

## Supplementary Materials for **Quasi-ballistic carbon nanotube array transistors with current density exceeding Si and GaAs**

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## SUPPLEMENTARY TEXT

### Full sweep transfer characteristics:

Example forward and backward sweep  $I_{DS} - V_{GS}$  characteristics are plotted in fig. S1, where the forward sweep corresponds to the same data plotted in Fig. 2A. The device exhibits hysteresis, the origin of which has previously been attributed to adsorbed water and oxygen molecules (64). Recent work has shown that hysteresis can be eliminated from CNT FETs via encapsulation (65). The forward sweep data presented in the main text are more representative of  $I_{DS} - V_{GS}$  characteristics in the absence of hysteresis (66).

### Lead resistance subtraction:

Parasitic lead and wire resistances ( $R_{leads}$ ) on the order of 100  $\Omega$  (accounting for 23% of the total measured resistance ( $R_{total}$ ) at the highest on-state conductance) are subtracted from  $R_{total}$  to obtain the intrinsic device resistance ( $R_{device}$ ) following

$$R_{device} = R_{total} - R_{leads} \quad (1)$$

To measure  $R_{leads}$ , a 30 nm thick, 5  $\mu\text{m}$  wide stripe of Au is deposited and patterned via e-beam lithography, spanning the FET channel from source to drain in order to intentionally electrically short-circuit the FETs. This measurement is performed on each device to account for small pattern-to-pattern differences in lithography and Pd electrode thickness. Example device lead structures and  $I$ - $V$  measurements for CNT FETs with different lead structure are shown in fig. S2A,B.

### Optical spectroscopy evidence of high semiconducting purity:

An absorbance spectrum of the initial dispersion of PFO-BPy sorted arc discharge CNTs in toluene is presented in fig. S3. Metallic CNT peaks, which are typically observed in the range of 630 – 810 nm are not detected (33). Previous reports on polyfluorene based sorting of CNTs have estimated the semiconducting CNT purity using the parameter  $\phi$ , which takes into account the area under the CNT absorbance peaks and the absorbance

background (33, 37). Ding *et al.* calibrated the  $\phi$  value showing that  $\phi = 0.38$  corresponds to 99% semiconducting purity (37). From our absorbance spectrum in fig. S3 we calculate a higher  $\phi$  of 0.428, which indicates the purity is higher than 99% semiconducting.

#### Effect of annealing on $I_{on}/I_{off}$ :

To determine whether annealing treatment selectively etches metallic versus semiconducting CNTs, we measure  $I_{on}/I_{off}$  and off-conductance ( $G_{off}$ ) and compare devices without and with 400°C annealing treatment in fig. S4. The  $I_{on}/I_{off}$  increases slightly from  $10^{4.7\pm 1.6}$  to  $10^{5.5\pm 1.1}$  when comparing devices fabricated without annealing to devices fabricated with 400°C annealing. However, the on-conductance increase caused by annealing the films has a significant contribution to the on/off ratio increase, thus  $G_{off}$  is a more direct way to compare the semiconducting purity before and after annealing.  $G_{off}$  is  $10^{-8.2\pm 1.8}$  without annealing and  $10^{-8.7\pm 1.2}$  with annealing. The similarity in  $G_{off}$  with and without annealing suggests the semiconducting purity is not altered significantly by vacuum annealing.

#### $I_{off}$ variability:

Figure S5A shows that  $I_{on}/I_{off}$  exponentially decreases and Fig. 2A shows that ambipolar behavior increases as a function of  $V_{DS}$  with a similar trend for three  $L_{ch}$  regimes used in this study. Recent work on single CNT FETs demonstrate that ambipolar behavior can be better controlled through design of asymmetric top gate (67), dual feedback gate (60), or wrap around gate (61) architectures to reduce tunneling currents at high  $V_{DS}$ . Future implementation of such gate architectures may enable similar electrostatic control over CNT array FETs (67). Further analysis of  $I_{on}/I_{off}$  is performed at a constant low  $V_{DS}$  of -100 mV to shed light on other possible sources of  $I_{on}/I_{off}$  variability and  $L_{ch}$  trends.

In fig. S5B,  $I_{on}$  and  $I_{off}$  are presented for 88 “rinsed+annealed” devices as a function of  $L_{ch}$ , where  $I_{on}$  and  $I_{off}$  are the absolute maximum and minimum current, respectively, extracted from the forward sweep of the  $I_{DS}-V_{GS}$  transfer curves measured at  $V_{DS} = -100$  mV.  $I_{off}$  variability is significant even at low  $V_{DS}$ , spanning ~4 orders of

magnitude, and increases at  $L_{ch}$  shorter than 250 nm. Several possibilities for the  $I_{off}$  variability, leading to  $I_{on}/I_{off}$  lower than  $10^5$  in 30% of devices, are discussed below.

(i) Metallic CNTs are one possible source of  $I_{off}$  variability where at shorter channel lengths there is an increased probability that any metallic CNTs that exist in the channel will directly span the source-drain gap. Based on previous  $I$ - $V$  measurements of single metallic CNTs at positive gate bias we predict a metallic short-circuit would result in a total  $I_{off}$  for the CNT arrays (measured at  $V_{DS} = -100$  mV) that is greater than  $1 \mu\text{A}$  ( $\sim 0.25 \mu\text{A} \mu\text{m}^{-1}$  for  $W_{ch} = 4 \mu\text{m}$ ) (68). A high  $I_{off}$  of  $0.3 \mu\text{A} \mu\text{m}^{-1}$  is observed in only 1 out of 88 devices measured, resulting in a low  $I_{on}/I_{off} = 300$  for this particular device. The remaining 87 devices have  $I_{off}$  lower than  $0.25 \mu\text{A} \mu\text{m}^{-1}$  and  $I_{on}/I_{off} > 10^3$ , suggesting that no metallic CNTs span the channel in these devices.

(ii) CNT-CNT interactions (e.g., due to CNT-crossing or interactions) in the channel may lead to electronic hybridization or dielectric screening effects that increase  $I_{off}$  (15, 17).

(iii) Another possible factor that may decrease  $I_{on}/I_{off}$  and increase the variability of  $I_{on}/I_{off}$  is the presence of inhomogeneously distributed fixed and mobile charged impurities at the CNT/SiO<sub>2</sub> interface (66, 69). Previous work has indicated that threshold voltage ( $V_T$ ) variability can be attributed to such inhomogeneity (70). In an array FET, if the  $V_T$  of each CNT in the array differs then this effect leads to increased subthreshold swing ( $SS = (d(\log_{10}(I_{DS})/dV_{GS}))^{-1}$ ) and decreased  $I_{on}/I_{off}$  for the ensemble array (17). Cao *et al.* have shown that at short channel lengths surface charge heterogeneity becomes even more prevalent, which further increases  $V_T$  variability (70). This is consistent with the increased  $I_{off}$  variability observed, here, however further studies will be required to decouple the separate contributions of (i), (ii), and (iii) to  $I_{on}/I_{off}$  variability and its  $L_{ch}$  dependence. In fig. S5C, we plot  $I_{on}/I_{off}$  vs.  $SS$  showing a negative correlation, which could also be attributed to (iii) (70). If inhomogeneously distributed fixed and mobile charged impurities or CNT-CNT interactions prove to be the origin of the  $I_{off}$  variability in the future then these non-idealities may be improved by implementing high-k dielectrics or by passivating the dielectric surface.

An estimate of the semiconducting CNT purity is obtained from (i), above. We assume that the one device with poor  $I_{on}/I_{off} < 10^3$  contains a single metallic CNT. The total number of CNTs that span  $S$ - $D$  electrodes is 10,208 CNTs (116 per device  $\times$  88 devices). From this analysis, we estimate a semiconducting CNT purity greater than 99.99%.

#### Postdeposition treatment effect on output characteristics:

$I_{DS} - V_{DS}$  curves for “as-deposited” and “rinsed+annealed” samples are plotted for  $V_{GS} - V_T = -2$  V in fig. S6A in order to qualitatively determine the effect of postdeposition treatments on the Pd-CNT injection barrier. Prior to treatment, an injection barrier is present in the “as-deposited” sample, which is evidenced by the highly non-linear  $I_{DS} - V_{DS}$  characteristics and increasing differential conductance ( $dI_{DS}/dV_{DS}$ ) at small  $V_{DS} > -0.6$  V, as shown in fig. S6B. The “rinsed+annealed” treatment significantly reduces this barrier leading to ohmic  $I_{DS} - V_{DS}$  characteristics and nearly constant  $dI_{DS}/dV_D$  at  $V_{DS} < -0.3$  V.

#### Protocols used to quantify atomic concentrations from XPS and optical absorbance spectra:

Optical absorbance measurements for CNT and PFO-BPy solutions are presented in fig. S7 where absorbance coefficients are log base 10. CNT concentration ( $c$ ) is calculated for the  $S_{22}$  transition envelope using the relationship

$$\frac{\int_{E_1}^{E_2} A_{S_{22}}^{meas} dE}{c} = \frac{\int_{E_1}^{E_2} A_{S_{22}}^{REF} dE}{c^{REF}} \quad (2)$$

where  $A_{S_{22}}^{meas}$  is the measured absorbance of our CNT sample, integrated over the energy range 1.0 –1.8 eV, and  $A_{S_{22}}^{REF}$  is the absorbance of a CNT sample of known concentration  $c^{REF}$  from Mistry *et al.* (33), integrated over the energy range 1.0 –1.8 eV. A PFO-BPy optical cross section ( $\sigma_{PFO-BPy}$ ) of  $1.1 \times 10^5$  cm<sup>2</sup> g<sup>-1</sup> was determined by measuring  $A$  at 350 nm for a series of dilutions of PFO-BPy reference samples of known concentration. To deconvolute the PFO-BPy absorbance in the CNT@PFO-BPy solution spectrum, we subtract off the contribution of bare CNTs in fig. S7A. The resulting concentrations of CNTs and PFO-BPy are 4.2 and 4.4  $\mu$ g ml<sup>-1</sup>, respectively, for the black-line spectrum shown in fig. S7A. Based on this analysis we determine a PFO-BPy:CNT mass ratio of 1.0, and a C:N atomic ratio of 41.

XPS is performed on FESA-deposited films at a packing density of 47 CNTs  $\mu\text{m}^{-1}$  (AFM shown in fig. S8).

Atomic surface densities ( $N_i$ ) for  $i=\text{C}$ , N, and Cl are quantified and summarized in table S1. The  $N_i$  are quantified using the  $\text{SiO}_2$  substrate as a standard for calibration and the relationship (71)

$$N_i = \frac{A_i S_{Si}}{A_{Si} S_i} \times (\sin(\theta) \times \lambda_{Si, SiO_2} \times \rho_{Si, SiO_2}) \quad (3)$$

where  $A_i$  is the measured peak-area, and  $A_{Si}$  is the area of the silicon background peak from the  $\text{SiO}_2$  substrate used for calibration. The instrument specific sensitivity factors for each element are  $S_C = 1$ ,  $S_N = 1.8$ ,  $S_{Cl} = 2.29$ , and  $S_{Si} = 0.82$ . The measurement take off angle ( $\theta$ ) is  $45^\circ$ , the inelastic mean free path of Si in  $\text{SiO}_2$  ( $\lambda_{Si, SiO_2}$ ) is  $2.96 \times 10^{-7}$  m, and the density of Si atoms in  $\text{SiO}_2$  ( $\rho_{Si, SiO_2}$ ) is  $2.28 \times 10^{22}$  atoms  $\text{cm}^{-3}$ . To check the accuracy of the  $N_i$  calculations, we convert the  $N_C = 1.05 \times 10^{15}$  C atoms  $\text{cm}^{-2}$  measured in the “rinsed+annealed” film into a CNT packing density. We subtract off the instrument background C atoms and the C atoms that can be attributed to the PFO-BPy backbone, which can be determined from the known ratio of N to C in the backbone of 1:11. We assume for simplicity that the CNTs are of a single chirality corresponding to the experimental mean diameter of  $\sim 1.5$  nm (diameter distribution measured in fig. S7B). Based on these assumptions, we estimate a packing density of 57 CNTs  $\mu\text{m}^{-1}$ , which is very close to the density of 47 CNT  $\mu\text{m}^{-1}$  measured via SEM imaging.

Reduction ratios ( $RR$ ) are reported in the main text and in table S1 to quantify the changes in  $N_i$  resulting from post-deposition treatments. These ratios are measured with respect to the “as-deposited” samples.

Control experiments are performed using chloroform ink without CNTs. When these inks are deposited on  $\text{SiO}_2/\text{Si}$  substrates via FESA, the inks do not leave a detectable trace of Cl. These control experiments suggest that the Cl detected in Fig. 3 arises from Cl residues that are directly bonded to the CNTs or PFO-BPy.

Temperature-dependent FTIR measurement of PFO-BPy side chains and backbone:

Thermal degradation of alkyl side chains from PFO-BPy is confirmed by the reduction of the alkyl C-H stretch peaks located at 2855 and 2926  $\text{cm}^{-1}$  in Fig. 3C. The spectra are presented as a function of temperature in fig. S9A. The C=C stretching modes are presented, as well, in fig. S9B. In fig. S9A it is evident that the alkyl C-H bonds start to decompose at a temperature in the range of 300 – 400°C. The  $\text{sp}^2$  C=C breathing mode is observed at 1450  $\text{cm}^{-1}$ , shown in fig. S9B, which originates from the PFO-BPy backbone (53). The peak at 1450  $\text{cm}^{-1}$  persists after annealing at 400°C, suggesting that the PFO-BPy backbone units remain partially intact following the “rinsed+annealed” post-deposition treatment at 400°C that is used to eliminate the side-chains.

#### On-conductance comparison of “annealed” versus “rinsed+annealed” treatments:

CNT films having undergone “annealed” and “rinsed+annealed” treatments are compared via electrical measurement in fig. S10 and XPS elemental quantification in table S1. The on-state conductance for FETs constructed from each treatment are similar at each channel length indicating both samples have similar channel and contact resistance. XPS quantities in table S1 show that “annealed” and “rinsed+annealed” samples have similar C, N and Cl reduction ratios. Thus, we conclude annealing is the dominant treatment and that high electronic performance can be achieved with both “annealed” and “rinsed+annealed” treatments.

#### Conductance per tube measured for FETs with $L_{\text{ch}} < 150$ nm:

On-state conductance and the number of CNTs spanning the *S-D* electrodes are measured for 13 CNT array FETs in order to obtain a conductance per tube measurement. These values are plotted in units of quantum conductance ( $2e^2/h$ ) in fig. S11. The SEM images from these short FETs are used to calculate an average packing density of 47 CNTs  $\mu\text{m}^{-1}$ .

### Contact resistance and mean free path extraction:

A transmission line measurement (54) is performed to extract  $2R_C$  and  $L_{MFP}$  for a set of 88 “rinsed+annealed” CNT array FETs in fig. S12A. There is scatter in the experimentally measured data, and the fit uncertainty in the each of the two parameters is correlated. Thus, to extract possible  $2R_C$  and  $L_{MFP}$  values, we start by fitting  $2R_C$  (y-intercept) and the uncertainty in  $2R_C$  ( $\sigma_{2R_C}$ ) as shown in fig. S12A. Next, we fix  $2R_C$  to three values (best fit  $2R_C$ , best fit  $2R_C - \sigma_{2R_C}$ , and best fit  $2R_C + \sigma_{2R_C}$ ) and then fit  $L_{MFP}$  (which is proportional to the inverse of the slope) for each  $2R_C$ . The best fit yields  $2R_C = 34$  k $\Omega$ -CNT and  $L_{MFP} = 42$  nm.  $\sigma_{2R_C}$  is 6 k $\Omega$ -CNT. For  $2R_C = 28$  k $\Omega$ -CNT, the fit  $L_{MFP} = 33$  nm. For  $2R_C = 40$  k $\Omega$ -CNT, the fit  $L_{MFP} = 56$  nm. These parameters are separately verified using the Y-function method (YFM) (details on YFM extraction method are provided in our previous studies of CNT array FETs (31)). We analyze 11 devices with diffusive transport (longer than 225 nm) where  $L_{ch} > L_{MFP}$ , with results shown in fig. S12B. From YFM analysis we obtain  $2R_C = 33 - 47$  k $\Omega$ -CNT, and  $L_{MFP} = 45 - 63$  nm, which agrees well with the ranges obtained from TLM.

Using a simple drift diffusion model, we approximate free carrier mobility for the CNT array

$$\mu = \frac{eL_{MFP}}{v_F m^*} \quad (4)$$

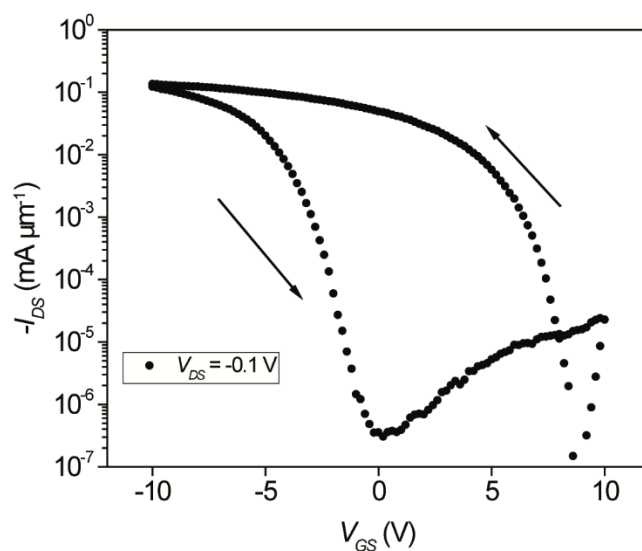
where  $e$  is the electron charge,  $v_F$  is the Fermi velocity in graphene equal to  $8 \times 10^5$  m s<sup>-1</sup>, and  $m^*$  is the diameter dependent ( $d = 1.5$  nm) hole effective mass  $m^* = \frac{2\hbar}{3dv_f}$  (72). Mobility ranges from 1130 - 1912 cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup> for the TLM fit  $L_{MFP}$  ranging from 33 to 56 nm.

### Polymer wrapped CNT diameter distribution:

Optical absorbance spectra of PFO-BPy sorted CNTs in chloroform are measured in order to estimate the diameter distribution from the  $S_{22}$  peak in fig. S7B. The  $S_{22}$  peak can be fit relatively well to a Gaussian function and from this fit we obtain a diameter distribution of  $1.5 \pm 0.2$  nm based on known chirality-dependent  $S_{22}$  energies from literature (73).

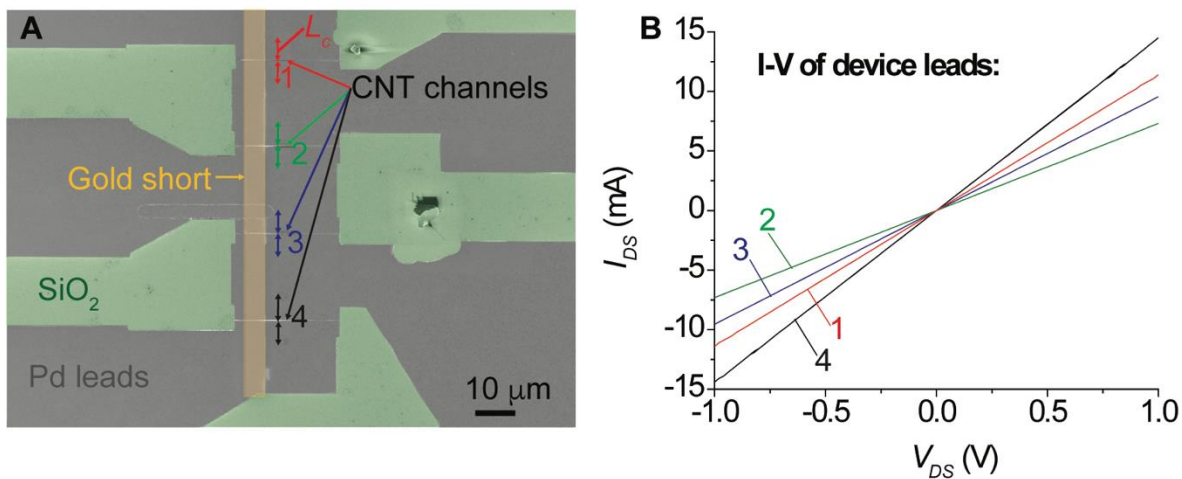


## SUPPLEMENTARY FIGURES AND TABLES

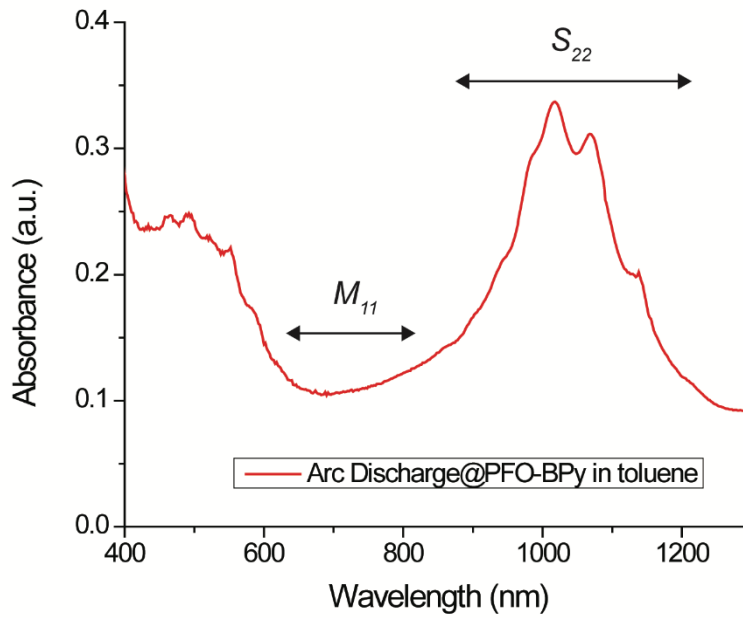


**fig. S1. Full sweep transfer characteristics.** Full hysteresis loop for CNT array device shown in Fig. 2A.

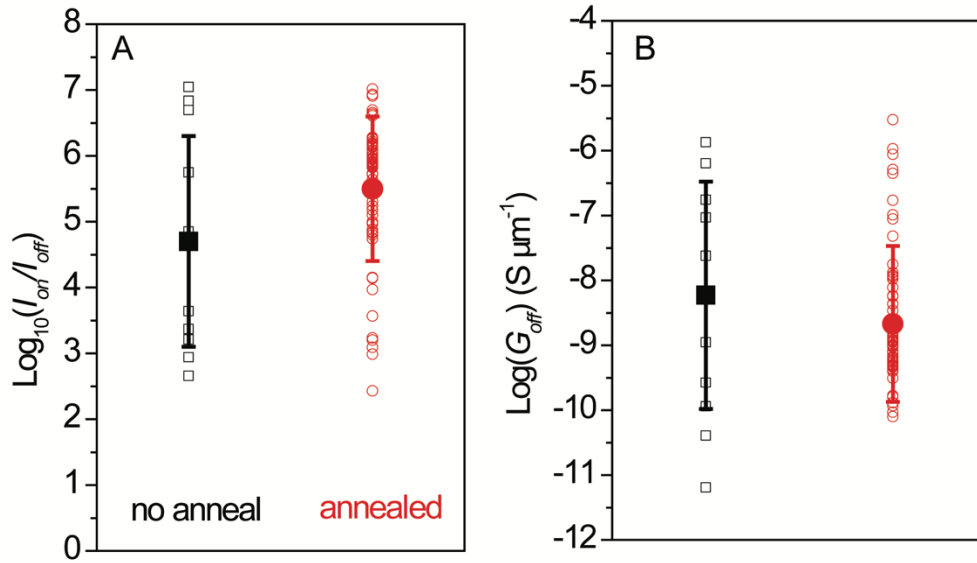
Arrows indicate directions of forward ( $V_{GS} = -10 \rightarrow 10$  V) and backward ( $V_{GS} = 10 \rightarrow -10$  V) sweeps.



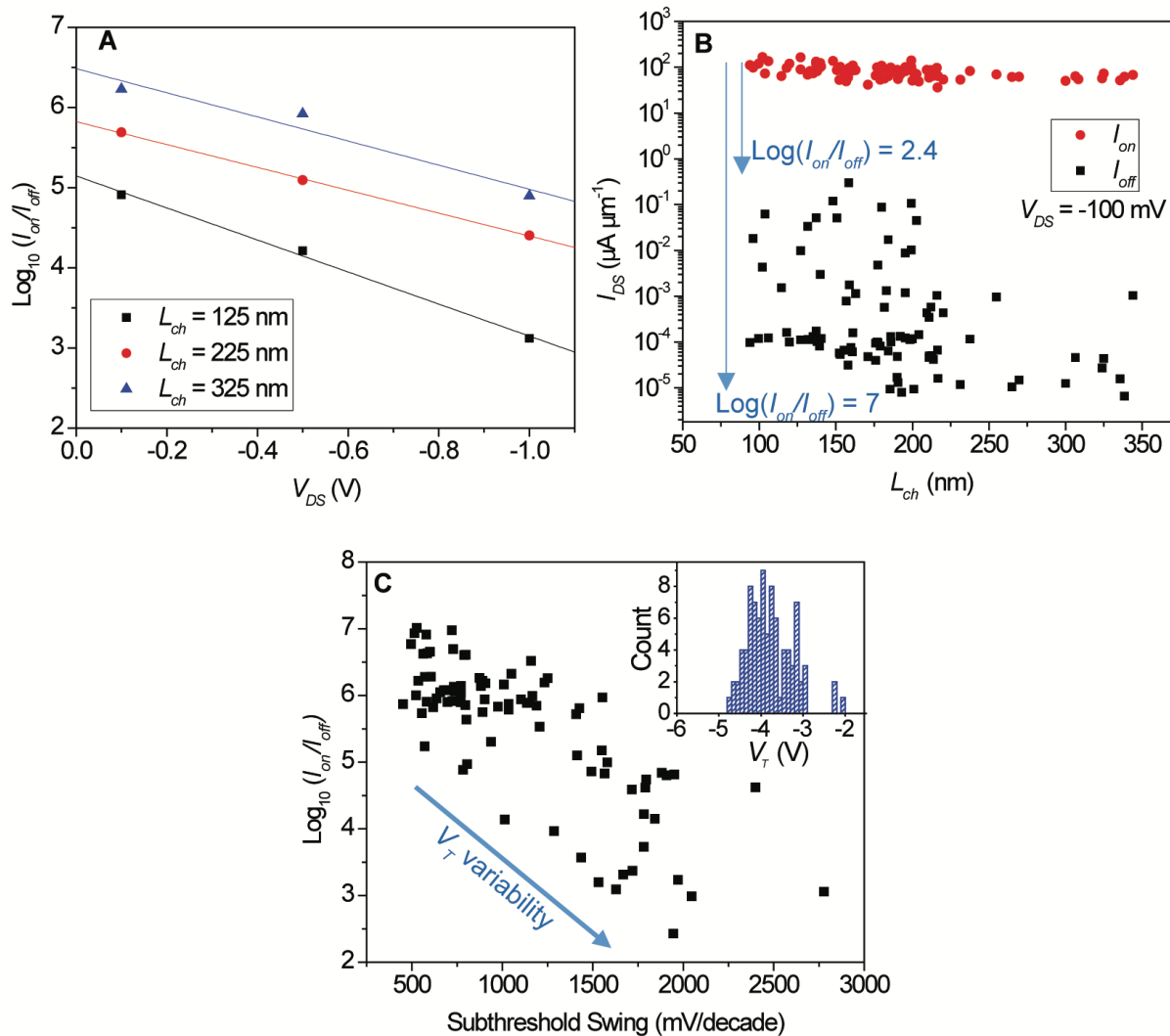
**fig. S2. Lead resistance subtraction.** (A) Top-view false-colored SEM image of CNT array FET lead- and electrode- structure with CNT channel regions and gold stripe short-circuit regions highlighted and labeled. (B)  $I$ - $V$  measurements of leads from different devices, showing some variation in the lead resistance.



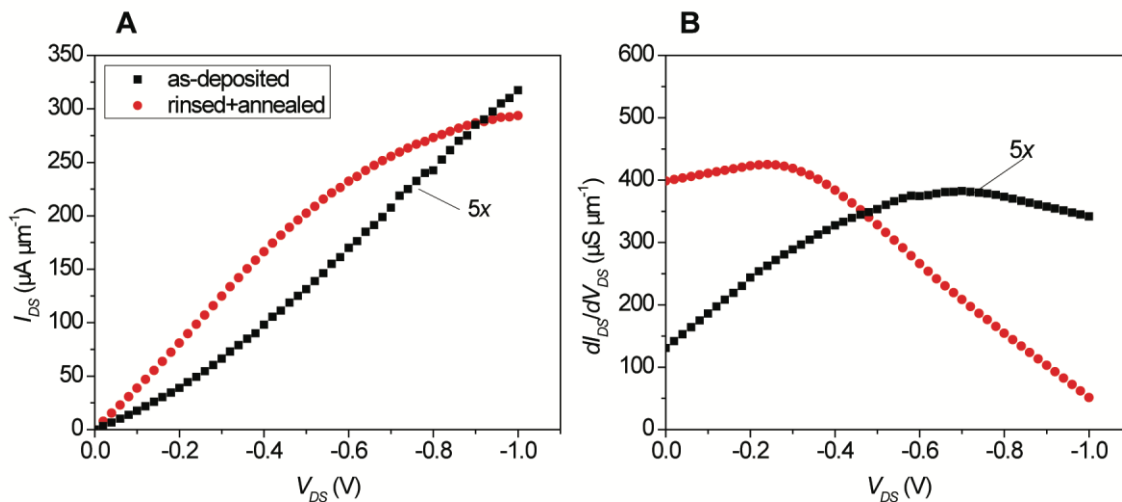
**fig. S3. Optical spectroscopy evidence of high semiconducting purity.** Absorbance spectrum of PFO-BPy sorted arc discharge CNTs in toluene. No metallic CNTs are detected as evidenced by the absence of a  $M_{11}$  peak in the range of 630 – 810 nm. To extract  $\phi$  values, wavelength was converted to energy prior to calculating CNT peak areas.



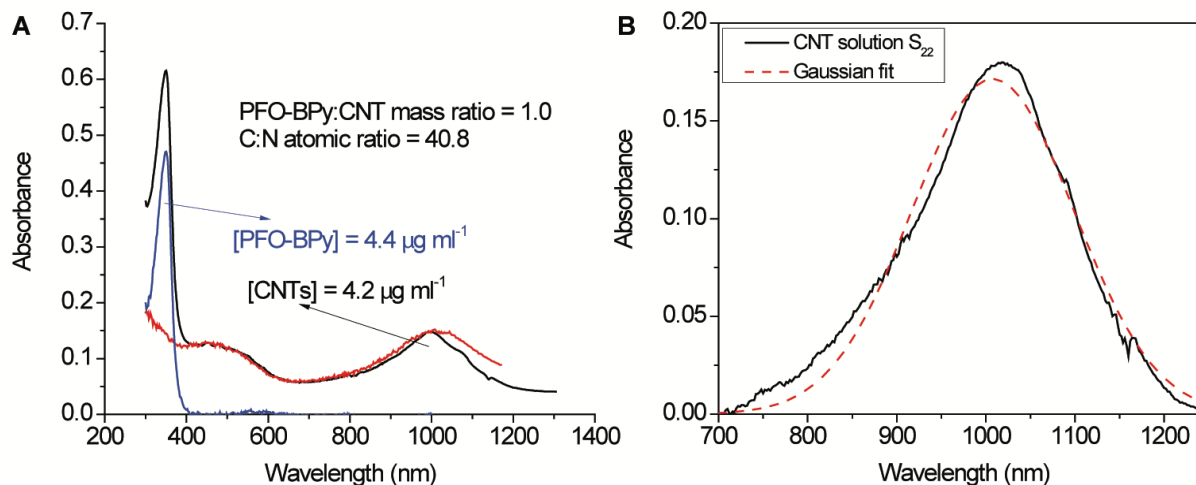
**fig. S4. Effect of annealing on  $I_{\text{on}}/I_{\text{off}}$ .** Comparison of the on/off ratio (A) and off-conductance (B) of FETs constructed using films without (black squares) versus with (red circles) 400°C annealing treatment. Individual devices are plotted as open symbols, where a total of 10 and 56 devices are measured for without and with annealing treatments, respectively. The mean and standard deviation are also plotted as filled symbols. To ensure a majority of CNTs span the channel, only devices with  $L_{\text{ch}}$  in the range of 150 – 250 nm are analyzed, with the same mean  $L_{\text{ch}}$  of ~200 nm for each treatment condition.



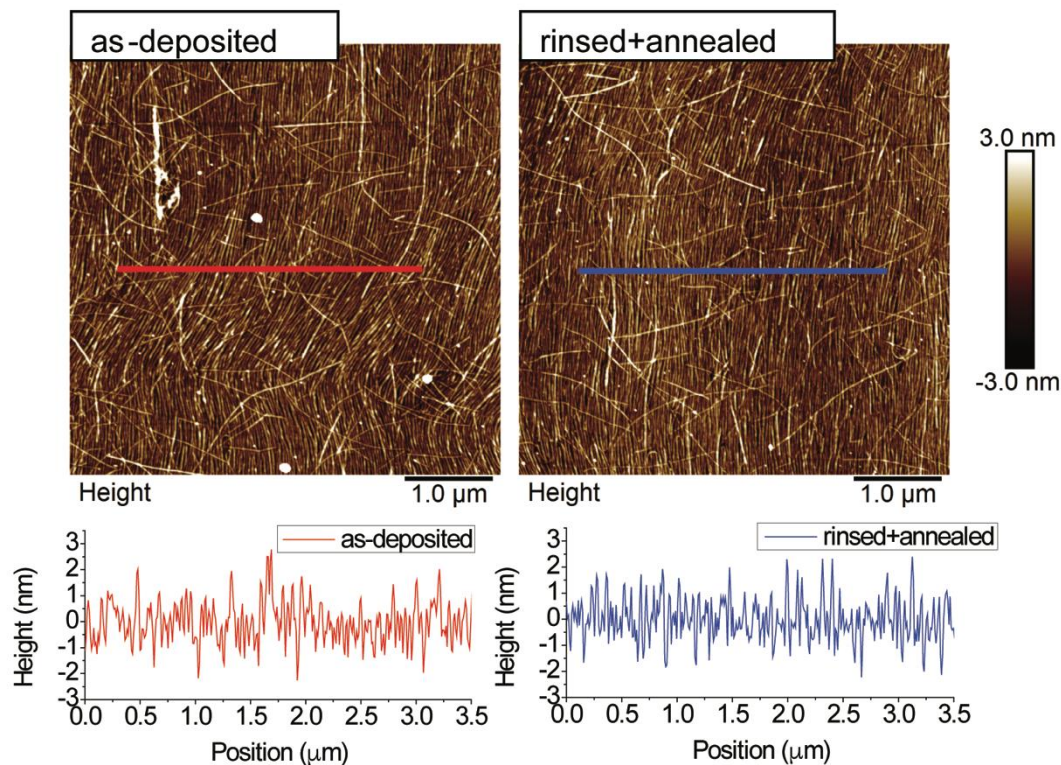
**fig. S5.  $I_{off}$  variability.** (A) Exponential decrease in  $I_{on}/I_{off}$  as a function of  $V_{DS}$  at various  $L_{ch}$  regimes.  $I_{on}/I_{off}$  is the mean calculated for  $> 7$  devices for  $L_{ch}$  ranges of 100-150 nm (125 nm), 200-250 (225 nm), and 300-350 nm (325 nm). (B) On and off state current for 88 “rinsed+annealed” devices at various  $L_{ch}$ . (C) Dependence of  $I_{on}/I_{off}$  on sub-threshold swing (SS).



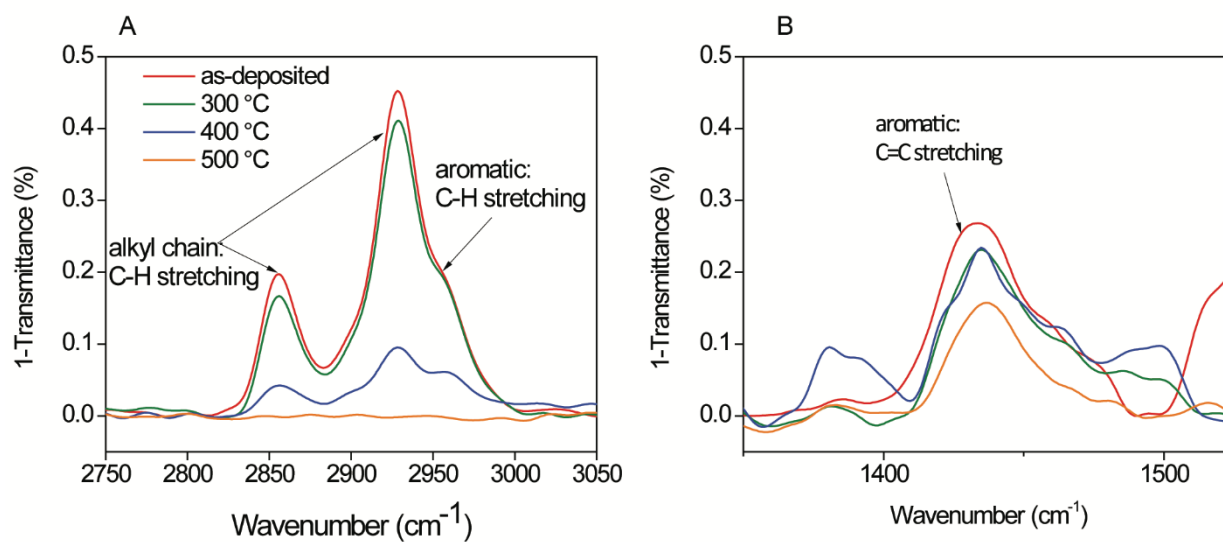
**fig. S6. Postdeposition treatment effect on output characteristics.**  $I_{DS} - V_{DS}$  measurement of “as-deposited” and “rinsed+annealed” CNT array FETs suggest the presence of a Schottky barrier at the Pd-CNT interface of the “as-deposited” CNT array FET.



**fig. S7. Optical absorbance spectra used to verify XPS atomic concentrations and CNT diameter distribution.** (A) Black curve is an optical absorbance spectrum of an arc discharge@PFO-BPy solution in chloroform measured in a 1 cm path length cuvette showing the polymer peak centered at 350 nm and the CNT  $S_{33}$  and  $S_{22}$ . Red curve is a drop-cast arc discharge@PFO-BPy film subject to 900 ° C annealing in carbon dioxide atmosphere to completely degrade polymer. Blue spectrum is the black minus the red curve. (B) Gaussian fit of the CNT  $S_{22}$  to obtain a diameter distribution of  $1.5 \pm 0.2$  nm.

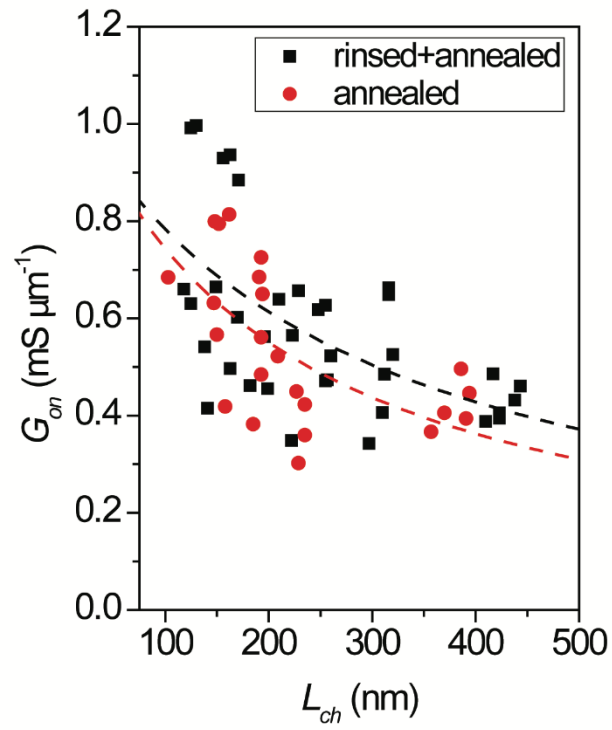


**fig. S8. Atomic force microscopy height measurement used to quantify packing density for XPS calibration.** AFM height maps of “as-deposited” and “rinsed+annealed” films demonstrates similar film morphology before and after treatments with average film height of ~3 nm.



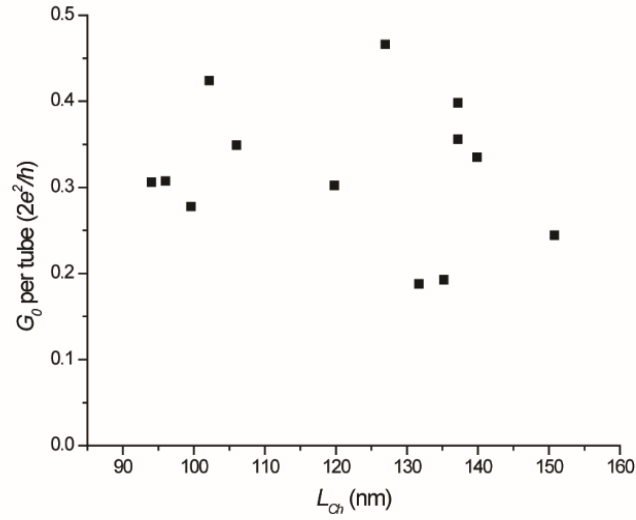
**fig. S9. Temperature-dependent FTIR measurement of PFO-BPy side chains and backbone.** (A) Temperature dependent FTIR measurement of the alkyl side chain C-H bonds from PFO-BPy. (B) FTIR spectra

of C=C breathing mode at  $1450\text{ cm}^{-1}$  for PFO-BPy films indicate the presence of aromatic carbon before and after the films are annealed and that the C=C bonds from carbon in the PFO-BPy backbone begin to degrade at  $500^\circ\text{C}$  annealing treatment.

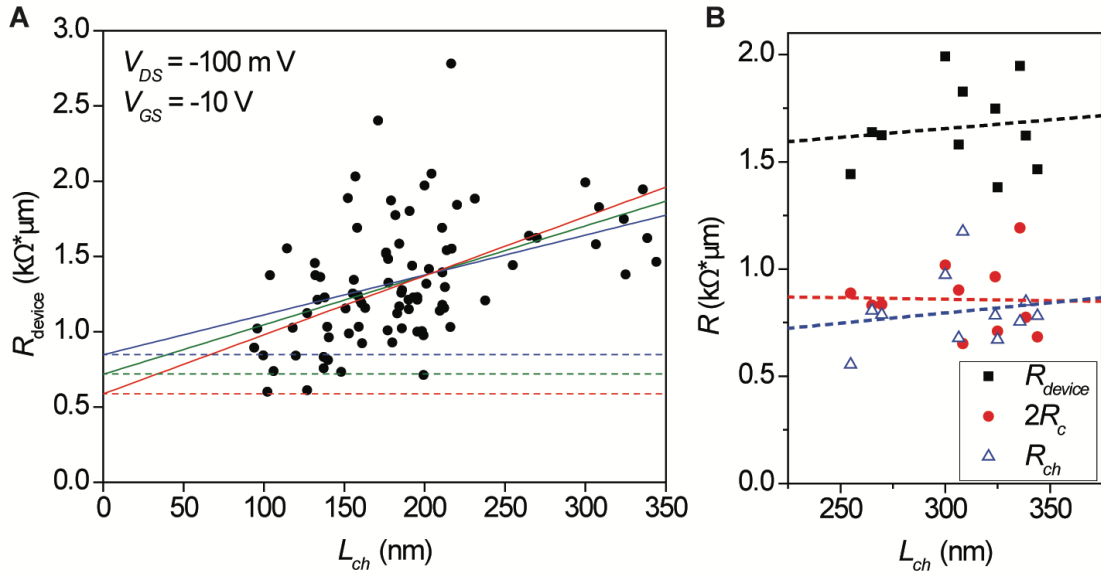


**fig. S10. On-conductance comparison of “annealed” versus “rinsed + annealed” treatments.** On-conductance is extracted from transfer characteristics at  $V_{DS} = -100$  mV and  $V_{GS} = -10$  V as a function of  $L_{ch}$  for “annealed” and “rinsed + annealed” surface treatments. Lead resistance is not subtracted from these data.





**fig. S11. Conductance per tube measured for FETs with  $L_{ch} < 150$  nm.** On-state conductance is normalized to the number of CNTs spanning  $S$ - $D$  electrodes for 13 devices.



**fig. S12. Contact resistance and mean free path extraction.** (A) Transmission line analysis of  $R_{device}$  vs.  $L_{ch}$  using equation (1) of the main text to extract  $2R_c$  and  $L_{MFP}$ . Red lines correspond to a lower bound for  $2R_c$  and  $L_{MFP}$ , green lines are the original fit, and blue lines correspond to the upper bound. (B) YFM extracted  $2R_c$  and channel resistance ( $R_{ch}$ ) for 11 separate devices as a function of channel length. Linear fits demonstrate that  $R_{ch}$  decreases with  $L_{ch}$  while  $2R_c$  is constant.

**table S1. Surface treatment XPS summary.** \*Non-negligible concentrations of carbon ( $4.0 \times 10^{14}$  atoms  $\text{cm}^{-2}$ ) and nitrogen ( $5.6 \times 10^{12}$  atoms  $\text{cm}^{-2}$ ) measured on blank substrates are accounted for and subtracted from the CNT array values for each element, respectively. Reduction percentages are reported in parentheses for “rinsed”, “annealed”, and “rinsed+annealed” treatments.

	Total C* ( $10^{15}$ atoms $\text{cm}^{-2}$ )	C from PFO-BPy inferred from N density ( $10^{15}$ atoms $\text{cm}^{-2}$ )	Total N* ( $10^{13}$ atoms $\text{cm}^{-2}$ )	Total Cl ( $10^{13}$ atoms $\text{cm}^{-2}$ )
		C from sources other than PFO-BPy ( $10^{15}$ atoms $\text{cm}^{-2}$ )		
<b>As-deposited</b>	$2.68 \pm 0.12$	$1.25 \pm 0.06$	$6.41 \pm 0.28$	$1.67 \pm 0.46$
		$1.43 \pm 0.10$		
<b>Rinsed</b>	$2.45 \pm 0.16$ (RR = $8.5 \pm 2.5\%$ )	$1.14 \pm 0.03$ (RR = $8.9 \pm 3.2\%$ )	$5.83 \pm 0.18$ (RR = $8.9 \pm 3.2\%$ )	$0.77 \pm 0.02$ (RR = $53.8 \pm 5.2\%$ )
		$1.31 \pm 0.12$ (RR = $8.0 \pm 2.3\%$ )		
<b>Annealed</b>	$1.83 \pm 0.18$ (RR = $31.8 \pm 5.4\%$ )	$0.61 \pm 0.36$ (RR = $50.7 \pm 0.9\%$ )	$5.57 \pm 0.33$ (RR = $12.6 \pm 1.7\%$ )	~0 (RR = ~100%)
		$1.22 \pm 0.14$ (RR = $15.1 \pm 10.0\%$ )		
<b>Rinsed + Annealed</b>	$1.62 \pm 0.05$ (RR = $33.7 \pm 2.5\%$ )	$0.53 \pm 0.01$ (RR = $57.8 \pm 1.4\%$ )	$4.8 \pm 0.09$ (RR = $25.1 \pm 2.53\%$ )	~0 (RR = ~100%)
		$1.1 \pm 0.04$ (RR = $23.1 \pm 2.9\%$ )		

**table S2. CNT FET performance comparisons.** Key metrics for CNT array FETs produced in this work compared to previous state-of-the-art and simulated/extrapolated CNT FETs (1, 17, 19, 74).

	Reference	Material / channel length (nm)	On-state conductance ( $\mu\text{S } \mu\text{m}^{-1}$ ) per CNT	Highest $I_{\text{on}}$ ( $\mu\text{A } \mu\text{m}^{-1}$ )	$I_{\text{on}}/I_{\text{off}}$	Density (CNTs $\mu\text{m}^{-1}$ )
Experimental	This work	Solution CNT array / 100	1700	900	$10^3 - 10^7$	50
			$0.47 G_0$			
	Shulaker, Stanford, <i>IEEE IEDM</i> 2014 [19]	CVD CNT array / 400	120	120	$10^4 - 10^5$	100
			$0.015 G_0$			
Cao, IBM, <i>Nature Nanotech.</i> 2013 [17]	Solution CNT array / 120	250	130	$10 - 10^3$	1000	
		$0.003 G_0$				
Hypothetical	Ouyang, Univ. of Florida, <i>Solid State Electron.</i> 2011 [74]	Simulated CNT array / 100	1200	720	$10^4$	50
			$0.15 G_0$			
	Franklin, IBM, <i>Nano Lett.</i> 2012 [1]	Hypothetical, extrapolated from single CNT FET data / 10	1300	630	$10^3$	200
			$0.04 G_0$			