molecular Dynamics Simulations of Phosphatidylcholine Membranes: a Comparative Force Field Study *J. Chem. Theo. Comput.* **2012**, *8*, 4593–4609.
- (13) Lupyan, D.; Mezei, M.; Logothetis, D. E.; Osman, R. A Molecular Dynamics Investigation of Lipid Bilayer Perturbation by PIP2. *Biophys. J.* **2010**, *98*, 240–247.
- (14) Kooijman, E. E.; King, K. E.; Gangoda, M.; Gericke, A. Ionization Properties of Phosphatidylinositol Polyphosphates in Mixed Model Membranes *Biochemistry* 2009, 48, 9360–9371.