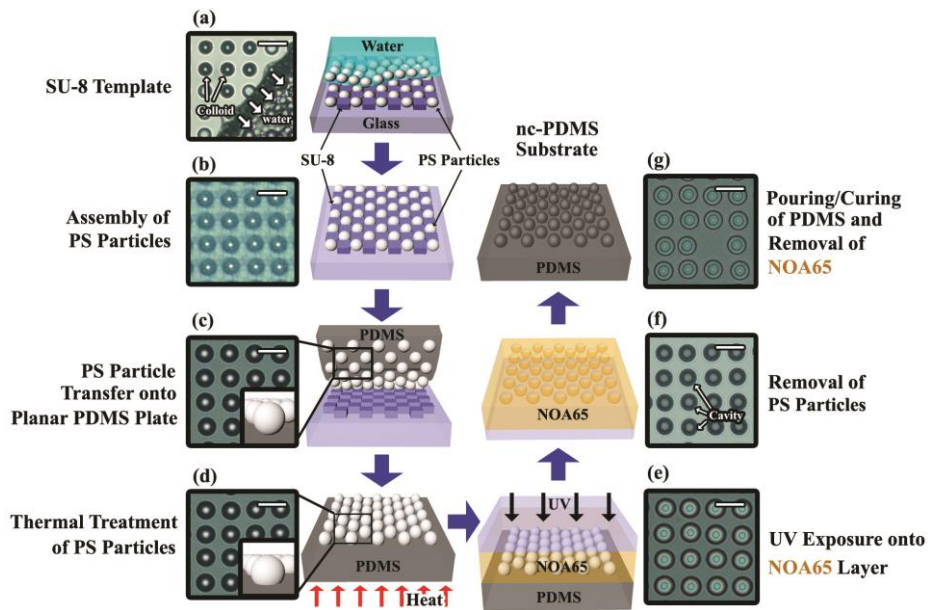
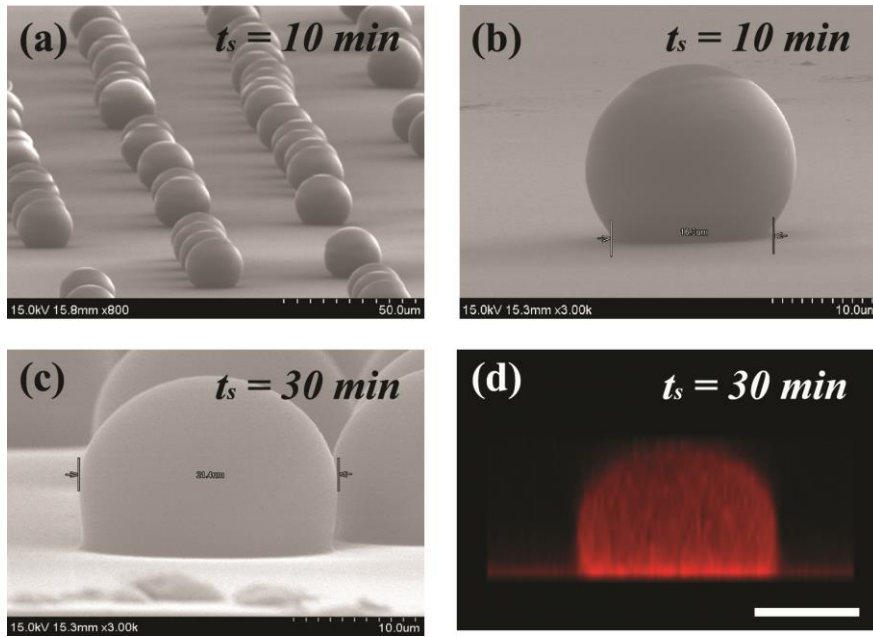


## Supplementary Information



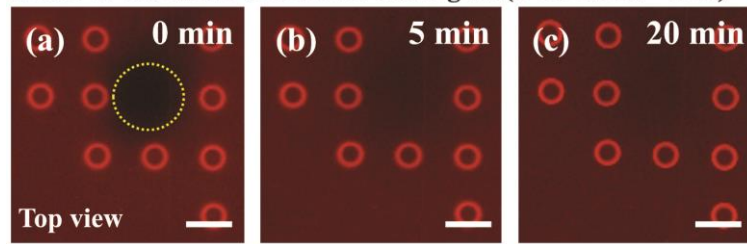
**Supplementary Figure 1. Fabrication processes of bud-mimicking PDMS patterns based on a combination of colloidal, soft, and photo-lithography a-d,** The polystyrene (PS) particles in colloid were assembled on a periodically patterned SU-8 template through a convective assembly process<sup>1, 2</sup>, transferred onto a PDMS plate, and treated thermally for softening and fixation. **e,** A photo-curable polymer (NOA65) layer was coated on the PDMS plate, a glass cover was placed on it, and the whole plate was then exposed to ultraviolet light for 2 min. **f,** After peeling off the PDMS plate, the PS particles were removed from the NOA65 layer in toluene for 1 min. **g,** The PDMS was poured over the NOA65 layer with the PS vacancies and then thermally cured for 3 h. The NOA65 layer was finally removed in a mixed solution for 1 day to leave out only the bud-mimicking patterns of the PDMS (the nc-PDMS substrate).

Note that after the plasma treatment, the nc-PDMS substrate exhibited significant resistance to hydrophobic recovery and stabilized the lipid membranes supported on the PDMS substrate for extended periods of time (up to several days). Scale bars: 50  $\mu\text{m}$

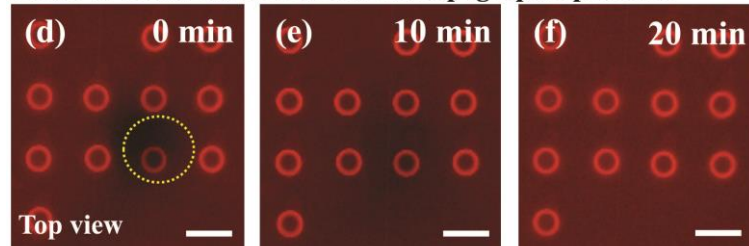


**Supplementary Figure 2. The SEM images of the bud-mimicking PDMS patterns and the confocal image (side-view) of the SLB covered on the PDMS pattern a-c, The SEM images of an array (a) and one unit (b) of the PMDS bud-pits produced at the softening time  $t_s = 10$  min (the bud-neck diameter  $\approx 16$   $\mu\text{m}$ ). The bud-neck of about 18  $\mu\text{m}$  in diameter was produced at  $t_s = 30$  min (c). d, The confocal image (side-view) of the SLB covering the PDMS pattern in (c). Scale bar: 10  $\mu\text{m}$**

**FRAP of TR on the SLB in the flat region (DOPC:TR = 99:1)**

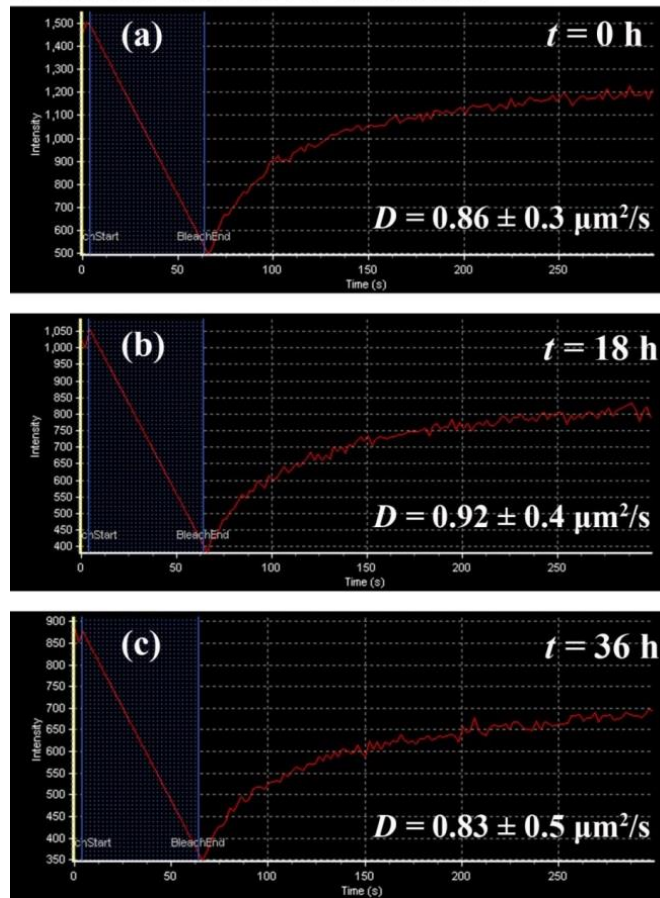


**FRAP of TR on the SLB across the topographic patterns**

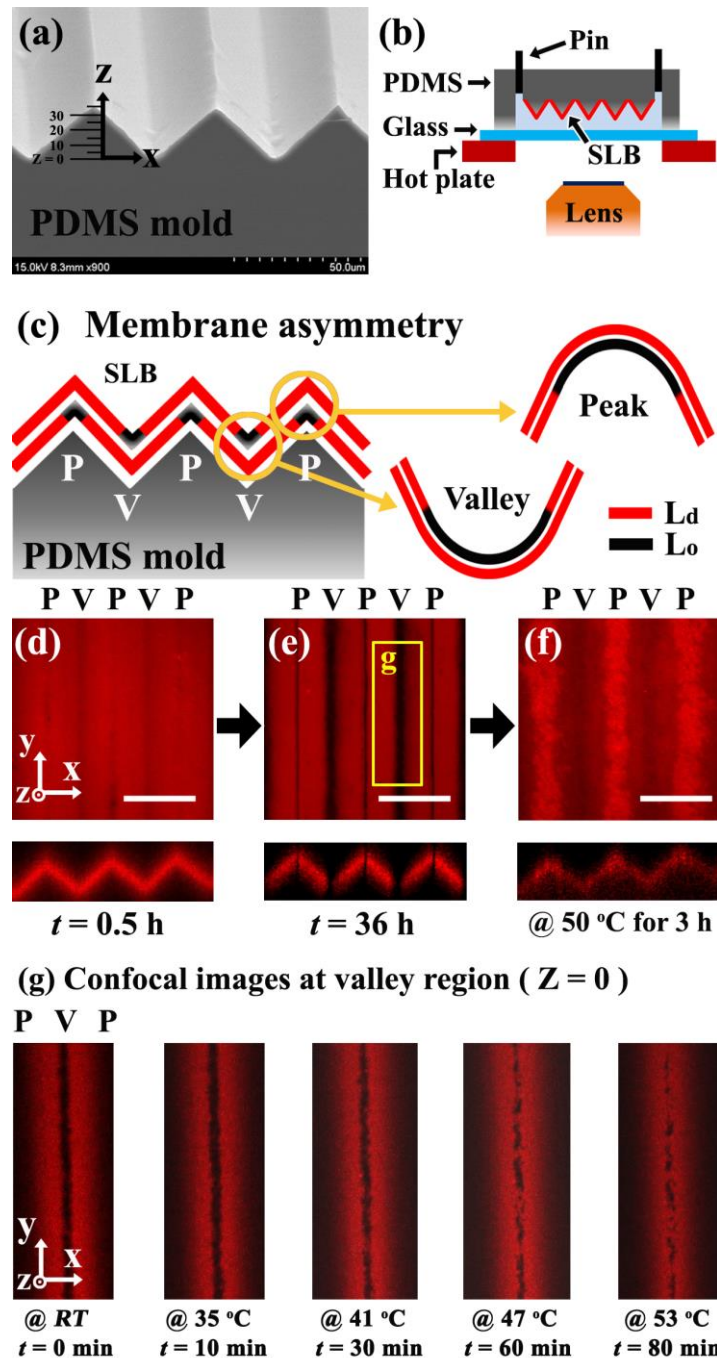


**Supplementary Figure 3. Lateral fluidity for the supported B1 membrane (DOPC: TR-DHPE=99:1) in the flat- and the bud-pit regions** The fluorescence recovery after photo-bleaching (FRAP)<sup>3</sup> studies were performed at the time of 3 days after the preparation of the SLB. **a-c**, The FRAP results in the flat region at the recovery time of 0, 5, and 20 min. **d-f**, The FRAP results across the topographic patterns with  $r \approx 100$  nm at the recovery time of 0, 10, and 20 min. The yellow dotted circle represents the photo-bleached area. Scale bars: 50  $\mu\text{m}$

### FRAP test with the B7 membrane

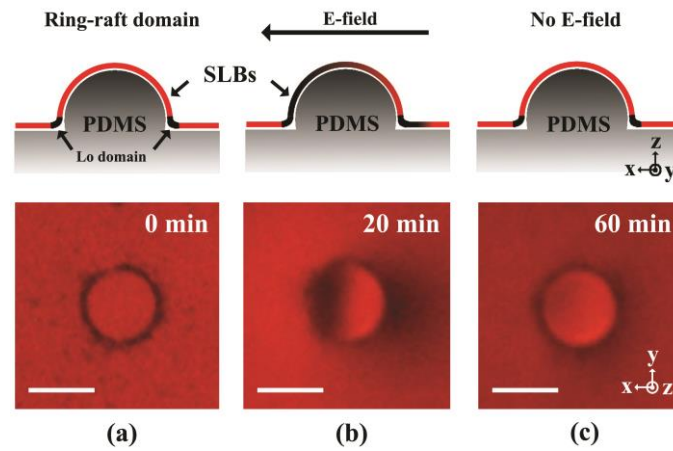


**Supplementary Figure 4. Diffusion coefficients in the ring-raft around the collar band at three elapsed times for the B7 membrane (DOPC:SPM:CHOL:TR=33:33:33:1) a-c,** The profiles of the fluorescence recovery after photo-bleaching (FRAP) for the B7 membrane supported on the oxidized nc-PDMS substrate at three elapsed times: (a) immediately after vesicle fusion ( $t = 0$  h); (b) 18 hours after vesicle fusion ( $t = 18$  h); and (c) 36 hours after vesicle fusion ( $t = 36$  h). Similar diffusion coefficients over the course of the elapsed time for 36 h ensure the long-term lateral fluidity and the structural contiguity of the SLB.

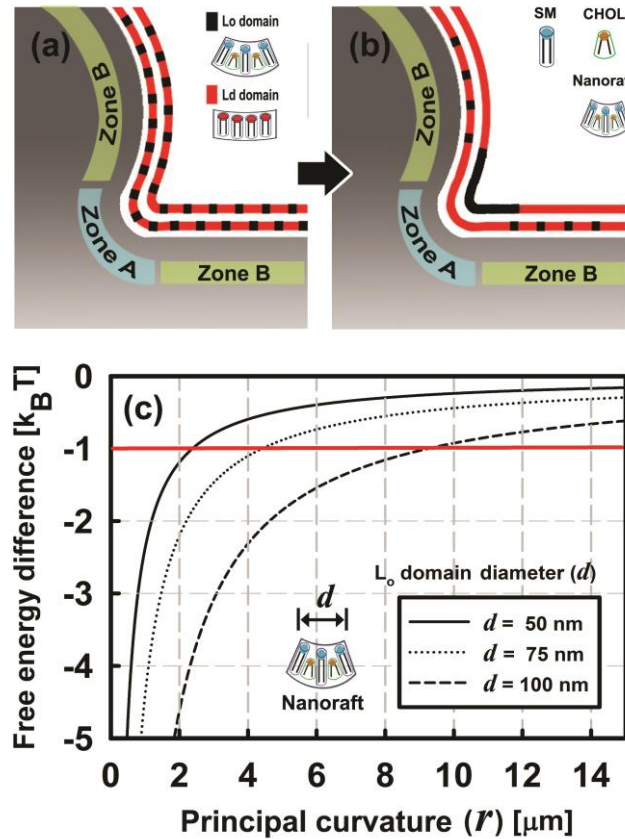


**Supplementary Figure 5. Reconstitution of the  $L_o$  domains in the prescribed regions with local curvatures and the thermal behavior of topographically structured SLB** In determining the thermal behavior of the curvature-driven  $L_o$  domain, the PDMS substrate with periodic wedges of the peaks (denoted by P's) and the valleys (denoted by V's) was used as a support for the SLB. The peak-to-peak distance was 50  $\mu\text{m}$ . Note that the wedge substrate represents the essential topographic features of the bud-neck regions and facilitates the observation of the  $L_o$  domains without being obscured by the bud-pit. The radius of curvature ( $r$ ) in the valley or the peak was comparable to that in the bud-neck ( $r = 100 - 300$  nm). The simplicity of the wedge substrate enables to directly characterize the thermal behavior of the  $L_o$  domains. **a**, The SEM image of the PDMS substrate with periodic wedges .

**b,c**, Schematic illustration of a custom-designed thermal stage mounted on a confocal fluorescence microscope (b) and reconstitution of the  $L_o$  domains in the prescribed regions of local curvatures together with the arrangement of the inner and outer leaflets in the valley and peak regions (c). The inter-leaflet compositional asymmetry may be considered for understanding more delicate interactions but it is beyond the scope of this work. **d,e**, The confocal images projected along the  $z$ -direction (top) and the cross-sectional images (bottom) obtained for the B7 membrane (DOPC:SPM:CHOL:TR=33:33:33:1) at the duration time of  $t = 0.5$  h (d) and 36 h (e) after vesicle fusion. Clearly, the TR-depleted  $L_o$  phase gradually decorated the edges and the raft domain was developed in the valleys and peaks as expected<sup>4</sup>. **f**, Upon heating the SLB sample at 50°C for 3 h, the fluorescence pattern became partially homogenized, indicating that the  $L_o$  domains, devoid of fluorescence, were mixed with the surrounding background and lost the preferential localization in the edges. **g**, The thermal and temporal behavior of the raft domain in the valley in the course of heating for 80 min (in a step of 10 min) from room temperature to 53°C. The thermal process of the pattern homogenization and domain dissolution was clearly observed. The images were acquired in a semi-confocal mode where a pinhole was open at the maximum of 800  $\mu\text{m}$  in diameter to obtain a large depth of field per image, allowing to keep the valleys in focus ( $Z = 0$   $\mu\text{m}$ , a yellow rectangular inset depicted in e). The observed temperature for the domain dissolution (about 41°C) was found to be somewhat higher than the bulk miscibility transition temperature (38°C). This may be attributed to the error in temperature control in our experimental setup and the effect of the substrate interface<sup>5</sup>. Scale bars: 50  $\mu\text{m}$



**Supplementary Figure 6. The closed ring-raft domain as a barrier to the lipid change between the bud membrane and the donor membrane a-c,** The curvature-driven localization of the  $L_o$  domain in the SLB was produced before the application of an electric field. The electric field  $E$  was applied along the  $x$ -axis from two Pt wires placed 1.5 cm apart. In the bottom, the epifluorescence micrographs observed at  $t = 0$  min (at the time of the application of  $E$ ; **a**),  $t = 20$  min (20 min under  $E = 45$  V/cm; **b**), and  $t = 60$  min (40 min after the removal of  $E$  at  $t = 20$  min; **c**) were shown. The closed topology of the ring-raft domain indeed serves as a barrier to the lipid exchange across the bud-neck between the bud membrane and the donor membrane. Scale bars: 20  $\mu$ m



**Supplementary Figure 7. A simple model for the preferential coarsening of raft domains at the bud-neck interface** The spontaneous curvature of the L<sub>o</sub> phase ( $c_{L_o} = -1/68 \text{ \AA}^{-1}$ )<sup>6</sup>, originated from the asymmetry in the lipid molecule<sup>7</sup>, has a more negative value than that of the L<sub>d</sub> phase ( $c_{L_d} = -1/160 \text{ \AA}^{-1}$ ) composed of the DOPC molecules<sup>8</sup> (shown in **a**). When the lipid membrane in the bud-neck region (Zone A) is under the structural deformation by the negative curvature of the underlying PDMS substrate, the L<sub>o</sub> domain experiences less severe elastic distortions than the L<sub>d</sub> domain. In other words, the curvature gradient of the negatively curved membrane provides one of the contributions to the driving force for the accumulation and coarsening of ‘nanorafts’<sup>9</sup> at the bud-neck interface. Together with the asymmetric distribution across two monolayer leaflets<sup>10</sup>, the curvature gradient leads to the formation of a ring-raft in the outer leaflet around the bud (shown in **b**). More quantitatively<sup>6</sup>, two limiting cases for the localization of the L<sub>o</sub> domain can be considered within the formalism of the Helfrich-type free energy per unit area<sup>11</sup>; (1) the preferential localization of the L<sub>o</sub> phase in Zone A (case I) and (2) the exclusion of the L<sub>o</sub> phase from Zone A (case II). In a simple model, the difference in the elastic energy per unit area ( $\Delta f$ ) between the two cases as a function of the principle curvature radius ( $r$ ) can be written as

$$\Delta f = f(\text{case I}) - f(\text{case II})$$



$$\begin{aligned}
&= \frac{1}{2} \left( \frac{1}{2} k_{Lo} (2c_{zone A} - 2c_{Lo})^2 + \frac{1}{2} k_{Ld} (2c_{zone B} - 2c_{Ld})^2 \right) \\
&\quad - \frac{1}{2} \left( \frac{1}{2} k_{Lo} (2c_{zone B} - 2c_{Lo})^2 + \frac{1}{2} k_{Ld} (2c_{zone A} - 2c_{Ld})^2 \right) \\
&= -\frac{2}{r_{zone A}} (k_{Lo} c_{Lo} - k_{Ld} c_{Ld}) \\
&\quad \text{( if } c_{zone A} \ll c_{Lo}, c_{zone A} \ll c_{Ld}, c_{zone B} = 0 \text{ )}
\end{aligned}$$

where  $c_{zone B}$  and  $c_{zone A}$  are the principle curvature in zone B and that in zone A, , respectively,  $k_{Ld}$  and  $k_{Lo}$  are the bending moduli, and  $c_{Ld}$  and  $c_{Lo}$  are the spontaneous curvatures. The free energy difference between the two cases is then given by  $\Delta f \times \pi(d/2)^2$  where  $d$  is the size of the unit raft (or nanoraft).

The free energy differences for three different sizes of the Lo domain were calculated as a function of the principal curvature radius ( $r$ ) (shown in **c**). The parameters<sup>12-14</sup> used for the calculations were  $c_{Ld} = -1/160 \text{ \AA}^{-1} = -62.5 \text{ \mu m}^{-1}$ ,  $c_{Lo} = -1/68 \text{ \AA}^{-1} = -147.1 \text{ \mu m}^{-1}$ ,  $k_{Ld} = 5 k_B T$ , and  $k_{Lo} = 6.25 k_B T$ . It was found that the coarsening of nanorafts is energetically favored at the bud-neck interface (zone A) in relative to zone B (planar surface) when the free energy difference exceeds thermal energy of  $1 k_B T$  (represented by the red line). For  $d = 50 \text{ nm}$ , the local radius of curvature below  $r \leq 2.38 \text{ \mu m}$  is small enough to accumulate the nanorafts in the bud-neck region. The critical value  $r$  was found to increase with increasing  $d$ . This might explain the coarsening process of nanorafts to produce a rather wide rim beyond the ring structure. In our case, the curvature radius is in the range of  $r \approx 100 - 900 \text{ nm}$  at the bud-neck interface from which the theoretical criterion for the spatial localization of raft units onto a negatively curved surface of the outer leaflet can be made. However, a complete picture of the curvature-based formalism where the interleaflet asymmetry is explicitly taken into account remains to be explored.

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