# **Supporting Material for Reversible Effects of Peptide Concentration and Lipid Composition on H-Ras Lipid Anchor Clustering**

X. Lin, Z. Li, and A. A. Gorfe\*

Department of Integrative Biology and Pharmacology, Medical School, The University of Texas Health Science Center at Houston, Houston, Texas 77030, United States

\* Corresponding author: E-mail: [alemayehu.g.abebe@uth.tmc.edu](mailto:alemayehu.g.abebe@uth.tmc.edu)

## **Contents**



## **1. Lipid domain formation and effect of bilayer asymmetry on tH clustering**

The simulations in the current work involved a bilayer of DPPC (dipalmitoylphosphatidylcholine), DLiPC (dilinoleoylphosphatidylcholine) and CHOL (cholesterol), with or without H-Ras lipid anchor (tH) embedded in one leaflet. To find a bilayer system capable of domain formation and to examine the effect of asymmetric tH insertion on clustering, we tested the three test systems listed in Table S1. As shown in Fig S1, simulation of these systems from a random initial mixture led to rapid (several  $\mu$ s) lipid de-mixing and formation of two striped domains (Fig S1a). One of these domains is enriched with DPPC and CHOL (liquid ordered domain,  $L_0$ ), while the other is enriched with DLiPC (liquid disordered domain,  $L_d$ ). The two domains differ by  $\sim 0.8$ nm in thickness (see Fig S1b).

System	Upper leaflet				Lower leaflet				Length $(\mu s)$
	<b>DPPC</b>	<b>DLiPC</b>	<b>CHOL</b>	tH	<b>DPPC</b>	<b>DLiPC</b>	<b>CHOL</b>	tH	
	480	288	192	$\theta$	480	288	192	$\boldsymbol{0}$	12
2	480	288	192	$\theta$	480	288	192	64	12
3	480	288	192	$\theta$	416	256	192	64	12

**Table S1: Summary of test simulations performed in this work&**

 $\overline{\text{a}}$  Simulation details are as described in the following section; system 2 is the same as the 20% CHOL system discussed in the main text and in Table S2.

Note that we inserted tH only to the lower leaflet in order to mimic Ras binding to the inner leaflet of plasma membrane. This may cause inter-leaflet area asymmetry and curvature. A detailed analysis of this issue in previous work<sup>1</sup> showed that the total area of the leaflet with 64 tH molecules (the maximum number of tH used in the current study) was  $563 \text{ nm}^2$  while that of the tH-free leaflet was 557 nm<sup>2</sup>, resulting in a difference of only 6 nm<sup>2</sup> or 1.1% (see Fig 7a in ref. 1). This is because (i) the maximum fraction of tH relative to the total number of molecules in the bilayer is only 3.2%, and (ii) area mismatch is partially relieved by cholesterol re-distribution between leaflets (see Fig 6a in ref. 1). In one of the current simulations containing 64 asymmetrically bound tH (system 2),  $74.6\pm7.3$ cholesterol molecules have transferred from the lower to the upper leaflet (see Fig. S2c).

Park *et al*<sup>2</sup> reported that area per lipid mismatch of up to 5% is tolerated in all-atom MD simulation of reasonably-sized (40-160 lipids per leaflet) bilayers. Assuming a direct correlation between total area and area per lipid, the  $\sim$ 1.1% inter-leaflet area mismatch in our system should not have significant effect on tH binding or clustering. Nevertheless, we directly tested the impact of area asymmetry on tH clustering by simulating a bilayer in which the overall surface area of the two leaflets was nearly the same. This was achieved by deleting 64 DPPC and 32 DLiPC lipids from the leaflet in which tH was bound. The number of deleted lipids was determined based on the work of Vogel *et al*.<sup>3,4</sup>, who found that the average cross-sectional area of a saturated lipid tail of N-Ras varied between  $0.236$  nm<sup>2</sup> and 0.354 nm<sup>2</sup>, depending on the chain length and saturation of the host bilayer. For a DPPC/CHOL

bilayer, the cross-secretion areas of N-Ras and DPPC lipid chains were similar  $(0.236 \text{ vs. } 0.227 \text{ nm}^2)$ ; table 1 of Vogel *et al*<sup>3</sup>). When the host bilayer was POPC, the cross-section of an N-Ras hydrocarbon tail is similar to that of POPC  $(0.333 \text{ vs. } 0.306 \text{ nm}^2)$ . Since tH predominantly localizes at the lipid domain boundary so that--on the main--its two saturated lipids interact with the DPPC/CHOL domain and its unsaturated farnesyl tail with the DLiPC domain<sup>5,6</sup>, we estimated that the total cross-section of 128 palmitoyl chains ( $64 \times 2$ ) would be roughly equal to the total area of 64 DPPC lipids ( $64 \times 2$  16:0 chains). Since the DLiPC-interacting farnesyl chain is more flexible, we estimated that the crosssection of 64 farnesyls would be equivalent to the area of 32 DLiPC lipids. Therefore, we deleted 64 DPPC and 32 DLiPC from the leaflet containing 64 tH molecules. Simulation of this system (system 3) yielded domain thicknesses (Fig. S1b), number of cholesterol/leaflet (Fig. S1c) and box area (Fig. S1d) that are comparable with those from a tH-free bilayer (system 1). More importantly, there was no significant difference between systems 3 and 2 in terms of lipid domain formation and structure (Fig. S1a, b), tH clustering profile (Fig. S1e, f) and even overall tH lateral distribution (Fig. S1a). These results demonstrate that, at least at the concentration regimes of the current work, asymmetric tH insertion faithfully models membrane binding of cellular Ras without adversely affecting cluster formation.



Figure S1. (a) Two-dimensional plots of average number density of DPPC (blue to red) and the location of tH molecules (black points); additional details can be found the in legend of Fig 1 in the main text. (b) Thicknesses of the  $L_0$  and  $L_d$  domains measured as the average inter-leaflet head-to-head distance based on the z-position of the PO4 beads of DPPC and DLiPC, respectively. To account for curvature and undulations, the thicknesses were calculated in 1nm-side slabs and averaged over the slabs. (c) Average number of cholesterol in each leaflet of the bilayer. (d) Average area of the simulation box in the x-y plane. (e-f) Time evolution of average cluster size (d) and probability distribution (f) of tH cluster sizes in systems 1 and 2. Except in (e), the last 4  $\mu$ s of the trajectory was used; standard deviations are shown as error bars in (b), (c) and (d).

## **2. Simulation Details**

In addition to the test systems referred to in the previous section, we simulated eight systems each in replicates of three (Table S2). In four of these the DPPC/DLiPC/CHOL ratio was kept fixed while the number of tH was progressively reduced as described in the main text. In the other four, the fraction of cholesterol was progressively decreased while keeping the number of peptides fixed at 64 (and always maintaining equal numbers of cholesterol in the two leaflets in the beginning). In every case, tH was asymmetrically bound to one leaflet, which, as shown above, has little impact on clustering. Following model building, each system was first (re)-equilibrated in two steps each of 100 ns duration. In the first step, tH, lipid and  $C_{60}$  (where applicable) molecules were position-restrained using a harmonic force constant of 1000 kJ  $mol^{-1}$   $nm^{-2}$ . In the second step, the position restraint was applied only on tH. Subsequently, all restraints were removed and production simulations commenced for the durations shown in Table S2. In each simulation, the Lennard-Jones potentials describing van der Waals interactions were smoothly shifted to zero between 0.9 nm and 1.2 nm, Coulomb electrostatic interactions were truncated at 1.2 nm, and a default dielectric constant of 15 was used<sup>7</sup>. Lipids, H-Ras lipid anchors (tH), water and ions were coupled separately to V-rescale heat baths<sup>8</sup> at  $T = 301K$  with a coupling constant  $\tau = 1 \text{ ps}$ . Constant number of particle, pressure and temperature (NPT ensemble) simulations were conducted at  $1 bar$  using semi-isotropic Parrinello-Rahman pressure coupling scheme<sup>9</sup> with a coupling constant  $\tau = 5 \rho s$  and a compressibility of  $3 \times 10^{-4}$  *bar*<sup>-1</sup>. The neighbor list for non-bonded interactions was updated every 10 steps and the time step was 20 fs. A summary of the simulations performed in this work is presented in Table S1. We have used the  $8-12 \mu s$  of each trajectory for the analysis of cluster size distribution as well as number of density of lipids, and the last  $32 \mu s$  of each 40  $\mu s$  trajectory for the analysis of long timescale processes such as molecular expulsion autocorrelation functions.





\*  $S_{a,b,c}$  represents the name of the system, where "a" is the number of tH molecules, "b" is the lipid fraction of cholesterol and "c" is the number of  $C_{60}$  nanoparticles. Simulation length is the effective time (4  $\times$  actual simulation length). The simulations were run for 12  $\mu$ s for cases where we were only interested in cluster size distribution or 40  $\mu$ s for cases where we were also interested in analyzing the compositional dynamics of tH nanoclusters whose relaxation time was relatively long (ref. 8, 9).

 $\frac{1}{2}$  Each system was simulated in replicates of three with the second and third runs being 12  $\mu$ s long. The results from the replicates were very similar (see main text) and therefore we used only one of the copies for most of the analysis.

#### **3. Time evolution of average cluster sizes**

Clustering was monitored by a Single-Linkage (SL) algorithm described previously<sup>10-12</sup>, with neighbors defined based on a distance cutoff of 7.5 Å; this cutoff was derived from an analysis of twodimensional pair distributions. Fig. S2 shows the time evolution of average cluster size.



**Figure S2**. Time evolution of average cluster sizes during simulations with different peptide (a) and cholesterol (b) concentrations, and in the presence of  $C_{60}$  (c).

#### **4. Error analysis**

The error bars in Fig. 1c and 2b (main text) were calculated using averaging with a single 500ns block size over the last 4  $\mu$ s of 12  $\mu$ s trajectories. To check if this block size is reasonable, we divided the 4  $\mu$ s trajectory into  $N_b$  blocks of progressively increasing size and evaluated the sampling error  $\sigma_b$  as

$$
\sigma_b = \sqrt{\frac{1}{N_b} \sum_{i=1}^{N_b} (x_i - \bar{x})^2}
$$
 (1)

where  $x_i$  and  $\bar{x}$  are averages over block size i<4µs and 4µs, respectively. As shown in Fig. S3, larger block sizes result in smaller errors. In most cases, the errors for different cluster sizes and simulation conditions plateaued at about 400-600ns. The relatively large scattering for the larger clusters at larger block sizes reflects the diminishing number of samples (smaller  $N_b$ ). We therefore chose a reasonable tradeoff of 500ns block size ( $N_b = 8$ ) for estimation of the errors shown in Figs. 1 and 2 of the main text and section 7 of the Supporting Material.



Figure S3. Error as a function of block size for different cluster sizes, calculated using equation 1 (lines represent exponential fits using the equation  $y = y_0 + A_1 e^{-(x-x_0)/t_1} + A_2 e^{-(x-x_0)/t_2} + A_3 e^{-(x-x_0)/t_3}$ .

#### **5. Stability of H-Ras lipid anchor nanoclusters**

In order to estimate the characteristic time scales of molecular exchange among nanoclusters or between nanoclusters and the non-clustered fraction, we used an autocorrelation function  $f(t)$  defined as:

$$
f(t) = \langle \frac{N(t_0) - N_{leave}(t_0 + t)}{N(t_0)} \rangle_{t_0, n}
$$
 (2)

where  $N(t_0)$  is the nanocluster size at time  $t_0$  ( $t_0$  representing the time point at which we begin following a given cluster),  $N_{leave}(t_0 + t)$  is the number of molecules that leave the given cluster after a time lapse of t, n is the number of nanoclusters of size  $N(t_0)$ . When  $t = 0$ ,  $N_{leave}(t_0 + t) = 0$ ,  $f(t) =$ 1, the composition of the nanocluster is identical to the initial state. When  $t \to \infty$ ,  $N_{leave}(t_0 + t) \to$  $N(t_0)$ ,  $f(t) \rightarrow 0$ , the composition of the nanocluster is completely different from the initial state. In other words,  $f(t)$  is the fraction of tH molecules that remain in the same initial cluster at time *t* relative to time  $t_0$  so that the speed with which  $f(t)$  drops from 1 to 0 can be regarded as the decay rate of the given cluster. Additional details can be found in our previous reports<sup>5, 6</sup>.

#### **6. Effect of C60 on dynamics of membrane domains and H-Ras lipid anchor nanoclusters**



Figure S4. (a) Side-view of the initial system setup of  $C_{60}$  and tH nanoclusters (C60 is in green; color scheme for the rest of the molecules is the same as in main text Fig. 1). (b) Bottom-view of the last snapshot. (c) Twodimensional normalized DPPC number density with the black points representing the location of tH. (d) Effect of  $C_{60}$  on tH nanocluster size distribution. e) Effect of  $C_{60}$  on tH nanocluster dynamics.



## **7. Comparison of tH clustering and dynamics upon depletion (current work) and increase (previous work) of peptide/cholesterol content.**

Figure S5. Reversible effects of decreasing and increasing peptide/cholesterol concentrations on tH nanocluster size distribution (a-d) and stability (e and f, for cluster size 6).



## **8. Results of two additional replicates of the simulations listed in Table 2.**

Figure S6. tH nanocluster size distribution from two replicates of each of the simulations listed in Table 2: (a, d), simulations with decreasing tH fraction; (b,c) simulations with decreasing cholesterol fraction; (c, f) simulations in the presence of the C60 nanoparticle. Note that the data in the two top left panels are also shown as inset in Figs. 1 and 2 of the main text.

### **Supporting References:**

- 1. Z. Li and A. Gorfe, *Soft Matter*, 2013, 9, 11249-11256.
- 2. S. Park, A. H. Beaven, J. B. Klauda and W. Im, *J. Chem. Theory Comput.*, 2015, **11**, 3466-3477.
- 3. A. Vogel, G. Reuther, K. Weise, G. Triola, J. Nikolaus, K. T. Tan, C. Nowak, A. Herrmann, H. Waldmann and R. Winter, *Angew. Chem. Int. Ed.*, 2009, 48, 8784-8787.
- 4. A. Vogel, C. P. Katzka, H. Waldmann, K. Arnold, M. F. Brown and D. Huster, *J. Am. Chem. Soc.*, 2005, **127**, 12263-12272.
- 5. L. Janosi, Z. Li, J. F. Hancock and A. A. Gorfe, *Proc. Natl. Acad. Sci. U. S. A.*, 2012, **109**, 8097-8102.
- 6. Z. Li, L. Janosi and A. A. Gorfe, *J. Am. Chem. Soc.*, 2012, **134**, 17278-17285.
- 7. S. J. Marrink, H. J. Risselada, S. Yefimov, D. P. Tieleman and A. H. De Vries, *J. Phys. Chem. B*, 2007, **111**, 7812-7824.
- 8. G. Bussi, D. Donadio and M. Parrinello, *J. Chem. Phys.*, 2007, **126**, 014101.
- 9. M. Parrinello and A. Rahman, *J. Appl. Phys.*, 1981, **52**, 7182-7190.
- 10. J. C. Gower and G. Ross, *J. R. Stat. Soc. Ser. C-Appl. Stat.*, 1969, **18**, 54-64.
- 11. Z. Li and E. E. Dormidontova, *Soft Matter*, 2011, **7**, 4179-4188.
- 12. M. V. Gudheti, N. M. Curthoys, T. J. Gould, D. Kim, M. S. Gunewardene, K. A. Gabor, J. A. Gosse, C. H. Kim, J. Zimmerberg and S. T. Hess, *Biophys. J.*, 2013, 104, 2182-2192.