## **Supporting Information**

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SI Text

## **RNA Construct and Reagents**

The model tetraloop-receptor RNA construct (Fig. 1) is assembled by annealing together three pieces of RNA that were purchased from IDT: (i) a DNA surface tether strand, 5'-biotin-CGC ACT CGT CTG AG-3'; (ii) a Cy3-labeled tetraloop-linkerhelix strand, 5'-Cy3-GGC GAA AGC C-PEG<sub>6</sub>-CGU GUC GUC CUA AGU CGG C-3'; and (iii) a Cy5-labeled helix strand, 5'-Cy5-GCC GAU AUG GAC GAC ACG CCC CUC AGA CGA GUG CG-3'. Sandwich-style sample holders were prepared by making a channel between a glass slide and a coverslip separated by two strips of double-sided tape. The channel is flushed with a 10:1 mixture of BSA:BSA-biotin at ~1 mg/mL total protein concentration, followed by streptavidin (~0.1 mg/mL) and lastly the tetraloop-receptor RNA construct at ~100 pM dilution. Just before imaging, the channel is flushed with an imaging solution that contains the following: (i) 50 mM Hepes buffer (pH 7.5), (ii) 0-100 mM NaCl and 0.1 mM EDTH, (iii) 2 mM Trolox (9hydroxy-2,5,7,8-tetramethylchroman-2-carboxylic acid), (iv) ~0.1 mg/mL protocatechuate 3,4-dioxygenase (PCD), (v) ~10 mM 3,4-dihydroxybenzoic acid (PCA) to catalytically remove oxygen, and (vi) PEG 8000 to achieve the desired solution conditions.

## Confocal Imaging

Surface immobilized molecules are imaged with a home-built confocal microscope that has been described in detail elsewhere (1, 2). Briefly, the samples are illuminated with 532-nm light from a pulsed ND:YAG laser (10 ps at 20 MHz) that is focused to a diffraction limited spot by a 1.2 N.A. water immersion objective. The total fluorescence from single dye-labeled RNA molecules is collected through the same objective, focused through a 50-µm pinhole, and split into Cy3 and Cy5 channels with dichroics (645 nm) just before collection by avalanche

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photodiodes. Background correction,  $E_{\text{FRET}}$  calculation, and dwell time analysis are performed with in-house written software.

As in previous studies, the kinetics of GAAA tetraloopreceptor folding is measured from single-molecule fluorescence intensity trajectories as follows (2-8). Folding and unfolding events are observed as anticorrelated intensity fluctuations between the Cy3 and Cy5 channels. Fluorescence intensities are then used to calculate the fluorescence resonance energy transfer efficiency values ( $E_{\text{FRET}}$ ), generating  $E_{\text{FRET}}$  trajectories, that reveal binary switching between high ( $E_{\rm FRET} \sim 0.7$ ) and low  $(E_{\text{FRET}} \sim 0.3)$  energy transfer efficiency (7). Crossings between high and low  $E_{\text{FRET}}$  define the time the construct spends in each state, or dwell times ( $\tau_{fold}$  and  $\tau_{unfold}$ ), which are then integrated with cumulative distribution functions (CDFs) to describe the decay out of the folded and unfolded states (9). Under typical illumination intensities (1-2 µW), each individual molecule survives long enough to record between 5 and 50 switching events before irreversibly photobleaching. To build up the 200-500 total switching events to accurately describe the single-exponential decay out of each state, ~10-50 individual molecules are monitored under one set of solution conditions. The aggregate switching events are then randomly sorted into three subdatasets and leastsquares fit to single-exponential functions yielding unimolecular rate constants for folding  $(k_{fold})$  and unfolding  $(k_{unfold})$  (10, 11). The average and SD of the three independent fits are reported on the plots. Typically, the fractional error on individual data points is  $\sim 10\%$  of the absolute value.

To control the sample temperature for thermodynamic measurements, the sample is heated in two ways: (*i*) simultaneous proportional-integral-derivative controlled Peltier heating of both stage and objective and (*ii*) an IR laser tuned to the first overtone of the OH stretch of water (1,455 nm) and focused to a ~17-µm spot coaxial with the visible laser excitation beam (7, 12). Both methods have been shown to provide temperature control to a precision and accuracy of  $\pm 0.1$  °C (6–8, 12).

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**Fig. S1.** Sample data analysis is shown. (*A*) Cy3 and Cy5 fluorescence time trajectories illustrate the anticorrelated intensity fluctuations due to folding and unfolding transitions. (*B*) The  $E_{FRET}$  time trajectory is calculated from fluorescence intensities revealing the binary two-state switching behavior. (*C*) The dwell times in the folded and unfolded states are used to construct the CDFs for folding and unfolding, which are then fit with single-exponential functions to extract  $k_{fold}$  and  $k_{unfold}$ .



**Fig. S2.** Temperature dependences of the measured rate constants are shown in aqueous (red) and 8% PEG 8000 (blue) solutions. In both the folding ( $k_{on}$ ) and unfolding ( $k_{off}$ ) processes, the lines are parallel, indicating no significant change in the transition state enthalpies for either process. The largest effect is the upward offset of  $\ln(k_{on})$  under crowded conditions. Similar to our analysis of the equilibrium results, this uniform upward shift in  $k_{on}$  corresponds to crowding-induced changes in the entropy rather than enthalpy. However, with both forward and backward single-molecule rate constants explicitly measured as a function of temperature, the data further reveal that this shift arises predominantly from a reduction in the entropic cost of forming the transition state from the unfolded state.