APPENDIX A: Analytical calculation of SPEN diffusion effects upon employing linearly-swept 180° pulses

This Appendix presents the results involved in calculating diffusion effects when employing 180° sweeps for performing the SPEN encoding, rather than the 90° chirp pulse used in the main text. Presenting such analysis is important, as 180° sweeps appear as the most promising way of decoupling the diffusion and the imaging gradients cross-coupling terms (*cf.* Fig. 4). By contrast to sequences using chirp 90° pulses to excite the spin packets, 180° sweeps in SPEN schemes are tuned to invert the spins in the transverse plane; excitation is thus created at the beginning of the experiment by a (usually slice-selective) homogeneous 90° pulse. As a result of this, the phase accumulated by spins during a gradient G_e of period T_e and a gradient G_a of duration T_a (*cf.* Fig. 2C)

$$K_{local}^{180^{\circ}-SPEN}(t,z) = \frac{d(\phi_{e}(z) + \phi_{a}(t,z))}{dz} = \frac{d}{dz} \begin{cases} \gamma G_{e}zt & 0 \le t \le t_{180}(z) \\ \phi_{rf} \left[t_{180}(z) \right] + \left[\phi_{rf} \left[t_{180}(z) \right] - \gamma G_{e}z \left[t_{180}(z) \right] \right] + \gamma G_{e}z \left[t - t_{180}(z) \right], \quad t_{180}(z) \le t \le T_{e} \\ 2\phi_{rf} \left[t_{180}(z) \right] - \gamma G_{e}z \left[t_{180}(z) - t \right] + \gamma G_{a}zt, \qquad T_{e} \le t \le T_{a} \end{cases}$$
(A1)

Expressions for the timing $t_{180}(z)$ when the adiabatic RF sweep reaches the resonance frequency of spin packet at a particular z coordinate, as well as for the phase ϕ_{RF} taken by the B₁ field at the time of this inversion, are similar to those given in the main text for the chirp 90° pulse. Taking the spatial derivatives involved in Eq. (A1) yields the relevant K_{local} wavenumbers,

$$K_{local}^{180^{\circ}-SPEN}(t,z) = \begin{cases} \gamma G_{e}t & 0 \le t \le t_{180}(z) \\ \gamma G_{e}[t-2t_{180}(z)], & t_{180}(z) \le t \le T_{e} \\ \gamma G_{e}[T_{e}-2t_{180}(z)] + \gamma G_{a}t, & T_{e} \le t \le T_{a} \end{cases}$$
(A2)

Substituting this expression into Eq. 8 provides the full argument of the exponential attenuation function,

$$A_{180^{\circ}-SPEN}(t,z) = \exp\left[-D\gamma^{2} \left(\frac{t_{180}^{3}(z)G_{e}^{2}}{3} + \frac{G_{e}^{2}(t_{180}^{3}(z) + (-2t_{180}(z) + T_{e})^{3})}{3} + \frac{(2t_{180}(z)G_{e} - G_{e}T_{e})^{3} + (-2t_{180}(z)G_{e} + G_{a}T_{a} + G_{e}T_{e})^{3}}{3G_{a}}\right)\right]$$
(A3)

These K_{local} and attenuation functions still have to account for the diffusion gradients G_d , and for the purging gradients G_{pr} that in 180°-encoded SPEN MRI are needed for shifting the symmetric phase parabola that the adiabatic sweep imparts, to one corner of the FOV (*cf.* Fig. 2C). Taking these additional factors into account, the total K_{local} wavenumber becomes

$$K_{local}(t,z) = \begin{cases} -\gamma G_{pr}t & 0 \le t \le T_{pr} \\ -\gamma G_{pr}T_{pr} + \gamma G_d(t - T_{pr}) & T_{pr} \le t \le T_{pr} + \delta \\ -\gamma G_{pr}T_{pr} + \gamma G_d \delta + \gamma G_e(t - (T_{pr} + \delta)) & T_{pr} + \delta \le t \le T_{pr} + \delta + t_{180}(z) \\ -\gamma G_{pr}T_{pr} + \gamma G_d \delta + \gamma G_e(t - 2t_{180}(z)) & T_{pr} + \delta + t_{180}(z) \le t \le T_{pr} + \delta + T_e \\ -\gamma G_{pr}T_{pr} + \gamma G_d \delta + \gamma G_e(T_e - 2t_{180}(z)) - \gamma G_d(t - (T_{pr} + \delta + T_e)) & T_{pr} + \delta + T_e \le t \le T_{pr} + \delta + \Delta \\ -\gamma G_{pr}T_{pr} + \gamma G_e(T_e - 2t_{180}(z)) + \gamma G_a t & T_{pr} + \delta + \Delta \le t \le T_{pr} + \delta + \Delta + T_a \end{cases}$$
(A4)

is

And the ensuing exponential attenuation function

$$A(t,z) = \exp \left[-D\gamma^2 \left[\frac{G_{pr}^2 T_{pr}^3}{3} + \frac{\delta (G_d^2 \delta^2 - 3G_d \delta G_{pr} T_{pr} + 3G_{pr}^2 T_{pr}^2)}{3G_e} + \frac{(G_d \delta + t_{180}(z)G_e - G_{pr} T_{pr})^3 + (-G_d \delta + G_{pr} T_{pr})^3}{3G_e} + \frac{(t_{180}(z)G_e - \delta (G_d + G_e) + (-G_e + G_{pr}) T_{pr})^3}{3G_e} + \frac{(G_d \delta - 2t_{180}(z)G_e - G_{pr} T_{pr} + G_e (\Delta + T_{pr}))^3}{3G_e} + \frac{(-2t_{180}(z)G_e + G_e T_e + G_d (-\Delta + 2\delta + T_e) - G_{pr} T_{pr})^3}{3G_d} + \frac{(2t_{180}(z)G_e + G_d (\Delta - \delta - T_e) - G_e T_e + G_{pr} T_{pr})^3}{3G_a} + \frac{(2t_{180}(z)G_e - G_e T_e + G_p T_{pr})^2}{3G_a} + \frac{(2t_{180}(z)G_e - G_e T_e + G_p T_{pr})^2}{3G_a} + \frac{(2t_{180}(z)G_e - G_e T_e + G_p T_{pr})^3}{3G_a} + \frac{(2t_{180}(z)G_e - G_e T_e + G_p T_{pr})^3}{3G_a} + \frac{(2t_{180}(z)G_e - G_e T_e + G_p T_{pr})^2}{3G_a} + \frac{(2t_{180}(z)G_e - G_e T_e + G_p T_{pr})^3}{3G_a} + \frac{(2t_{18$$

(A5)

APPENDIX B: Experimental validation of exact vs. PGSE-based b-value effects in dSPEN sequence

Figure B.1 assesses experimentally the efficiency of all the DW pulse sequences variants explored, with a series of single-scan 2D MRI comparisons conducted for water samples. Diffusion-sensitizing gradients were applied separately along the readout, the low-bandwidth (phase-encode or SPEN) and the slice-selection directions, such that the final ADC map is a geometric mean of the three directions. In all cases, a high consistency is evidenced by the diffusion coefficient values arising over the entire phantom sample for all five sequences –regardless of the diffusion-gradient's encoding axis– once the correct *b*-values have been computed and are taken into account for computing the corresponding maps.



Figure B1. ADC maps derived from b^{Exact_-} (top) and from b^{PGSE} -values (bottom) for a neat H₂O sample scanned by the five pulse sequences described in Fig. 2. PGSE parameters were δ = 3ms, Δ = 14ms (except for Fig. 1D where Δ = 11.4ms), $0.8 \le G_d \le 3.1$ G/cm gradients applied along all three measurements directions (readout, PE/SPEN and slice-selection). The maps represent an average of these three directions. The mean diffusion values arising over the entire phantom from b^{PGSE_-} / b^{Exact} -values are: 2.16±0.03 / 1.9±0.03 (A); 2.15±0.03 / 1.99±0.03 (B); 2.03±0.03 / 1.97±0.03 (C); 1.96.03±0.03 / 1.96±0.03 (D); 1.99±0.03 / 1.98±0.03 (E) (all values x10⁻³ mm²/s). Minor features disturbing the flatness of the images reflect artifacts arising from the various ultrafast imaging modalities. Common scanning parameters: square FOV = 30x30 mm², nominal resolution = 0.4×0.4 mm², 2 mm slice. Other gradient and timing values: (A) T₉₀=T₁₈₀=2ms, T_a=21ms, total scan duration = 51.5ms; (B) T₉₀=T_a=21ms, G_e=1.2 G/cm, G_a=4.5 G/cm, total scan duration = 51ms; (D) T₉₀=2ms, T₁₈₀=T_a/2=10.5ms, G_e=0.8 G/cm, G_a=3 G/cm, total scan duration = 65ms; (E) T₉₀=2ms, T₁₈₀=T_a/2=10.5ms, G_e=0.8 G/cm, G_a=3 G/cm, total scan duration = 62ms.