Supplemental Material

General expressions for $R_{1\rho}$ relaxation for *N*-site chemical exchange **and the special case of linear chains**

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I. Reduction of the 12×12 **linear four-site mechanism evolution matrix to a** 3×3 **matrix**

The evolution matrix for the linear four-site exchange mechanism is:

$$
\mathbf{L} + \mathbf{K} = \begin{pmatrix} \mathbf{L}_A & 0 & 0 & 0 \\ 0 & \mathbf{L}_B & 0 & 0 \\ 0 & 0 & \mathbf{L}_C & 0 \\ 0 & 0 & 0 & \mathbf{L}_D \end{pmatrix} + \begin{pmatrix} -k_{12} & k_{21} & 0 & 0 \\ k_{12} & -k_{21} - k_{13} & k_{31} & 0 \\ 0 & k_{13} & -k_{31} - k_{34} & k_{43} \\ 0 & 0 & k_{34} & -k_{43} \end{pmatrix} \otimes \mathbf{I}
$$
 (S1)

The first application of the matrix determinant lemma yields:

$$
0 = |L + K - \lambda I|
$$

\n
$$
= |L_A - \lambda I||L_B - \lambda I||L_C - \lambda I||L_D - \lambda I|
$$

\n
$$
\times |I + \begin{pmatrix} -k_{12} & k_{21} & 0 & 0 \\ 0 & -k_{13} & k_{31} & 0 \\ 0 & 0 & -k_{34} & k_{33} \end{pmatrix} \begin{pmatrix} (L_A - \lambda I)^{-1} & 0 & 0 & 0 \\ 0 & (L_B - \lambda I)^{-1} & 0 & 0 \\ 0 & 0 & (L_C - \lambda I)^{-1} & 0 \\ 0 & 0 & 0 & (L_D - \lambda I)^{-1} \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ -1 & 1 & 0 \\ 0 & -1 & 1 \\ 0 & 0 & -1 \end{pmatrix}
$$

\n
$$
= |L_A - \lambda I||L_B - \lambda I||L_C - \lambda I||L_D - \lambda I|
$$

\n
$$
\times \begin{pmatrix} I - k_{12}(L_A - \lambda I)^{-1} - k_{21}(L_A - \lambda I)^{-1} & k_{21}(L_B - \lambda I)^{-1} & k_{31}(L_C - \lambda I)^{-1} & k_{31}(L_C - \lambda I)^{-1} \\ 0 & k_{13}(L_B - \lambda I)^{-1} & I - k_{13}(L_B - \lambda I)^{-1} & -k_{34}(L_C - \lambda I)^{-1} - k_{33}(L_D - \lambda I)^{-1} \end{pmatrix}
$$

\n
$$
= |L_A - \lambda I||L_B - \lambda I||L_C - \lambda I||L_D - \lambda I|
$$

\n
$$
\begin{pmatrix} I - k_{12}(L_A - \lambda I)^{-1} & 0 & 0 \\ 0 & I & 0 & 0 \\ 0 & 0 & 0 & -k_{43}(L_D - \lambda I)^{-1} \end{pmatrix}
$$

\n
$$
\times \begin{pmatrix} I - k_{21}(L_B - \lambda I)^{-1} & k_{21}(L_B - \lambda I)^{-1} & 0 \\ k_{33}(L_B - \lambda I)^{-1} & I - k_{13}(L_B - \lambda I)^{-1} & k_{31}(L_C - \lambda I)^{-1} \\ 0 & k_{34}(L_C - \lambda I)^{-1} & -k_{34}(L_C - \lambda I)^{-1
$$

The second application of the matrix determinant lemma yields:

$$
0 = |L_A - \lambda I||L_B - \lambda I||L_C - \lambda I||L_D - \lambda I||I - k_{12}(L_A - \lambda I)^{-1}||I - k_{43}(L_D - \lambda I)^{-1}|
$$

\n
$$
\times \begin{bmatrix}\nI + \begin{bmatrix}\n-(L_B - \lambda I)^{-1} & (L_B - \lambda I)^{-1} & 0 \\
0 & (L_C - \lambda I)^{-1} & -(L_C - \lambda I)^{-1}\n\end{bmatrix} \\
\times \begin{bmatrix}\nI - k_{12}(L_A - \lambda I)^{-1} & 0 & 0 \\
0 & I & 0 \\
0 & 0 & I - k_{43}(L_D - \lambda I)^{-1}\n\end{bmatrix}\n\end{bmatrix} + k_{13} - k_{14} - \lambda I||L_D - \lambda I||L_D - \lambda I||I - k_{12}(L_A - \lambda I)^{-1}||I - k_{43}(L_D - \lambda I)^{-1}|
$$

\n
$$
\times \begin{bmatrix}\nI - k_{21}(L_B - \lambda I)||L_C - \lambda I||L_D - \lambda I||I - k_{12}(L_A - \lambda I)^{-1}||I - k_{43}(L_D - \lambda I)^{-1} \\
\cdot & -k_{13}(L_C - \lambda I)^{-1} & I - k_{31}(L_C - \lambda I)^{-1} \\
\cdot & -k_{13}(L_C - \lambda I)^{-1} & I - k_{31}(L_C - \lambda I)^{-1} - k_{34}(L_C - \lambda I)^{-1} \\
\cdot & -k_{13}(L_C - \lambda I)^{-1} & I - k_{31}(L_C - \lambda I)^{-1} - k_{34}(L_C - \lambda I)^{-1} \\
\cdot & -k_{12}(L_B - \lambda I||L_D - \lambda I||I - k_{12}(L_A - \lambda I)^{-1}||I - k_{43}(L_D - \lambda I)^{-1}|\n\end{bmatrix}
$$

\n
$$
\times \begin{bmatrix}\nI - k_{21}(L_B - \lambda I) - (I - k_{12}(L_A - \lambda I)^{-1}) & 0 \\
\cdot & 0 & I - k_{34}(L_C - \lambda I)^{-1} - (I - k_{43}(L_D - \lambda I)^{-1} \\
\cdot & 0 & I - k_{43}(L_C - \lambda I)^{-1} - (I - k_{43}(L_D - \lambda I)^{-1} \\
\cdot & 0 & I - k_{43}(L_C - \lambda I)^{-1} -
$$

The third and last application of the matrix determinant lemma yields:

$$
0 = |L_A - \lambda I||L_B - \lambda I||L_D - \lambda I||L_D - \lambda I||I - k_{12}(L_A - \lambda I)^{-1}||I - k_{43}(L_D - \lambda I)^{-1}|
$$

\n
$$
\times |I - k_{21}(L_B - \lambda I)^{-1}(I - k_{12}(L_A - \lambda I)^{-1})||I - k_{34}(L_C - \lambda I)^{-1}(I - k_{43}(L_D - \lambda I)^{-1})^{-1}|
$$

\n
$$
|I - (k_{13} k_{31})\n\begin{pmatrix}\n(I - k_{21}(L_B - \lambda I)^{-1}(I - k_{12}(L_A - \lambda I)^{-1})^{-1} \\
0\n\end{pmatrix}^{-1}\n\begin{pmatrix}\nI - k_{34}(L_C - \lambda I)^{-1}(I - k_{43}(L_D - \lambda I)^{-1})^{-1} \\
0\n\end{pmatrix}^{-1}
$$

\n
$$
= |L_A - \lambda I||L_B - \lambda I||L_C - \lambda I||L_D - \lambda I||I - k_{12}(L_B - \lambda I)^{-1}||I - k_{43}(L_D - \lambda I)^{-1}|
$$

\n
$$
\times |I - k_{21}(L_B - \lambda I)^{-1}(I - k_{12}(L_A - \lambda I)^{-1})^{-1}||I - k_{34}(L_C - \lambda I)^{-1}|
$$

\n
$$
= |L_A - \lambda I||L_B - \lambda I||L_C - \lambda I||L_D - \lambda I||I - k_{12}(L_B - \lambda I)^{-1}||I - k_{43}(L_D - \lambda I)^{-1}|
$$

$$
\times \begin{bmatrix}\nI - k_{13} \left(I - k_{21} (L_{B} - \lambda I)^{-1} \left(I - k_{12} (L_{A} - \lambda I)^{-1} \right)^{-1} \right)\n- k_{31} \left(I - k_{34} (L_{C} - \lambda I)^{-1} \right)\n\end{bmatrix}\n- k_{13} \left(I - k_{14} (L_{C} - \lambda I)^{-1} \right)\n\begin{bmatrix}\nI - k_{13} (L_{C} - \lambda I)^{-1} \left(I - k_{43} (L_{D} - \lambda I)^{-1} \right)^{-1}\n\end{bmatrix}\n\begin{bmatrix}\n(L_{B} - \lambda I)\left(I - k_{12} (L_{A} - \lambda I)^{-1} \right)^{-1}\n\end{bmatrix}\n\begin{bmatrix}\n(L_{B} - \lambda I)\left(I - k_{12} (L_{A} - \lambda I)^{-1} \right)^{-1}\n\end{bmatrix}\n\begin{bmatrix}\nL_{D} - \lambda I\left(I - k_{13} (L_{D} - \lambda I)^{-1} \right)^{-1}\n\end{bmatrix}\n\begin{bmatrix}\nI - k_{13} (L_{B} - \lambda I)^{-1} \left(I - k_{12} (L_{A} - \lambda I)^{-1} \right)^{-1}\n\end{bmatrix}\n\begin{bmatrix}\nI - k_{13} (L_{B} - \lambda I)^{-1} \left(I - k_{12} (L_{A} - \lambda I)^{-1} \right)^{-1}\n\end{bmatrix}\n- k_{31} \left(I - k_{31} (L_{C} - \lambda I)^{-1} \left(I - k_{31} (L_{D} - \lambda I)^{-1} \right)^{-1}\n\end{bmatrix}\n\begin{bmatrix}\nL_{B} - \lambda I\left(I\right)\n\end{bmatrix}\n\begin{bmatrix}\nI - k_{13} (L_{C} - \lambda I)^{-1} \left(I - k_{12} (L_{A} - \lambda I)^{-1} \right)^{-1}\n\end{bmatrix}\n\begin{bmatrix}\nL_{B} - \lambda I\left(I\right)^{-1} \left(L_{B} - \lambda I\right)^{-1}\n\end{bmatrix}\n\begin{bmatrix}\nL_{A} - \lambda I - k_{13} I \left((L_{B} - \lambda I)^{-1} (L_{B} - \lambda I)^{-1} \right)^{-1}\
$$

Expanding the above expression and collecting powers of *λ* yields:

$$
0 = \mathbf{L}'_{A}\mathbf{L}'_{B}\mathbf{L}'_{C}\mathbf{L}'_{D} - k_{34}k_{43}\mathbf{L}'_{A}\mathbf{L}'_{B} - k_{13}k_{31}\mathbf{L}'_{C}\mathbf{L}'_{D} + k_{12}k_{21}k_{34}k_{43} - \lambda \begin{Bmatrix} \mathbf{L}'_{A}\mathbf{L}'_{B}\mathbf{L}'_{C} + \mathbf{L}'_{A}\mathbf{L}'_{B}\mathbf{L}'_{D} + \mathbf{L}'_{A}\mathbf{L}'_{C}\mathbf{L}'_{D} + \mathbf{L}'_{B}\mathbf{L}'_{C}\mathbf{L}'_{D} \end{Bmatrix} - (k_{13}k_{31} + k_{34}k_{43})\mathbf{L}'_{A} - k_{34}k_{43}\mathbf{L}'_{B} - k_{13}k_{31}\mathbf{L}'_{C} - (k_{12}k_{21} + k_{13}k_{31})\mathbf{L}'_{D} \end{Bmatrix}
$$
\n
$$
+ \lambda^{2} \left\{ \sum_{\substack{i,j\\j>i}} \mathbf{L}'_{i}\mathbf{L}'_{j} - k_{12}k_{21} - k_{13}k_{31} - k_{34}k_{43} \right\} + \lambda^{3} \left\{ \mathbf{L}'_{A} + \mathbf{L}'_{B} + \mathbf{L}'_{C} + \mathbf{L}'_{D} \right\} + \lambda^{4}
$$
\n(S5)

Keeping only the three lowest order terms in *λ* yields Eq. (42).

II. Woodbury approximation for the triangular three-site scheme

In the following, "first" refers to the first order approximation, "second" to the second order approximation. The first order approximation of *Rex* in a triangular three-site exchange situation can be written as (Eq. 48)

$$
R_{ex}^{triang, first} = R_{ex}^{linear, first} \left[\frac{1}{1 + \sin^2 \theta R_{ex}^{linear, first} Tr\{\mathbf{Z}\}} \right]
$$
 (S6)

The inverse of R_{ex} is:

$$
\frac{1}{R_{ex}^{triang, first}} = \frac{1}{R_{ex}^{linear, first}} \left[1 + \sin^2 \theta R_{ex}^{linear, first} Tr\left\{ \mathbf{Z} \right\} \right]
$$
(S7)

The absolute value of differences of reciprocal R_{ex} values can be used to obtain a new approximation for $R_{ex}^{triangle, exact}$, which we call $R_{ex}^{triangle, new}$.

$$
\left| \frac{1}{R_{ex}^{triang, new}} - \frac{1}{R_{ex}^{triang, exact}} \right| < \left| \frac{1}{R_{ex}^{triang, first}} - \frac{1}{R_{ex}^{triang, exact}} \right|
$$
\n
$$
\Leftrightarrow \left| \frac{1}{R_a} \left[1 + \sin^2 \theta R_b Tr \{ \mathbf{Z} \} \right] - \frac{1}{R_{ex}^{triang, exact}} \right| < \left| \frac{1}{R_{ex}^{triang, first}} - \frac{1}{R_{ex}^{triang, exact}} \right|
$$
\n(S8)

in which R_a and R_b must be determined and in which vertical bars denote absolute values. The two cases in which $R_{ex}^{triang, new}$ is smaller or larger than $R_{ex}^{triang, exact}$ must be considered separately. Note that $R_{ex}^{triang, exact} > R_{ex}^{triang, first}$ and $R_{ex}^{linear, exact} > R_{ex}^{linear, first}$.

In the first case,
$$
R_{ex}^{triangle, new} < R_{ex}^{triangle, exact}
$$
 and $R_{ex}^{triangle, first} < R_{ex}^{triangle, exact}$:
\n
$$
\frac{1}{R_a} \left[1 + \sin^2 \theta R_b Tr\{\mathbf{Z}\}\right] - \frac{1}{R_{ex}^{triangle, exact}} < \frac{1}{R_{ex}^{triangle, first}} - \frac{1}{R_{ex}^{triangle, exact}}
$$
\n
$$
\Leftrightarrow \frac{1}{R_a} \left[1 + \sin^2 \theta R_b Tr\{\mathbf{Z}\}\right] < \frac{1}{R_{ex}^{triangle, first}}
$$
\n
$$
(S9)
$$

Setting R_a to $R_{ex}^{linear, exact}$ and R_b to $R_{ex}^{linear, first}$ yields

$$
\frac{1}{R_{ex}^{\text{linear},\text{exact}}} < \frac{1}{R_{ex}^{\text{linear},\text{first}}} \Leftrightarrow R_{ex}^{\text{linear},\text{exact}} > R_{ex}^{\text{linear},\text{first}} \tag{S10}
$$

which is true as noted above. Setting both R_a and R_b to $R_{ex}^{linear, exact}$ usually violates the condition $R_{ex}^{triang,new} < R_{ex}^{triang,exact}$, as shown empirically.

In the second case,
$$
R_{ex}^{triang, new} > R_{ex}^{triang, exact}
$$
 and $R_{ex}^{triang, first} < R_{ex}^{triang, exact}$:
\n
$$
-\frac{1}{R_a} \left[1 + \sin^2 \theta R_b Tr\{\mathbf{Z}\}\right] + \frac{1}{R_{ex}^{triang, exact}} < \frac{1}{R_{ex}^{triang, first}} - \frac{1}{R_{ex}^{triang, exact}}
$$
\n
$$
\Leftrightarrow \frac{2}{R_{ex}^{triang, exact}} - \frac{1}{R_{ex}^{triang, first}} < \frac{1}{R_a} \left[1 + \sin^2 \theta R_b Tr\{\mathbf{Z}\}\right]
$$
\n
$$
\Leftrightarrow \frac{2}{R_{ex}^{triang, exact}} - \frac{1}{R_{ex}^{triang, first}} < \frac{1}{R_a} + \frac{R_b}{R_a} \left[\frac{1}{R_{ex}^{triang, first}} - \frac{1}{R_{ex}^{linear, first}}\right]
$$
\n(S11)

Again, setting $R_a = R_{ex}^{linear, exact}$ and $R_b = R_{ex}^{linear, first}$ yields:

$$
\frac{2}{R_{ex}^{triang, exact}} - \frac{1}{R_{ex}^{triang, first}} < \frac{1}{R_{ex}^{linear, exact}} + \frac{R_{ex}^{linear, first}}{R_{ex}^{linear, exact} R_{ex}^{triang, first}} - \frac{1}{R_{ex}^{linear, exact}}
$$
\n
$$
\Leftrightarrow \frac{2}{R_{ex}^{triang, exact}} < \frac{R_{ex}^{linear, first}}{R_{ex}^{linear, exact} R_{ex}^{triang, first}} + \frac{1}{R_{ex}^{triang, first}}
$$
\n
$$
\Leftrightarrow \frac{R_{ex}^{triang, first}}{R_{ex}^{triang, first}} < \frac{R_{ex}^{linear, first}}{R_{ex}^{linear, first}} + \frac{1}{R_{ex}^{triang, first}}
$$
\n
$$
\Leftrightarrow \frac{R_{ex}^{triang, first}}{R_{ex}^{triang, exact}} < \frac{R_{ex}^{linear, first}}{2R_{ex}^{linear, exact}} + \frac{1}{2}
$$
\n(512)

which establishes a boundary condition that is frequently, but not always met. In contrast, setting $a = R_{ex}^{linear, exact}$ and $b = R_{ex}^{linear, exact}$ lead to a boundary condition, which is, as empirically shown, usually not met:

$$
\frac{2}{R_{ex}^{triang, exact}} - \frac{1}{R_{ex}^{triang, first}} < \frac{1}{R_{ex}^{linear, exact}} + \frac{1}{R_{ex}^{triang, first}} - \frac{1}{R_{ex}^{linear, first}}
$$
\n
$$
\Leftrightarrow \frac{2}{R_{ex}^{triang, exact}} + \frac{1}{R_{ex}^{linear, first}} < \frac{2}{R_{ex}^{triang, first}} + \frac{1}{R_{ex}^{linear, exact}}
$$
\n
$$
\Leftrightarrow 2\left(\frac{1}{R_{ex}^{triang, exact}} - \frac{1}{R_{ex}^{triang, first}}\right) > \frac{1}{R_{ex}^{linear, exact}} - \frac{1}{R_{ex}^{linear, first}}
$$
\n(S13)

The new approximation for the triangular three-site scheme performs better than the first-order approximation if the boundary condition Eq. S12 is met.

As a result of the above considerations, the new approximation for R_{ex} is:

$$
R_{ex}^{triang,new} = R_{ex}^{linear,exact} \left[\frac{1}{1 + \sin^2 \theta R_{ex}^{linear,first} Tr\{\mathbf{Z}\}} \right]
$$
(S14)

and $R_{ex}^{linear, exact}$ can be substituted by $R_{ex}^{linear, second}$, giving the Woodbury approximation:

$$
R_{ex}^{triang, Woodbury} = R_{ex}^{linear, second} \left[\frac{1}{1 + \sin^2 \theta R_{ex}^{linear, first} Tr\{\mathbf{Z}\}} \right]
$$
(S15)

III. Fitting three-state models

Aspects of using the above expressions for fitting to experimental data are described below for three-state models; extension to more complex topologies are straightforward. The exchange contribution to relaxation is

$$
-\lambda = \sin^2 \theta R_{ex} = -1/\operatorname{Tr} \left\{ \left(\mathbf{L} + \mathbf{K} \right)^{-1} \right\} \tag{S16}
$$

and for a triangular three-state model,

$$
\mathbf{L} = \left[\begin{array}{ccc} \mathbf{L}_A & \mathbf{0} & \mathbf{0} \\ \mathbf{0} & \mathbf{L}_B & \mathbf{0} \\ \mathbf{0} & \mathbf{0} & \mathbf{L}_C \end{array} \right] \tag{S17}
$$

in which

$$
L_{A} = \begin{bmatrix} 0 & -\delta_{A} & 0 \\ \delta_{A} & 0 & -\omega_{1} \\ 0 & \omega_{1} & 0 \end{bmatrix}
$$

\n
$$
= \begin{bmatrix} 0 & -\Delta\Omega + p_{B}\Omega_{BA} + (1 - p_{A} - p_{B})\Omega_{CA} & 0 \\ \Delta\Omega - p_{B}\Omega_{BA} - (1 - p_{A} - p_{B})\Omega_{CA} & 0 & -\omega_{1} \\ 0 & 0 & \omega_{1} & 0 \end{bmatrix}
$$

\n
$$
L_{B} = \begin{bmatrix} 0 & -\delta_{B} & 0 \\ \delta_{B} & 0 & -\omega_{1} \\ 0 & \omega_{1} & 0 \end{bmatrix}
$$

\n
$$
= \begin{bmatrix} 0 & -\Delta\Omega - (1 - p_{B})\Omega_{BA} + (1 - p_{A} - p_{B})\Omega_{CA} & 0 \\ \Delta\Omega + (1 - p_{B})\Omega_{BA} - (1 - p_{A} - p_{B})\Omega_{CA} & 0 & -\omega_{1} \\ 0 & 0 & \omega_{1} & 0 \end{bmatrix}
$$

\n
$$
L_{C} = \begin{bmatrix} 0 & -\delta_{C} & 0 \\ \delta_{C} & 0 & -\omega_{1} \\ 0 & \omega_{1} & 0 \end{bmatrix}
$$

\n
$$
= \begin{bmatrix} 0 & -\delta_{C} & 0 \\ \delta_{C} & 0 & -\omega_{1} \\ 0 & \omega_{1} & 0 \end{bmatrix}
$$

\n
$$
= \begin{bmatrix} 0 & -\Delta\Omega + p_{B}\Omega_{BA} - (p_{A} + p_{B})\Omega_{CA} & 0 \\ 0 & \omega_{1} & 0 \end{bmatrix}
$$

$$
\mathbf{K} = \begin{bmatrix} -k_{12} - k_{13} & k_{21} & k_{31} \\ k_{12} & -k_{21} - k_{23} & k_{32} \\ k_{13} & k_{23} & -k_{31} - k_{32} \end{bmatrix} \otimes \mathbf{I}
$$

\n
$$
= \begin{bmatrix} -\frac{p_{B}k_{AB}}{p_{A} + p_{B}} - \frac{(1 - p_{A} - p_{B})k_{AC}}{1 - p_{B}} & \frac{p_{A}k_{AB}}{p_{A} + p_{B}} & \frac{p_{A}k_{AC}}{1 - p_{B}} \\ \frac{p_{B}k_{AB}}{p_{A} + p_{B}} & -\frac{p_{A}k_{AB}}{p_{A} + p_{B}} - \frac{(1 - p_{A} - p_{B})k_{BC}}{1 - p_{A}} & \frac{p_{B}k_{BC}}{1 - p_{A}} \end{bmatrix} \otimes \mathbf{I}
$$

\n(S19)
\n
$$
\frac{(1 - p_{A} - p_{B})k_{AC}}{1 - p_{B}} \qquad \frac{(1 - p_{A} - p_{B})k_{BC}}{1 - p_{A}} - \frac{p_{A}k_{AC}}{1 - p_{B}} - \frac{p_{B}k_{BC}}{1 - p_{B}} \end{bmatrix}
$$

 $\Omega_{mA} = \Omega_m - \Omega_A$, $k_{AB} = k_{12} + k_{21}$, $k_{AC} = k_{13} + k_{31}$, and $k_{BC} = k_{23} + k_{32}$.

In the above equations, the known quantities are $\Delta\Omega$ and ω_1 ; the parameters to be optimized are p_A , p_B , Ω_{AB} , Ω_{AC} , k_{AB} , k_{AC} , and k_{BC} . Use of these variables incorporates the constraints that $p_A + p_B + p_C = 1$ and that the change in free-energy around the reaction cycle must be zero (which constrains k_{23}/k_{32} given the other rate constants). The exchange matrix for the linear three-site model in which A is the middle state (A exchanges with B and C) is obtained by setting $k_{BC} = 0$. The exchange matrix for the linear three-site model in which A is the end state (A exchanges with B and B exchanges with C) is obtained by setting $k_{AC} = 0$. In practice, independent estimates of some of these parameters, such as limiting chemical shifts, may be necessary to reduce the complexity of data fitting. The derivative of $\sin^2\theta R_{ex}$ with respect to a parameter *x* is:

$$
\frac{d\left(\sin^2\theta R_{ex}\right)}{dx} = \frac{Tr\left\{\left(\mathbf{L} + \mathbf{K}\right)^{-2} \frac{d\left(\mathbf{L} + \mathbf{K}\right)}{dx}\right\}}{Tr^2\left\{\left(\mathbf{L} + \mathbf{K}\right)^{-1}\right\}}
$$
(S20)

which is needed for gradient-based optimization methods. The derivative on the right-hand-side of Eq. S20 is obtained analytically from Eqs. S16-S19. For example,

$$
\frac{d(\mathbf{L} + \mathbf{K})}{d\Omega_{BA}} = \begin{bmatrix} \frac{d\mathbf{L}_A}{d\Omega_{BA}} & \mathbf{0} & \mathbf{0} \\ \mathbf{0} & \frac{d\mathbf{L}_B}{d\Omega_{BA}} & \mathbf{0} \\ \mathbf{0} & \mathbf{0} & \frac{d\mathbf{L}_C}{d\Omega_{BA}} \end{bmatrix}
$$
(S21)

$$
\frac{d\mathbf{L}_A}{d\Omega_{BA}} = \begin{bmatrix} 0 & -p_B & 0 \\ p_B & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}
$$

$$
\frac{d\mathbf{L}_B}{d\Omega_{BA}} = \begin{bmatrix} 0 & (1 - p_B) & 0 \\ -(1 - p_B) & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}
$$
(S22)

$$
\frac{d\mathbf{L}_C}{d\Omega_{BA}} = \begin{bmatrix} 0 & -p_B & 0 \\ p_B & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}
$$

As another example,

$$
\frac{d(\mathbf{L} + \mathbf{K})}{dk_{AB}} = \frac{1}{p_A + p_B} \begin{bmatrix} -p_B & p_A & 0 \\ p_B & -p_A & 0 \\ 0 & 0 & 0 \end{bmatrix} \otimes \mathbf{I}
$$
(S23)

Note that Equations S16 and S20 are not expanded analytically (that is, algebraic expressions for the matrix inverses are not needed); rather, numerical values are calculated for each combination of parameters directly from these equations.

IV. Supplemental Figure Captions

Figure S1. $R_{1\rho}$ magnetization decay for two-site chemical exchange. Calculations were performed by numerical evaluation of the integrated Bloch-McConnell equation for each decay component in Eq. 6, using $a_i = d | \mathbf{u}_i \rangle \langle \mathbf{u}_i^{-1} | \mathbf{M}(0)$. Parameters used for the calculations were $p_A =$ $0.5, p_B = 0.5, \omega_1 = 100 \text{ s}^{-1}, \Omega_B - \Omega_A = 1000 \text{ s}^{-1}, \Delta \Omega = 500 \text{ s}^{-1}, R_1 = 1.5 \text{ s}^{-1}, R_2 = 11 \text{ s}^{-1}$ and $k_{ex} = k_{12} + k_{12} = 1.5 \text{ s}^{-1}$ k_{21} = 1000 s⁻¹. Magnetization decay for the components associated with the (a) least negative, real eigenvalue and (b) the second real eigenvalue. (c) and (d) Real-valued magnetization decays obtained by addition of pairs of components for complex conjugate eigenvalues.

Figure S2. Offset dependence of the R_{ex} contribution to $R_{1\rho}$ for three-site exchange in the strong field limit. (Solid) Numerical calculation of $R_{ex} = -\lambda/\sin^2\theta$ from the least negative real eigenvalue of the 9×9 evolution matrix, (dashed) calculation from the first-order approximation from Eqs. 11 , and (dotted) calculation from the second-order approximation from Eqs. 29 and 40. To center the graph, parameters were chosen so that $\delta_A = \Delta \Omega$. Parameters used for the calculations were $p_A = 0.85$, $p_B = 0.10$, $p_C = 0.05$, $k_{12} + k_{21} = 1550$ s⁻¹, $k_{13} + k_{31} = 2500$ s⁻¹, $\omega_1 = 1250$ s⁻¹, $\Omega_B - \Omega_A$ $= 750 \text{ s}^{-1}, \Omega_C - \Omega_A = -1500 \text{ s}^{-1}.$

Figure S3. Offset dependence of the R_{ex} contribution to $R_{1\rho}$ for linear three-site exchange in which the sites exchange in different exchange regimes. (Solid) Numerical calculation of $R_{ex} = \lambda/\sin^2\theta$ from the least negative real eigenvalue of the 9 × 9 evolution matrix, (dashed) calculation from the first-order approximation from Eqs. 11, and (dotted) calculation from the second-order approximation from Eqs. 29 and 40. The inset exemplifies a region in which the results of the calculations differ. To center the graph, parameters were chosen so that $\delta_A = \Delta \Omega$. Parameters used for the calculations were (in partial analogy to Fig. 2) $p_A = 0.95$, $p_B = 0.035$, $p_C = 0.015$, $k_{12} + k_{21} = 200 \text{ s}^{-1}, k_{13} + k_{31} = 5000 \text{ s}^{-1}, \omega_1 = 500 \text{ s}^{-1}, \Omega_B - \Omega_A = 1500 \text{ s}^{-1}, \Omega_C - \Omega_A = -3500 \text{ s}^{-1}.$

Figure S4. Second-order approximations of $R_{1\rho}$ linear three-site exchange at different dominant site populations. (Left) Actual $R_{1\rho}$ plots, (Right) absolute difference between exact numerical and approximate results. (Solid) Numerical calculation of the least negative real eigenvalue of the 9×9 evolution matrix, (dashed) calculation from the second-order approximation from Eqs. 29 and 40, using *Reff* from Eq. 13, (dotted) calculation from the second-order approximation from Eqs. 29 and 40, using R_{eff} from Eq. 12, (dashed-dotted) calculation from the second-order approximation from Eqs. 29 and 40, but replacing \mathbf{L}'_k with \mathbf{L}'_k + **R** in Eq. 40 and not using R_{eff} . To center the graph, parameters were chosen so that $\delta_A = \Delta \Omega$. Parameters used for calculation: p_B $=p_c = (1-p_A)/2$, $k_{12}+k_{21} = 100$ s⁻¹, $k_{13}+k_{31} = 100$ s⁻¹, $\omega_1 = 50$ s⁻¹, $\Omega_B - \Omega_A = -300$ s⁻¹, $\Omega_C - \Omega_A =$ 300 s^{-1} , $R_1 = 1$ s^{-1} , $R_2 = 6$ s^{-1} . (a) $p_A = 0.33$; (b) $p_A = 0.9$.

Figure S5. R_{ex} contribution to $R_{1\rho}$ in triangular 3-state chemical exchange. To center the graph, parameters were chosen so that $\delta_A = \Delta\Omega$. Parameters used for all calculations were $p_A = 0.85$, p_B $= 0.1, p_c = 0.05, k_{12} + k_{21} = 50 \text{ s}^{-1}, k_{13} + k_{31} = 2000 \text{ s}^{-1}, \Omega_B - \Omega_A = -1000 \text{ s}^{-1}, \Omega_C - \Omega_A = 2000 \text{ s}^{-1}, \omega_1 = 0.000 \text{ s}^{-1}$ 500 s–1 . Left: Triangular 3-state chemical exchange scheme and approximations of the *Rex* contribution. (Solid) Numerical calculation of $R_{ex} = -\lambda/\sin^2\theta$ from the least negative real eigenvalue of the 9×9 evolution matrix, (dashed) calculation from the first order approximation (Eq. 11), (dotted) calculation from the Woodbury approximation (Eq. 50), (dashed-dotted) calculation from a less well performing variation of the Woodbury approximation in which all first order terms are replaced by second order terms; additional parameter: $k_{23} + k_{32} = 700 \text{ s}^{-1}$. Right: Difference in error (approximation – numerical value) between the Woodbury approximation and the first order approximation. Values greater than 0 mean that the first order approximation is more accurate than the Woodbury approximation. The value of $k_{23}+k_{32}$ has been varied. Smaller rate constants for $k_{23}+k_{32}$, relative to the other kinetic rate constants, give a more accurate Woodbury approximation.

10/20/16

Figure S6. Approximations for the pseudo-kite 4-site scheme at different population levels. To center the graph, parameters were chosen so that $\delta_A = \Delta \Omega$. Parameters used for all calculations were $p_B = 10/16 \cdot (1 - p_A)$, $p_C = 5/16 \cdot (1 - p_A)$, $p_D = 1/16 \cdot (1 - p_A)$, $k_{21} + k_{12} = 140 \text{ s}^{-1}$, $k_{31} + k_{13} = 350 \text{ s}^{-1}$, $k_{43} + k_{34} = 700 \text{ s}^{-1}, k_{41} + k_{14} \text{ (non-linear fragment)} = 350 \text{ s}^{-1}, \Omega_B - \Omega_A = -850 \text{ s}^{-1}, \Omega_C - \Omega_A = 2550 \text{ s}^{-1}$ $\Omega_D - \Omega_A = -4250 \text{ s}^{-1}$, $\omega_1 = 350 \text{ s}^{-1}$. Left: 2-dimensional plots showing the performance of the approximations at three different population scenarios. (Solid) Numerical calculation of $R_{ex} = \lambda$ /sin² θ from the least negative real eigenvalue of the 12 × 12 evolution matrix of the pseudo-kite four-site scheme, (dashed) calculation from the first order approximation from Eq. 11, (dotted) Woodbury approximation (Eq. 50). Right: 3-dimensional plots showing the performance of the approximations at three different population scenarios ($0.2 \le p_A \le 0.99$). The absolute differences between the numerical result and the first-order approximation (top) and the Woodbury approximation are plotted.

Figure S7. Use of pseudo-sites to calculate the Woodbury approximation for complex schemes. Top: A pseudo-kite 4-site scheme (grey: non-linear component of the scheme) can be collapsed to a linear three-site scheme. In this example, D* is the pseudo-site which has the same chemical shift as site A. The population in D^* can be set low and pseudo-exchange rates between A and D can be set to very large numbers. Middle: Equivalence of the linear three-site scheme and the pseudo-kite 4-site scheme. (Solid) Numerical calculation of $R_{ex} = -\lambda/\sin^2\theta$ from the least negative real eigenvalue of the 12×12 evolution matrix of the pseudo-kite four-site scheme, (dashed – overlapping with solid) numerical calculation of the 9×9 evolution matrix of the linear threesite scheme, (dotted) Woodbury approximation (Eq. 50) of the pseudo-kite four-site scheme. To center the graph, parameters were chosen so that $\delta_A = \Delta \Omega$. Parameters used for all calculations were $p_B = 0.1$, $p_C = 0.05$, $k_{12} + k_{21} = 10 \text{ s}^{-1}$, $k_{12} + k_{21} = 500 \text{ s}^{-1}$, $\omega_I = 100 \text{ s}^{-1}$, $\Omega_B - \Omega_A = -500 \text{ s}^{-1}$, $\Omega_C \Omega_A = 200 \text{ s}^{-1}$. Additional parameters for three-site scheme: $p_A = 0.85$. Additional parameters for four-site scheme: $p_A = 0.84999$, $p_D = 0.00001$, $k_{14} + k_{41} = 10000$ s⁻¹, $k_{34} + k_{43} = 10000$ s⁻¹, $\Omega_D - \Omega_A =$

 $0 s⁻¹$. Bottom: A pseudo-kite 5-site system can be used to obtain the Woodbury approximation for a star four-site system (grey: non-linear component of the scheme).

