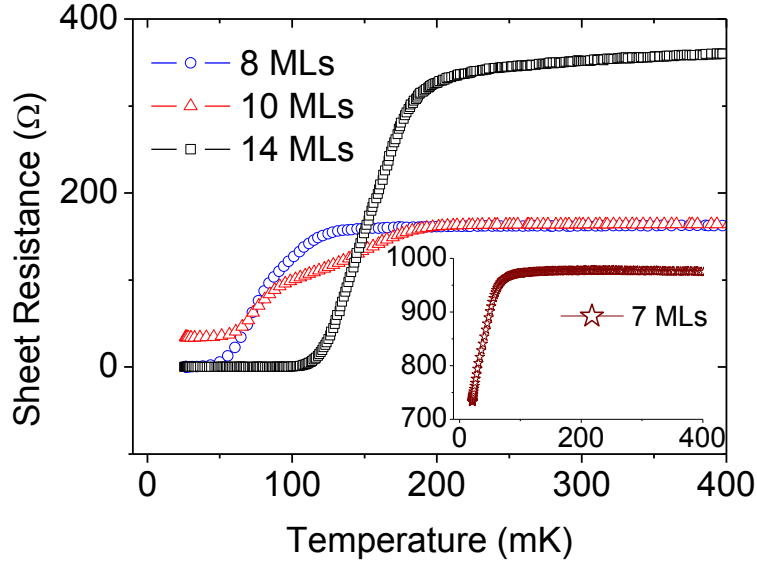


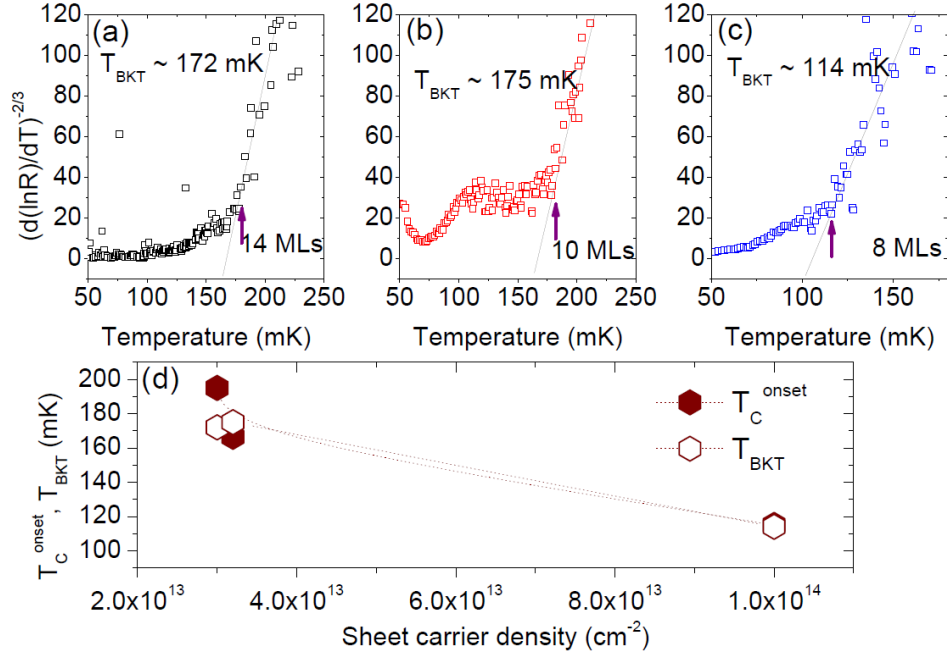
**Supplementary Figure 1| Out-of-plane upper critical fields at (001) and (110) LaAlO<sub>3</sub>/SrTiO<sub>3</sub> interfaces.**

Sheet resistance of (a) the (110)-interface with  $t = 14$  MLs and (b) the (001)-interface with  $t = 10$  MLs, under magnetic fields applied normal to the interface. The field values are indicated in the panels of Supplementary Figure 1.



**Supplementary Figure 2 | Superconductive transitions at the (110) LaAlO<sub>3</sub>/SrTiO<sub>3</sub> interface.**

We defined the onset of the transition as the temperature where the resistance falls to 90% of its normal state value (measured at  $T = 400$  mK). We observed that the superconductive transition temperature  $T_C$  of (110)-oriented interfaces was strongly dependent on the sample, being the highest for  $t = 14$  MLs ( $T_C \approx 195$  mK), whereas lower values were observed for  $t = 10$  MLs ( $T_C \approx 166.5$  mK),  $t = 8$  MLs ( $T_C \approx 115$  mK) and  $t = 7$  MLs ( $T_C \approx 45$  mK) (Supplementary Figure 1). Samples with  $t = 14$  MLs and  $t = 8$  MLs exhibited the narrowest transition widths  $\Delta T_C \approx 40$  mK (defined between 20% and 80% of the normal state resistance), comparable to the  $\Delta T_C$  observed in superconductive (001)-LaAlO<sub>3</sub>/SrTiO<sub>3</sub> as well as (001)-LaTiO<sub>3</sub>/SrTiO<sub>3</sub><sup>1-3</sup> interfaces. The other samples exhibited significantly broader transitions, indicating more inhomogeneous electronic states. This observation is more perceptible in sample  $t = 10$  MLs that exhibits a complex transition to superconductivity with two kinks.



**Supplementary Figure 3| Analysis of the 2D-superconductivity transition: the Berezinskii-Kosterlitz-Thouless (BKT) model.**

In the 2D limit, it is well established that superconductivity is suppressed by vortex-antivortex pair unbinding by thermal fluctuations, so that the transition to the 2D superconductive state belongs to the Berezinskii-Kosterlitz-Thouless (BKT) universality class. In this scenario, above the transition temperature  $T_{BKT}$  the free motion of vortices triggers the transition from the 2D superconductive towards the normal state. Within a narrow region around  $T_{BKT}$  the BKT theory predicts specific laws for the temperature dependence of the resistance  $R$ , giving a quantifiable criterion to assess the dimensional character of the superconductive state. In particular, it is expected that  $[\partial \ln R / \partial T]^{-2/3}$  should scale linearly with  $(T - T_{BKT})$  for a certain region around  $T_{BKT}$ .<sup>4</sup> [As seen in panels a-c, similarly to (001)- LAO/STO interfaces, the  $[\partial \ln R / \partial T]^{-2/3}$  data corresponding to (110) interfaces fall approximately on a straight line above a certain temperature that –following the same protocol as described in the analogous analysis of (001)-interfaces<sup>1</sup> – we identify as the  $T_{BKT}$  transition temperature. Therefore, the temperature  $T_{BKT}$  is at the boundary between the two regimes (superconductive/normal) giving rise to a prominent kink in the  $[\partial \ln R / \partial T]^{-2/3}$  plot. We note that the values of  $T_{BKT}$  thus obtained are in the same range as those found in (001)-LaAlO<sub>3</sub>/SrTiO<sub>3</sub>. We find that the  $T_{BKT}$  values correlate with  $T_C$  (panel d), both being strongly dependent on the sheet carrier density  $n_{\text{sheet}}$  in the normal state, a feature that is reminiscent of the strong dependence of the bulk superconductive transition on carrier density.

## Supplementary References

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