Quantum and All-Atom Molecular Dynamics Simulations of Protonation and Divalent Ion Binding to Phosphatidylinositol 4,5- Bisphosphate (PIP₂)

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Supporting Information

Using the Gaussian '09 program, the molecular geometry of PIP_2 was obtained using the Hartree-Fock scheme in conjunction with the $6-31+G(d)$ basis set. The initial structure contained a proton placed equidistant between oxygen atoms from the 4-phosphate and 5phosphate groups. The final Z-matrix is as follows:

Figure S1. Relative energies of protonation isomers. The relative stabilities of protonated $PIP₂$ isomers were computed with Gaussian 09 at the HF/3-21G level of theory with implicit water solvent using a Polarizable Continuum Model¹⁻³. Relaxed dihedral potential energy scans were used to identify energetic minima among isomers. In the lowest energy structure (middle panel), the proton is shared between the 4 and 5 phosphates. The distances from the proton to the 5 and 4-phosphate oxygens are 1.11 Å and 1.30 Å, respectively. This shared state is 15.4 kcal/mol and 22.0 kcal/mol more stable than localization to the 5-phosphate (left panel) or the 4-phosphate (right panel), respectively.

Our results are obtained from all-atom molecular dynamics simulations of membrane bilayers of 800 total lipids containing PIP2 in the presence of phosphatidylcholine (PC), phosphatidylserine (PS), phosphatidylethanolamine (PE), and cholesterol with 150 mM NaCl in a solvent box containing approximately 80,000 waters. The outer leaflet contains 300 POPC [3-palmitoyl-2-oleoyl-D-glycero-1-phosphatidylcholine] molecules and 100 cholesterol molecules. The inner leaflet contains 200 DOPE [2,3-dioleyl-D-glycero-1 phosphatidylethanolamine] molecules, 100 cholesterol molecules, 60 DOPS [2,3-dioleyl-D-glycero-1-phosphatidylserine] molecules, 40 SA-PIP2 [1-stearoyl-2-arachadonyl-*sn*glycero-1-phosphatidylinositol-4,5-bisphosphate] molecules. The lipids were packed at approximately 65 A^2 per lipid and randomly seeded with PIP₂ typically about 10% of the total lipids on one leaflet. The system was equilibrated for 10 ns before the production run, for which the analysis is presented below. We used the CHARMM36 lipid forcefield in the GROMACS simulation package⁴⁻⁷ to prepare membrane bilayers in an NPT ensemble and have simulated for 31 ns so far, using a 2 fs timestep. Pressure is maintained using the Parrinello-Rahman barostat $8,9$, coupled independently (semiisotropic) in directions normal and perpendicular to the bilayer with a time constant of 2.0 ps and a pressure of 1 bar. Electrostatics are treated with the particle mesh Ewald $(PME)^{10}$ method with cubic periodic boundaries, using a real space cutoff of 1.2 nm and a Fourier grid spacing of 0.16 nm with cubic interpolation. Temperature coupling is accomplished using the thermostat due to Bussi¹¹, in which velocity is rescaled with a stochastic term. These simulations employ the CHARMM-specific TIP3P (i.e., TIPS3P) water model (with van der Waals interactions on the water hydrogen atoms) and van der Waals interactions which are smoothly switched off between 0.8 and 1.2 nm^{12} .

Histograms were computed (Figure S2) for the QM/MM simulations presented in the main text (a single PIP2 molecule in a water sphere) and for each lipid every 20 frames of the classical all-atom trajectory. The head-tail angle was measured as described in the main text. Our QM/MM and classical bilayer simulation results quantitatively agree with each other, but disagree if we employ the parameters of Lupyan, et al.

Figure S2. A. The probability distribution for 80 PIP₂ lipids in one leaflet of a bilayer also containing phosphatidylserine, phosphatidylethanolamine, and cholesterol. The histogram was averaged over each lipid and every 20 frames of the all-atom simulation for 15 ns using CHARMM36 lipid parameters. The mean head-tail angle is 104 degrees. **B.** The probability distribution of the angle between the PIP_2 headgroup and the PIP_2 acyl chains as measured from an isolated PIP_2 molecule in a water sphere over 10 ps of a QM/MM simulation. Sodium ions were present for charge neutrality but fixed to be \sim 10 Å away from the PIP2. The mean head-tail angle is 103 degrees. **C.** The probability distribution from a simulation of a bilayer containing pure PIP_2 using the same system parameters as panel A, but with the parameters of Lupyan, et al. used instead of the CHARMM36 lipid parameters¹³. The average head-tail angle is 162 degrees.

We computed the average PIP_2 molecular area in membrane bilayer systems to compare with the QM/MM results presented in the main text. Since there was only a single molecule in the QM/MM simulations, we used the separation of the phosphate groups as a proxy for molecular area. To calculate the molecular area in the bilayer simulations, we projected the center of gravity of each lipid onto a plane and created a Delaunay triangulation. We summed the area created by forming a polygon from the bisector of each line meeting a vertex corresponding to PIP₂. There are subtleties in defining the average molecular area of lipids in a heterogeneous bilayer simulation and these methods are still

evolving.

Figure S3. A. The probability distribution for 80 PIP₂ lipids in a bilayer also containing phosphatidylcholine, phosphatidylserine, and phosphatidylethanolamine. **B**. The probability distribution of the approximate molecular area of PIP_2 as measured from an isolated PIP_2 molecule in a water sphere over 10 ps of a QM/MM simulation. Sodium were present for charge neutrality but fixed to be \sim 10 Å away from the PIP₂.

Figure S4. The free energy landscape for dissociating a proton from the 5-phosphate group of PIP2. Full deprotonation in the presence of magnesium has not been achieved. Insets show magnesium (in red) between the two phosphomonoester groups of PIP_2 and the proton to be dissociated in black. The initial state of the QM/MM 2D umbrella sampling is on the left, with proton H51 bound to oxygen O51. The inset in the middle corresponds to the energy well on the contour plot (blue), where proton H51 is shared between oxygen O51 and oxygen O4P. The inset on the right corresponds to the final window of the trajectory, where proton H51 is shared between oxygen O4P of $PIP₂$ and a QM water (in blue).

It is possible to calculate pKa values from the free energy plots presented in the main text. To do this, we consider the canonical definition of the association constant, K_a , between a ligand, L, and its binding partner, R, in terms of the ratio of the probabilities of being not bound or bound. In this case, we take the ligand to be a proton and its binding partner to be a PIP_2 molecule. So we write,

$$
K_a = \frac{1}{[L]} \times \frac{\mathcal{P}_{scalar}}{\mathcal{P}_{unbound}}
$$

where P_{bound} and P_{unburn} can be determined by integrating over the degrees of freedom and the potentials of mean force. pKa is related to K using the following relationship: using the standard concentration of 1 mol/liter = $1/1661$ Å³.

We can perform the pKa calculation in two ways. In the first method, we integrate the bound state until $r = r^*$ which is the point where the proton leaves the domain of influence of the oxygen to which it was originally bound. This is judged by viewing the simulation trajectories and looking at the minima of the potentials of mean force. The unbound state is integrated from $r = r^*$ until the final point in the energy landscape.

In the second method, we integrate the bound state until $r = r^*$ as before. To evaluate the unbound state, we integrate from $r = r^*$ until the next metastable minimum at $r = r^{**}$. In this technique, events that we categorize as dissociation, such as proton transfer to other atoms on PIP2, may not correspond to dissociation events judged from experiments. Values for r^* and r^{**} and the results of these calculations are shown in Figure S5. Error

bars in the pKa calculation represent the outcome of changing the value of r^{\star} by 0.25 Å or approximately the width of a well in the potential of mean force.

We stress that the absolute value of these numbers should not be compared to the results of experiments due to the approximations necessary in obtaining absolute free energy curves from the potentials of mean force calculated along a single reaction coordinate. Kooijman et al. calculate the pKa of PIP_2 in lipid vesicles, which is different from our case in several respects (e.g., the effect of a negative surface charge density of the vesicle membrane) so a direct comparison is not strictly valid¹⁴. Moreover, the pKa values reported by Kooijman et al. correspond to $PI(3,5)P_2$ and $PI(3,4,5)P_2$ which could have different ionization properties due to the positioning of the phosphate groups. In fact, for $PI(4,5)P_2$, Kooijman et al. note that a biphasic pH-dependent ionization behavior that cannot be explained by a Henderson-Hasselbach equation was obtained. Our calculations are consistent with this; our free energy landscape does not simply represent two states (bound or unbound) but displays intermediate metastable states. These states correspond to sharing of the proton between the vicinal phosphomonoester groups. Hence, our pKa should be regarded as an apparent pKa for proton transfer from the host oxygen (O51 in the main text) to the next energetically favorable state (rather than a free unbound state).

The subtleties caused by a complex (rugged) free energy landscape and nonindependence of proton binding to vicinal sites cause problems in defining an absolute pKa; moreover, the apparent pKa is subject to the definitions of vicinal binding sites as metastable states (or our choice of $r^{\star\star}$). Another complexity is a computational one. In the QM/MM potentials of mean force, the energy landscape is only accurate for short

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excursions from the initial bound state along the reaction coordinate. For longer excursions, the wave functions from the QM region are spread over larger volumes causing an imbalance in basis set superposition. Also for these large distances along the reaction coordinate, close interaction between classical charges and the QM wave function become significant.

In general, we find that the pKa for protonation of PIP_2 in the presence of calcium is the smallest, implying strong dissociation of the proton. The pKa in the presence of magnesium is similar to that of a pure water environment, where we find protonation of PIP₂ on the 5-phosphate group to be the most stable conformation.

Figure S5. The potentials of mean force for proton dissociation from PIP_2 in a water sphere with either calcium, magnesium, or pure water. In the latter case, sodium ions are present for charge neutrality but confined \sim 10 Å away from the PIP₂.

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