

## Linear stability analysis of the within-season system

We consider a within-season system derived from Eqs. (1) and (3). In such a system, we first neglect the equations expressing the time evolution of statoblasts, because statoblasts only affect other state variables in the transition between two seasons. Moreover, in order to define a disease-free equilibrium, we neglect the equation related to the time evolution of susceptible bryozoan biomass and susceptible fish abundance and treat the terms  $B_S$  and  $F_S$  appearing in other equations as parameters. All parameters are assumed as constant in time. Hence, this system reads:

$$\frac{dB_C}{d\tau} = r [1 - \rho(B_S + B_C + B_O)] B_C - (d_{CO} + \psi)B_C + d_{OC}B_O + B_S Z_F^*; \quad (\text{A1a})$$

$$\frac{dB_O}{d\tau} = r_O [1 - \rho(B_S + B_C + B_O)] B_O + d_{CO}B_C - (d_{OC} + \psi)B_O; \quad (\text{A1b})$$

$$\frac{dF_E}{d\tau} = -(\mu_F + h)F_E + F_S Z_B^*; \quad (\text{A1c})$$

$$\frac{dF_I}{d\tau} = (1 - \epsilon)hF_E - (\mu_F + a + \gamma)F_I; \quad (\text{A1d})$$

$$\frac{dF_C}{d\tau} = \epsilon hF_E + \gamma F_I - (\mu_F + \zeta)F_C; \quad (\text{A1e})$$

$$\frac{dZ_B}{d\tau} = \pi_B^* B_O - \mu_Z Z_B^*; \quad (\text{A1f})$$

$$\frac{dZ_F}{d\tau} = \pi_F^* (F_I + \kappa F_C) - \mu_Z Z_F^*. \quad (\text{A1g})$$

The disease-free equilibrium (DFE) of the system (A1) is  $\mathbf{x}_{DF} = \{B_C; B_O; F_E; F_I; F_C; Z_B^*; Z_F^*\} = \mathbf{0}$ .

We now linearise (A1) around  $\mathbf{x}_{DF}$  and calculate its Jacobian:

$$\mathbf{J}_{\mathbf{DF}} = \begin{bmatrix} J_{DF,11} & d_{OC} & 0 & 0 & 0 & 0 & B_S \\ d_{CO} & J_{DF,22} & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & -h - \mu_F & 0 & 0 & F_S & 0 \\ 0 & 0 & (1 - \epsilon)h & -a - \gamma - \mu_F & 0 & 0 & 0 \\ 0 & 0 & \epsilon h & \gamma & -\zeta - \mu_F & 0 & 0 \\ 0 & \pi_B^* & 0 & 0 & 0 & -\mu_Z & 0 \\ 0 & 0 & 0 & \pi_F^* & \kappa \pi_F^* & 0 & -\mu_Z \end{bmatrix}.$$

where  $J_{DF,11} = r(1 - \rho B_S) - \psi - d_{CO}$  and  $J_{DF,22} = r_O(1 - \rho B_S) - \psi - d_{OC}$ .

According to a classical result of dynamical system theory [1], the spectral abscissa  $\alpha_s(\mathbf{J}_{\mathbf{DF}})$  of the Jacobian of (A1) expresses the speed at which a perturbation of the equilibrium state will propagate; if  $\alpha_s(\mathbf{J}_{\mathbf{DF}}) < 0$ , the DFE is stable.

The determinant of  $\mathbf{J}_{\mathbf{DF}}$  reads:

$$\det(\mathbf{J}_{\mathbf{DF}}) = J_1 - J_2 \mathcal{T}$$

where

$$J_1 = B_S F_S d_{CO} h \pi_B^* \pi_F^* [\epsilon \kappa a + \kappa \gamma + (1 - \epsilon) \zeta + (1 + \epsilon \kappa - \epsilon) \mu_F];$$

$$J_2 = (h + \mu_F)(a + \gamma + \mu_F)(\zeta + \mu_F) \mu_Z^2;$$

$$\mathcal{T} = \psi^2 + [d_{CO} + d_{OC} - (r + r_O)(1 - \rho B_S)] \psi - (1 - \rho B_S) [r_O d_{CO} + r d_{OC} + r r_O(1 - \rho B_S)].$$

Note that  $J_1$  and  $J_2$  are non-negative by construction. Given that the order of  $\mathbf{J}_{\mathbf{DF}}$  is odd, a switch in the sign of  $\alpha_s(\mathbf{J}_{\mathbf{DF}})$  from negative to positive corresponds to a switch in the sign of  $\det(\mathbf{J}_{\mathbf{DF}})$  from negative to positive. Also,  $\det(\mathbf{J}_{\mathbf{DF}}) > 0$  is a sufficient condition for the instability of the  $\mathbf{x}_{\mathbf{DF}}$ . This condition is always true when  $\mathcal{T} < 0$ : in this case, the parasite can spread within the bryozoan population even in the absence of the fish population (i.e. when  $F_S = 0$  and thus  $J_1 = 0$ ).

Indeed, it is straightforward to show that  $\mathcal{T}$  is the determinant of the Jacobian of a system that considers only bryozoans (i.e. Eqs. (A1a) with  $Z_F^* = 0$  and (A1b)), with  $\mathcal{T} < 0$  representing the instability criterion for the DFE of such a system. If instead  $\det(\mathbf{J}_{\text{DF}}) > 0$  with  $\mathcal{T} > 0$ ,  $\mathbf{x}_{\text{DF}}$  is still unstable, but in this case the parasite needs to cycle between the two hosts to invade the system.

A condition of incipient instability can be derived when  $\det(\mathbf{J}_{\text{DF}}) = 0$ , which leads to an expression for the reproduction number  $\mathcal{R}$ :

$$\mathcal{R} = \frac{J_1}{J_2 \mathcal{T}}. \quad (\text{A2})$$

The DFE is unstable when  $\mathcal{R}$  is larger than unity. Note that the expression (A2) is only valid if  $\mathcal{T} > 0$ ; failing that, the DFE is always unstable and  $\mathcal{R}$  can not be defined.

Figure A1 identifies regions of stability and instability for the disease-free equilibrium in the parameter space.

## References

- [1] Luenberger, David G. (1979). Introduction to dynamic systems: theory, models, and applications. Wiley. ISBN 978-0-471-02594-8.

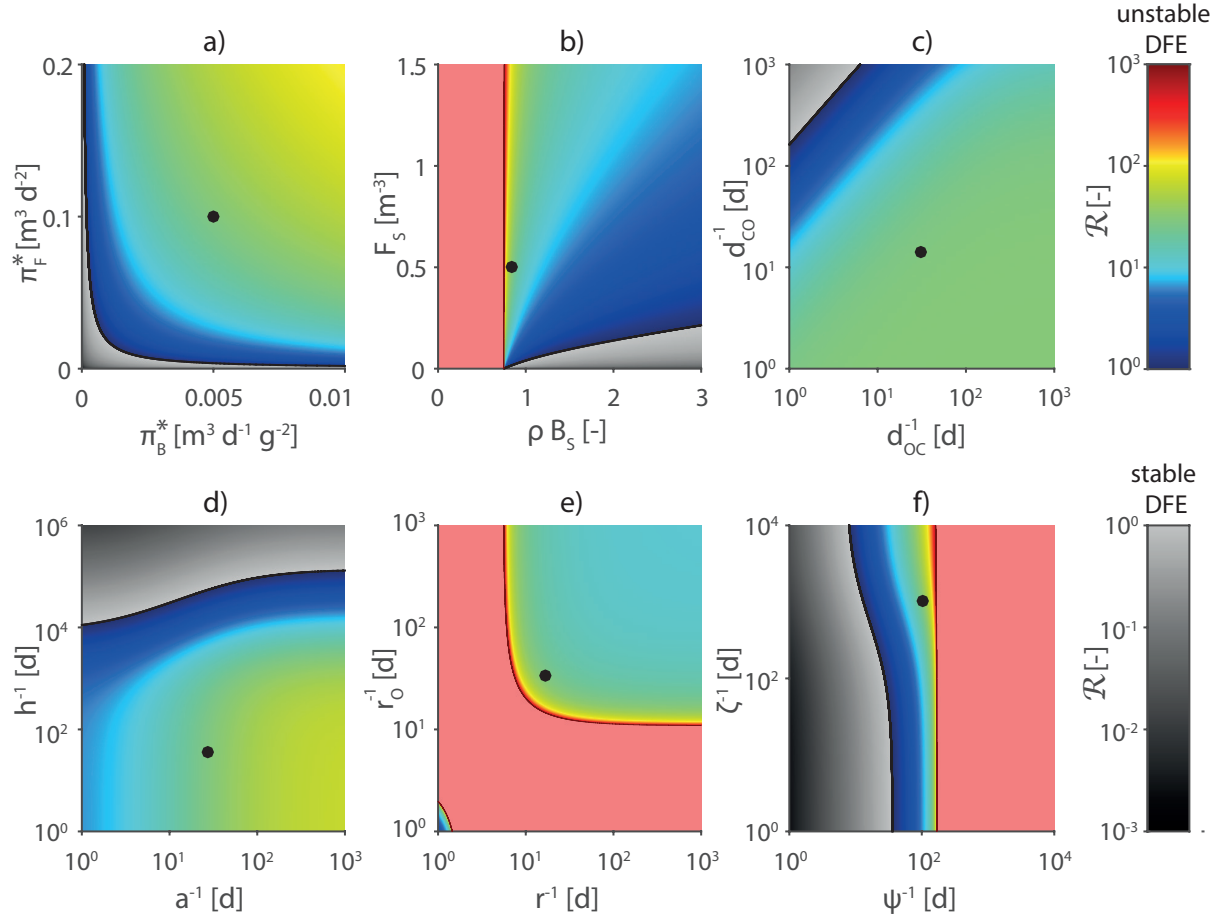


Figure A1: Values of the reproductive number  $\mathcal{R}$  of system (A1) as a function of the model parameters. Two parameters are varied at a time, the others are assumed as in Table 3. Regions where  $\mathcal{T} < 0$  (DFE unstable) are coloured in light red. Black dots refer to the reference parameter set. All rates are expressed as mean times (i.e., by their inverse). With regards to temperature-dependent parameters, their value at 15 °C is displayed.