# **Simulation of avascular tumor growth by agent-based game model involving phenotype-phenotype interactions: Supplementary**

Yong Chen<sup>1,2</sup>, Hengtong Wang<sup>3</sup>, Jiangang Zhang<sup>4</sup>, Ke Chen<sup>5</sup>, and Yuming Li<sup>1,+</sup>

<sup>1</sup>Key Laboratory of Digestive System Tumors of Gansu Province, Lanzhou University, Lanzhou 730000, China <sup>2</sup>Center of Soft Matter Physics and its Applications, Beihang University, Beijing 100191, China

<sup>3</sup> School of Physics and Information Technology, Shaanxi Normal University, Xi'an 710119, China

4 Institute of Pathology, Lanzhou University, Lanzhou 730000, China

<sup>5</sup>Key Laboratory of Genome Sciences and Information, Beijing Institute of Genomics, Chinese Academy of Sciences, Beijing 100101, China

## **ABSTRACT**





#### **Supplementary Figures**



Figure S1. An example of renormalized proliferative and invasive probabilities for a cell in the spatial grid following by Eqs. (3-4). Here, the living cell count  $N_T = 3$ , the shape parameters  $\theta_p = \theta_i = 0.2$ . The phenotypic parameters in Table 1 are  $\alpha_{pp} = -0.1$ ,  $\beta_{ii} = 0.1$ ,  $\alpha_{pi} = -0.02$ , and  $\beta_{ip} = 0.02$ . The blue circled line is the proliferative probability and the green one is of invasive cell. The red dashed line denotes the nutrient criterion of necrotic phenotype  $\phi_c = 0.02$ .



Figure S2. Snapshots of necrotic cells in growing tumor based on our simulation at different times, (a)  $t = 50 \tau$ , (b)  $t = 100 \tau$ , (c)  $t = 150 \tau$ , (d)  $t = 200 \tau$ . Here, the living cells and ECM have been removed. Parameters are the same as those in Figure 3.



Figure S3. The average radial distance  $\langle r \rangle$  again time *t*.  $\langle r \rangle$  increases linearly with *t*. It means the tumor grows with constant radial velocity *v*. Note that *v* decreases with increasing the inhibition degree  $\alpha_{pp}$  except for  $\alpha_{pp} = -0.9$ . Here, the parameter  $\beta_{ii} = 0.1$  and the others are in Table S1.



Figure S3. The average roughness  $\langle R \rangle$  again time *t*.  $\langle R \rangle$  increases linearly with *t*. It means the surface roughness grows with constant velocity which increases with the inhibition degree  $\alpha_{pp}$ . Here, the parameter  $\beta_{ii} = 0.1$  and the others are in Table S1.

#### Text S1. The numerical scheme for Eqs. (1-2).

In this study, the nutrient diffusion equations Eq. (1) is a nonlinear parabolic partial differential equation in two-dimensional space under the conditions Eqs. (2). In the case of one-dimensional space, the similar general parabolic equation is<sup>[1](#page-4-1)</sup>

<span id="page-3-0"></span>
$$
\partial_t u = a \partial_{xx} U + b u, \qquad 0 < x < 1, \quad t > 0. \tag{S.1}
$$

The initial and boundary conditions are

$$
u(x,0) = g(x), \t 0 \le x \le 1 u(0,t) = \psi(t), \t t \ge 0 u(1,t) = \phi(t), \t t \ge 0.
$$

For this class of numerical problem, the differences of Eq. [\(S.1\)](#page-3-0) are

$$
\partial_t u\Big|_{j}^{n} = \frac{u_j^{n+1} - u_j^{n}}{\tau} + O(\tau)
$$
  

$$
\partial_{xx} u\Big|_{j}^{n} = \frac{1}{h} \Big( \partial_x u\Big|_{j+1/2}^{n} - \partial_x u\Big|_{j-1/2}^{n} \Big) + O(h^2)
$$
  

$$
\partial_x u\Big|_{j-1/2}^{n} = \partial_x u\Big|_{j-1/2}^{n+1} - \tau \partial_x \partial_t u\Big|_{j-1/2}^{t_n + \theta \tau}, \qquad 0 \le \theta \le 1.
$$

Here,  $u_j^n$  denotes the value of function *u* at  $(x_j, t_n)$ .  $\tau$  and *h* represent the time and spatial step size, respectively. Substituting the above equations to Eq.  $(S.1)$ , we have

$$
\frac{u_j^{n+1} - u_j^n}{\tau} = \frac{a}{h} \left( \partial_x u \Big|_{j+1/2}^n - \partial_x u \Big|_{j-1/2}^n \right) + b u_j^n + O(\tau + h^2).
$$
 (S.2)

Considering the central differences as below

$$
\partial_x u \Big|_{j+1/2}^n = \frac{1}{h} (u_{j+1}^n - u_j^n) + O(h^2),
$$
  

$$
\partial_x u \Big|_{j-1/2}^n = \frac{1}{h} (u_j^{n+1} - u_{j-1}^{n+1}) + O(h^2),
$$

a highly efficient quasi-Saul'ev difference scheme for one-dimensional Eq.  $(S.1)$  is deduced,  $1-5$  $1-5$ 

<span id="page-3-1"></span>
$$
\frac{u_j^{n+1} - u_j^n}{\tau} = \frac{a}{h^2} \left( u_{j+1}^n - u_j^n - u_j^{n+1} + u_{j-1}^{n+1} \right) + bu_j^n + O(\tau + h^2).
$$
 (S.3)

A reduced form of two-dimensional parabolic equation used in our simulation could be described by

$$
\partial_t u = a \left( \partial_{xx} + \partial_{yy} \right) u + bu. \tag{S.4}
$$

With extending the difference scheme of Eq.  $(S.3)$ , we deduce the numerical scheme,

$$
\frac{u_{j,l}^{n+1} - u_{j,l}^n}{\tau} = \frac{a}{h^2} \left( u_{j+1,l}^n - u_{j,l}^n - u_{j,l}^{n+1} + u_{j-1,l}^{n+1} \right) + \frac{a}{h^2} \left( u_{j,l+1}^n - u_{j,l}^n - u_{j,l}^{n+1} + u_{j,l-1}^{n+1} \right) + bu_j^n + O(\tau + h^2).
$$
 (S.5)

After sorting out the above equation, we obtain the final iterative scheme performed in our simulations,

$$
u_{j,l}^{n+1} = \frac{1+b\tau - 2\lambda}{1+2\lambda} u_{j,l}^n + \frac{\lambda}{1+2\lambda} \left( u_{j+1,l}^n + u_{j,l+1}^n + u_{j-1,l}^{n+1} + u_{j,l-1}^{n+1} \right),
$$
\n(S.6)

where  $\lambda = a\tau/h^2$ .

### <span id="page-4-0"></span>**References**

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