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Effective Stress Data. We have compiled effective stress data of both sedimentary and crystalline rocks (Fig. 1). Unlike hard crystalline rocks, the softer sedimentary rocks are usually noncritically stressed. The data displayed in Fig. 1 is detailed in Table S1.

Model of the Upper Crust. To test whether sedimentary formations are likely to become critically stressed or not, we built a model of the upper crust to simulate the stress buildup induced by the strain of plate tectonics (Fig. S1). We impose a strain rate typical of plate tectonics, i.e., 10^{-17} s⁻¹ (1) at the left boundary of the model. As a result of this strain, the upper crust, which we consider that its initial stress state is isotropic, will undergo an increase of the horizontal stresses. Apart from the isotropic stress state, the initial conditions are hydrostatic pressure and a geothermal gradient of 30 °C/km, with a surface temperature of 10 °C. Both the pressure and temperature are maintained constant throughout the simulation. In this model, we consider that the shallower 2.5 km are sedimentary rocks and below them, there is the crystalline basement.

We assume that our model of the upper crust behaves elastically. We deem that this rheology is appropriate because the elastic plate thickness is usually thicker than the 16 km considered in our model (1–3). The stiffness of the rocks is a function of depth and of the rock type. We chose the stiffness of the sedimentary rocks to be in agreement with the compressibility values given by Neuzil (4) for sedimentary rocks (Fig. S2). The stiffness of the crystalline basement is based on values given by Brudy et al. (5) for the gneiss at the deep well KTB, Germany. We assume that the stiffness increases with depth and that the upper part of the crystalline basement is softer than the deeper part because it could have been exposed to alteration (6). The resulting vertical profile of the Young's modulus as a function of depth is shown in Fig. S3.

Despite the simplicity of the model, it illustrates the concept that sedimentary rocks, which are softer than the crystalline basement, become less stressed than the crystalline basement. We solve this thermo-hydro-mechanical problem using the finite element code CODE_BRIGHT (7, 8). The model is discretized with structured quadrilateral elements. The mesh is regular, with an element size of 2 km in the horizontal direction and 100 m and 500 m in the vertical direction for the sedimentary rocks and the crystalline basement, respectively. We have ensured that further refinement of the mesh does not affect the results.

Model of CO₂ Injection. For studying the overpressure and caprock stability evolution as a result of $CO₂$ injection, we consider an idealized horizontal reservoir overlaid and underlain by a lowpermeability and high entry pressure formation (Fig. S4). The reservoir thickness is of 50 m and the thickness of the seals on top and below of the reservoir is 100 m. The top of the reservoir is placed at a depth of 1,500 m. The model is completed by an upper aquifer that extends up to 500-m depth and by a basal aquifer that reaches a depth of 2,500 m. We assume that the upper 500-m-thick overburden has such a low shear stiffness that it does not need to be included in the model. To avoid boundary effects on the hydromechanical behavior of the model, the model extends laterally 100 km. The model is axisymmetric to represent a vertical well, which has a radius of 0.15 m.

The hydromechanical properties of the rocks that form the model are listed in Tables S2 and S3. Table S2 displays the values of the base case, and Table S3 includes reasonable ranges of values of the stiffness and the porosity of the rocks. The reservoir corresponds to a permeable limestone and the seals to a shale (9). The upper and basal aquifers have the same hydraulic properties, but the basal aquifer is stiffer because of its higher confining pressure.

The initial conditions are hydrostatic pressure and a constant temperature of 60 °C, which corresponds to the mean temperature of the reservoir. We assume isothermal conditions and that both horizontal stresses are equal. The initial stress field displays a relationship between horizontal and vertical effective stresses of $\sigma'_{h_0} = 0.5\sigma'_{v_0}$, where σ'_{h_0} is the initial horizontal effective stress, $\sigma'_{v_0} = 0.013z$ MPa is the initial vertical effective stress and z is depth. This relationship between the vertical and the horizontal effective stresses is equivalent to a mobilized friction coefficient of 0.35, which is in the lower limit of the friction coefficients of sedimentary rocks of $CO₂$ storage sites (Table S1). As a first step, a steady-state calculation is carried out to ensure consistent initial conditions in equilibrium for the pressure and stress fields.

The hydraulic boundary conditions are constant hydrostatic pressure at the outer boundary and no flow at the other boundaries. At the injection well, a mass flow rate of 2 Mt/y of $CO₂$ is uniformly injected along the whole thickness of the reservoir for 2,000 d. As for the mechanical boundary conditions, we apply a constant lithostatic stress on the upper boundary and no displacement perpendicular to the other boundaries.

The hydromechanical response of the rocks to $CO₂$ injection in a deep saline reservoir is modeled using the fully coupled finite element numerical code CODE_BRIGHT (7, 8), which was extended to simulate $CO₂$ injection by Vilarrasa et al. (10). The model is discretized with structured quadrilateral elements. The mesh is finer close to the injection well in the reservoir and seals and becomes coarser further away. We have ensured that further refinement of the mesh does not change the results.

We discuss here the sensitivity of overpressure evolution to variations of Young's modulus and porosity within the ranges of Table S3 (Fig. S5). Overpressure is virtually insensitive to changes in the stiffness and porosity of the upper and basal aquifers and to the porosity of the caprock. However, caprock stiffness does affect overpressure, because it affects the storativity of the reservoir (11). Stiff caprocks reduce the effective storage coefficient of the reservoir and therefore, pressure buildup is higher than for soft caprocks. Similarly, stiff reservoirs yield a higher overpressure than soft reservoirs. Finally, high porosity reservoirs induce a higher overpressure than low porosity reservoirs. This may seem paradoxical, but reflects that permeability has not been changed. A small porosity implies that the $CO₂$ plume is large, which will cause a low overpressure because $CO₂$ viscosity is small, which reduces the overall resistance to flow.

CO₂ Dissolution Calculations. Dissolution is sensitive to two factors: resident volume of water (porosity, ϕ , times thickness, b) and vertical permeability, k. Dissolution may also be affected by aquifer bottom slope, as the dense $CO₂$ rich brine will tend to sink downslope, but this effect is neglected here. Permeability controls the dissolution rate, m, whereas the resident volume controls the maximum dissolution capacity, M . Both M and m are proportional to CO_2 solubility, χ . The maximum dissolution capacity is given by

$$
M = \phi b \rho \chi \tag{S1}
$$

where ρ is water density. The solubility is highly sensitive to pressure, temperature and brine salinity. Pool et al. (12) provide values of solubility (mass fraction) for 0.5 M water (salinity comparable to that of seawater) that range from 4.1% (100 bars, 55 °C) to 5.5% (200 bars, 30 °C, corresponding to cold $CO₂$ injection or storage below the ocean). Adopting a storage formation with $b =$ 100 m, $\phi = 0.1$, $\rho = 1,000$ kg/m³, and $\chi = 0.049$, the maximum dissolution capacity results in $M = 0.49$ Mt/km², or a capacity of 19.6 Mt for a 40 km² CO₂ plume, which suggests that storage 19.6 Mt for a 40 km² $CO₂$ plume, which suggests that storage capacity for dissolved $CO₂$ is significant in deep saline formations. Different expressions have been derived for the dissolution rate

(13, 14). Expressed in the same form, they can be written as

$$
m = \alpha \chi \rho u_b, \nu_b = k \epsilon \rho g / \mu,
$$
 [S2]

where u_b is the flux of a dense fluid sinking in a light fluid when the density difference is $\Delta \rho$, g is gravity, μ is water viscosity and

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 $\varepsilon(\Delta\rho/\rho)$ is the buoyancy factor. Pool et al. (12) also provide values for this factor, which ranges around 0.8–1.6%, for the conditions adopted for solubility above. The factor α was numerically evaluated as 0.0175 by Pau et al. (14) and as 0.0142 by Hidalgo and Carrera (13).

It is clear that the critical factor in the evaluation of m is the vertical permeability. If $\chi = 0.49$, $\varepsilon = 0.012$, $\mu = 5.9 \cdot 10^{-4}$ Pa·s, and $k = 2 \cdot 10^{-12}$ m², then $m = 0.01$ Mt/y/km², or 0.4 Mt/y for a 40-km²
plume which is significant. For example, if we analyze a storage plume, which is significant. For example, if we analyze a storage site with a high permeability as Sleipner, where the $CO₂$ plume extends 12 km² after having injected 12 Mt of $CO₂$ in 12 y (15), the dissolution rate would be of 0.12 Mt/y, which is a 12% of the injection rate. Obviously, the dissolution rate would drop to 5.10^{-4} Mt/y/km² if $k = 10^{-13}$ m², which is negligible and supports the generalized percention that dissolution is only relevant over the generalized perception that dissolution is only relevant over long periods of time.

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Fig. S1. Schematic representation of the upper crust model (not to scale). The upper 2.5 km represent sedimentary rocks, and the lower 13.5 km characterize the crystalline basement. The initial stress state is isotropic, but the horizontal stresses increase with time as a result of the strain rate typical of plate tectonics.

Fig. S2. Rock compressibility of the sedimentary rocks considered in our study. The shaded region indicates the range of compressibility of sedimentary rocks given by Neuzil (1).

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Fig. S3. Young's modulus as a function of depth of the sedimentary and crystalline rocks considered in the upper crust model. The sedimentary rocks are softer that the crystalline basement.

Fig. S4. Schematic representation of the CO₂ injection model (not to scale). 2 Mt/y of CO₂ are injected through a vertical well in a 50-m-thick reservoir overlaid and underlain by a seal.

Fig. S5. Effect of subsurface heterogeneity on overpressure evolution.

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Table S1. Effective stress data of sedimentary and crystalline rocks shown in Fig. 1

 $σ₁$ is the maximum effective stress, $σ₃$ is the minimum effective stress and $μ$ is the mobilized friction coefficient. -, no units.

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-, no dimensions.

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Table S3. Reasonable ranges of stiffness and porosity of the deep sedimentary rocks considered in the model

-, no dimensions.