

## Article

# A Simple Model of Multivalent Adhesion and Its Application to Influenza Infection

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**ABSTRACT** Adhesion between biological surfaces, which is typically the result of molecular binding between receptors on one surface and ligands on another, plays a fundamental role in biology and is key to the infection mechanisms of certain viruses, including influenza. The physiological outcome of adhesion depends on both the number of bound cells (or viruses, or other biological particles) and the properties of the adhesion interface that is formed, including the equilibrium number of receptor-ligand connections. Here, we introduce a quantitative model for biological adhesion by adapting thermodynamic models developed for the related problem of multivalent molecular binding. In our model, adhesion affinity is approximated by a simple, analytical expression involving the numbers of ligands and receptors at the interface. Our model contains only two fitting parameters and is simple to interpret. When applied to the adhesion between the hemagglutinin ligands on influenza viruses and the sialic acid receptors on biosensors or on host cells, our model generates adhesion affinities consistent with experimental measurements performed over a range of numbers of receptors, and provides a semiquantitative estimate of the affinity range of the hemagglutinin-sialic acid interaction necessary for the influenza virus to successfully infect host cells. The model also provides a quantitative explanation for the experimental finding that a mutant avian virus gained transmissibility in mammals despite the mutations conferring only a less than twofold increase in the affinity of its hemagglutinin for mammalian receptors: the model predicts an order-of-magnitude improvement in adhesion to mammalian cells. We also extend our model to describe the competitive inhibition of adhesion: the model predicts that hemagglutinin inhibitors of relatively modest affinity can dramatically reduce influenza virus adhesion to host cells, suggesting that such inhibitors, if discovered, may be viable therapeutic agents against influenza.

## INTRODUCTION

Adhesion between two cells, or between a virus and a cell, mediates diverse biological phenomena including pathogen recognition of host cells (1–3), cell trafficking (4–7), and cell signaling (8,9). Such adhesion typically involves the simultaneous binding of multiple copies of ligands on one of the contact surfaces with multiple copies of receptors on the other (10,11), and at biological interfaces there are often hundreds to thousands of such ligand-receptor pairs. Influenza virus, for example, in its first step of infection, attaches to vertebrate host cells through the binding between multiple hemagglutinin (HA) trimers on the viral surface and multiple copies of the *N*-acetylneuraminic acid (also known as sialic acid (SA)) moiety on the host-cell surface (12). Although the binding between a single pair of HA and SA molecules is weak (equilibrium dissociation constant  $K_D \sim 1$  mM) (13), the large number of HA and SA molecules present at the contact surface collectively ensures that the virus adheres to the host cell with high affinity. This

strong adhesion is thus a consequence of multivalent (also referred to as polyvalent) interactions, where valency refers to the number of simultaneous connections between one kind of particle and another.

The physiological outcome of adhesion depends on both the number of bound cells (or viruses, or other biological particles) and the properties of the adhesion interface that is formed, such as the contact surface area, the distance of cell-cell separation, and the equilibrium number of receptor-ligand connections. Thermodynamic models have been developed to study the properties of the individual adhesion interface (14–18). Such models, however, do not calculate the fraction of bound cells in a given condition (15). Full characterization of adhesion requires models that can estimate the affinity of adhesion and thus the population of bound cells, which can now be measured experimentally. In this study we develop such a model.

Many biological processes involve modulating the strength of multivalent adhesion, either by changing the number of available receptors or ligands at the interface or by changing the binding affinity of each receptor-ligand pair. Quantitative characterization of such modulation requires a model that accurately estimates the adhesion affinity from the parameters of the underlying receptor-ligand binding.

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Several empirical models of multivalent adhesion have been reported in the literature (11). One widely used approach approximates the dissociation constant of adhesion,  $K_{D,ad}$ , as  $(K_D)^m$ , where  $m$  is referred to as the multiplicity, the value of which depends on the number of adhesive receptor and ligand molecules at the contact surface (19). To our knowledge, however, no theoretical model has been proposed to estimate  $m$  from the number of receptors,  $N_R$ , the number of ligands,  $N_L$ , and  $K_D$ : It can only be measured empirically at discrete, static numbers of receptors and/or ligands. As suggested above, it is often necessary to consider adhesion for changing numbers of receptors and ligands and for a range of binding affinities. The empirical approach is typically difficult to apply in these problems.

Thermodynamic models have been developed for the related problem of binding between molecules with multiple binding sites, such as in multivalent host-guest binding (20–25). These models relate the avidity of multivalent binding to the numbers of available binding sites on the multivalent molecules (receptors and ligands or host and guest molecules) and to the  $K_D$  of individual pairs of binding sites. These models have enjoyed great success treating molecular systems in which there are  $< \sim 10$  binding site pairs, but they have not been applied to adhesion between larger biological particles. Adhesion between larger particles (e.g., between a virus and a host cell) often involves a much larger contact interface, which incurs nonspecific interactions, such as electrostatic repulsion between membranes, that can significantly affect the adhesion affinity.

In this work, we extend previous thermodynamic models of multivalent molecular binding to propose a simple model of the phenomenon of multivalent adhesion, in part by introducing a fitting parameter to account for the contribution of nonspecific interactions. We also introduce approximations appropriate for the large number of adhesive connections, resulting in a simple equation that estimates  $K_{D,ad}$  from  $K_D$ ,  $N_R$ , and  $N_L$ . Our model establishes a clear relationship between  $K_{D,ad}$  and the equilibrium number of receptor-ligand connections, the latter being a key quantity in the groundbreaking model developed by Bell and co-workers (14–18).

To demonstrate the utility of this approach, we use our model to study adhesion of influenza virus to SA-conjugated biosensors and to host cells. The influenza virus exploits two mechanisms to modulate its adhesion to host cells: 1) changing the number of receptors to spread from cell to cell (26,27), and 2) changing the affinity of receptor-ligand binding to adapt from avian to mammalian hosts (28–30). We apply our model to the analysis of both mechanisms of modulation. Our model accurately reproduces the experimentally measured relationship between adhesion and the number of available SA molecules conjugated to biosensors, and quantitatively explains the observed

sensitivity of adhesion affinity to small changes in the HA-SA binding affinity. Our model also provides a quantitative technique for analyzing the functional interplay between HA-SA binding affinity and the catalytic efficiency of neuraminidase, and yields a reasonable estimate for the range of affinity of HA-SA binding that leads to a viable virus.

We also extend our model to the study of competitive inhibition of multivalent adhesion. A competitive inhibitor of the adhesive molecules—against either the receptor or the ligand—reduces the number of available connection sites, thus weakening or abolishing adhesion (31). The efficacy of such an inhibitor depends not only on its affinity for its molecular target, but also on how the adhesion affinity changes with respect to the numbers of available receptors and ligands. Our model permits a quantitative analysis of the competitive inhibition of multivalent adhesion. When applied to influenza, our model indicates that it may be possible to effectively inhibit the adhesion of influenza virus to host cells using a molecule with only modest affinity for HA. It might thus be feasible to take existing modest inhibitors of HA, or discover new ones, and develop them as therapeutic agents against influenza infection.

## MATERIALS AND METHODS

Please see the [Supporting Material](#).

## RESULTS AND DISCUSSION

### Theoretical model

Our model estimates  $K_{D,ad}$  of adhesion from the  $K_D$  of individual receptor-ligand binding and the numbers of receptors,  $N_R$ , and ligands,  $N_L$ , at the adhesion interface. The result, which we derive later in this section, is

$$K_{D,ad} = K_0 \exp \left( N_R \ln \frac{N_R - \bar{n}}{N_R} + N_L \ln \frac{N_L - \bar{n}}{N_L} + \bar{n} \right), \quad (1)$$

where  $\bar{n}$  is the equilibrium number of connections formed between receptors and ligands, and is obtained from the equilibrium equation

$$\frac{(N_R - \bar{n})(N_L - \bar{n})}{\bar{n}} = K_D V_{\text{eff}}. \quad (2)$$

Here,  $K_0$  is the 0-valency dissociation constant between the two surfaces when no connections are formed, and  $V_{\text{eff}}$  is the effective molar volume in which receptor-ligand binding occurs.  $K_0$  and  $V_{\text{eff}}$  are difficult to measure directly, but can be estimated by fitting to adhesion measurements at varying values of  $N_R$  and/or  $N_L$ . Typically,  $K_D$  can be measured by standard binding assays or calorimetry, and  $N_R$  and  $N_L$  can be measured by electron microscopy or fluorescence. In

the specific case of HA-SA-mediated adhesion by the influenza virus, the HA-SA binding affinity has been measured by nuclear magnetic resonance (13) and micro-scale thermophoresis (MST) (19), and the number of sialic acids on the cell surface has been measured by cleavage experiments (32).

We first derive Eq. 1 by the method of statistical mechanics, and then give a simple interpretation of the result at the end. At the adhesion interface, individual connections between receptors and ligands form and break in a thermodynamic equilibrium, and the number of connections at a given instant,  $n$ , fluctuates between 0 and  $\min(N_R, N_L)$ . We assume that the connections form and break independently and that the Gibbs free energy change of forming a connection between a given ligand and receptor is  $\Delta G = RT \ln(K_D V_{\text{eff}})$ , with  $R$  being the gas constant and  $T$  the temperature. The molar volume  $V_{\text{eff}}$  accounts for the adjustment of the reaction condition, because the binding between the receptors and the ligands is restrained in the narrow confine of the adhesive interface, and both the receptors and the ligands are tethered to their respective surfaces. Our model in this form ignores potential cooperativity in receptor-ligand binding (25), in which the binding free energy of a given receptor-ligand pair depends on the number of existing connections. Such cooperativity, however, can be included in our model at the cost of extra parameters, as discussed in a subsequent section.

At thermal equilibrium, the probability that  $n$  connections are formed is given by

$$\begin{aligned} p_n &= Z^{-1} \frac{1}{N_R!} \frac{1}{N_L!} \exp\left(-\frac{n\Delta G}{RT}\right) \binom{N_R}{n} \binom{N_L}{n} n! \\ &= Z^{-1} \exp\left(-\frac{n\Delta G}{RT}\right) \frac{1}{(N_R - n)!(N_L - n)!n!}, \end{aligned} \quad (3)$$

where

$$Z = \sum_{n=0}^{\min(N_R, N_L)} \exp\left(-\frac{n\Delta G}{RT}\right) \frac{1}{(N_R - n)!(N_L - n)!n!} \quad (4)$$

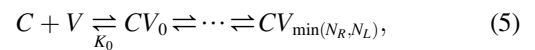
is the grand canonical partition function. In the first line of Eq. 3, the binomials

$$\binom{N_R}{n} \binom{N_L}{n}$$

count the numbers of ways to choose  $n$  receptors out of a total  $N_R$  and  $n$  ligands out of a total  $N_L$  to form  $n$  connections. The value  $n!$  stands for the number of ways to pair the  $n$  chosen receptors and ligands. The factor  $1/(N_R!N_L!)$  takes into account that the receptors are indistinguishable from one another and so are the ligands. In situations where the receptors are attached to fixed locations on a surface (such as experiments in which the receptors are teth-

ered to a surface through streptavidin-biotin linkers), thus becoming distinguishable,  $N_R!$  should be removed from the denominator. The value  $p_n$ , however, remains unchanged. Equation 3 is an extension of the Klotz equation, first developed for the binding of ligands to receptors with multiple binding sites (33). The exact number of receptors and ligands at the contact surface varies among individual biological particles; in our model we take  $N_R$  to be the average number of receptors on receptor-bearing particles, and  $N_L$  to be the average number of ligands on ligand-bearing particles.

The above receptor-ligand binding equilibrium is part of the overall equilibrium of adhesion between two biological surfaces  $C$  and  $V$ ,



where  $CV_n$  stands for the microstate of the adhesion surface where  $n$  connections are formed between the receptors and the ligands. The 0-valency dissociation constant,  $K_0$ , accounts for the contribution of nonspecific interactions to the adhesion affinity.

The dissociation constant of adhesion is thus given by

$$\begin{aligned} K_{D,\text{ad}} &= \frac{[C][V]}{[CV]} = \frac{[C][V]}{\sum_{n=0}^{\min(N_R, N_L)} [CV_n]} \\ &= \frac{[C][V]}{[CV_0]} p_0 = K_0 p_0. \end{aligned} \quad (6)$$

Given  $K_D$ ,  $N_R$ ,  $N_L$ ,  $K_0$ , and  $V_{\text{eff}}$ , one can calculate  $K_{D,\text{ad}}$  from Eqs. 3, 4, and 6 using a computer program. When  $N_R$  and  $N_L$  are large, however, we can introduce approximations that greatly simplify the expression of  $K_{D,\text{ad}}$  and remove the factorials in the equation.

Under conditions where reasonably large numbers of connections form, the grand canonical partition function  $Z$  can be well approximated by the maximum term in the summand,

$$Z \approx \exp\left(-\frac{\bar{n}\Delta G}{RT}\right) \frac{1}{(N_R - \bar{n})!(N_L - \bar{n})!\bar{n}!}, \quad (7)$$

where  $\bar{n}$  is the number at which  $p_n$  is at its maximum. The value  $\bar{n}$  thus satisfies  $p_{\bar{n}}/p_{\bar{n}+1} \approx 1$ , so

$$\frac{(N_R - \bar{n})(N_L - \bar{n})}{\bar{n}} \approx \exp\left(\frac{\Delta G}{RT}\right) = K_D V_{\text{eff}}, \quad (8)$$

which implies that  $\bar{n}$  is the equilibrium number of connections formed between the receptors and the ligands confined within the adhesion interface. Applying the Stirling

approximation  $\ln x! = x \ln x - x + \ln p_0 = -\ln(Z N_R! N_L!)$ , we have

$$\begin{aligned}
\ln p_0 &= -\ln Z - \ln N_R! - \ln N_L! \\
&\approx \frac{\bar{n}\Delta G}{RT} + (N_R - \bar{n})\ln(N_R - \bar{n}) - (N_R - \bar{n}) \\
&\quad + (N_L - \bar{n})\ln(N_L - \bar{n}) - (N_L - \bar{n}) + \bar{n}\ln\bar{n} \\
&\quad - \bar{n} - N_R \ln N_R + N_R - N_L \ln N_L + N_L \\
&= N_R \ln \frac{N_R - \bar{n}}{N_R} + N_L \ln \frac{N_L - \bar{n}}{N_L} + \bar{n} \\
&\quad + \bar{n} \left( \frac{\Delta G}{RT} - \ln \frac{(N_R - \bar{n})(N_L - \bar{n})}{\bar{n}} \right) \\
&\approx N_R \ln \frac{N_R - \bar{n}}{N_R} + N_L \ln \frac{N_L - \bar{n}}{N_L} + \bar{n},
\end{aligned} \tag{9}$$

where the last line is due to Eq. 8. Substituting  $p_0$  above in Eq. 6, we arrive at our expression for  $K_{D,\text{ad}}$  in Eq. 1.

Equations 1 and 2 can also be derived by considering the following imaginary thermodynamic process (because the free energy change of adhesion is the difference between two state functions and is independent of the connecting thermodynamic path). The reaction of adhesion can be approximately decomposed into two steps: 1) the attachment of the two surfaces without the formation of any receptor-ligand connections, and 2) the establishment of the equilibrium of receptor-ligand binding at the interface. The total free energy of adhesion is thus

$$RT \ln K_{D,\text{ad}} = RT \ln K_0 + \Delta G_{\bar{n}}, \tag{10}$$

where  $\Delta G_{\bar{n}}$  is the free energy of forming  $\bar{n}$  receptor-ligand connections, and is given by

$$\Delta G_{\bar{n}} = \bar{n}\Delta G - RT \ln \left( \binom{N_R}{\bar{n}} \binom{N_L}{\bar{n}} \bar{n}! \right), \tag{11}$$

where the second term on the right-hand side is the entropic contribution due to the number of ways to form  $\bar{n}$  connections out of  $N_R$  receptors and  $N_L$  ligands. Equations 10 and 11, with the Stirling approximation and the same algebra as in Eq. 9, lead to Eq. 1.

Finally, we derive an approximate expression for the so-called 0-valency dissociation constant  $K_0$ , which is valid when the contribution of nonspecific interactions to this quantity can be neglected. Suppose that the total numbers of receptors and ligands exposed on the respective biological particles are  $M_R$  and  $M_L$ , out of which the  $N_R$  receptors and the  $N_L$  ligands that form the connections are the subsets that happen to lie within the interface. The value  $K_0$  can be related to the fractions  $N_R/M_R$  and  $N_L/M_L$ , and the effective molar volume  $V_{\text{eff}}$ , as follows.

The reaction to form one connection between the two biological particles is



If we ignore the nonspecific interactions between the two biological particles, however, this monovalent reaction is equivalent to the reaction of individual receptor-ligand binding, where the receptor and ligand concentrations are  $M_R$  and  $M_L$  times the concentrations of the respective biological particles. As a result, we have

$$\frac{[C][V]}{[CV_1]} \frac{([R]/M_R)([L]/M_L)}{[RL]} = \frac{K_D}{M_R M_L} \tag{13}$$

and

$$K_0 = \frac{[C][V]}{[CV_0]} = \frac{[C][V]}{[CV_1]} \frac{[CV_1]}{[CV_0]} = \frac{K_D}{M_R M_L} \frac{p_1}{p_0} = \frac{N_R}{M_R} \frac{N_L}{M_L} V_{\text{eff}}^{-1}. \tag{14}$$

If the fractions  $N_R/M_R$  and  $N_L/M_L$  are known,  $V_{\text{eff}}$  becomes the only free parameter in our model. In this work, we treat  $K_0$  as a free fitting parameter, and we check its estimated value against sensible guesses of the fractions.

We have so far developed the model assuming that one particle of  $V$  adheres to one particle of  $C$ . Different types of biological particles, however, often have disparate sizes. If  $C$  is much larger than  $V$ , many particles of  $V$  can adhere to one particle of  $C$ , each occupying a nonoverlapping site of adhesion and making a number of adhesive molecules on  $C$ , which we can call  $M^*$ , unavailable to other particles of  $V$ . If the total number of adhesive molecules on one particle of  $C$  is  $M$ , the number of particles of  $V$  adhering to one particle of  $C$  at adhesion equilibrium can be shown to be

$$\bar{v} = \frac{S}{1 + S K'_{D,\text{ad}}/[V]}, \tag{15}$$

where  $S = M/M^*$  is the maximum number of particles of  $V$  that can be accommodated on the surface of one particle of  $C$ . Equation 15 is equivalent to the Langmuir equation,  $S$  being the number of sites of adhesion, and  $K'_{D,\text{ad}} \equiv S K_{D,\text{ad}}$  being the per-site dissociation constant of adhesion. Correspondingly, the per-site 0-valency dissociation constant is  $K'_0 = S K_0$ . Clearly,  $K'_{D,\text{ad}}/K'_0 = K_{D,\text{ad}}/K_0$ , and within our model this ratio still follows from Eq. 1. The value  $K_0$  follows from Eq. 14; note that  $K'_0$  is also given by Eq. 14 if we reinterpret  $M_R$  (or  $M_L$ , if the ligands rather than receptors lie on  $C$ ) as the number of adhesive molecules per site of adhesion (i.e., the number of adhesive molecules on  $C$  made unavailable to other particles of  $V$  by adhesion of one particle of  $V$ ). Below, we use the more convenient per-site dissociation constants  $K'_0$  and  $K'_{D,\text{ad}}$ , but we switch notation by dropping the prime symbols.

Our model in its simple form (Eqs. 1 and 2) assumes that only one type of receptor and one type of ligand are present at the interface, and that the binding affinity between any receptor-ligand pair is the same. Real cells, however, may present diverse sets of different receptor types, each with a

different binding affinity for the ligand. A vertebrate cell, for example, may present a range of different glycans that terminate with the SA moiety, and each type of glycan may have a different binding affinity for the HA molecule (34–36). It is straightforward to extend our model to account for such heterogeneity. If the number of receptors with binding affinity between  $K$  and  $K + dK$  is  $N(K)dK$ , the dissociation constant of adhesion in our model is given by

$$K_{D,ad} = K_0 \exp \left( \int N(K) \ln \frac{N(K) - \bar{n}(K)}{N(K)} dK + N_L \ln \frac{N_L - \bar{n}}{N_L} + \bar{n} \right), \quad (16)$$

where  $\bar{n}(K)dK$ , satisfying the equilibrium equation

$$\frac{(N(K) - \bar{n}(K))(N_L - \bar{n})}{\bar{n}(K)} = KV_{\text{eff}}, \quad (17)$$

is the equilibrium number of receptor-ligand connections with binding affinity between  $K$  and  $K + dK$ , and  $\bar{n} = \int \bar{n}(K)dK$  is the total equilibrium number of connections formed between all receptors and the ligand molecules. The value  $\bar{n}$  can be determined from the following equation:

$$\bar{n} = \int N(K) \left( 1 + \frac{KV_{\text{eff}}}{N_L - \bar{n}} \right)^{-1} dK. \quad (18)$$

Equations 16–18 (derived in the [Supporting Material](#)) describe a cell's properties of adhesion once the distribution of binding affinity  $N(K)$  is known for the receptors on the cell. Equations 1 and 2 are a special case of Eqs. 16–18, with  $N(K) = \delta(K - K_D)$ . Our model, thus generalized, can be used to analyze the variability in adhesion affinity in a population of cells due to the fluctuations in the expression levels of different receptors.

### Properties of the model

The expression of  $p_n$  and the partition function  $Z$  (Eqs. 3 and 4, without the approximations introduced in our model) can also be derived by equations of mass action, as has been done in the development of previous models of multivalent molecular binding (23), if the combinatorial factors of pairing any receptor molecule with any ligand molecule are properly accounted for. The exact expression of  $Z$  (Eq. 4) can be used as is, but our approximation leads to a simple analytical expression of the adhesion affinity (Eq. 1), making apparent its relationship to the other variables. The expression in Eq. 1 makes it easy to include  $N_R$  and  $N_L$ , which are discrete variables in the exact solution, as continuous parameters in fitting to experimental adhesion data, when their values are unknown. It also relates the adhesion affinity to the equilibrium number of receptor-ligand connections,  $\bar{n}$ , in a simple formula.

Although Eq. 1 is derived by applying the Stirling approximation at large values of  $N_R$ ,  $N_L$ , and  $\bar{n}$ , Eq. 1 remains a good approximation even at small values of these quantities. [Fig. 1 A](#) plots the ratio  $p_0/p_{0,\text{exact}}$  as  $N_R$  is varied between 1 and 20,000, where  $p_{0,\text{exact}}$  is given by Eq. 3. In the case of weak receptor-ligand affinities ( $K_D \sim 1$  mM), as in the case of cell-influenza adhesion, the ratio remains very close to 1 throughout this range of  $N_R$ . At stronger affinity,  $K_D = 0.001$  mM, our model can overestimate the adhesion affinity by a factor of up to  $\sim 2$ , which may still be acceptable for most applications.

The strength of adhesion can be modulated by changing the available number of receptors (or ligands) at the adhesion interface. Consider the adhesion between viruses and cells: the fraction of bound virus particles is

$$f = \frac{1}{1 + K_{D,ad}/[\text{cell}]}, \quad (19)$$

where  $[\text{cell}]$  is the concentration of the cells. The value  $f$  changes with the adhesion affinity, and thus with the number of receptors  $N_R$ . [Fig. 2 A](#) plots  $f$  against  $N_R$ , holding the number of ligands,  $N_L$ , at various fixed values. These adhesion isotherms are steep at large values of  $N_L$ . By solving Eq. 19 for  $N_R$  for any given  $f$ , our model can be used to estimate  $N_{R,f}$ , the number of receptors at which that fraction of the virus particles is bound ([Fig. 2 B](#)). At high values of  $N_L$ , the

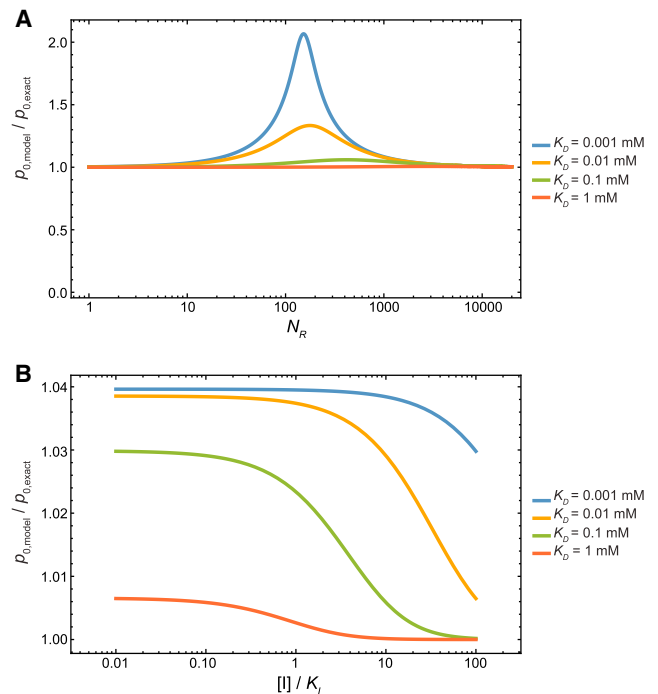
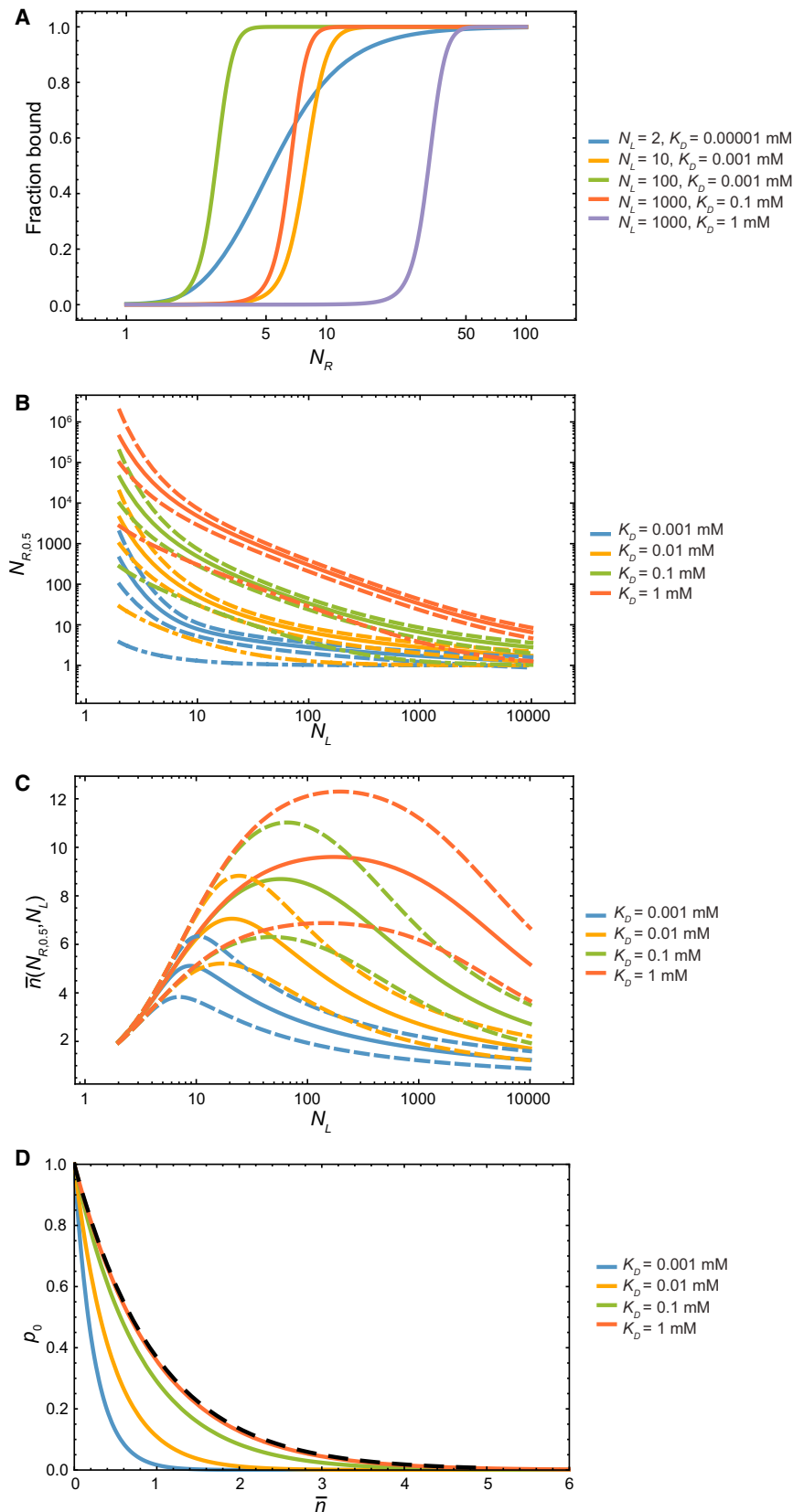


FIGURE 1 The accuracy of our model. (A) The ratio  $p_0/p_{0,\text{exact}}$  as  $N_R$  is varied between 1 and 20,000, for different receptor-ligand affinities. The other model parameters are  $N_L = 150$ ,  $V_{\text{eff}} = 2.7 \times 10^3 \text{ mM}^{-1}$ . (B) The ratio  $p_0/p_{0,\text{exact}}$  as the concentration of an inhibitor of the ligand is varied. The parameters  $N_L$  and  $V_{\text{eff}}$  have the same values as in (A), and  $N_R = 2000$ .



**FIGURE 2** Modulation of multivalent adhesion by changing the number of receptors at the adhesion interface. (A) The fraction of bound virus particles as a function of the number of receptors at the adhesion interface. (B)  $N_{R,0.5}$  (solid lines),  $N_{R,0.95}$  (dashed lines above), and  $N_{R,0.05}$  (dashed lines below) values (see text) as a function of the number of ligands at the adhesion surface. The gap between the two dashed lines narrows as  $N_L$  increases, indicating that at higher valencies it takes a smaller increase in the number of receptors to produce the same increase in the adhesion affinity. The number of receptors below which the equilibrium number of connections becomes no greater than 1 ( $\bar{n} \leq 1$ ) is also shown (dash-and-dot lines). (C) The equilibrium number of connections,  $\bar{n}$ , when the fraction of bound virus particles is 0.5 (solid lines), 0.95 (dashed lines), and 0.05 (dashed lines below). (D) The probability,  $p_0$ , of spontaneously breaking all receptor-ligand connections plotted against the equilibrium number of connections,  $\bar{n}$ , at varying numbers of receptors,  $N_R$ , while the number of ligands is held constant, at  $N_L = 150$ . The black dashed curve corresponds to the limit  $p_0 = \exp(-\bar{n})$ . Each color of lines (B–D) corresponds to a different receptor-ligand binding affinity  $K_D$ . The parameters used here are  $[\text{cell}] = 4.0 \times 10^{-10} \text{ mM}$ ,  $V_{\text{eff}} = 2.7 \times 10^3 \text{ mM}^{-1}$ , and  $K_0 = 1.0 \times 10^{-5} \text{ mM}$ ; the derivation of these parameter values is explained in the discussion of cell-influenza adhesion.

bound fractions increase quickly from 5 to 95%, with little change in the number of receptors.

To break apart an adhesion interface, it is not necessary to reduce the equilibrium number of receptor-ligand connections to  $\bar{n} = 0$ ; when  $\bar{n}$  becomes small enough, thermal fluctuations can spontaneously break all connections—so that the instantaneous number of connections is 0—which abolishes the adhesion. Without specific receptor-ligand connections, cells (and viruses) tend to repel each other, due to similar charges on their surfaces (15), such that  $K_D/[\text{cell}] \gg 1$ . Consequently, adhesion can be disrupted even when  $p_0$  is still very small, and a few receptor-ligand connections persist in binding equilibrium at the interface. At different values of  $K_D$  and  $N_L$ ,  $\bar{n}$  can take different values when half of the virus particles are bound (Fig. 2 C). Conversely, the number of receptors at which most of the virus particles are unbound ( $f = 0.05$ ) differs from the value below which  $\bar{n} \leq 1$  (Fig. 2 B). Our model will thus predict different and likely more accurate conditions under which adhesion is abolished than previous models, which consider that condition to occur when  $\bar{n} = 0$  (14).

If the affinity between the receptor and the ligand is sufficiently weak, the equilibrium number of connections is small (i.e.,  $\bar{n} \ll \min(N_R, N_L)$ ), and Eq. 9 can be well approximated by  $p_0 \approx \exp(-\bar{n})$ . Under such conditions, the adhesion affinity is  $(K_{D,\text{ad}})^{-1} \propto \exp(\bar{n}) \approx 1 + \bar{n} \approx \bar{n}$ , thus approximately agreeing with a previous suggestion that took adhesion affinity to be proportional to  $\bar{n}$  (16). At moderate receptor-ligand affinities, however,  $p_0$ , and thus the adhesion affinity, cannot be estimated from  $\bar{n}$  alone (Fig. 2 D).

### Adhesion of influenza virus

We now apply our model to the analysis of the binding of influenza virus. The influenza virus uses its surface glycoprotein HA to bind to terminal SA moieties of carbohydrate chains on the host-cell surface. SA is frequently referred to as the receptor, and HA as the ligand. HA of avian influenza (bird flu) strains preferentially binds to SA that is in  $\alpha 2,3$  linkage to its neighboring sugar (galactose) in carbohydrate chains (characteristic of cells in the avian enteric tracts), whereas HA of influenza strains transmissible in mammals, including humans, needs to bind to SA in  $\alpha 2,6$  linkage to galactose (characteristic of human trachea airway epithelia) (28,29). Mutations in the HA sequence can cause a change in its binding preference from  $\alpha 2,3$  linkage to  $\alpha 2,6$  linkage, enabling transmissibility of the virus from birds to mammals (30), as was demonstrated in highly pathogenic H5N1 influenza strains (37,38).

#### Adhesion to SA-conjugated biosensors

We will first use our model to analyze the adhesion of influenza virus to SA-conjugated biosensors. This simplified experimental system, which measures the binding of free vi-

rus to immobilized SA, has been used to give insights relevant to the attachment of influenza virus to vertebrate cells and, in particular, to help address the question of the species specificity of different viral strains.

Recently, Xiong et al. (19) measured the binding affinities between HA in various influenza strains and SA with  $\alpha 2,3$  and  $\alpha 2,6$  linkages. The authors also measured the adhesion affinities of the whole viruses of these strains (except for a ferret-transmissible mutant of an H5N1 virus) to SA attached to biosensors, and they demonstrated that a modest change in the receptor-ligand binding affinity is magnified to a much larger change in the viral adhesion to the cell surface. Using the empirical model  $K_{D,\text{ad}} = (K_D)^m$ , where the value of  $m$  depends on the SA concentration, and where its value at each SA concentration is estimated from  $\ln(K_{D,\text{ad}})/\ln(K_D)$  using data at multiple HA concentrations, the authors predicted the binding isotherms of the ferret transmissible mutant virus and suggested that the small changes in the HA-SA binding affinities were sufficient to switch the binding preference of the virus from avian receptors to human receptors.

We analyzed the adhesion measurements by Xiong et al. (19) using Eq. 1. To validate our model, we fit the two parameters,  $K_0$  and  $V_{\text{eff}}$ , to three of the measured adhesion isotherms and used the parameters to predict the remaining three measured isotherms. These isotherms relate the amount of attached viral particles and the relative sugar loading (i.e., the amount of SA as a fraction of the amount at maximum loading) on the surface. The fractional saturation of the biosensor by the virus is given by

$$f = \frac{1}{1 + K_{D,\text{ad}}/[\text{virus}]}, \quad (20)$$

where  $[\text{virus}] = 10^{-7}$  mM is the solution concentration of the virus used in the experiment, and  $K_{D,\text{ad}}$  varies according to Eq. 1 with respect to the relative sugar loading. The values of the fitted parameters are summarized in Table 1, and the resulting adhesion isotherms are shown in Fig. 3 A. To compare our model to the empirical model  $K_{D,\text{ad}} = (K_D)^m$ , which is used by Xiong et al. (19) in their analysis of the data, we computed the multiplicity  $m = \ln K_{D,\text{ad}}/\ln K_D$  at different relative sugar loading values for each HA-SA pair, using the  $K_{D,\text{ad}}$  value estimated from the experiment (by Eq. 19) and from our model (by Eq. 1). The results are shown in Fig. 3 B. As expected, the multiplicity varies with both the relative sugar loading and the individual  $K_D$ ; this variation is reasonably well captured by our model.

The  $K_D$  values of HA-SA binding, used in the fitting and prediction, are taken from the MST measurements by Xiong et al. (19) and are listed in Table 2. To further assess the quality of our model, we took the  $K_0$  and  $V_{\text{eff}}$  values from the fitting and refit the dissociation constant,  $K_D$ , of each individual HA-SA binding to the corresponding adhesion

**TABLE 1** Values of the model parameters for influenza virus-host cell adhesion

$N_{\text{HA}}$	$N_{\text{SA}}$	$K_0$ (mM)	$V_{\text{eff}}$ (mM <sup>-1</sup> )
$3 \times 50$	2000	$1.0 \pm 0.5 \times 10^{-5}$	$2.7 \pm 0.3 \times 10^3$

The parameters  $K_0$  and  $V_{\text{eff}}$  are fitted to three experimentally measured isotherms.  $N_{\text{HA}}$  is the number of HA monomers, and  $N_{\text{SA}}$  is the number of SA receptors at maximum sugar loading. The number of ligands in our model is then  $N_L = N_{\text{HA}}$ , and the number of receptors  $N_R = N_{\text{SA}} \times$  relative sugar loading. In principle,  $N_{\text{HA}}$  and  $N_{\text{SA}}$  can be treated as fitting parameters as well. The fitting problem then becomes underrestrained, however, given the small amount of isotherm data, and there are large statistical uncertainties in the fitted  $N_{\text{HA}}$ ,  $N_{\text{SA}}$ , and  $K_0$  parameters. Instead, we used rough estimates of  $N_{\text{HA}}$  and  $N_{\text{SA}}$  from previous publications. It has been estimated (19) that 50 HA trimers per virus participate in forming adhesive connections, and thus  $N_{\text{HA}} \approx 3 \times 50$ . The estimate for  $N_{\text{SA}}$  ranges from 600 to 4500 (19,32,39). We used an intermediate value. Different values of  $N_{\text{SA}}$  used in fitting result in compensatory values of  $V_{\text{eff}}$ , but the consequent isotherms are little changed. More accurate measurements of  $N_{\text{HA}}$  and  $N_{\text{SA}}$  can help determine more reliably the parameters  $K_0$  and  $V_{\text{eff}}$ . We have also fit the parameters  $K_0$  and  $V_{\text{eff}}$  using the exact solution (Eqs. 3 and 4), obtaining the same values.

isotherm (Fig. 3 C). The  $K_D$  values obtained from the adhesion isotherms are in good agreement with the values determined from MST (Table 2). This suggests that our model may be used to estimate the binding affinities from adhesion experiments. Such an approach may be advantageous when the adhesion affinity is easier to measure than the receptor-ligand binding affinity (e.g., when the receptor-ligand binding is too weak to be detected by the binding assays).

We can check that the fitted value of  $K_0$  is reasonable by using Eq. 14. Assuming that 10% of the virus surface is in contact with the cell surface (19) (i.e.,  $N_{\text{HA}}/M_{\text{HA}} \approx 0.1$ ), we can estimate the fraction of the SA molecules per adhesion site that lie within contact interface to be  $N_{\text{SA}}/M_{\text{SA}} = K_0 V_{\text{eff}} M_{\text{HA}}/N_{\text{HA}} \approx 0.3$ , a reasonable estimate for the SA-loaded surfaces considered in this work, which are relatively flat compared to the virus—more of the available SA molecules per site of adhesion on the cell thus lie within the interface.

#### Adhesion to cells

Our model can be used to analyze the implications of multivalent adhesion in the biological function of influenza virus and the competitive inhibition of influenza adhesion needed to effectively reduce influenza infection. Unlike biosensors, which are conjugated with a pure form of SA-terminated oligosaccharide, a real vertebrate cell presents a diverse set of SA-terminated glycans with different affinities for HA molecules. To characterize the adhesion of influenza virus to real cells, we ideally need the distribution of binding affinities,  $N(K)$ , of the surface glycans to the HA molecules as input for the model. We also ideally need to fit the values of  $K_0$  and  $V_{\text{eff}}$  to a series of measured adhesion affinities between the cells and the influenza virus. The necessary experimental data are not available, however, and so we make

some reasonable assumptions and approximations to enable us to apply our model to cell-influenza adhesion. Although some quantitative predictions of our model will improve as new experimental data yield more realistic parameters, some of the predictions will remain essentially unchanged with respect to changes in experimental inputs.

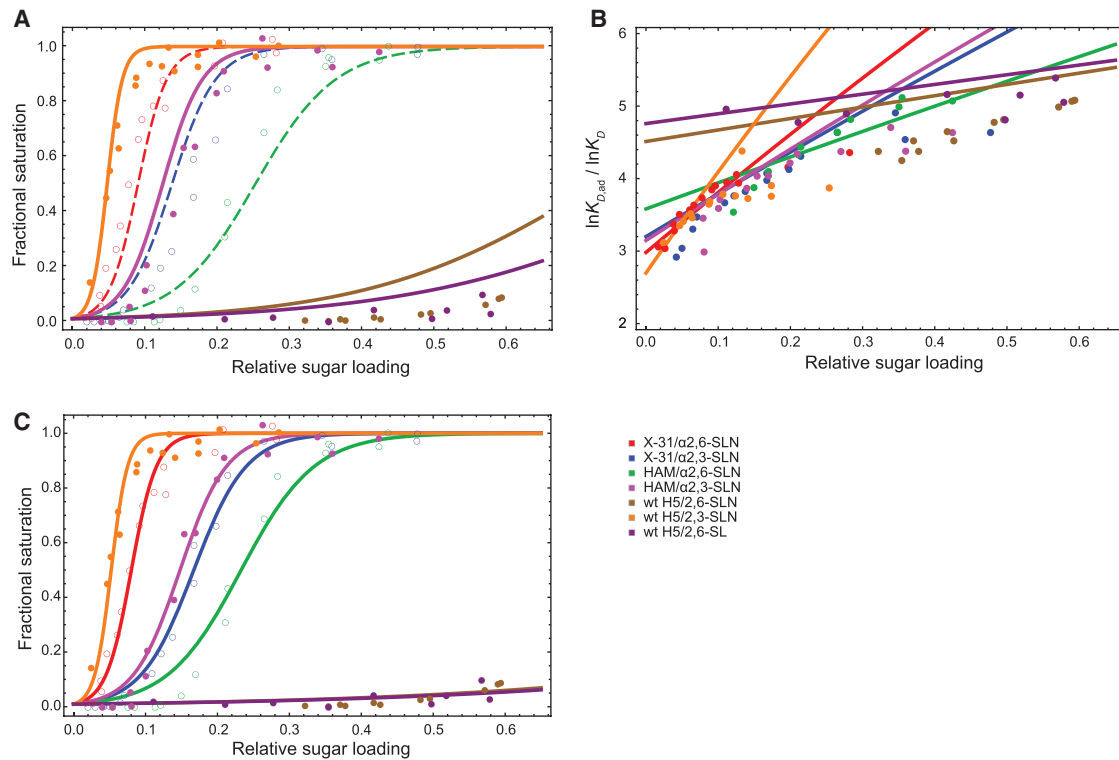
In the following examples, our model is applied to determine the condition at which adhesion is dramatically altered (e.g., the value of  $K_D$  above which the adhesion is effectively abolished). Here, we discuss why such a condition is insensitive to the value of  $K_0$ , which may be different between adhesion to cells and adhesion to biosensors because of the difference in the nonspecific interactions of the different adhesion surfaces. According to Eq. 19, the population of bound virus changes rapidly around  $K_{D,\text{ad}} \approx [\text{cell}]$ . The predicted parameter  $\theta$ —which may be  $K_D$ ,  $N_R$ ,  $N_L$ , or the concentration of hemagglutinin inhibitor (see below)—corresponding to this condition thus changes with  $K_0$  as  $d\theta/d\ln K_0 = -(d\ln K_{D,\text{ad}}/d\theta)^{-1}$ . In our model, because  $d\ln K_{D,\text{ad}}/d\theta$  is large (e.g., a small change in  $K_D$  can lead to orders-of-magnitude change in  $K_{D,\text{ad}}$ , as discussed below),  $d\theta/d\ln K_0$  is small. This insensitivity to the value of  $K_0$  will be demonstrated in the later discussion on the viable range of HA-SA binding. We will thus continue to use its value determined above from the biosensor experiments.

The parameter  $V_{\text{eff}}$  appears as a scaling factor of  $K_D$ , adjusting the affinity for the confined geometry in which receptor binding takes place. As such, the predictions of our model would remain valid, and the adhesion curves would look identical, for a different value of  $V_{\text{eff}}$ , if we simply replace the binding affinity  $K_D$  with  $K'_D = K_D V_{\text{eff}}/V_{\text{eff}}$ . The value of  $V_{\text{eff}}$  in Table 1 will be used in the subsequent discussions, but it will be straightforward to update the conclusions should new experimental measurements on cell-to-influenza virus adhesion permit the determination of a more accurate value of  $V_{\text{eff}}$  for cells.

Experimental data for the distribution of binding affinity  $N(K)$  is not yet available. We will first apply our model using the assumption that all SA molecules on a cell bind to the HA molecule with the same binding affinity; later, we examine the sensitivity of the results to changes in the distribution of binding affinity.

Multivalency amplifies the weak binding between HA and SA to achieve strong affinity between the influenza virus and the host cell. Assuming that the host cells are closely packed against one another, with an individual diameter of 20  $\mu\text{m}$ , the cell density is then approximately  $[\text{cell}] = 4 \times 10^{-10}$  mM. We can estimate the fraction of virus particles that are bound to the host cells using Eq. 19. The red curve in Fig. 4 shows how the fraction of virus particles bound to host cells changes with HA-SA binding affinity. Even when the affinity is as weak as  $K_D = K_{D,\text{max}} \approx 10$  mM ( $K_{D,\text{max}}$  is determined by solving for  $K_D$  in Eq. 19 with  $f = 0.5$ ), the majority of virus particles are still bound to host cells. The fraction of bound virus particles increases steeply





**FIGURE 3** The agreement between our model and experimental measurements of adhesion. (A) Isotherms of adhesion between influenza virus and biosensors loaded with SA. (Lines) Isotherms predicted by our model; (circles) experimental results. Different colors correspond to different combinations of viral strains and  $\alpha$ 2,3-SA or  $\alpha$ 2,6-SA. (Dashed lines and open circles) The three isotherms used for fitting the model parameters; (solid lines) predictions of isotherms for which experimental measurements are available. (B) Multiplicity of adhesion,  $m = \ln K_{D,ad} / \ln K_D$ , as a function of the relative sugar loading for different HA-SA pairs. The variation of  $m$  across different relative sugar-loading values is reasonably well captured by our model. (C) Computed and measured isotherms of adhesion, with  $K_D$  values refitted to each individual isotherm, using the parameter values of  $K_0$  and  $V_{eff}$  in Table 1. The fractional saturation predicted by our model is in good agreement with the experimental measurements, with fitted  $K_D$  values close to those measured by MST. X-31, HAM, and wild-type (wt) H5 are different influenza strains, each with a different HA sequence and thus different binding affinities for sialic acids (19) (Table 2).  $\alpha$ 2,3-SLN ( $\alpha$ 2,3-linked sialyl lactosamine),  $\alpha$ 2,6-SLN ( $\alpha$ 2,6-linked sialyl lactosamine), and  $\alpha$ 2,6-SL ( $\alpha$ 2,6-linked sialyl lactose) are different terminal sialic acid moieties.

(by >100-fold) as  $K_D$  decreases from 20 to 5 mM, allowing the virus to gain affinity to and thus to infect new host species by small changes in its HA sequence.

A nascent influenza virus emerging from a host cell must detach itself from that cell to spread the infection to other cells. To overcome the strong multivalent adhesion, the virus uses neuraminidase to cleave the SA moieties at the adhesion interface (26,27), reducing the number of SA mol-

**TABLE 2** The dissociation constants,  $K_D$ , of HA-SA binding for various viral strains and SA

Viral Strain	SA	$K_D$ (mM) (MST)	$K_D$ (mM) (Refitted)
X-31	$\alpha$ 2,6-SLN	$2.1 \pm 0.3$	$1.85 \pm 0.04$
X-31	$\alpha$ 2,3-SLN	$3.2 \pm 0.6$	$3.88 \pm 0.07$
HAM	$\alpha$ 2,6-SLN	$5.9 \pm 0.7$	$5.4 \pm 0.2$
HAM	$\alpha$ 2,3-SLN	$2.9 \pm 0.3$	$3.44 \pm 0.08$
wt H5	$\alpha$ 2,6-SLN	$17 \pm 3$	$34 \pm 4$
wt H5	$\alpha$ 2,3-SLN	$1.1 \pm 0.2$	$1.20 \pm 0.03$
wt H5	$\alpha$ 2,6-SL	$21 \pm 6$	$37 \pm 6$

The last column lists the  $K_D$  values estimated by refitting them to the virus adhesion isotherms, using the  $K_0$  and  $V_{eff}$  parameters shown in Table 1.

ecules available and thus the affinity of adhesion. The stronger the binding between HA and SA, the more SA neuraminidase needs to cleave for effective release of the virus. We define a quantity  $R_{50}$ , which is the factor by which the population of SA molecules on the surface needs to be reduced for half of the virus particles to be unbound. The blue curve in Fig. 4 plots  $R_{50}$  against  $K_D$ . At  $K_D = K_{D,min} \approx 0.08$  mM,  $R_{50} = 0.01$  ( $K_{D,min}$  is determined by solving for  $K_D$  in Eq. 19 with  $f = 0.5$  and  $N_R = R_{50} N_{SA} = 0.01 N_{SA}$ ), which implies that neuraminidase will need to remove 99% of all SA moieties to release half of the virus particles from the host cell. The enzymatic efficiency of neuraminidase thus imposes a lower limit of viable  $K_D$  in HA-SA binding. On the other hand, the virus can compensate for inefficient neuraminidase, and develop resistance to neuraminidase inhibitors (40), with mutations in HA that decrease HA-SA affinity (41,42).

Our model suggests that  $K_D$  falls within a range that satisfies the dual requirements that the virus be able to attach to host cells and, with the aid of neuraminidase, to detach from them. Values of the parameters—namely the percentage of

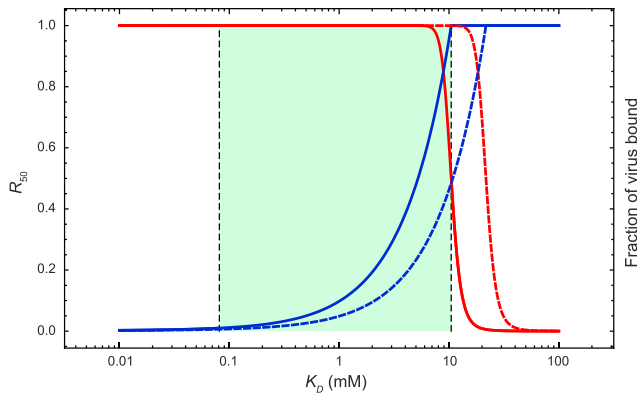


FIGURE 4 Dependence of virus adhesion on HA-SA binding affinity. (Red curve) Fraction of virus particles bound to host cells versus  $K_D$  of HA-SA binding. (Dark-blue curve)  $R_{50}$  (see text) versus  $K_D$ . (Light-green region) Interval of viable  $K_D$  values such that the virus can 1) attach sufficiently to and infect cells and 2) release itself, with the help of neuraminidase, from the cells to spread the infection. The upper bound of the region corresponds to the  $K_D$  value at which half of the viruses attach to cells; the lower bound corresponds to the  $K_D$  value at which 99% of all SA moieties must be removed in order for half of the viruses to detach from the cells. (Solid curves) Results from Eqs. 1 and 2, which assume that all SA receptors on the cell bind to the HA molecule with the same affinity  $K_D$ ; (dashed curves) results from Eqs. 16–18, which assume that the binding free energy is uniformly distributed between  $RT\ln(K_D/10)$  and  $RT\ln(K_D10)$ .

SA moieties neuraminidase can remove before the destruction of the immobilized virus on the cell surface and the fraction of nascent viral particles that need to detach for successful establishment of influenza within a host—are unfortunately not known, but using the speculative values discussed above yields a permissible range of  $K_D$  from our model of  $K_{D,\min} = 0.08$  to  $K_{D,\max} = 10$  mM. This estimated range is in line with what has been measured experimentally (12,19). The specific range estimate can vary depending on the assumed values of the parameters in our model. This dependence on the parameters—including the discrete  $N_R$  and  $N_L$ —is easy to analyze: differentiating Eq. 19 using the chain rule, then substituting the numerical values of the current parameters, we obtain

$$d \ln k_{D,\min} = 1.0 d \ln N_{\text{HA}} + 1.2 d \ln N_{\text{SA}} - d \ln V_{\text{eff}} - 0.12 d \ln K_0 + 0.12 d \ln [\text{cell}] \quad (21)$$

and

$$d \ln k_{D,\max} = 1.0 d \ln N_{\text{HA}} + 1.0 d \ln N_{\text{SA}} - d \ln V_{\text{eff}} - 0.10 d \ln K_0 + 0.10 d \ln [\text{cell}], \quad (22)$$

which suggest that our estimated range of viable  $K_D$  has relative errors comparable to those in the estimates of  $N_{\text{HA}}$  and  $N_{\text{SA}}$ , but is insensitive to the parameters  $K_0$  and  $[\text{cell}]$ . Increasing  $K_0$  by a factor of 1000, for example, changes the estimated viable range by only a factor of  $\sim 2$ , to  $K_{D,\min} = 0.04$  mM and  $K_{D,\max} = 6$  mM. Although the predicted

values of  $K_{D,\min}$  and  $K_{D,\max}$  depend on  $V_{\text{eff}}$ , the width of the viable range (i.e.,  $K_{D,\max}/K_{D,\min}$ ) does not.

Different cells will have different SA densities on their surfaces, and thus different numbers of receptors at the interface. The fraction of bound virus particles, however, changes steeply over a narrow range of  $N_R$ , and remains essentially constant at  $\sim 0$  for  $N_R$  below this range and at  $\sim 1$  for  $N_R$  above this range (Fig. 2 A). This allows us to divide the cells into two populations, one with an  $N_R$  above the midpoint of the transition range and the other with an  $N_R$  below it, and to consider only the adhesion of virus to the first population. The adhesion can thus be analyzed using a cell concentration adjusted by the fraction of the first population. As shown above in Eqs. 21 and 22, the predictions of our model are insensitive to the value of  $[\text{cell}]$ , suggesting that the predictions of our model will remain mostly valid despite the variability in SA densities among different cells. We present in the Supporting Material a detailed analysis of viral adhesion to a population of cells with varying numbers of interface receptors;  $K_{D,\max}$  obtained from this analysis is essentially the same as that calculated for adhesion to cells with a uniform number of interface receptors.

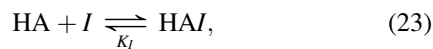
The above analysis assumes that all SA molecules on the cell surface bind to the HA molecule with the same binding affinity  $K_D$ . On the surface of a real cell, SA moieties in different glycans may bind to HA with different affinities  $K_D$ . We now discuss the effect of such heterogeneity.

We examine a simple, concrete example, in which the HA-SA binding affinity spans two orders of magnitude from  $K_{\min}$  to  $K_{\max} = 100 K_{\min}$ , the binding free energy is uniformly distributed from  $RT\ln K_{\min}$  to  $RT\ln K_{\max}$ , and the total number of SA molecules is  $N_{\text{SA}}$ . The distribution is then  $N(K) = N_{\text{SA}}(\ln(K_{\max}/K_{\min}))^{-1} K^{-1}$ . Defining the mean affinity  $K'_D \equiv \sqrt{K_{\min}K_{\max}}$ , the value of  $K'_D$  then determines the distribution  $N(K)$  (e.g.,  $K'_D = 1$  mM signifies that the binding affinity ranges from  $K_{\min} = 0.1$  mM to  $K_{\max} = 10$  mM, with the binding free energy uniformly distributed in the corresponding range). We can compute the fraction of bound viral particles and the  $R_{50}$  parameter at different values of  $K'_D$  (the dashed blue and red curves, respectively, in Fig. 4). Compared to the case of homogeneous receptors, the viable range of the mean affinity is shifted by a factor of  $\sim 2$ , but the width of the range is similar ( $K'_{D,\max}/K'_{D,\min} = 137$  versus  $K_{D,\max}/K_{D,\min} = 129$ ). It will be straightforward to incorporate more realistic estimates of  $N(K)$  in our model as innovative experimental techniques, such as glycan microarrays (43), more comprehensively characterize the HA-binding properties of diverse glycans on real cells.

HA of avian influenza preferentially binds SA of  $\alpha 2,3$  linkage, whereas HA of human influenza preferentially binds SA of  $\alpha 2,6$  linkage. Due to multivalency in the adhesion, a small difference in the binding affinity can be amplified to a big difference in biological adhesion. The X-31 HA (an extensively studied HA construct), for instance, binds

$\alpha 2,6$  SA with  $K_D = 2.1 \pm 0.3$  mM, and it binds  $\alpha 2,3$  SA with  $K_D = 3.2 \pm 0.6$  mM (13). The ability of the whole virus to adhere to cells, in contrast, can differ drastically between cells with surface SA of  $\alpha 2,3$  linkage and those of  $\alpha 2,6$  linkage: Assuming the parameters in Table 1, the ratio of the adhesion constants is  $K_{D,ad}^{(\alpha 2,3)} / K_{D,ad}^{(\alpha 2,6)} = 10^{6 \pm 3}$ . As another example, the wild-type HA of the avian strain A/Vietnam/1194/2004 binds  $\alpha 2,6$  SA with  $K_D = 17 \pm 3$  mM (19). A mutant HA was identified in a strain derived from A/Vietnam/1194/2004 that is transmissible in ferrets. This mutant HA, which contains a handful of point mutations, binds  $\alpha 2,6$  SA with  $K_D = 12 \pm 2.5$  mM (19). According to our model, this 1.4-fold increase in the binding affinity will lead to a much larger increase in the adhesion affinity, as  $K_{D,ad}^{(wild-type)} / K_{D,ad}^{(mutant)} = 10^{1.1 \pm 0.9}$ , and thus potentially a much higher fraction of virus particles bound to the mammalian cells. Our model thus quantitatively predicts that small changes in the HA-SA binding affinities might be sufficient to switch the binding preference of the virus from avian cells to human cells, and that a small gain in the hemagglutinin's affinity for human receptors may sufficiently increase the virus's adhesion to human cells to enable human infection.

There has been a longstanding interest in developing HA inhibitors as potential therapeutic agents against influenza (31,44,45). A competitive inhibitor  $I$  of HA forms complexes  $HAI$  in the reaction



where  $K_I = [HA][I]/[HAI]$  is the equilibrium dissociation constant, and reduces the number of available HA by the number of complexes formed,  $N_{HAI}$ . The presence of the inhibitor can be included in our model through the following set of simultaneous equations for  $N_{HAI}$  and  $\bar{n}$ :

$$\frac{(N_{HA} - N_{HAI} - \bar{n})(N_{SA} - \bar{n})}{\bar{n}} = K_D V_{eff} \quad (24)$$

$$\frac{(N_{HA} - N_{HAI} - \bar{n})}{N_{HAI}} = K_I / [I]$$

and

$$K_{D,ad} = K_0 \exp \left( N_{HA} \ln \frac{N_{HA} - N_{HAI} - \bar{n}}{N_{HA} / \left( 1 + \frac{[I]}{K_I} \right)} + N_{SA} \ln \frac{N_{SA} - \bar{n}}{N_{SA}} + \bar{n} \right). \quad (25)$$

Solving for  $N_{HAI}$  in terms of  $N_{HA} - \bar{n}$  and  $K_I / [I]$  from the second line of Eq. 24, and substituting it into Eq. 25, we

can show that  $K_{D,ad}$  has the same expression as in Eq. 1, but with the value of  $\bar{n}$  determined by Eq. 24.

The above equations can be derived from the partition function of the system

$$Z = \frac{1}{N_{HA}! N_{SA}!} \sum_{m=0}^{N_{HA}} \binom{N_{HA}}{m} \exp \left( -\frac{m \Delta \mu}{RT} \right) \times \sum_{n=0}^{\min(N_{HA}-m, N_{SA})} \binom{N_{HA}-m}{n} \binom{N_{SA}}{n} \exp \left( -\frac{n \Delta G}{RT} \right), \quad (26)$$

where  $\exp(-(\Delta \mu / RT)) = [I] / K_I$  is the excess chemical potential of the HA-inhibitor complex. In the above, we assume that the inhibitor is in excess such that the free inhibitor concentration  $[I]$  does not depend on the number of HA-inhibitor complexes formed. The probability of spontaneously breaking all HA-SA connections (i.e.,  $n = 0$ ) is then

$$p_0 = Z^{-1} \frac{1}{N_{HA}! N_{SA}!} \sum_{m=0}^{N_{HA}} \binom{N_{HA}}{m} \exp \left( -\frac{m \Delta \mu}{RT} \right) \quad (27)$$

$$= Z^{-1} \frac{1}{N_{HA}! N_{SA}!} (1 + [I] / K_I)^{N_{HA}}.$$

By replacing  $Z$  with the largest term in the summand of Eq. 26, which corresponds to  $m = N_{HAI}$  and  $n = \bar{n}$  in Eq. 24, and applying the Stirling approximation, we arrive at Eq. 25. The approximation is quite accurate: the ratio of the approximate value and the exact value of  $p_0$  stays close to 1 across a large range of inhibitor concentrations (Fig. 1 B).

Inhibition of the multivalent adhesion depends on inhibitor concentration in a very different manner than inhibition of monovalent binding, as shown in Fig. 5, in which the bound fraction of virus particles, normalized by the maximum bound fraction (at  $[I] = 0$ ), is plotted against the inhibitor concentration  $[I] / K_I$ . The normalized fraction of HA bound to SA at different inhibitor concentrations, which is given by  $f_{bound} / f_{max} = 1 / (1 + [I] / K_I)$ , if the inhibitor binds to HA with much higher affinity than SA (i.e.,  $K_I \ll K_D$ ), is also shown to contrast the difference between inhibition of monovalent binding and of multivalent adhesion. Whereas the normalized fraction of HA bound with SA decreases gradually with increasing  $[I] / K_I$ , reaching half at  $[I] = K_I$ , the normalized fraction of bound virus drops precipitously at different  $[I] / K_I$  values, depending on the affinity between HA and SA. In viral strains with high HA-SA affinity  $K_D = 0.5$  mM, adhesion is half-inhibited only when  $[I] / K_I = 20$ . On the other hand, once  $[I] / K_I$  exceeds that value, the normalized fraction steeply decreases, reaching  $f_{bound} / f_{max} = 0.001$  at  $[I] / K_I = 66$ . The inhibitor concentrations  $[I] / K_I$  required to reduce the virus adhesion to given  $f_{bound} / f_{max}$  values are computed for different HA-SA affinities and shown in Fig. 6.

Our analysis above suggests that inhibitors of multivalent adhesion should be assessed differently than inhibitors of

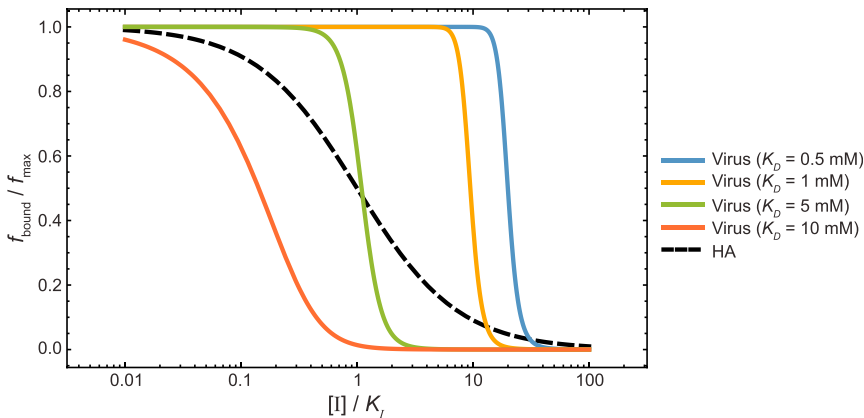


FIGURE 5 Inhibition of multivalent adhesion and of monovalent binding. (Solid lines) The normalized fractions of bound viruses at different inhibitor concentrations for viruses of different HA-SA binding affinities. (Dashed line) Fraction of HA bound to SA at different inhibitor concentrations.

monovalent binding. Consider, for example, the challenges facing the search for effective HA inhibitors: the receptor-binding pocket of HA is very shallow, making it difficult for any small molecule to achieve high-affinity (i.e., nanomolar) binding. It has been hypothesized that a strong HA inhibitor is required to counter the strong adhesion due to multivalency, yet to our knowledge no quantitative estimate has been made as to the inhibitor potency required for effective inhibition of adhesion. According to our model, a relatively weak HA inhibitor with good bioavailability can effectively inhibit virus attachment. At a typical plasma concentration of  $[I] = 10 \mu\text{M}$ , for example, an inhibitor of submicromolar affinity ( $K_I = 0.1 \mu\text{M}$ ) is predicted to reduce the virus binding by  $\sim 3000$ -fold, even for viral strains with HA of high SA-binding affinities ( $K_D \sim 0.5 \text{ mM}$ ).

### The role of binding cooperativity in multivalent adhesion

Equation 1 is predicated on the constancy of the affinity of individual receptor-ligand binding throughout the adhesion

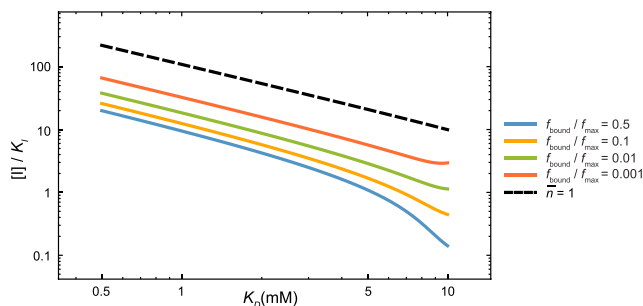


FIGURE 6 Inhibitor concentrations required to reduce the virus adhesion to given normalized fractions for different HA-SA affinities. Even when the affinity of receptor-ligand binding is high, a moderate inhibitor concentration relative to the inhibition constant is sufficient to drastically reduce viral adhesion. (Dashed line) Inhibitor concentration required to reduce the equilibrium number of connections,  $\bar{n}$ , to  $< 1$ . The inhibitor concentration required for  $\bar{n} \leq 1$  is greater than that required to reduce the fraction of bound virus by 1000-fold, which suggests that adhesion can be effectively inhibited before  $\bar{n} \approx 0$ .

interface. The binding affinity of a receptor-ligand pair, however, may vary with the position of the receptor-ligand pair within the interface, and it may also be affected by existing receptor-ligand connections. The latter dependence is often referred to as cooperativity. Here we discuss the implications of cooperativity and position-dependent affinity in multivalent adhesion.

In general, the free energy change of forming  $n$  labeled and distinguishable receptor-ligand connections can be written as

$$\exp\left(-\frac{G(n)}{RT}\right) = \int \prod_{i=1}^n dr_i \exp(-g(r_1, r_2, \dots, r_n)/RT), \quad (28)$$

where  $r_i$  is the position of the  $i$ th receptor-ligand pair, and  $g(r_1, r_2, \dots, r_n)$  is the free energy change of forming  $n$  connections in the corresponding positions. (We note that because Eq. 28 refers to distinguishable receptor-ligand connections, there is no factor of  $1/n!$  on the right-hand side.) Without better experimental and theoretical characterization of the microscopic details of the adhesion interface, the function  $G(n)$  is not known for realistic biological adhesion interfaces, and it is necessary to make some approximations in our analysis.

The study of adhesion is most often concerned with the conditions under which adhesion is established or abolished, corresponding to conditions where the bound fraction is  $f \approx 0.5$ . We will show that the effect of cooperativity on such conditions diminishes as the numbers of receptors and ligands at the adhesion interface become large. Under such conditions, the probability,  $p_0$ , of spontaneously breaking all receptor-ligand connections cannot be vanishingly small, and, according to the above analysis, the equilibrium number of connections  $\bar{n}$  must be small (Fig. 2). If  $N_R$  and  $N_L$  are large,  $\bar{n} \ll \min(N_R, N_L)$ , the small number of formed receptor-ligand pairs will almost never be close to each other, and thus the binding affinity of one pair will not be affected by other pairs. In such cases, we can ignore the multibody terms in  $g$ , and approximate it as

$$g(r_1, r_2, \dots, r_n) \approx \sum_{i=1}^n g_1(r_i) \quad (29)$$

and

$$G(n) \approx -nRT \ln \int dr \exp(-g_1(r)/RT) \equiv n\Delta G, \quad (30)$$

where  $\Delta G = -RT \ln(\int dr \exp(-g_1(r)/RT))$  is a constant. The details of  $g_1(r)$  determine  $V_{\text{eff}}$ , the value of which is obtained by parameter fitting in our model. Equation 1, and thus all the quantitative results of our model, remain unchanged when the receptor-ligand affinity varies with the position in the adhesion interface.

To quantify the effect of cooperativity, we consider a special and commonplace case, which is pertinent to the study of influenza-cell adhesion: cooperativity within oligomeric receptors or ligands. The effect of such cooperativity on the equilibrium number of receptor-ligand connections at the cell-cell adhesion interface has been previously investigated (18); here we study its effect on adhesion affinity. HA is a homotrimer. Each monomer can individually bind an SA molecule, and its affinity may depend on the number of SA molecules already bound to the same HA trimer. If the binding of  $k = 1, 2, 3$  SA molecules to the same HA trimer corresponds to the free energy change of  $\Delta G_k$ , the probability of breaking all connections is given by

$$\ln p_0 = N_{\text{HA}_3} \ln \frac{N_{\text{HA}_3} - \bar{n}_1 - \bar{n}_2 - \bar{n}_3}{N_{\text{HA}_3}} + N_{\text{SA}} \ln \frac{N_{\text{SA}} - \bar{n}}{N_{\text{SA}}} + \bar{n}, \quad (31)$$

where  $\bar{n}_k$  is the equilibrium number of HA trimers bound to  $k = 1, 2, 3$  SA molecules,  $\bar{n} = \bar{n}_1 + 2\bar{n}_2 + 3\bar{n}_3$  is the equilibrium number of SA molecules bound to HA, and  $N_{\text{HA}_3} = N_{\text{HA}}/3$  is the total number of HA trimers. The values of  $\bar{n}_{k=1,2,3}$  are given by the following equilibrium condition:

$$\begin{aligned} \frac{3(N_{\text{HA}_3} - \bar{n}_1 - \bar{n}_2 - \bar{n}_3)(N_{\text{SA}} - \bar{n})}{\bar{n}_1} &= \exp\left(\frac{\Delta G_1}{RT}\right) \\ \frac{3(N_{\text{HA}_3} - \bar{n}_1 - \bar{n}_2 - \bar{n}_3)(N_{\text{SA}} - \bar{n})^2}{\bar{n}_2} &= \exp\left(\frac{\Delta G_2}{RT}\right) \\ \frac{(N_{\text{HA}_3} - \bar{n}_1 - \bar{n}_2 - \bar{n}_3)(N_{\text{SA}} - \bar{n})^3}{\bar{n}_3} &= \exp\left(\frac{\Delta G_3}{RT}\right) \\ \bar{n} &= \bar{n}_1 + 2\bar{n}_2 + 3\bar{n}_3. \end{aligned} \quad (32)$$

Equations 31 and 32, derived in the [Supporting Material](#), account for the cooperativity in the binding of SA molecules to different monomers within a HA trimer. In the absence of cooperativity,  $\Delta G_3/3 = \Delta G_2/2 = \Delta G_1 = \Delta G$ , and Eqs. 31 and 32 can be shown to reduce to Eq. 1.

We introduced into our model the cooperativity between SA molecules binding to the same HA trimer by the coefficient  $\gamma$ , such that the binding affinity of SA to the first

HA monomer is  $K_D V_{\text{eff}}$ , but the subsequent binding affinities to the remaining two HA monomers in the same trimer are  $K_D V_{\text{eff}} \gamma$  (i.e.,  $\exp(\Delta G_1/RT) = K_D V_{\text{eff}}$ ,  $\exp((\Delta G_2 - \Delta G_1)/RT) = \exp((\Delta G_3 - \Delta G_2)/RT) = K_D V_{\text{eff}} \gamma$ ). Attempting to derive the parameter  $\gamma$  by fitting the model to the experimental isotherms yields large statistical uncertainty on its value (the standard error for  $\ln \gamma$  is 600), suggesting that these experimental isotherms do not provide sufficient information to infer the presence or the extent of the cooperativity embodied by our model. The fitted model (parameters given in the caption of [Fig. 7 A](#)) exhibits almost identical isotherms, as does the model without cooperativity ([Fig. 7 A](#)).

To explore the extent to which cooperativity can affect the adhesion isotherms, we computed  $R_{50}$  as a function of  $\gamma$ , for various  $N_{\text{HA}_3}$  and  $K_D$  values, holding the other parameters in our model to the corresponding values in [Table 1](#) ([Fig. 7 B](#)). As expected, the effect of cooperativity diminishes as the number of HA trimers increases; when  $N_{\text{HA}_3} = 1000$ , the effect of cooperativity is negligible. At the assumed value  $N_{\text{HA}_3} = 50$  for the influenza virus, cooperativity has a noticeable effect:  $R_{50} = 0.030$  for  $\gamma = 0.01$ ,  $R_{50} = 0.069$  for  $\gamma = 0.1$ , and  $R_{50} = 0.098$  for  $\gamma = 1$  ( $K_D = 1$  mM). The adhesion isotherms of the various viral strains are plotted for  $\gamma = 0.1$  as dash-and-dot lines in [Fig. 7 A](#).

Cooperativity within the trimer can be similarly considered in the inhibition of adhesion. Denoting  $\bar{m}_i$  as the equilibrium number of HA trimers that have  $i = 0, 1, 2, 3$  inhibitors bound, and  $\bar{n}_{ij}$  as the equilibrium number of HA trimers that have  $i$  inhibitors and  $j$  SA molecules bound, the probability of breaking all HA-SA connections is

$$\ln p_0 \approx N_{\text{HA}_3} \ln \frac{\bar{m}_0 - \bar{n}_{01} - \bar{n}_{02} - \bar{n}_{03}}{N_{\text{HA}_3}} + N_{\text{SA}} \ln \frac{N_{\text{SA}} - \bar{n}}{N_{\text{SA}}} + \bar{n} \quad (33)$$

with the following equilibrium condition:

$$\binom{3}{i} \frac{\bar{m}_0 - \sum_{j=1}^3 \bar{n}_{0j}}{\bar{m}_i - \sum_{j=1}^{3-i} \bar{n}_{ij}} = \left(\frac{K_I}{[I]}\right)^i, \quad i = 1, 2, 3$$

$$\binom{3-i}{j} \frac{\bar{m}_i - \sum_{k=1}^{3-i} \bar{n}_{ik}}{\bar{n}_{ij}} = \frac{\exp\left(\frac{\Delta G_j}{RT}\right)}{(N_{\text{SA}} - \bar{n})^j},$$

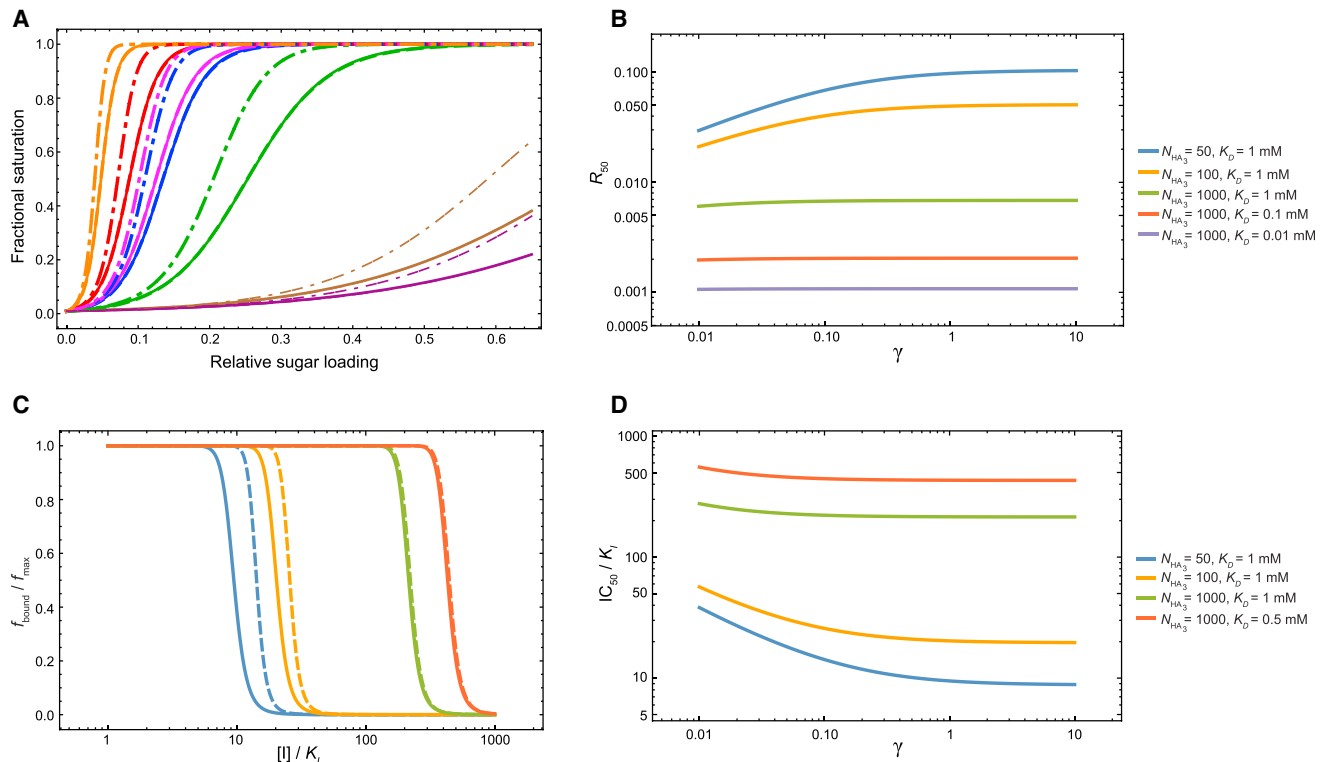
$$i = 0, 1, 2, 3; \quad j = 1, \dots, 3 - i$$

$$\bar{m}_0 + \bar{m}_1 + \bar{m}_2 + \bar{m}_3 = N_{\text{HA}_3}$$

$$\sum_{i=0}^3 \sum_{j=1}^{3-i} j \bar{n}_{ij} = \bar{n}.$$

(34)

The derivation of Eqs. 33 and 34 is outlined in the [Supporting Material](#).



**FIGURE 7** The effect of binding cooperativity within the same HA trimer on the isotherms and the inhibition of adhesion. (A) Adhesion isotherms for the viral strains considered in this work. (Solid lines) Model without cooperativity (i.e.,  $\gamma = 1$ ), with the fit parameters in Table 1. (Dashed lines) Model refit with cooperativity, which yielded a value of  $\gamma$  with large statistical uncertainty. The parameters from the refitting are  $V_{\text{eff}} = 2.6 \times 10^3 \text{ mM}^{-1}$ ,  $K_0 = 1.2 \times 10^{-5}$ , and  $\gamma = 72$ . (Dash-and-dot lines) Isotherms generated with the parameters in Table 1, but with  $\gamma = 0.1$ . (B) The dependence of  $R_{50}$  on the cooperativity parameter  $\gamma$ . (C) The effect of cooperativity on the inhibition of adhesion. (Solid lines) Inhibition of adhesion if there is no cooperativity (i.e.,  $\gamma = 1$ ); (dashed lines) inhibition of adhesion if the binding affinities of the second and third SA molecules are 10-fold stronger than that of the first SA molecule to the same HA trimer ( $\gamma = 0.1$ ). (D) The dependence of  $IC_{50}$  values ( $[I] / K_I$ ) at which the fraction of bound cells is half that in absence of the inhibitor) on cooperativity. The values of  $N_{SA}$ ,  $V_{\text{eff}}$ , and  $K_0$  in Table 1 are used for (B)–(D). In (C), the cell concentration is as estimated in the text ( $[c_{\text{cell}}] = 4 \times 10^{-10} \text{ mM}$ ). (B–D) Effect of cooperativity diminishes with the number of HA trimers.

The binding cooperativity shifts the inhibition curve to higher inhibitor concentrations: given the parameters in Table 1, with  $N_{HA_3} = N_{HA}/3 = 50$ , a typical HA-SA binding affinity of  $K_D = 1 \text{ mM}$ , and a cooperativity coefficient  $\gamma = 0.1$ , the inhibitor concentration necessary to reduce the adhesion by half is 1.5 times the concentration necessary if there is no cooperativity (Fig. 7 C). This shift, however, all but vanishes at large numbers of trimers (Fig. 7, C and D).

In all the discussions above, we have assumed that the HA inhibitor has the same inhibition constant against HA anywhere in the adhesion interface that it does against HA in the solution, which is likely a valid assumption for small-molecule HA inhibitors. If the inhibitor binding to HA is diminished or enhanced inside the adhesion interface, our model would underestimate or overestimate, respectively, the inhibitor concentration required to inhibit virus-cell adhesion.

## CONCLUSIONS

Multivalent adhesion at biological surfaces often involves a very large number of simultaneous receptor-ligand connec-

tions, commonly on the order of hundreds to thousands. The adhesion affinity is the result of both receptor-ligand binding and nonspecific interactions. We have proposed a model for the adhesion affinity that takes into account both contributions. Our model gives a simple expression for the adhesion affinity by taking advantage of approximations appropriate for large numbers of receptors and ligands. We have shown that these approximations are accurate for a wide range of parameters common in biological adhesion.

We have applied our model to the analysis of the adhesion between influenza virus and SA-conjugated biosensors, and between influenza virus and vertebrate cells. Compared to experimental measurements, our model has reasonable quantitative accuracy using just two fitting parameters. Our model quantifies the high amplification that results from the high valency characteristic of many examples of biological adhesion: the high valency not only amplifies the weak receptor-ligand binding to produce high-affinity adhesion but also turns gradual changes in the conditions (e.g., receptor-ligand affinity, the number of receptors or ligands, and the concentration of competitive inhibitors) into

abrupt transitions in adhesion. Analysis of the relationship between adhesion and receptor-ligand affinity in our model provides quantitative support for the notion that the range of observed HA-SA affinities is a natural consequence of the functional requirements for the influenza virus. When used to study the inhibition of adhesion by competitive inhibitors of HA, our model suggests that it might be feasible to inhibit adhesion with a relatively weak (0.1  $\mu\text{M}$  in dissociation constant) HA inhibitor. Our model may aid the analysis of experiments and the design of therapeutic strategies involving cell-cell and cell-virus adhesion in general.

## SUPPORTING MATERIAL

Supporting Material is available at [http://www.biophysj.org/biophysj/supplemental/S0006-3495\(15\)01122-4](http://www.biophysj.org/biophysj/supplemental/S0006-3495(15)01122-4).

## AUTHOR CONTRIBUTIONS

H.X. and D.E.S. designed research; H.X. performed research and analyzed results; and H.X. and D.E.S. wrote the article.

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# Supporting Material

## A Simple Model of Multivalent Adhesion and Its Application to Influenza Infection

### Derivation of Eqs. 16–18

We present here the outline of the derivation of Eqs. 16–18. These equations estimate the affinity of adhesion mediated by a range of different receptors, each with a different binding affinity to the ligand. Suppose that there are  $N_i$  receptors of type  $i = 1, 2, \dots, m$  with binding free energy  $\Delta G_i = RT \ln K_i' \equiv RT \ln(K_i V_{\text{eff}})$ . We can write down the grand canonical partition function of the system as

$$Z = \frac{1}{\prod_i N_i! N_L!} \sum_{\{n_i\}} \binom{N_i}{n_i} \binom{N_L}{n_1 n_2 \dots n_m} \prod_i n_i! \exp\left(-\frac{1}{RT} \sum_i n_i \Delta G_i\right). \quad (\text{S1})$$

Following the derivation of Eqs. 1 and 2 in the text, we obtain the equilibrium condition

$$\frac{(N_i - \bar{n}_i)(N_L - \sum_i \bar{n}_i)}{\bar{n}_i} \approx K_i V_{\text{eff}}, \quad i = 1, 2, \dots, \quad (\text{S2})$$

and the probability of breaking all connections (i.e.,  $n_i = 0 \forall i$ )

$$\begin{aligned} \ln p_0 &= \ln \frac{1}{\prod_i N_i! N_L!} - \ln Z \\ &\approx \sum_i N_i \ln \frac{N_i - \bar{n}_i}{N_i} + N_L \ln \frac{N_L - \sum_i \bar{n}_i}{N_L} + \sum_i \bar{n}_i. \end{aligned} \quad (\text{S3})$$

Defining  $\bar{n} \equiv \sum_i \bar{n}_i$ , and taking the continuous limit by replacing the summation with the integral, we obtain Eqs. 16 and 17.

From Eq. S2, we have

$$\bar{n}_i = N_i \left( 1 + \frac{K_i V_{\text{eff}}}{N_L - \bar{n}} \right)^{-1}, i = 1, 2, \dots, \quad (\text{S4})$$

Summing over  $i$ , and taking the continuous limit, we obtain Eq. 18 in the text.

## Derivation of Eqs. 31 and 32

We derive here Eqs. 31 and 32, which account for the cooperativity in the binding of SA molecules to monomers within a HA trimer. If the binding of  $k = 1, 2, 3$  SA molecules to the same HA trimer corresponds to the free energy change of  $\Delta G_k$ , the partition function of the system is

$$\begin{aligned} Z &= \frac{1}{3^{N_{\text{HA}_3}} N_{\text{HA}_3}! N_{\text{SA}}!} \sum_{n_1, n_2, n_3} \binom{N_{\text{HA}_3}}{n_1 \ n_2 \ n_3} \binom{N_{\text{SA}}}{n} n! \binom{3}{1}^{n_1} \binom{3}{2}^{n_2} \binom{3}{3}^{n_3} \cdot \\ &\quad \exp(-(n_1 \Delta G_1 + n_2 \Delta G_2 + n_3 \Delta G_3)/RT) \\ &= \sum_{n_1, n_2, n_3} \frac{3^{n_1 + n_2 - N_{\text{HA}_3}}}{(N_{\text{HA}_3} - n_1 - n_2 - n_3)! n_1! n_2! n_3! (N_{\text{SA}} - n)!} \exp(-(n_1 \Delta G_1 + n_2 \Delta G_2 + n_3 \Delta G_3)/RT), \quad (\text{S5}) \end{aligned}$$

where  $\binom{N_{\text{HA}_3}}{n_1 \ n_2 \ n_3}$  is the number of ways to choose  $n_k$  HA trimers to bind  $k = 1, 2, 3$  SA molecules,  $\binom{3}{k}$  accounts for the number of ways of choosing  $k$  monomers out of each HA trimer, and  $n = n_1 + 2n_2 + 3n_3$  is the total number of SA molecules bound to HA.  $N_{\text{HA}_3} = N_{\text{HA}} / 3$  is the total number of HA trimers.

As before, we approximate  $Z$  with the largest term in the summand, which is given by

$$\frac{p_{n_1+1}}{p_{n_1}} = \frac{3^{(N_{\text{HA}_3} - n_1 - n_2 - n_3)(N_{\text{SA}} - n)}}{n_1 + 1} \exp\left(-\frac{\Delta G_1}{RT}\right)$$

$$\begin{aligned}
&\approx \frac{3(N_{\text{HA}_3} - n_1 - n_2 - n_3)(N_{\text{SA}} - n)}{n_1} \exp\left(-\frac{\Delta G_1}{RT}\right) \approx 1 \\
\frac{p_{n_2+1}}{p_{n_2}} &= \frac{3(N_{\text{HA}_3} - n_1 - n_2 - n_3)(N_{\text{SA}} - n)(N_{\text{SA}} - n - 1)}{n_2 + 1} \exp\left(-\frac{\Delta G_2}{RT}\right) \\
&\approx \frac{3(N_{\text{HA}_3} - n_1 - n_2 - n_3)(N_{\text{SA}} - n)^2}{n_2} \exp\left(-\frac{\Delta G_2}{RT}\right) \approx 1 \\
\frac{p_{n_3+1}}{p_{n_3}} &= \frac{(N_{\text{HA}_3} - n_1 - n_2 - n_3)(N_{\text{SA}} - n)(N_{\text{SA}} - n - 1)(N_{\text{SA}} - n - 2)}{n_3 + 1} \exp\left(-\frac{\Delta G_3}{RT}\right) \\
&\approx \frac{(N_{\text{HA}_3} - n_1 - n_2 - n_3)(N_{\text{SA}} - n)^3}{n_3} \exp\left(-\frac{\Delta G_3}{RT}\right) \approx 1. \tag{S6}
\end{aligned}$$

The largest term thus corresponds to the equilibrium numbers of HA trimers bound to 1, 2, or 3 SA molecules, satisfying

$$\begin{aligned}
\frac{3(N_{\text{HA}_3} - \bar{n}_1 - \bar{n}_2 - \bar{n}_3)(N_{\text{SA}} - \bar{n})}{\bar{n}_1} &\approx \exp\left(\frac{\Delta G_1}{RT}\right) \\
\frac{3(N_{\text{HA}_3} - \bar{n}_1 - \bar{n}_2 - \bar{n}_3)(N_{\text{SA}} - \bar{n})^2}{\bar{n}_2} &\approx \exp\left(\frac{\Delta G_2}{RT}\right) \\
\frac{(N_{\text{HA}_3} - \bar{n}_1 - \bar{n}_2 - \bar{n}_3)(N_{\text{SA}} - \bar{n})^3}{\bar{n}_3} &\approx \exp\left(\frac{\Delta G_3}{RT}\right) \\
\bar{n} &= \bar{n}_1 + 2\bar{n}_2 + 3\bar{n}_3. \tag{S7}
\end{aligned}$$

The probability of breaking all connections (i.e.,  $n_1 = n_2 = n_3 = 0$ ), is

$$\begin{aligned}
\ln p_0 &= \ln \frac{1}{3^{N_{\text{HA}_3}} N_{\text{HA}_3}! N_{\text{SA}}!} - \ln Z \\
&\approx N_{\text{HA}_3} \ln \frac{N_{\text{HA}_3} - \bar{n}_1 - \bar{n}_2 - \bar{n}_3}{N_{\text{HA}_3}} + N_{\text{SA}} \ln \frac{N_{\text{SA}} - \bar{n}}{N_{\text{SA}}} + \bar{n} \\
&\quad - \bar{n}_1 \left( \ln \frac{3(N_{\text{HA}_3} - \bar{n}_1 - \bar{n}_2 - \bar{n}_3)(N_{\text{SA}} - \bar{n})}{\bar{n}_1} - \frac{\Delta G_1}{RT} \right)
\end{aligned}$$

$$\begin{aligned}
& -\bar{n}_2 \left( \ln \frac{3^{(N_{\text{HA}_3} - \bar{n}_1 - \bar{n}_2 - \bar{n}_3)} (N_{\text{SA}} - \bar{n})^2}{\bar{n}_2} - \frac{\Delta G_2}{RT} \right) \\
& -\bar{n}_3 \left( \ln \frac{(N_{\text{HA}_3} - \bar{n}_1 - \bar{n}_2 - \bar{n}_3) (N_{\text{SA}} - \bar{n})^3}{\bar{n}_2} - \frac{\Delta G_3}{RT} \right) \\
& \approx N_{\text{HA}_3} \ln \frac{N_{\text{HA}_3} - \bar{n}_1 - \bar{n}_2 - \bar{n}_3}{N_{\text{HA}_3}} + N_{\text{SA}} \ln \frac{N_{\text{SA}} - \bar{n}}{N_{\text{SA}}} + \bar{n}.
\end{aligned} \tag{S8}$$

Eqs. S7 and S8 are Eqs. 32 and 31 in the text.

### Derivation of Eqs. 33 and 34

We outline here the derivation of Eqs. 33 and 34, which describe the inhibition of adhesion when there is cooperativity in the binding of SA molecules to different monomers within a HA trimer.

Denoting  $m_i$  as the number of HA trimers that have  $i = 0, 1, 2, 3$  inhibitors bound, and  $n_{ij}$  as the number HA trimers that have  $i$  inhibitors and  $j$  SA molecules bound, we can write the partition function of the system as

$$\begin{aligned}
Z &= \frac{1}{3^{N_{\text{HA}_3}} N_{\text{HA}_3}! N_{\text{SA}}!} \sum_{\{m_i\}} \binom{N_{\text{HA}_3}}{m_1 \ m_2 \ m_3} \binom{3}{1}^{m_1} \binom{3}{2}^{m_2} \exp \left( -\frac{(m_1 + 2m_2 + 3m_3)\Delta\mu}{RT} \right) \cdot \\
& \sum_{\{n_{ij}\}} \binom{m_0}{n_{01} \ n_{02} \ n_{03}} \binom{m_1}{n_{11} \ n_{12}} \binom{m_2}{n_{21}} \cdot \binom{3}{1}^{n_{01}} \binom{2}{1}^{n_{11}} \binom{3}{2}^{n_{02}} \binom{N_{\text{SA}}}{n} n! \cdot \\
& \exp \left( -\frac{(n_{01} + n_{11} + n_{21})\Delta G_1 + (n_{02} + n_{12})\Delta G_2 + n_{03}\Delta G_3}{RT} \right).
\end{aligned} \tag{S9}$$

In Eq. S9, we assume that the inhibitors are in excess of the HA, such that the solution concentration of inhibitor  $[I]$  does not change with  $\{m_i\}$ , and we assume that the affinity of

inhibitor-HA binding does not depend on the presence or absence of other inhibitor-HA or SA-HA complexes.

The largest term in the summand corresponds to the equilibrium condition

$$\binom{3}{i} \frac{\bar{m}_0 - \sum_{j=1}^3 \bar{n}_{0j}}{\bar{m}_i - \sum_{j=1}^{3-i} \bar{n}_{ij}} = \left(\frac{K_I}{[I]}\right)^i, i = 1, 2, 3$$

$$\binom{3-i}{j} \frac{\bar{m}_i - \sum_{k=1}^{3-i} \bar{n}_{ik}}{\bar{n}_{ij}} = \frac{\exp\left(\frac{\Delta G_j}{RT}\right)}{(N_{SA} - \bar{n})^j}, i = 0, 1, 2, 3; j = 1, \dots, 3-i$$

$$\bar{m}_0 + \bar{m}_1 + \bar{m}_2 + \bar{m}_3 = N_{HA_3}$$

$$\sum_{i=0}^3 \sum_{j=1}^{3-i} j \bar{n}_{ij} = \bar{n}, \quad (\text{S10})$$

and the probability of breaking all HA-SA connections (i.e.,  $n_{ij} = 0$  for  $j > 0$ ), is

$$\ln p_0 = \ln \left( \frac{1}{3^{N_{HA_3}} N_{HA_3}! N_{SA}!} \sum_{\{m_i\}} \binom{N_{HA_3}}{m_1 m_2 m_3} \binom{3}{1}^{m_1} \binom{3}{2}^{m_2} \exp\left(-\frac{(m_1 + 2m_2 + 3m_3)\Delta\mu}{RT}\right) \right) - \ln Z$$

$$= \ln \left( \frac{1}{3^{N_{HA_3}} N_{HA_3}! N_{SA}!} \left(1 + \exp\left(-\frac{\Delta\mu}{RT}\right)\right)^{3N_{HA_3}} \right) - \ln Z$$

$$\approx N_{HA_3} \ln \frac{\bar{m}_0 - \bar{n}_{01} - \bar{n}_{02} - \bar{n}_{03}}{N_{HA_3} / \left(1 + \frac{[I]}{K_I}\right)^3} + N_{SA} \ln \frac{N_{SA} - \bar{n}}{N_{SA}} + \bar{n}. \quad (\text{S11})$$

The last line is the result of replacing  $Z$  by the largest term in the summand, applying the Stirling approximation, and substituting in the equilibrium condition in Eq. S10. Eqs. S10 and S11 are Eqs. 34 and 33 in the text.

## Adhesion of influenza virus to a cell population with variable numbers of receptors

Different cells have different SA densities on their surfaces, and thus different numbers of receptors at the interface. Here we analyze the effect of this variability on adhesion. Assuming that the expression and surface presentation of each individual receptor molecule is independent of other receptor molecules on the surface of a cell, the number of receptors at the interface in a cell population follows the Poisson distribution:

$$\chi(N_R) = \bar{N}_R^{N_R} \exp(-\bar{N}_R) / N_R!, \quad (\text{S12})$$

where  $\bar{N}_R$  is the average number of interface receptors in the cell population. The viral concentration bound to cells with  $N_R$  number of receptors,  $[CV](N_R)$ , satisfies

$$\frac{[CV](N_R)}{[V]} = [C](N_R)K_{D,ad}^{-1}(N_R) = [\text{cell}]\chi(N_R)K_{D,ad}^{-1}(N_R). \quad (\text{S13})$$

The total fraction of viral particles bound to cells is thus given by

$$f = \frac{\sum_{N_R} [CV](N_R)}{[V] + \sum_{N_R} [CV](N_R)} = \left( 1 + \frac{1}{[\text{cell}] \sum_{N_R} \chi(N_R) K_{D,ad}^{-1}(N_R)} \right)^{-1}. \quad (\text{S14})$$

Typically,  $K_{D,ad}^{-1}(N_R)$  increases rapidly with  $N_R$  around  $\bar{N}_R$ , while  $\chi(N_R)$  decreases rapidly with  $N_R > \bar{N}_R$ , and their product sharply peaks at a maximum near  $N_{R,max} \approx \bar{N}_R$ . The summation over  $N_R$  in Eq. S14 can thus be well approximated by the term  $\chi(N_{R,max})K_{D,ad}^{-1}(N_{R,max})$ , and Eq. S14 then reduces to Eq. 19, but with  $[\text{cell}]$  scaled by  $\chi(N_{R,max})$ . This suggests that the conclusions drawn for adhesion to an individual cell remain largely unchanged for adhesion to a population of cells with variable receptor densities.

To explore the quantitative difference resulting from the variability in  $N_R$ , we computed the  $K_{D,max}$  that corresponds to the HA-SA affinity at which half of the viral particles are bound (see

the discussion on the viable range of binding affinity in the main text), assuming the value  $\bar{N}_R = N_{SA}$  (in Table 1), and assuming the other parameter values in Table 1. Solving Eq. S14 for  $K_D$  with  $f = 0.5$  yields  $K_{D,\max} = 10.5$  mM, which is essentially the same as the value obtained for adhesion to cells with a uniform number of receptors (see main text).