### **Supplemental Material**

#### General expressions for $R_{1\rho}$ relaxation for *N*-site chemical exchange and the special case of linear chains

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### I. Reduction of the $12 \times 12$ linear four-site mechanism evolution matrix to a $3 \times 3$ matrix

The evolution matrix for the linear four-site exchange mechanism is:

$$\mathbf{L} + \mathbf{K} = \begin{pmatrix} \mathbf{L}_{A} & 0 & 0 & 0 \\ 0 & \mathbf{L}_{B} & 0 & 0 \\ 0 & 0 & \mathbf{L}_{C} & 0 \\ 0 & 0 & 0 & \mathbf{L}_{D} \end{pmatrix} + \begin{pmatrix} -k_{12} & k_{21} & 0 & 0 \\ k_{12} & -k_{21} - k_{13} & k_{31} & 0 \\ 0 & k_{13} & -k_{31} - k_{34} & k_{43} \\ 0 & 0 & k_{34} & -k_{43} \end{pmatrix} \otimes \mathbf{I}$$
(S1)

The first application of the matrix determinant lemma yields:

$$\begin{split} 0 = |\mathbf{L} + \mathbf{K} - \lambda \mathbf{I}| \\ = |\mathbf{L}_{A} - \lambda \mathbf{I}| |\mathbf{L}_{B} - \lambda \mathbf{I}| |\mathbf{L}_{C} - \lambda \mathbf{I}| |\mathbf{L}_{D} - \lambda \mathbf{I}| \\ \times |\mathbf{I} + \begin{pmatrix} -k_{12} & k_{21} & 0 & 0 \\ 0 & -k_{13} & k_{31} & 0 \\ 0 & 0 & -k_{34} & k_{43} \end{pmatrix} \begin{pmatrix} (\mathbf{L}_{A} - \lambda \mathbf{I})^{-1} & 0 & 0 & 0 \\ 0 & (\mathbf{L}_{C} - \lambda \mathbf{I})^{-1} & 0 & 0 \\ 0 & 0 & (\mathbf{L}_{C} - \lambda \mathbf{I})^{-1} & 0 \\ 0 & 0 & 0 & (\mathbf{L}_{D} - \lambda \mathbf{I})^{-1} \end{pmatrix} | \\ = |\mathbf{L}_{A} - \lambda \mathbf{I}| |\mathbf{L}_{B} - \lambda \mathbf{I}| |\mathbf{L}_{C} - \lambda \mathbf{I}| |\mathbf{L}_{D} - \lambda \mathbf{I}| \\ \times \left[ \begin{pmatrix} \mathbf{I} - k_{12} (\mathbf{L}_{A} - \lambda \mathbf{I})^{-1} & k_{21} (\mathbf{L}_{A} - \lambda \mathbf{I})^{-1} & k_{21} (\mathbf{L}_{B} - \lambda \mathbf{I})^{-1} & 0 \\ 0 & 0 & 0 & (\mathbf{L}_{D} - \lambda \mathbf{I})^{-1} \end{pmatrix} \right] \\ = |\mathbf{L}_{A} - \lambda \mathbf{I}| |\mathbf{L}_{B} - \lambda \mathbf{I}| |\mathbf{L}_{C} - \lambda \mathbf{I}| |\mathbf{L}_{D} - \lambda \mathbf{I}| \\ & 0 & k_{34} (\mathbf{L}_{C} - \lambda \mathbf{I})^{-1} & -k_{33} (\mathbf{L}_{C} - \lambda \mathbf{I})^{-1} & k_{31} (\mathbf{L}_{C} - \lambda \mathbf{I})^{-1} \\ & 0 & k_{34} (\mathbf{L}_{C} - \lambda \mathbf{I})^{-1} & -k_{43} (\mathbf{L}_{C} - \lambda \mathbf{I})^{-1} - k_{43} (\mathbf{L}_{D} - \lambda \mathbf{I})^{-1} \\ & = |\mathbf{L}_{A} - \lambda \mathbf{I}| |\mathbf{L}_{B} - \lambda \mathbf{I}| |\mathbf{L}_{C} - \lambda \mathbf{I}| |\mathbf{L}_{D} - \lambda \mathbf{I}| \\ & = |\mathbf{L}_{A} - \lambda \mathbf{I}| |\mathbf{L}_{B} - \lambda \mathbf{I}| |\mathbf{L}_{C} - \lambda \mathbf{I}| |\mathbf{L}_{D} - \lambda \mathbf{I}| \\ & + \begin{pmatrix} \mathbf{I} - k_{12} (\mathbf{L}_{A} - \lambda \mathbf{I})^{-1} & 0 & 0 \\ 0 & \mathbf{I} & 0 \\ 0 & 0 & \mathbf{I} - k_{43} (\mathbf{L}_{D} - \lambda \mathbf{I}|^{-1} \end{pmatrix} \\ & + \begin{pmatrix} \mathbf{I} - k_{21} (\mathbf{L}_{B} - \lambda \mathbf{I}) |\mathbf{L}_{D} - \lambda \mathbf{I}| \\ & + \begin{pmatrix} \mathbf{I} - k_{21} (\mathbf{L}_{B} - \lambda \mathbf{I})^{-1} & k_{21} (\mathbf{L}_{B} - \lambda \mathbf{I})^{-1} \\ 0 & k_{34} (\mathbf{L}_{C} - \lambda \mathbf{I})^{-1} & -k_{34} (\mathbf{L}_{C} - \lambda \mathbf{I})^{-1} \end{pmatrix} \\ \end{pmatrix}$$
(S2)

The second application of the matrix determinant lemma yields:

$$0 = |\mathbf{L}_{A} - \lambda \mathbf{I}||\mathbf{L}_{B} - \lambda \mathbf{I}||\mathbf{L}_{C} - \lambda \mathbf{I}||\mathbf{L}_{D} - \lambda \mathbf{I}||\mathbf{I} - k_{12}(\mathbf{L}_{A} - \lambda \mathbf{I})^{-1}||\mathbf{I} - k_{43}(\mathbf{L}_{D} - \lambda \mathbf{I})^{-1}||$$

$$\times \left| \mathbf{I} + \left( -(\mathbf{L}_{B} - \lambda \mathbf{I})^{-1} - (\mathbf{L}_{B} - \lambda \mathbf{I})^{-1} - (\mathbf{L}_{C} - \lambda \mathbf{I})^{-1} - (\mathbf{L}_{C} - \lambda \mathbf{I})^{-1} - (\mathbf{L}_{C} - \lambda \mathbf{I})^{-1} \right) \right| \left| \mathbf{I} - k_{43}(\mathbf{L}_{D} - \lambda \mathbf{I})^{-1} - (\mathbf{L}_{C} - \lambda \mathbf{I})^{-1} - (\mathbf{L}_{C}$$

The third and last application of the matrix determinant lemma yields:

$$0 = |\mathbf{L}_{A} - \lambda \mathbf{I}||\mathbf{L}_{B} - \lambda \mathbf{I}||\mathbf{L}_{C} - \lambda \mathbf{I}||\mathbf{L}_{D} - \lambda \mathbf{I}||\mathbf{I} - k_{12}(\mathbf{L}_{A} - \lambda \mathbf{I})^{-1}||\mathbf{I} - k_{43}(\mathbf{L}_{D} - \lambda \mathbf{I})^{-1}||$$

$$\times |\mathbf{I} - k_{21}(\mathbf{L}_{B} - \lambda \mathbf{I})^{-1}(\mathbf{I} - k_{12}(\mathbf{L}_{A} - \lambda \mathbf{I})^{-1})^{-1}||\mathbf{I} - k_{34}(\mathbf{L}_{C} - \lambda \mathbf{I})^{-1}(\mathbf{I} - k_{43}(\mathbf{L}_{D} - \lambda \mathbf{I})^{-1})^{-1}||$$

$$\times |\mathbf{I} - (k_{13} - k_{31})||\mathbf{L}_{B} - \lambda \mathbf{I}||\mathbf{L}_{B} - \lambda \mathbf{I}|^{-1}(\mathbf{I} - k_{12}(\mathbf{L}_{A} - \lambda \mathbf{I})^{-1})^{-1}||\mathbf{I} - k_{34}(\mathbf{L}_{C} - \lambda \mathbf{I})^{-1}||^{-1}||$$

$$\times |\mathbf{L}_{B} - \lambda \mathbf{I}||\mathbf{L}_{C} - \lambda \mathbf{I}||\mathbf{L}_{D} - \lambda \mathbf{I}||\mathbf{I} - k_{12}(\mathbf{L}_{B} - \lambda \mathbf{I})^{-1}||\mathbf{I} - k_{43}(\mathbf{L}_{D} - \lambda \mathbf{I})^{-1}||^{-1}||$$

$$= |\mathbf{L}_{A} - \lambda \mathbf{I}||\mathbf{L}_{B} - \lambda \mathbf{I}||\mathbf{L}_{C} - \lambda \mathbf{I}||\mathbf{I} - k_{12}(\mathbf{L}_{B} - \lambda \mathbf{I})^{-1}||\mathbf{I} - k_{43}(\mathbf{L}_{D} - \lambda \mathbf{I})^{-1}|||$$

$$\times |\mathbf{I} - k_{21}(\mathbf{L}_{B} - \lambda \mathbf{I})||\mathbf{L}_{C} - \lambda \mathbf{I}||\mathbf{I} - k_{12}(\mathbf{L}_{B} - \lambda \mathbf{I})^{-1}|||\mathbf{I} - k_{43}(\mathbf{L}_{D} - \lambda \mathbf{I})^{-1}|||$$

$$\times \left| \mathbf{I} - k_{13} \left( \mathbf{I} - k_{21} \left( \mathbf{L}_{B} - \lambda \mathbf{I} \right)^{-1} \left( \mathbf{I} - k_{12} \left( \mathbf{L}_{A} - \lambda \mathbf{I} \right)^{-1} \right)^{-1} \left( \mathbf{L}_{B} - \lambda \mathbf{I} \right)^{-1} \right|^{-1} \left| \mathbf{L}_{33} \left( \mathbf{I} - k_{34} \left( \mathbf{L}_{C} - \lambda \mathbf{I} \right)^{-1} \left( \mathbf{I} - k_{43} \left( \mathbf{L}_{D} - \lambda \mathbf{I} \right)^{-1} \right)^{-1} \left( \mathbf{L}_{C} - \lambda \mathbf{I} \right)^{-1} \right)^{-1} \left| \left| \left( \mathbf{L}_{A} - \lambda \mathbf{I} \right) \left( \mathbf{I} - k_{12} \left( \mathbf{L}_{A} - \lambda \mathbf{I} \right)^{-1} \right)^{-1} \left\| \left( \mathbf{L}_{D} - \lambda \mathbf{I} \right) \left( \mathbf{I} - k_{43} \left( \mathbf{L}_{D} - \lambda \mathbf{I} \right)^{-1} \right)^{-1} \right)^{-1} \right|^{-1} \right|^{-1} \left| \left| \left( \mathbf{L}_{B} - \lambda \mathbf{I} \right)^{-1} \right|^{-1} \right|^{-1} \right|^{-1} \left| \left( \mathbf{L}_{B} - \lambda \mathbf{I} \right)^{-1} \right|^{-1} \right|^{-1} \right|^{-1} \left| \left( \mathbf{L}_{B} - \lambda \mathbf{I} \right)^{-1} \right|^{-1} \left| \left( \mathbf{L}_{B} - \lambda \mathbf{I} \right)^{-1} \right|^{-1} \right|^{-1} \left| \left( \mathbf{L}_{B} - \lambda \mathbf{I} \right)^{-1} \right|^{-1} \right|^{-1} \left| \left( \mathbf{L}_{B} - \lambda \mathbf{I} \right)^{-1} \right|^{-1} \right|^{-1} \left| \left( \mathbf{L}_{B} - \lambda \mathbf{I} \right)^{-1} \right|^{-1} \left| \left( \mathbf{L}_{B} - \lambda \mathbf{I} \right)^{-1} \right|^{-1} \right|^{-1} \left| \left( \mathbf{L}_{B} - \lambda \mathbf{I} \right)^{-1} \right|^{-1} \right|^{-1} \left| \left( \mathbf{L}_{B} - \lambda \mathbf{I} \right)^{-1} \right|^{-1} \right|^{-1} \left| \left( \mathbf{L}_{B} - \lambda \mathbf{I} \right)^{-1} \right|^{-1} \right|^{-1} \left| \left( \mathbf{L}_{B} - \lambda \mathbf{I} \right)^{-1} \right|^{-1} \left| \left( \mathbf{L}_{B} - \lambda \mathbf{I} \right)^{-1} \right|^{-1} \right|^{-1} \left| \left( \mathbf{L}_{B} - \lambda \mathbf{I} \right)^{-1} \right|^{-1} \right|^{-1} \left| \left( \mathbf{L}_{B} - \lambda \mathbf{I} \right)^{-1} \right|^{-1} \right|^{-1} \left| \left( \mathbf{L}_{B} - \lambda \mathbf{I} \right)^{-1} \right|^{-1} \right|^{-1} \left| \left( \mathbf{L}_{B} - \lambda \mathbf{I} \right)^{-1} \right|^{-1} \left| \left( \mathbf{L}_{B} - \lambda \mathbf{I} \right)^{-1} \right|^{-1} \right|^{-1} \left| \left( \mathbf{L}_{B} - \lambda \mathbf{I} \right)^{-1} \right|^{-1} \right|^{-1} \left| \left( \mathbf{L}_{B} - \lambda \mathbf{I} \right)^{-1} \right|^{-1} \left| \left( \mathbf{L}_{B} - \lambda \mathbf{I} \right)^{-1} \right|^{-1} \right|^{-1} \left| \left( \mathbf{L}_{B} - \lambda \mathbf{I} \right)^{-1} \right|^{-1} \left$$

Expanding the above expression and collecting powers of  $\lambda$  yields:

$$0 = \mathbf{L}'_{A}\mathbf{L}'_{B}\mathbf{L}'_{C}\mathbf{L}'_{D} - k_{34}k_{43}\mathbf{L}'_{A}\mathbf{L}'_{B} - k_{13}k_{31}\mathbf{L}'_{C}\mathbf{L}'_{D} + k_{12}k_{21}k_{34}k_{43} -\lambda \left\{ \begin{aligned} \mathbf{L}'_{A}\mathbf{L}'_{B}\mathbf{L}'_{C} + \mathbf{L}'_{A}\mathbf{L}'_{B}\mathbf{L}'_{D} + \mathbf{L}'_{A}\mathbf{L}'_{C}\mathbf{L}'_{D} + \mathbf{L}'_{B}\mathbf{L}'_{C}\mathbf{L}'_{D} \\ -(k_{13}k_{31} + k_{34}k_{43})\mathbf{L}'_{A} - k_{34}k_{43}\mathbf{L}'_{B} - k_{13}k_{31}\mathbf{L}'_{C} - (k_{12}k_{21} + k_{13}k_{31})\mathbf{L}'_{D} \end{aligned} \right\} +\lambda^{2} \left\{ \sum_{\substack{i,j\\j>i}} \mathbf{L}'_{i}\mathbf{L}'_{j} - k_{12}k_{21} - k_{13}k_{31} - k_{34}k_{43} \\ +\lambda^{3} \left\{ \mathbf{L}'_{A} + \mathbf{L}'_{B} + \mathbf{L}'_{C} + \mathbf{L}'_{D} \right\} + \lambda^{4} \end{aligned}$$
(S5)

Keeping only the three lowest order terms in  $\lambda$  yields Eq. (42).

#### II. Woodbury approximation for the triangular three-site scheme

In the following, "first" refers to the first order approximation, "second" to the second order approximation. The first order approximation of  $R_{ex}$  in a triangular three-site exchange situation can be written as (Eq. 48)

$$R_{ex}^{triang,first} = R_{ex}^{linear,first} \left[ \frac{1}{1 + \sin^2 \theta R_{ex}^{linear,first} Tr\{\mathbf{Z}\}} \right]$$
(S6)

The inverse of  $R_{ex}$  is:

$$\frac{1}{R_{ex}^{triang,first}} = \frac{1}{R_{ex}^{linear,first}} \left[ 1 + \sin^2 \theta R_{ex}^{linear,first} Tr \left\{ \mathbf{Z} \right\} \right]$$
(S7)

The absolute value of differences of reciprocal  $R_{ex}$  values can be used to obtain a new approximation for  $R_{ex}^{triang,exact}$ , which we call  $R_{ex}^{triang,new}$ :

$$\left|\frac{1}{R_{ex}^{triang,new}} - \frac{1}{R_{ex}^{triang,exact}}\right| < \left|\frac{1}{R_{ex}^{triang,first}} - \frac{1}{R_{ex}^{triang,exact}}\right|$$

$$\Leftrightarrow \left|\frac{1}{R_{a}}\left[1 + \sin^{2}\theta R_{b}Tr\left\{\mathbf{Z}\right\}\right] - \frac{1}{R_{ex}^{triang,exact}}\right| < \left|\frac{1}{R_{ex}^{triang,first}} - \frac{1}{R_{ex}^{triang,exact}}\right|$$
(S8)

in which  $R_a$  and  $R_b$  must be determined and in which vertical bars denote absolute values. The two cases in which  $R_{ex}^{triang,new}$  is smaller or larger than  $R_{ex}^{triang,exact}$  must be considered separately. Note that  $R_{ex}^{triang,exact} > R_{ex}^{triang,first}$  and  $R_{ex}^{linear,exact} > R_{ex}^{linear,first}$ .

In the first case, 
$$R_{ex}^{triang,new} < R_{ex}^{triang,exact}$$
 and  $R_{ex}^{triang,first} < R_{ex}^{triang,exact}$ :  

$$\frac{1}{R_a} \left[ 1 + \sin^2 \theta R_b Tr \left\{ \mathbf{Z} \right\} \right] - \frac{1}{R_{ex}^{triang,exact}} < \frac{1}{R_{ex}^{triang,first}} - \frac{1}{R_{ex}^{triang,exact}}$$

$$\Leftrightarrow \frac{1}{R_a} \left[ 1 + \sin^2 \theta R_b Tr \left\{ \mathbf{Z} \right\} \right] < \frac{1}{R_{ex}^{triang,first}}$$
(S9)

Setting  $R_a$  to  $R_{ex}^{linear, exact}$  and  $R_b$  to  $R_{ex}^{linear, first}$  yields

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$$\frac{1}{R_{ex}^{linear,exact}} < \frac{1}{R_{ex}^{linear,first}} \Leftrightarrow R_{ex}^{linear,exact} > R_{ex}^{linear,first}$$
(S10)

which is true as noted above. Setting both  $R_a$  and  $R_b$  to  $R_{ex}^{linear,exact}$  usually violates the condition  $R_{ex}^{triang,new} < R_{ex}^{triang,exact}$ , as shown empirically.

In the second case, 
$$R_{ex}^{triang,new} > R_{ex}^{triang,exact}$$
 and  $R_{ex}^{triang,first} < R_{ex}^{triang,exact}$  :  

$$-\frac{1}{R_a} \left[ 1 + \sin^2 \theta R_b Tr \left\{ \mathbf{Z} \right\} \right] + \frac{1}{R_{ex}^{triang,exact}} < \frac{1}{R_{ex}^{triang,first}} - \frac{1}{R_{ex}^{triang,exact}}$$

$$\Leftrightarrow \frac{2}{R_{ex}^{triang,exact}} - \frac{1}{R_{ex}^{triang,first}} < \frac{1}{R_a} \left[ 1 + \sin^2 \theta R_b Tr \left\{ \mathbf{Z} \right\} \right]$$

$$\Leftrightarrow \frac{2}{R_{ex}^{triang,exact}} - \frac{1}{R_{ex}^{triang,first}} < \frac{1}{R_a} \left[ 1 + \frac{R_b}{R_a} \left[ \frac{1}{R_{ex}^{triang,first}} - \frac{1}{R_{ex}^{triang,first}} \right]$$
(S11)

Again, setting  $R_a = R_{ex}^{linear, exact}$  and  $R_b = R_{ex}^{linear, first}$  yields:

$$\frac{2}{R_{ex}^{triang,exact}} - \frac{1}{R_{ex}^{triang,first}} < \frac{1}{R_{ex}^{linear,exact}} + \frac{R_{ex}^{linear,first}}{R_{ex}^{linear,exact}R_{ex}^{triang,first}} - \frac{1}{R_{ex}^{linear,exact}}$$

$$\Leftrightarrow \frac{2}{R_{ex}^{triang,exact}} < \frac{R_{ex}^{linear,first}}{R_{ex}^{linear,exact}R_{ex}^{triang,first}} + \frac{1}{R_{ex}^{triang,first}}$$

$$\Leftrightarrow \frac{R_{ex}^{triang,first}}{R_{ex}^{triang,exact}} < \frac{R_{ex}^{linear,first}}{2R_{ex}^{linear,exact}} + \frac{1}{2}$$
(S12)

which establishes a boundary condition that is frequently, but not always met. In contrast, setting  $a = R_{ex}^{linear,exact}$  and  $b = R_{ex}^{linear,exact}$  lead to a boundary condition, which is, as empirically shown, usually not met:

$$\frac{2}{R_{ex}^{triang,exact}} - \frac{1}{R_{ex}^{triang,first}} < \frac{1}{R_{ex}^{linear,exact}} + \frac{1}{R_{ex}^{triang,first}} - \frac{1}{R_{ex}^{linear,first}}$$

$$\Leftrightarrow \frac{2}{R_{ex}^{triang,exact}} + \frac{1}{R_{ex}^{linear,first}} < \frac{2}{R_{ex}^{triang,first}} + \frac{1}{R_{ex}^{linear,exact}}$$

$$\Leftrightarrow 2\left(\frac{1}{R_{ex}^{triang,exact}} - \frac{1}{R_{ex}^{triang,first}}\right) > \frac{1}{R_{ex}^{linear,exact}} - \frac{1}{R_{ex}^{linear,first}}$$
(S13)

The new approximation for the triangular three-site scheme performs better than the first-order approximation if the boundary condition Eq. S12 is met.

As a result of the above considerations, the new approximation for  $R_{ex}$  is:

$$R_{ex}^{triang,new} = R_{ex}^{linear,exact} \left[ \frac{1}{1 + \sin^2 \theta R_{ex}^{linear,first} Tr\{\mathbf{Z}\}} \right]$$
(S14)

and  $R_{ex}^{linear,exact}$  can be substituted by  $R_{ex}^{linear,second}$ , giving the Woodbury approximation:

$$R_{ex}^{triang,Woodbury} = R_{ex}^{linear,second} \left[ \frac{1}{1 + \sin^2 \theta R_{ex}^{linear,first} Tr\{\mathbf{Z}\}} \right]$$
(S15)

### **III.** Fitting three-state models

Aspects of using the above expressions for fitting to experimental data are described below for three-state models; extension to more complex topologies are straightforward. The exchange contribution to relaxation is

$$-\lambda = \sin^2 \theta R_{ex} = -1/Tr\left\{ \left( \mathbf{L} + \mathbf{K} \right)^{-1} \right\}$$
(S16)

and for a triangular three-state model,

$$\mathbf{L} = \begin{bmatrix} \mathbf{L}_{A} & \mathbf{0} & \mathbf{0} \\ \mathbf{0} & \mathbf{L}_{B} & \mathbf{0} \\ \mathbf{0} & \mathbf{0} & \mathbf{L}_{C} \end{bmatrix}$$
(S17)

in which

$$\begin{split} L_{A} &= \begin{bmatrix} 0 & -\delta_{A} & 0 \\ \delta_{A} & 0 & -\omega_{1} \\ 0 & \omega_{1} & 0 \end{bmatrix} \\ &= \begin{bmatrix} 0 & -\Delta\Omega + p_{B}\Omega_{BA} + (1 - p_{A} - p_{B})\Omega_{CA} & 0 \\ \Delta\Omega - p_{B}\Omega_{BA} - (1 - p_{A} - p_{B})\Omega_{CA} & 0 & -\omega_{1} \\ 0 & \omega_{1} & 0 \end{bmatrix} \\ L_{B} &= \begin{bmatrix} 0 & -\delta_{B} & 0 \\ \delta_{B} & 0 & -\omega_{1} \\ 0 & \omega_{1} & 0 \end{bmatrix} \\ &= \begin{bmatrix} 0 & -\delta_{B} & 0 \\ \delta_{B} & 0 & -\omega_{1} \\ 0 & \omega_{1} & 0 \end{bmatrix} \\ L_{C} &= \begin{bmatrix} 0 & -\Delta\Omega - (1 - p_{B})\Omega_{BA} + (1 - p_{A} - p_{B})\Omega_{CA} & 0 \\ \Delta\Omega + (1 - p_{B})\Omega_{BA} - (1 - p_{A} - p_{B})\Omega_{CA} & 0 & -\omega_{1} \\ 0 & \omega_{1} & 0 \end{bmatrix} \\ L_{C} &= \begin{bmatrix} 0 & -\delta_{C} & 0 \\ \delta_{C} & 0 & -\omega_{1} \\ 0 & \omega_{1} & 0 \end{bmatrix} \\ &= \begin{bmatrix} 0 & -\Delta\Omega + p_{B}\Omega_{BA} - (p_{A} + p_{B})\Omega_{CA} & 0 \\ \Delta\Omega - p_{B}\Omega_{BA} + (p_{A} + p_{B})\Omega_{CA} & 0 & -\omega_{1} \\ 0 & \omega_{1} & 0 \end{bmatrix} \end{split}$$

$$\mathbf{K} = \begin{bmatrix} -k_{12} - k_{13} & k_{21} & k_{31} \\ k_{12} & -k_{21} - k_{23} & k_{32} \\ k_{13} & k_{23} & -k_{31} - k_{32} \end{bmatrix} \otimes \mathbf{I}$$

$$= \begin{bmatrix} -\frac{p_B k_{AB}}{p_A + p_B} - \frac{(1 - p_A - p_B) k_{AC}}{1 - p_B} & \frac{p_A k_{AB}}{p_A + p_B} & \frac{p_A k_{AC}}{1 - p_B} \\ \frac{p_B k_{AB}}{p_A + p_B} - \frac{-\frac{p_A k_{AB}}{p_A + p_B} - \frac{(1 - p_A - p_B) k_{BC}}{1 - p_A} & \frac{p_B k_{BC}}{1 - p_A} \\ \frac{(1 - p_A - p_B) k_{AC}}{1 - p_B} & \frac{(1 - p_A - p_B) k_{BC}}{1 - p_A} & -\frac{p_A k_{AC}}{1 - p_B} \end{bmatrix} \otimes \mathbf{I}$$
(S19)

 $\Omega_{mA} = \Omega_m - \Omega_A, \ k_{AB} = k_{12} + k_{21}, \ k_{AC} = k_{13} + k_{31}, \ \text{and} \ k_{BC} = k_{23} + k_{32}.$ 

In the above equations, the known quantities are  $\Delta\Omega$  and  $\omega_1$ ; the parameters to be optimized are  $p_A$ ,  $p_B$ ,  $\Omega_{AB}$ ,  $\Omega_{AC}$ ,  $k_{AB}$ ,  $k_{AC}$ , and  $k_{BC}$ . Use of these variables incorporates the constraints that  $p_A + p_B + p_C = 1$  and that the change in free-energy around the reaction cycle must be zero (which constrains  $k_{23}/k_{32}$  given the other rate constants). The exchange matrix for the linear three-site model in which A is the middle state (A exchanges with B and C) is obtained by setting  $k_{BC} = 0$ . The exchange matrix for the linear three-site model in which A is the end state (A exchanges with B and B exchanges with C) is obtained by setting  $k_{AC} = 0$ . In practice, independent estimates of some of these parameters, such as limiting chemical shifts, may be necessary to reduce the complexity of data fitting. The derivative of  $\sin^2\theta R_{ex}$  with respect to a parameter x is:

$$\frac{d\left(\sin^{2}\theta R_{ex}\right)}{dx} = \frac{Tr\left\{\left(\mathbf{L} + \mathbf{K}\right)^{-2} \frac{d\left(\mathbf{L} + \mathbf{K}\right)}{dx}\right\}}{Tr^{2}\left\{\left(\mathbf{L} + \mathbf{K}\right)^{-1}\right\}}$$
(S20)

which is needed for gradient-based optimization methods. The derivative on the right-hand-side of Eq. S20 is obtained analytically from Eqs. S16-S19. For example,

$$\frac{d(\mathbf{L}+\mathbf{K})}{d\Omega_{BA}} = \begin{bmatrix} \frac{d\mathbf{L}_{A}}{d\Omega_{BA}} & \mathbf{0} & \mathbf{0} \\ \mathbf{0} & \frac{d\mathbf{L}_{B}}{d\Omega_{BA}} & \mathbf{0} \\ \mathbf{0} & \mathbf{0} & \frac{d\mathbf{L}_{C}}{d\Omega_{BA}} \end{bmatrix}$$
(S21)  
$$\frac{d\mathbf{L}_{A}}{d\Omega_{BA}} = \begin{bmatrix} \mathbf{0} & -p_{B} & \mathbf{0} \\ p_{B} & \mathbf{0} & \mathbf{0} \\ \mathbf{0} & \mathbf{0} & \mathbf{0} \end{bmatrix}$$
$$\frac{d\mathbf{L}_{B}}{d\Omega_{BA}} = \begin{bmatrix} \mathbf{0} & (1-p_{B}) & \mathbf{0} \\ -(1-p_{B}) & \mathbf{0} & \mathbf{0} \\ \mathbf{0} & \mathbf{0} & \mathbf{0} \end{bmatrix}$$
(S22)  
$$\frac{d\mathbf{L}_{C}}{d\Omega_{BA}} = \begin{bmatrix} \mathbf{0} & -p_{B} & \mathbf{0} \\ -(1-p_{B}) & \mathbf{0} & \mathbf{0} \\ \mathbf{0} & \mathbf{0} & \mathbf{0} \end{bmatrix}$$

As another example,

$$\frac{d(\mathbf{L} + \mathbf{K})}{dk_{AB}} = \frac{1}{p_A + p_B} \begin{bmatrix} -p_B & p_A & 0\\ p_B & -p_A & 0\\ 0 & 0 & 0 \end{bmatrix} \otimes \mathbf{I}$$
(S23)

Note that Equations S16 and S20 are not expanded analytically (that is, algebraic expressions for the matrix inverses are not needed); rather, numerical values are calculated for each combination of parameters directly from these equations.

#### **IV. Supplemental Figure Captions**

Figure S1.  $R_{1\rho}$  magnetization decay for two-site chemical exchange. Calculations were performed by numerical evaluation of the integrated Bloch-McConnell equation for each decay component in Eq. 6, using  $a_i = \mathbf{d} |\mathbf{u}_i\rangle \langle \mathbf{u}_i^{-1} | \mathbf{M}(0)$ . Parameters used for the calculations were  $p_A =$  $0.5, p_B = 0.5, \omega_1 = 100 \text{ s}^{-1}, \Omega_B - \Omega_A = 1000 \text{ s}^{-1}, \Delta\Omega = 500 \text{ s}^{-1}, R_1 = 1.5 \text{ s}^{-1}, R_2 = 11 \text{ s}^{-1}$  and  $k_{ex} = k_{12} + k_{21} = 1000 \text{ s}^{-1}$ . Magnetization decay for the components associated with the (a) least negative, real eigenvalue and (b) the second real eigenvalue. (c) and (d) Real-valued magnetization decays obtained by addition of pairs of components for complex conjugate eigenvalues.

Figure S2. Offset dependence of the  $R_{ex}$  contribution to  $R_{1\rho}$  for three-site exchange in the strong field limit. (Solid) Numerical calculation of  $R_{ex} = -\lambda/\sin^2\theta$  from the least negative real eigenvalue of the 9 × 9 evolution matrix, (dashed) calculation from the first-order approximation from Eqs. 11 , and (dotted) calculation from the second-order approximation from Eqs. 29 and 40. To center the graph, parameters were chosen so that  $\delta_A = \Delta\Omega$ . Parameters used for the calculations were  $p_A = 0.85$ ,  $p_B = 0.10$ ,  $p_C = 0.05$ ,  $k_{12}+k_{21} = 1550$  s<sup>-1</sup>,  $k_{13}+k_{31} = 2500$  s<sup>-1</sup>,  $\omega_1 = 1250$  s<sup>-1</sup>,  $\Omega_B - \Omega_A = 750$  s<sup>-1</sup>,  $\Omega_C - \Omega_A = -1500$  s<sup>-1</sup>.

Figure S3. Offset dependence of the  $R_{ex}$  contribution to  $R_{1\rho}$  for linear three-site exchange in which the sites exchange in different exchange regimes. (Solid) Numerical calculation of  $R_{ex} = -\lambda/\sin^2\theta$  from the least negative real eigenvalue of the 9 × 9 evolution matrix, (dashed) calculation from the first-order approximation from Eqs. 11, and (dotted) calculation from the second-order approximation from Eqs. 29 and 40. The inset exemplifies a region in which the results of the calculations differ. To center the graph, parameters were chosen so that  $\delta_A = \Delta\Omega$ . Parameters used for the calculations were (in partial analogy to Fig. 2)  $p_A = 0.95$ ,  $p_B = 0.035$ ,  $p_C=0.015$ ,  $k_{12}+k_{21} = 200 \text{ s}^{-1}$ ,  $k_{13}+k_{31} = 5000 \text{ s}^{-1}$ ,  $\omega_1 = 500 \text{ s}^{-1}$ ,  $\Omega_B - \Omega_A = 1500 \text{ s}^{-1}$ ,  $\Omega_C - \Omega_A = -3500 \text{ s}^{-1}$ .

Figure S4. Second-order approximations of  $R_{1\rho}$  linear three-site exchange at different dominant site populations. (Left) Actual  $R_{1\rho}$  plots, (Right) absolute difference between exact numerical and approximate results. (Solid) Numerical calculation of the least negative real eigenvalue of the 9 × 9 evolution matrix, (dashed) calculation from the second-order approximation from Eqs. 29 and 40, using  $R_{eff}$  from Eq. 13, (dotted) calculation from the second-order approximation from Eqs. 29 and 40, using  $R_{eff}$  from Eq. 12, (dashed-dotted) calculation from the second-order approximation from Eqs. 29 and 40, but replacing  $\mathbf{L}'_k$  with  $\mathbf{L}'_k + \mathbf{R}$  in Eq. 40 and not using  $R_{eff}$ . To center the graph, parameters were chosen so that  $\delta_A = \Delta \Omega$ . Parameters used for calculation:  $p_B$  $= p_C = (1-p_A)/2$ ,  $k_{12}+k_{21} = 100$  s<sup>-1</sup>,  $k_{13}+k_{31} = 100$  s<sup>-1</sup>,  $\omega_1 = 50$  s<sup>-1</sup>,  $\Omega_B - \Omega_A = -300$  s<sup>-1</sup>,  $\Omega_C - \Omega_A = 300$  s<sup>-1</sup>,  $R_1 = 1$  s<sup>-1</sup>,  $R_2 = 6$  s<sup>-1</sup>. (a)  $p_A = 0.33$ ; (b)  $p_A = 0.9$ .

Figure S5.  $R_{ex}$  contribution to  $R_{1\rho}$  in triangular 3-state chemical exchange. To center the graph, parameters were chosen so that  $\delta_A = \Delta \Omega$ . Parameters used for all calculations were  $p_A = 0.85$ ,  $p_B = 0.1$ ,  $p_C = 0.05$ ,  $k_{12}+k_{21} = 50 \text{ s}^{-1}$ ,  $k_{13}+k_{31} = 2000 \text{ s}^{-1}$ ,  $\Omega_B - \Omega_A = -1000 \text{ s}^{-1}$ ,  $\Omega_C - \Omega_A = 2000 \text{ s}^{-1}$ ,  $\omega_1 = 500 \text{ s}^{-1}$ . Left: Triangular 3-state chemical exchange scheme and approximations of the  $R_{ex}$  contribution. (Solid) Numerical calculation of  $R_{ex} = -\lambda/\sin^2\theta$  from the least negative real eigenvalue of the  $9 \times 9$  evolution matrix, (dashed) calculation from the first order approximation (Eq. 11), (dotted) calculation from the Woodbury approximation (Eq. 50), (dashed-dotted) calculation from a less well performing variation of the Woodbury approximation in which all first order terms are replaced by second order terms; additional parameter:  $k_{23}+k_{32} = 700 \text{ s}^{-1}$ . Right: Difference in error (approximation – numerical value) between the Woodbury approximation and the first order approximation. Values greater than 0 mean that the first order approximation is more accurate than the Woodbury approximation. The value of  $k_{23}+k_{32}$  has been varied. Smaller rate constants for  $k_{23}+k_{32}$ , relative to the other kinetic rate constants, give a more accurate Woodbury approximation.

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Figure S6. Approximations for the pseudo-kite 4-site scheme at different population levels. To center the graph, parameters were chosen so that  $\delta_A = \Delta \Omega$ . Parameters used for all calculations were  $p_B = 10/16 \cdot (1 - p_A)$ ,  $p_C = 5/16 \cdot (1 - p_A)$ ,  $p_D = 1/16 \cdot (1 - p_A)$ ,  $k_{21}+k_{12} = 140 \text{ s}^{-1}$ ,  $k_{31}+k_{13} = 350 \text{ s}^{-1}$ ,  $k_{43}+k_{34} = 700 \text{ s}^{-1}$ ,  $k_{41}+k_{14}$  (non-linear fragment) =  $350 \text{ s}^{-1}$ ,  $\Omega_B - \Omega_A = -850 \text{ s}^{-1}$ ,  $\Omega_C - \Omega_A = 2550 \text{ s}^{-1}$ ,  $\Omega_D - \Omega_A = -4250 \text{ s}^{-1}$ ,  $\omega_1 = 350 \text{ s}^{-1}$ . Left: 2-dimensional plots showing the performance of the approximations at three different population scenarios. (Solid) Numerical calculation of  $R_{ex} = -\lambda/\sin^2\theta$  from the least negative real eigenvalue of the  $12 \times 12$  evolution matrix of the pseudo-kite four-site scheme, (dashed) calculation from the first order approximation from Eq. 11, (dotted) Woodbury approximation (Eq. 50). Right: 3-dimensional plots showing the performance of the approximations at three different population scenarios ( $0.2 \le p_A \le 0.99$ ). The absolute differences between the numerical result and the first-order approximation (top) and the Woodbury approximation are plotted.

Figure S7. Use of pseudo-sites to calculate the Woodbury approximation for complex schemes. Top: A pseudo-kite 4-site scheme (grey: non-linear component of the scheme) can be collapsed to a linear three-site scheme. In this example, D\* is the pseudo-site which has the same chemical shift as site A. The population in D\* can be set low and pseudo-exchange rates between A and D can be set to very large numbers. Middle: Equivalence of the linear three-site scheme and the pseudo-kite 4-site scheme. (Solid) Numerical calculation of  $R_{ex} = -\lambda/\sin^2\theta$  from the least negative real eigenvalue of the 12 × 12 evolution matrix of the pseudo-kite four-site scheme, (dashed – overlapping with solid) numerical calculation of the 9 × 9 evolution matrix of the linear three-site scheme. To center the graph, parameters were chosen so that  $\delta_A = \Delta\Omega$ . Parameters used for all calculations were  $p_B = 0.1$ ,  $p_C = 0.05$ ,  $k_{12}+k_{21} = 10$  s<sup>-1</sup>,  $k_{12}+k_{21} = 500$  s<sup>-1</sup>,  $\omega_l = 100$  s<sup>-1</sup>,  $\Omega_B - \Omega_A = -500$  s<sup>-1</sup>,  $\Omega_C - \Omega_A = 200$  s<sup>-1</sup>. Additional parameters for three-site scheme:  $p_A = 0.84999$ ,  $p_D = 0.00001$ ,  $k_{14}+k_{41} = 10000$  s<sup>-1</sup>,  $k_{34}+k_{43} = 10000$  s<sup>-1</sup>,  $\Omega_D - \Omega_A = -500$  s<sup>-1</sup>,  $\Omega_A -$  0 s<sup>-1</sup>. Bottom: A pseudo-kite 5-site system can be used to obtain the Woodbury approximation for a star four-site system (grey: non-linear component of the scheme).





















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