

## Supporting Information

### Modulation of electronic properties from stacking orders and spin-orbit coupling for 3R-type MoS<sub>2</sub>

Xiaofeng Fan<sup>a,b,\*</sup>, W. T. Zheng<sup>b</sup>, Jer-Lai Kuo<sup>d</sup>, David J. Singh<sup>c</sup>, C.Q. Sun<sup>a</sup> and W. Zhu<sup>a,†</sup>

a. NOVITAS, School of Electrical and Electronic Engineering, Nanyang Technological University, 50 Nanyang Avenue, Singapore 639798, Republic of Singapore

b. College of Materials Science and Engineering, Jilin University, Changchun 130012, China

c. Department of Physics and Astronomy, University of Missouri, Columbia, Missouri 65211-7010, USA

d. Institute of Atomic and Molecular Sciences, Academia Sinica, Taipei, 10617, Taiwan

\*, † Correspondence and requests for materials should be addressed, xffan@jlu.edu.cn (X. F. Fan); ewzhu@ntu.edu.sg (W. Zhu)

**Fig. S1**

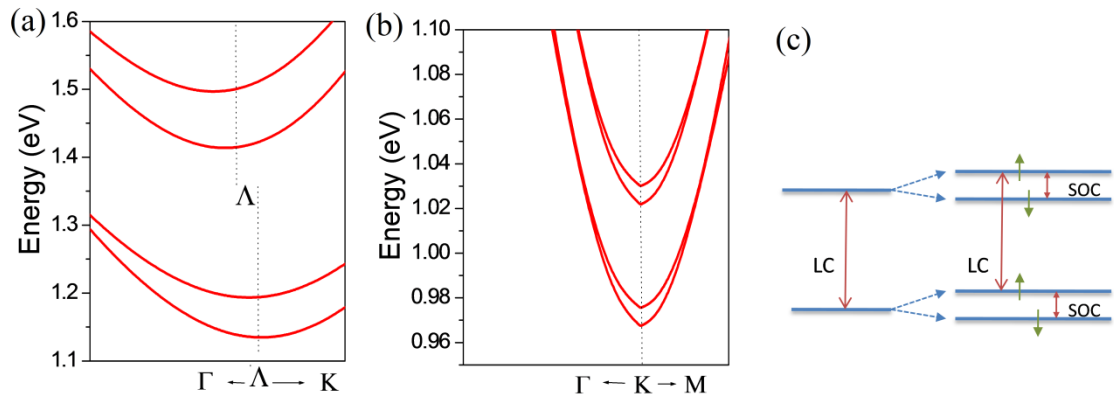


Fig. S1 Band structures of 2-layer MoS<sub>2</sub> with ab-stacking near the bottom of conduction band at  $\Lambda$  (a) and K point (b), and schematic of conduction band splitting at  $\Lambda$  and K point due to spin-orbit coupling and layer's coupling (c)

Fig. S2

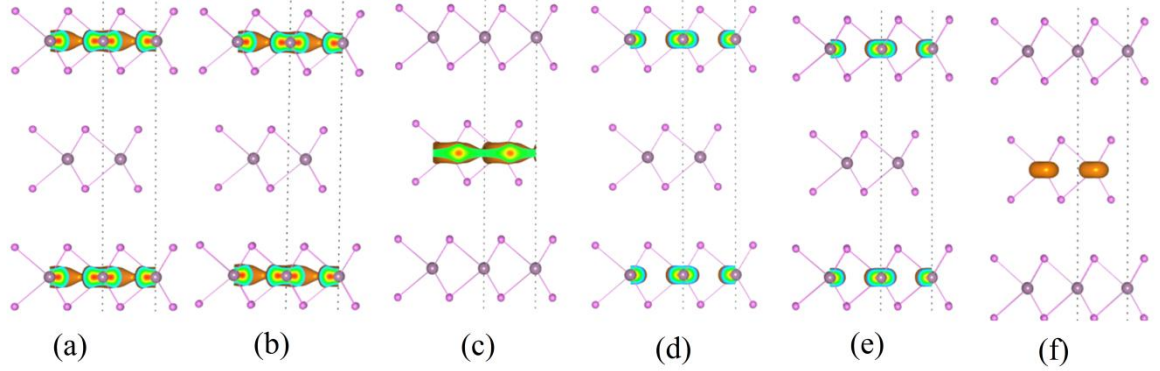


Fig. S2 Isosurface of band-decomposed charge density of six states of 3-layer MoS<sub>2</sub> with aba stacking order, at the top of valance band of K point (in Fig. 5a) including the states  $\sim|1+3\uparrow\rangle-1$ (a),  $\sim|1+3\uparrow\rangle-2$ (b),  $\sim|2\uparrow\rangle$ (c),  $\sim|1+3\downarrow\rangle-1$ (d),  $\sim|1+3\downarrow\rangle-2$ (e) and  $\sim|2\downarrow\rangle$  (f).

The spin-up states  $\sim|1+3\uparrow\rangle-1$  and  $\sim|1+3\uparrow\rangle-2$  are almost degenerate. The spin-up states  $\sim|1+3\downarrow\rangle-1$  and  $\sim|1+3\downarrow\rangle-2$  are almost degenerate. This may due to the being of mirror symmetry of aba stacking.

Fig. S3

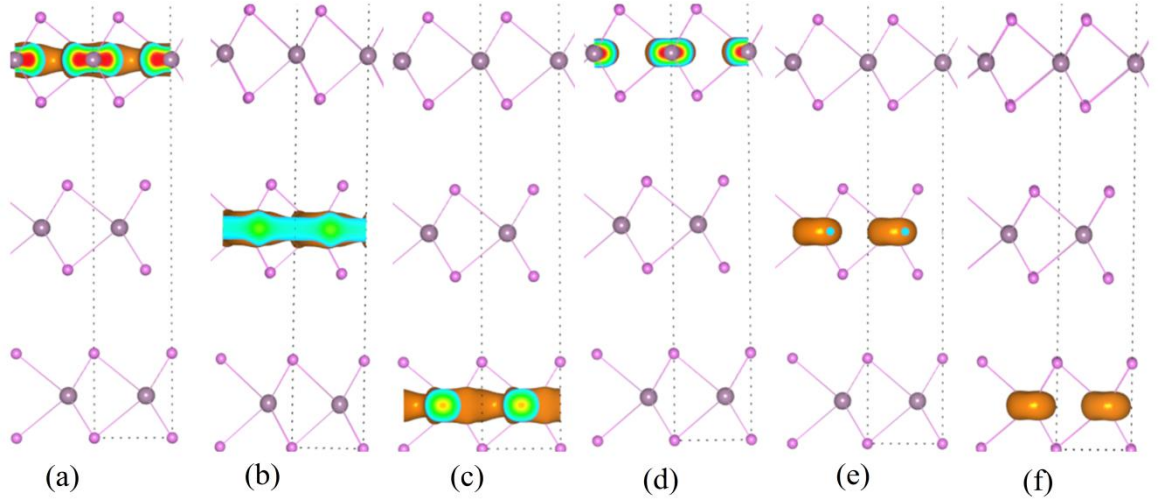


Fig. S3 Isosurface of band-decomposed charge density of six states of 3-layer MoS<sub>2</sub> with aba stacking order, at the top of valance band of K point (in Fig. 5a) including the states  $\sim|1\uparrow\rangle$  (a),  $\sim|2\uparrow\rangle$  (b),  $\sim|3\uparrow\rangle$  (c),  $\sim|1\downarrow\rangle$  (d),  $\sim|2\downarrow\rangle$  (e) and  $\sim|3\downarrow\rangle$  (f).

**Fig. S4**

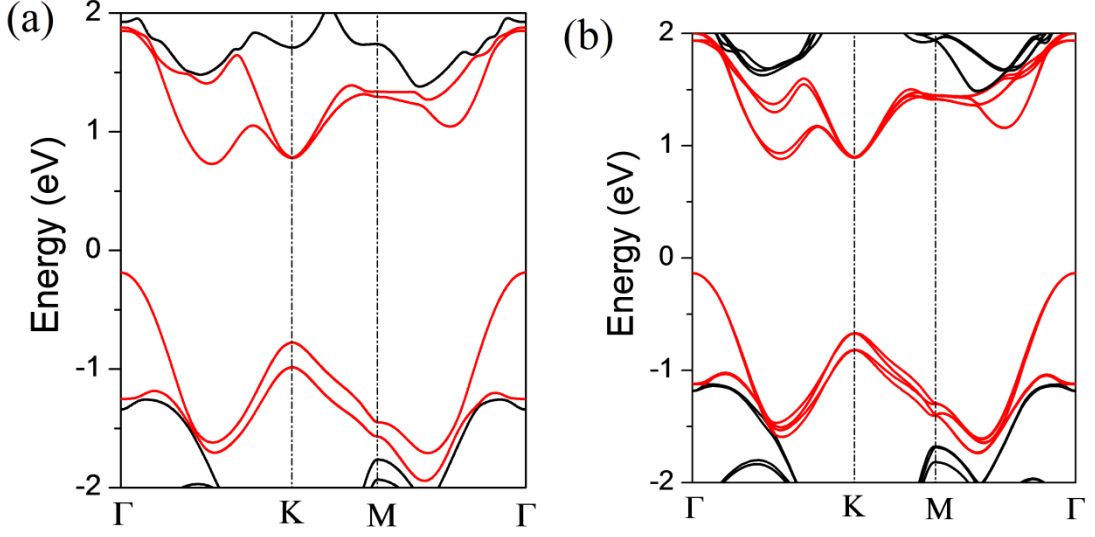


Fig. S4 Band structures of bulk 2H-MoS<sub>2</sub> (a) and 3R-MoS<sub>2</sub> (b) with spin-orbit coupling.

The band splitting of 2H-MoS<sub>2</sub> at CB-K is 208 meV and that of 3R-MoS<sub>2</sub> at CB-K is 149 meV.

### Model Hamiltonian

The model Hamiltonian for tree-layer MoS<sub>2</sub>

$$H(k) = \begin{pmatrix} E_0(k) & M_1(k) & M_2(k) \\ M_1^*(k) & E_0(k) & M_1(k) \\ M_2^*(k) & M_1^*(k) & E_0(k) \end{pmatrix}$$

The model Hamiltonian for four-layer MoS<sub>2</sub>

$$H(k) = \begin{pmatrix} E_0(k) & M_1(k) & M_2(k) & M_3(k) \\ M_1^*(k) & E_0(k) & M_1(k) & M_2(k) \\ M_2^*(k) & M_1^*(k) & E_0(k) & M_1(k) \\ M_3^*(k) & M_2^*(k) & M_1^*(k) & E_0(k) \end{pmatrix}$$

The model Hamiltonian is used to analyze the splitting of band splitting near band gap and understand the layer's coupling.  $M_1(k)$ ,  $M_2(k)$  and  $M_3(k)$  represent the coupling parameters of nearest-neighbor, second near-neighbor and third near-neighbor layer-coupling with same k point and same energy level , respectively.

Table S1. Band splitting values of the top of valance bands at  $\Gamma$  and K points (VB- $\Gamma$  and VB-K) and the bottom of conduction bands at  $\Lambda$  and K points (CB- $\Lambda$  and CB-K) for multi-layer 3R-type MoS<sub>2</sub> including 2L-ab, 3L-aba, 3L-abc, 4L-abab and 4L-abca with the different layers and stacking orders, calculated by DFT and fitted by model Hamiltonian (TM). Note that  $\Delta_1$ ,  $\Delta_2$  and  $\Delta_3$  represent the splitting values between the nearest-neighbor energy levels and  $\Delta$  is for total splitting value.

Stacking way	VB- $\Gamma$ (eV)		CB- $\Lambda$ (eV)		CB-K( $10^{-2}$ eV)		VB-K( $10^{-2}$ eV)	
	DFT	TM	DFT	TM	DFT	TM	DFT	TM
2L-ab								
$\Delta$	0.6389	0.6389	0.3058	0.3058	5.43	5.43	6.11	6.11
3L-aba								
$\Delta_1$	0.2732	0.2757	0.2699	0.2698	5.5370	5.5340	0.7900	0.7920
$\Delta_2$	0.5324	0.5307	0.1377	0.1378	0.1100	0.1100	6.0420	6.0420
$\Delta$	0.8057	0.8604	0.4076	0.4076	5.6470	5.6440	6.8320	6.8340
3L-abc								
$\Delta_1$	0.3052	0.3058	0.2278	0.2266	3.6910	3.6950	5.4900	5.4930
$\Delta_2$	0.5228	0.5218	0.1926	0.1936	5.3490	5.3450	4.2960	4.2930
$\Delta$	0.8281	0.8276	0.4204	0.4202	9.0400	9.0400	9.7860	9.7860
4L-abab								
$\Delta_1$	0.1707	0.1708	0.1572	0.1580	0.9750	0.9750	0.8110	0.8040
$\Delta_2$	0.3067	0.3066	0.2001	0.2006	4.2560	4.2220	5.1840	5.1760
$\Delta_3$	0.4214	0.4204	0.1017	0.1004	0.7360	0.7750	1.1500	1.1640
$\Delta$	0.8988	0.8978	0.4591	0.4590	5.9670	5.9720	7.1450	7.1440
4L-abca								
$\Delta_1$	0.1810	0.1809	0.1700	0.1698	4.0870	4.0730	5.6170	5.6640
$\Delta_2$	0.3037	0.3030	0.1798	0.1808	4.4640	4.4090	4.4890	4.4760
$\Delta_3$	0.4317	0.4319	0.1357	0.1357	5.4900	5.5470	4.6540	4.6480
$\Delta$	0.9164	0.9159	0.4855	0.4862	14.0410	14.0290	14.7600	14.7880

It is obvious that the values of band splitting from model Hamiltonian are consistent with that from DFT.