#### **Supplementary Information**

# **The intrinsic defect structure of exfoliated MoS2 single layers revealed by Scanning Tunneling Microscopy**

*Péter Vancsó\*, Gábor Zsolt Magda\*, János Pető, J.Y. Noh, Yong-Sung Kim, Chanyong Hwang, László P. Biró, and Levente Tapasztó*

# **I. Sample preparation**

The  $MoS<sub>2</sub>$  single layers have been exfoliated onto Au (111) substrates using a novel mechanical exfoliation technique developed by us, and described elsewhere [1], yielding single layers of several hundreds of microns lateral size. To ensure the high structural quality of the investigated samples we used synthetic bulk  $MoS<sub>2</sub>$  crystals (2DSemiconductors), which we found of higher quality than naturally occurring crystals investigated by us. The exfoliation process involves only a mild sonication step of a few seconds, which is not expected to induce any additional defect sites as samples sonicated for longer times did not display a higher defect concentration. The single layer nature of such flakes has been confirmed by a combination of optical microscopy, Raman spectroscopy as well as STM measurements of step height relative to the gold substrate near the  $MoS<sub>2</sub>$  flake edges.

### **II. Scanning Tunneling Microscopy and Spectroscopy characterization**

STM investigations have been performed on a Nanoscope E instrument, at room temperature, under ambient and in pure  $N_2$  atmosphere, using an atmospheric hood, in order to protect the samples from oxidation during scanning with higher bias voltages  $(>1V)$ . Landing the STM tip on a single layer  $MoS<sub>2</sub>$  flake has been achieved by guiding the tip under an optical microscope. This had been made possible by the large (>100 microns) lateral size of the exfoliated single layer flakes. Typical imaging parameters for atomic resolution were  $5 - 50$  $mV$ , and  $1 - 3 nA$ . To clearly resolve the atomic structure of point defects, ideal tip conditions and the fine tuning of the imaging parameters was needed. In contrast, to the apparent resolution of the hexagonal lattice of the top S atoms, the high quality images clearly resolving individual point defects could not be easily achieved. Typical STM images of point defect in  $MoS<sub>2</sub>$  single layers under non-ideal imaging conditions are shown in Fig S1a.



**Fig. S1**. *STM image of point defects in MoS<sup>2</sup> single layers acquired under non-ideal imaging conditions (10 mV, 2nA), when the Fermi level is located away from the defect states (a) and under ideal imaging conditions (50mV, 2 nA), in a different sample area (b).*

Reliable tunneling spectroscopy over a single point defect could not be achieved due to the thermal drift at the room temperature. However, reproducible tunneling spectra could be acquired near the defect sites. A typical spectrum is shown in Fig S2.a.



**Fig. S2**. *Tunneling spectra acquired on MoS<sup>2</sup> single layer deposited on Au (111) substrate.*

The tunneling spectra display a band gap of about 1.8 eV in excellent agreement with previous results. However, the density of states within the gap has a small but finite value, which we attribute to the presence of the Au substrate. It is also apparent that the Fermi level (vertical blue line) is close to the valence band maximum While the spectrum in Fig S2 a, has been predominantly observed, in a few areas spectra with significantly shifted Fermi level position have also been observed (Fig. S2 b), evidencing the spatial inhomogeneity of the doping, which we attribute to the variation of the interaction strength between the  $MoS<sub>2</sub>$  layer and the Au (111) substrate, due to impurity intercalation in restricted areas.

#### **III. Density Functional Calculations**

The DFT calculations were performed within the framework of local density approximation (LDA) using the Vienna *ab initio* simulation package (VASP) [2,3]. Projector augmented wave (PAW) pseudo-potentials [4,5] were used and the kinetic energy cut-off for the plane wave expansion was 350 eV. We applied  $8\times 8$  surface unit cell (192 host atoms) where the inplane lattice constant (distance between neighbor S atoms) was 3.124 Å from our previous LDA calculations [6] and the vacuum thickness was 15 Å. The atomic structures were relaxed until the Hellmann-Feynman forces are less than 0.02 eV/Å. The Γ- point was used for the Brillouin-zone summation. STM images were simulated with the simple Tersoff-Hamann approximation [7] using the calculated local density of states.

#### **IV. Sulfur Divacancies**

The atomic structures of  $V_S$  and  $V_{2S}$  are compared in Fig. S3 a and b. With the  $V_{2S}$ , the surrounding Mo atoms get closer to each other, as shown by the Mo-Mo bonding (red lines) in Fig. S3b. The defect level hybridization due to the binding of two  $V_S$ 's is schematically drawn in Fig. S3c. The bonding and anti-bonding states of  $(a_1-a_1)$  and  $(a_1-a_1)^*$  are generated by the hybridization of the singlet  $a<sub>l</sub>$  states, while four (*e-e*) states are generated by the double *e* state hybridization in a divacancy. The  $(a_1-a_1)$  bonding state is very close to the host VB and coupled with the host VB sates. The  $(a<sub>1</sub>-a<sub>1</sub>)$ <sup>\*</sup> antibonding state is found to be characteristically observed in STM for the  $V_{2S}$  defect.



**Figure S3.** Atomic structure and schematic bandstructure of disulfur vacancy  $(V_{2S})$ .

# **References**

**.** 

 $2$  G. Kresse, J. Hafner, Ab initio molecular-dynamics simulation of the liquid-metal amorphous-semiconductor transition in germanium. Phys. Rev. B 1994; 49(20):14251–14269.

 $3$  G. Kresse, J. Furthmüller, Efficient iterative schemes for ab initio total-energy calculations using a plane-wave basis set. Phys. Rev. B 1996; 54(16):11169–11186.

<sup>4</sup> G. Kresse, D. Joubert, From ultrasoft pseudopotentials to the projector augmented-wave method. Phys. Rev. B 1999; 59(3):1758–1775.

5 P.E. Blöchl, Projector augmented-wave method. Phys. Rev. B 1994; 50(24):17953–17979.

 $^6$  J.-Y. Noh, H. Kim, Y.-S. Kim, Stability and electronic structures of native defects in single-layer MoS2. Phys. Rev. B 2014; 89:205417.

 $^7$  J. Tersoff, D.R. Hamann, Theory and Application for the Scanning Tunneling Microscope. Phys. Rev. Lett. 1983; 50(25):1998–2001.

 $1$  G. Zs. Magda, J. Pető, G. Dobrik, C. Hwang, L. P. Biró, L. Tapasztó, Exfoliation of large-area transition metal chalcogenide single layers. Sci Rep. 2015; 5:14714.