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Supplemental Information

Opposing Intermolecular Tuning of Ca²⁺ Affinity for Calmodulin by Neurogranin and CaMKII Peptides

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Supporting Text

I. Sample preparation

I.1 Intact Ng

Ng is a 78-residue intrinsically disordered protein with IQ motif [\(1,](#page-62-0) [2\)](#page-62-1) composed of hydrophobic and basic amino acids. Available experimental measurement for residual structure of Ng is for the full-length protein Ng from mouse [\(3\)](#page-62-2), therefore, the sequence is used in the modeling and parameterization of Ng, as shown below,

1 10 20 30 40 50 MDCCTESACSKPDDDILDIPLDDPGANAAAAKIQASFRGHMARKKIKSGE 60 70 CGRKGPGPGGPGGAGGARGGAGGGPSGD 78

Following the work by Hoffman et al. [\(4\)](#page-62-3), we divided the Ng protein into several units, including the acidic N terminal (residues 13-25), the IQ motif (residues 26-49) and the poly-Glycine C terminal (residues 50-78). The underscored sequence stands for the IQ motif which partially retains residual structure (residues 25-42) in the unbound state. To reproduce the residual structure, replica exchange method (REM) [\(5,](#page-62-4) [6\)](#page-62-5) was employed for calculation of average helicity of the residual structure and the backbone model-free order parameter S^2 .

I.2 Ng13-49 peptide

Hoffman et al. showed from calorimetry and NMR experiments that the IQ motif alone does not reproduce similar Ng-mediated affinity of Ca^{2+} for Ca^{2+} -free CaM (apoCaM), or the pattern of intermolecular interaction between Ng and apoCaM [\(4\)](#page-62-3). A combination of the acidic region in the N-terminal and the IQ motif yields the minimum composition of Ng for its function.

Therefore the Ng13-49 peptide is used to study the interaction with apoCaM (PDB ID: 1CFD) that allows the comparison of the results from computer simulations with several experimental measurements, including dissociation constant of apoCaM and Ng_{13-49} , the affinity of Ca^{2+} to CaM, and the changes in the chemical shifts of apoCaM upon Ng₁₃₋₄₉ binding [\(4\)](#page-62-3). The sequence of the Ng13-49 peptide is provided below,

20 30 40 13 DDDILDIPLDDPGANAAAAKIQASFRGHMARKKIKSG 49

II. Coarse-grained protein or peptide models

II. 1 Hamiltonian and parametrization of the coarse-grained protein models

As described above, Ng is partly ordered in solution. The fragment G25-A42 has a fraction of 22% and 28% residual structure calculated from the C_{α} and the H_α chemical shifts in the nuclear magnetic resonance (NMR) experiment [\(3\)](#page-62-2), respectively. The rest part of the peptide remains unstructured. In order to compare with results from several experiments $(3, 4)$ $(3, 4)$, peptide Ng₁₃₋₄₉ comprising the IQ domain with adjacent acidic region as well as the full sequence Ng were used in the simulations (please find the sequences in session I from the *Supporting Information*.).

We adopted a side-chain- C_{α} coarse-grained model to represent protein/peptides developed by Cheung et al [\(7\)](#page-62-6). In this model, each amino acid (except glycine, which is represented by a C_α bead) is represented by two beads: the C_α bead is located at the C_α position to represent the backbone atoms; the side-chain bead is located at the center of mass of the sidechain atoms to represent the side-chain atoms. The use of side-chain- C_{α} model enables us to capture large structural fluctuations of the proteins at a large timescale that warrants a wide search of the phase space.

The total potential energy E of Ng or Ng₁₃₋₄₉ is given by,

$$
E^{Ng} = E_{bond} + E_{angle} + E_{dihedral} + E_{chirality} + E_{elec} + E_{HB} + E_{LI}
$$
 Eq. (S1)

The bonded interactions include the bond energy (E_{bond}) , the bond-angle energy (E_{angle}) , and the dihedral-angle energy $(E_{dihedral})$ the L-isomer restraint of chirality $(E_{chirality})$. E_{bond} , E_{angle} , and *Echirality* constrain the bond length, bond angle and side-chain chirality through harmonic potentials. The formula can be found in our previous work [\(8\)](#page-62-7). The equilibrium bond length, bond angle and side-chain chirality parameters were taken from the crystal structure of apoCaM and IQ motif of Ng (PDB ID: 405E), and those of the missing segment were obtained from the structures predicted by the Sparks-X protein structure prediction server [\(9\)](#page-62-8). Since these terms are mainly constrained by chemical rules, they do not vary significantly by the conformations.

The electrostatic energy (E_{elec}) between each two beads (C_{α} or side-chain bead) *i* and *j* with partial charges is described by Debye-Hückel potential [\(10\)](#page-62-9) to include the screening effect of the electrolyte solution.

$$
E_{elec}^{ij} = \frac{q_i q_j}{4\pi \epsilon_r \epsilon_0 r_{ij}} \exp\left(-r_{ij} / \sqrt{\frac{\epsilon_r \epsilon_0 k_B T}{2e^2 I}}\right)
$$
 Eq. (S2)

where q_i (q_j) is the partial charges on bead *i* (*j*); *e* is the elementary charge (see the method of generating the partial charges in the next section II.2 from the *Supporting Information*); ϵ_r is the relative dielectric constant and is set to 80 for aqueous solution; ϵ_0 is the permittivity of the vacuum; r_{ij} is the distance between beads *i* and *j*; k_B is Boltzmann constant; T = 1.1 ε / k_B is the temperature; $I = 0.15$ M is the ionic strength.

For the backbone hydrogen bonding interactions, we adopted an angular‐dependent function that captures directional properties of backbone hydrogen bonds [\(7\)](#page-62-6),

$$
E_{HB}^{ij} = \sum_{|j-i|>3} A(\rho) \varepsilon \left[\left(\frac{\rho_{ij}^0}{r_{ij}} \right)^6 - 2 \left(\frac{\rho_{ij}^0}{r_{ij}} \right)^{12} \right]
$$
 Eq. (S3)

$$
A(\rho) = \frac{1}{\left[1 + (1 - \cos^2 \rho)(1 - \frac{\cos \rho}{\cos \rho \alpha})\right]^2}
$$
 Eq. (S4)

where the strength $\varepsilon = 0.6$ kcal/mol; r_{ij} is the distance between backbone beads *i* and *j*; and $\rho_{ij}^0 =$ 4.6 Å is the typical length of a hydrogen bond; $A(\rho)$ measures the structural alignment of two interacting strands; ρ is the pseudo dihedral angle between backbone beads of the two interacting strands [\(7\)](#page-62-6); ρ_{α} is the pseudo dihedral angle of a canonical helical turn, 0.466 rad. *A*(*ρ*) preserves the tendency to form β-strands (*ρ* = 0 or π) or α-helices (*ρ* = *ρ_α*) as it is maximized to 1 in either of the two cases.

For the non-bonded interactions (E_{LJ}) , between C_{α} bead *i* from the backbone (bb) and side-chain (sc) bead *j*, a pure repulsive interaction was considered,

$$
E_{LJ}^{sc-bb} = \varepsilon \left(\frac{\rho_{ij}^0}{r_{ij}}\right)^{12} \qquad \text{Eq. (S5)}
$$

where the strength $\varepsilon = 0.6$ kcal/mol; i $\rho_{ij}^0 = 0.9(\rho_i^0 + \rho_j^0)$, ρ_i^0 is the size of backbone bead *i*, 0.5 $σ (σ = 3.8 Å)$ and $ρ_j^0$ is the size of side-chain bead *j*, which is the van der Waals radius of the side-chain. 0.9 is a scaling factor to remove clashes between bulky side-chains. For the nonbonded interactions between side-chain bead *i* and side-chain bead *j*, we applied the Lennard-Jones potential,

$$
E_{LJ}^{sc-sc} = \varepsilon_{ij} \left[\left(\frac{\rho_{ij}^0}{r_{ij}} \right)^{12} - 2 \left(\frac{\rho_{ij}^0}{r_{ij}} \right)^6 \right]
$$
 Eq. (S6)

where ε_{ij} is solvent mediated interactions between the involved amino acids obtained from Betancourt-Thirumalai's study [\(11\)](#page-62-10). $\rho_{ij}^0 = 0.9(\rho_i^0 + \rho_j^0)$, $\rho_{i(j)}^0$ is the van der Waals radius of side-chain bead i (j). r_{ij} is the distance between beads i and j in Eqs. S3, S5 and S6. Since the electrostatic interactions were explicitly included in the Hamiltonian, the Lennard-Jones potential coefficient ε_{ij} was rescaled as described in our previous work [\(8\)](#page-62-7).

To best describe the intrinsic disorder nature of Ng or Ng13-49, we employed a backbone dihedral angle potential that is sequence-specific but independent on protein topology. For this we used the model introduced by Karanicolas and Brooks [\(12\)](#page-62-11), or the KB model. The dihedral angle between four adjacent α -carbons depends on the backbone dihedral angles of the two middle residues. Brooks' group produced a probability distribution for 400 possible ordered pairs of amino acid residues from a survey of the Protein Data Bank and thus related the probability distribution to potential energy ignoring the entropy contribution. The dihedral angle potential presents two minima corresponding to local α-helical and β-strand geometries. The statistical potential is modeled as a 4-term cosine series,

$$
E_{\text{dihedral}}^{\text{ijkl}} = \varepsilon_{\text{KB}} \sum_{n=1}^{4} K_{ij,n} \left[1 + \cos(n\varphi_{ijkl} - \delta_{jk,n}) \right] \tag{S7}
$$

where φ_{ijkl} is the dihedral angle formed by four consecutive C_α beads *i, j, k, l* with beads *j* and *k* in the middle, $K_{jk,n}$ and $\delta_{jk,n}$ are statistically determined constants. ε_{KB} is a factor to adjust the strength in relative to other interactions in the current model.

We tested a total of three types of dihedral potentials for a comparative study. The first dihedral angle potential is a sequence-based model as described above. The second dihedral angle potential is a structure-based potential, or SB model, where the structure of a specific segment of Ng obtained from the crystal structure (PDB code: 4E50 for Ng) was used as the reference (please see Table S4 for the residues in this segment). The SB model is composed of two-term cosine-series,

$$
E_{dihedral}^{ijkl} = \varepsilon_{SB} \sum_{n=1,3} k_{\varphi}^{n} \left[1 - \cos \left(n \times \left(\varphi_{ijkl} - \varphi_{ijkl}^{0} \right) \right) \right]
$$
 Eq. (S8)

where *i, j, k* and *l* are four consecutive C_α beads. φ_{ijkl} is the dihedral angle formed by those four beads. The equilibration position of the dihedral angle φ_{ijkl}^0 for this specific segment G25-A42 was taken from the crystal structure (PDB code: 4E50 for Ng). $k_{\varphi}^1 = 2k_{\varphi}^3 = 2\varepsilon$. ε_{SB} is used to adjust the barrier of the dihedral angle potential. For the unstructured segment $\varepsilon_{SB} = 0$.

The third dihedral potential is the hybrid of the two models, or Hybrid model, as shown in Table S4. We replaced the dihedral angle potential of the unstructured segment (beyond G25- A42) from the SB model with a KB model.

Different barrier values of the structure-based dihedral potential ε_{SB} and the statistical dihedral potential ϵ_{KB} were tested. Fig. S7 shows the helicity of the fragment G25-A42 decreases with temperature in all the cases. At the temperature $T = 1.1 \epsilon/k_B$, for the SB model, the helicity fell within the experimental range when $\varepsilon_{SB} = 0.3$ (Fig. S7A); for the Hybrid model, the helical fraction fell within the experimental range with when $\varepsilon_{SB} = 0.3$, and ε_{KB} did not make any influence in the range from 1.8 to 2.2 (Fig. S7B); for the KB model, the helical fraction was best matched when ε_{KB} is in the range from 1.5 to 2.2 (Fig. S7C).

To determine which best represents the properties of Ng among all the three models, we further investigated them at the residue level. Therefore, for those models with parameters that matched the overall helicity (Fig. S7), we computed the nuclear magnetic resonance (NMR) model free order parameter S^2 for the backbone beads from our coarse-grained simulations (please find the details of calculation in V.1 from the *Supporting Information*), and compared with data from the NMR experiments [\(3\)](#page-62-2). The computed S^2 (Table S5) positively correlated with the experimental values in all cases. Interestingly, for the KB model, the correlation was overall higher than the SB and Hybrid model, especially for $\varepsilon_{KB} = 1.8$. Because Ng protein is intrinsically disordered [\(13\)](#page-62-12), The KB dihedral angle potential of no bias to any specific structure enables sampling of a broad spectrum of Ng conformations, whereas the SB model and the Hybrid model of full or partial bias to a specific structure limit the flexibility to explore more conformations. Therefore, for modeling intrinsically disordered protein Ng, we adopted the sequence-based KB dihedral potential model. $\varepsilon_{KB} = 1.8$ was used in the following study of apoCaM-Ng13-49 binding simulations.

II.2 Determination of partial charges for the coarse-grained models

II.2.a Neurogranin protein

We adopted a multi-scale method developed by Cheung group [\(14\)](#page-62-13) to assign partial charges to the intact Ng protein. Firstly, we ran REMD simulations without electro-static interactions (please find the details in III.1 in the *Supporting Information*) and obtained the free energy surface $F(\Delta, \chi)$ as a function of asphericity Δ and overlap function χ . Δ measures the shape of the protein: it is like a rod or a sphere when $\Delta = 1$ or 0, respectively [\(15\)](#page-62-14). χ measures the similarity to the reference structure [\(16\)](#page-62-15). Using the free energy, we thus selected about 400-600 frames from the REMD simulations of Ng protein through importance sampling [\(17\)](#page-63-0). We reconstructed the coarse-grained structures into atomistic protein models [\(17\)](#page-63-0), used H++ server [\(18\)](#page-63-1) to predict the protonation states and computed the partial charges using the semi-empirical

quantum chemistry program MOPAC [\(19\)](#page-63-2). We collected the partial charges on the backbone (side-chain) atoms as the charge for the C_α (sidechain) bead. We then obtained the averaged partial charges on the C_{α} and side-chain beads over all the structures. We repeated the same process for all the three models shown in Table S4 and the three sets of charges we obtained were highly similar to each other (the correlation coefficients were \sim 1.00). We therefore used only one set of charges for all the three models. The charges on Ng protein are provided in Table S9.

II.2.b Apo-calmodulin and neurogranin peptide system

We followed the same procedures as for the Ng protein. The free energy surface $F(d_{COM}, Z_{inter})$ was obtained from US simulations of apoCaM and Ng₁₃₋₄₉ without electrostatic interactions by using WHAM [\(20,](#page-63-3) [21\)](#page-63-4) (please see details about the US simulations in III.2 in the *Supporting Information*). d_{COM} is the distance between center of mass of apoCaM and center of mass of $Ng₁₃₋₄₉$, and Z_{inter} is the total number of intermolecular contacts (please see the definition of a contact in V.3 from the *Supporting Information*). Two groups of partial charges were generated for apoCaM and Ng13-49 according to the experimental conditions: the fluorescence experiments for measurement of the dissociation constant of apoCaM and Ng_{13-49} were conducted at pH = 7.2 and at ionic strength $I = 0.15$ M; the NMR experiments for determining the change in chemical shifts of apoCaM upon binding Ng₁₃₋₄₉ were conducted at $pH = 6.3$ and at ionic strength I = 0.10 M. The protonation states of histidine residues were determined by the H++ server [\(18\)](#page-63-1) before using MOPAC [\(19\)](#page-63-2). The calculated partial charges of apoCaM and Ng13-49 in the above two conditions were provided in Table S10-S13*.*

III. Coarse-grained Molecular Simulations

III.1 Replica Exchange Molecular Dynamics simulations

To study dynamics of unbound Ng protein, which has a rugged energy landscape because of its intrinsic disorder, we applied replica exchange molecular dynamics (REMD) [\(5,](#page-62-4) [6\)](#page-62-5) to enhance the sampling. For the three dihedral angle models, a range of barriers of dihedral angel potential were investigated: for the SB (structure-based) model, $\varepsilon_{SB} = 0.1, 0.2, 0.3, 0.4, 0.5, 0.6$ and 0.7; for the Hybrid model, $\varepsilon_{SB} = 0.3$, 0.4 for the part with residual structure (G25-A42), in combination with $\varepsilon_{KB} = 1.8$, 2.0 and 2.2 for the rest; for the KB model, $\varepsilon_{KB} = 1.0$, 1.5, 1.8, 2.0, 2.2, 2.5 and 3.0. For each of the models, 24 replicas were distributed at temperatures $T = 0.80$, 0.82, 0.83, 0.85, 0.87, 0.88, 0.90, 0.93, 0.97, 1.00, 1.03, 1.07, 1.10, 1.13, 1.17, 1.20, 1.23, 1.27, 1.30, 1.33, 1.40, 1.47, 1.57, 1.67 ε / k_B to produce ample exchanges between the replicas (ε is the reduced energy unit = 0.6 kcal/mol, and k_B is the Boltzmann constant). The average acceptance ratio of exchanging among the neighboring replicas ranges from 0.57 to 0.77. An exchange between neighboring replicas was attempted every 100,000 time steps. In order to explore more conformation space, in each REMD simulation, a total of 6 sets of distinct initial configurations were used, making up a total of 1200 exchange attempts for each model.

III.2 Umbrella sampling simulations

Umbrella sampling (US) method [\(22\)](#page-63-5) was used to explore the thermodynamic properties of the apoCaM-Ng system. The distance between the center of mass of apoCaM and the center of mass of Ng₁₃₋₄₉ d_{COM} was used as the reaction coordinate. d_{COM} was restrained by a harmonic force with spring constant 66.7 ε/σ^2 ($\sigma = 3.8$ Å is reduced length unit). The equilibrium positions of the harmonic force range from 0.2 σ to 20.0 σ with a bin size of 0.2 σ, making up a total of 100

windows. For each window, 10,000,000 time-steps of constrained molecular dynamic simulation were carried out.

To generate the initial structures at each window of d_{COM} for the US simulations, molecular dynamics simulations were carried out for the apoCaM-Ng13-49 complex from the bound form (PDB code: 4E50, in the coarse-grained model) at a high temperature $T = 1.33 \epsilon / k_B$ for dissociation ($\varepsilon = 0.6$ kcal/mol is reduced energy unit, and k_B is Boltzmann constant). A total of 5 sets of initial configurations for the following umbrella sampling simulations were generated from 5 dissociation trajectories.

US simulations were then performed without electrostatic interactions for generating ensemble of structures to generate ensemble of structures for partial charge determination (see 2.3 for the details). After we obtained the partial charges for apoCaM and Ng_{13-49} , US simulations were carried out with charge-charge interactions to study how the strength of nonelectrostatic intermolecular interactions influences the binding affinity and to determine the optimal strength by comparing with the experiments [\(4\)](#page-62-3). Using this strength of non-electrostatic intermolecular interactions, US simulations were carried up at appropriate experimental conditions to study the binding thermodynamic properties of apoCaM-Ng13-49 system. Each set of US simulations were performed from 5 different initial structures at the temperature $T = 1.1$ ε/k_B . We analyzed the data using Weighted Histogram Analysis Method (WHAM) [\(20,](#page-63-3) [21\)](#page-63-4).

IV. Calculation of difference in binding free energy using Jarzynski's equality from steered molecular dynamics simulations

IV.1 Selection of initial atomistic structures

holoCaM-Ng13-49 complex: We employed coarse-grained molecular simulations to efficiently sample a broad ensemble of complex structures. We used the experimental measurements as a guide to strategically select several structures from the major cluster of coarse-grained complex structures. According to the crystal structure of apoCaM and IQ motif of Ng (NgIQ) (PDB ID: 4E50), the two EF-hand motifs from cCaM are open (EF-hand angles ≈ 101) and the two EFhand motifs from nCaM are closed (EF-hand angles \approx 129) (Table S3). The EF-hand angles were computed as the angle between the vectors that define the orientation of the two helices in the EF-hand motif; the vectors were defined from the center of mass of the C_{α} beads of the first 4 residues to the center of mass of the C_{α} beads of the last 4 residues in a helix of the EF-hand motif. Moreover, NgIQ binds to cCaM in the crystal structure and the NMR experiments [\(4\)](#page-62-3) showed that Ng13-49 mainly interacts with cCaM. Therefore, we select the 4 structures from the simulations by following the criteria: (i) EF-hand motifs from cCaM must be open (EF-hand angles range from 85̊ to 105̊) and EF-hand motifs from nCaM must be closed (EF-hand angles are greater than 125°; (ii) cCaM has more interactions with Ng₁₃₋₄₉ than nCaM (Z_c – Z_n > 0). Using SCAAL method [\(17\)](#page-63-0), we reconstructed all-atomistic structures of apoCaM-Ng13-49 from the selected sidechain-C_α configurations, as shown in Fig. S8. Four Ca²⁺ ions were added and their positions were estimated as the center of mass of sidechain beads of the corresponding Ca^{2+} co-ordination residues. The Ca^{2+} positions were further optimized by energy minimization on the all-atomistic structures of $Ca^{2+}-CaM-Ng_{13-49}$ (holoCaM-Ng₁₃₋₄₉) using a gradient descent algorithm provided by the molecular dynamics software package GROMACS (version 5.0.6) [\(23\)](#page-63-6) with AMBER99SB-ILDN force-field [\(24\)](#page-63-7). The tolerance for the maximum force was set to 500 kJ / mol / nm to remove clashes between atoms.

holoCaM-CaMKII complex: We used the crystal structure of calcium-bound CaM-CaMKII complex (PDB ID: 1CDM) for Ca^{2+} binding free energy calculations.

holoCaM: We used the crystal structure of holoCaM (PDB ID: 1CLL) for assessing the Ca^{2+} binding free energy in the absence of the CaMBT as a reference.

IV.2 Protonation of the initial structures for pulling simulations

The charge distribution impacts accurate estimation of the binding free energy of the Ca^{2+} ions. We first took into consideration of the pH and ionic strength by using H++ server [\(18\)](#page-63-1) to predict the protonation states of the all-atomistic structure. We input the following parameters to $H++$ as used in the experiment: $pH = 7.4$, ionic strength I = 0.15 M, the external dielectric constant is 78.4 and the internal dielectric constant is 4. After we obtained the protonated states of all the residues, we performed an energy minimization using steepest descent algorithm and the tolerance for the maximum force was set to 500 kJ / mol / nm. Then we assigned partial charges according to the geometry of the proteins (input is the protonated structure in PDB format) by using a semi-empirical quantum chemistry program MOPAC [\(19\)](#page-63-2). We applied those protonated protein structures and partial charges for further all-atomistic calculations.

IV.3 Steered molecular dynamics simulations

Classical molecular dynamics simulations were carried out using GROMACS molecular dynamics package (version 5.0.6) [\(23\)](#page-63-6). We used the AMBER99SB-ILDN force field [\(24\)](#page-63-7) except the charge assignment. The partial charges were generated from MOPAC as explained in the

previous session. The rigid three-site TIP3P model [\(25\)](#page-63-8) was used for water molecules. We neutralized the system with $Na⁺$ and Cl⁻ ions and maintained an ionic strength of 150 mM.

The size of the box is about 10 nm x 10 nm x 10 nm. The proteins are placed at least 1 nm away from the edges of the cubic box. Periodic boundary conditions were employed to mimic the macroscopic settings for electrolytes. Electrostatic interactions between periodic images were treated using the particle mesh Ewald approach [\(26\)](#page-63-9), with a grid size of 0.16 nm, fourth-order cubic interpolation and a tolerance of 10^{-5} . A cutoff of 10 Å was used for van der Waals interactions and real space Coulomb interactions as well as for updating neighbor lists.

We adopted a gradient descent algorithm for energy minimization. Then we gradually heated the system temperature by 50 K per 0.1 ns to 298.15 K using NVT ensemble. We carried out 1 ns NPT equilibration with heavy atoms of the proteins (including the four Ca^{2+} ions) fixed. The proteins as well as the Ca^{2+} ions were afterwards released and were further equilibrated for another 1 ns. All NPT simulations maintained a constant pressure of 1 bar and temperature of 298.15 K using the Parrinello-Rahman barostat [\(27\)](#page-63-10). The bond lengths in proteins were constrained using the LINCS algorithm of Hess [\(28\)](#page-63-11). The equation of motion was integrated using a 2-fs time steps. As discussed in the main text, the positions of the Ca^{2+} ions in the bound state change during the minimization and equilibration stage, which yields inaccurate estimation of the binding free energy, therefore, we froze the positions of the Ca^{2+} ions as well as the backbone heavy atoms of the protein (or protein complex) during these preparation stage. They were free to move afterwards in the pulling simulations.

For each of the initial structures including 1 structure for each of holoCaM and holoCaM-CaMKII and four selected structure of holoCaM-Ng₁₃₋₄₉, I pulled the Ca²⁺ from site III and IV independently to 2 nm away where the interaction between the Ca^{2+} and the corresponding Ca^{2+}

binding loop is negligible. During the pulling simulations, the positions of the proteins may shift and this would cause inaccurate estimation of the distance between Ca^{2+} and the corresponding Ca^{2+} binding loop. Therefore, we fixed the C_α atom of the 100th or 136th residue in CaM, which showed smallest RMSD in a separate equilibration simulation, while pulling Ca^{2+} from the binding site III or IV, respectively. The force constant k and speed of pulling ν of the reference position are described in the next section. The direction of the pulling force was randomly assigned and pointed away from the center of mass of CaM to avoid clashes between CaM and the Ca²⁺. The pulling direction was selected if the angle Ω between the pulling vector and the vector connecting center of mass of CaM and center of mass of the corresponding Ca^{2+} binding loop was within 90 degrees since $\Omega > 90$ leads to a large work (as shown in Fig. S12). The displacement of the Ca^{2+} ion and the pulling forces were output every 20 fs for calculation of the work as shown in Eq. S9. The coordinates and velocities of the system were saved every 1 ps.

The setup of the pulling simulations is illustrated in Fig. S9. The Ca^{2+} is constrained to a reference position that is moving along \vec{x} direction at a speed of $|\vec{v}| = 0.001$ nm/ps. The force constant of the spring *k* is set to 10000 kJ / mol / nm² to guarantee that the Ca^{2+} strictly follows the reference position. Therefore, the pulling force is calculated as,

$$
\vec{f}(t) = -k(\vec{x}(t) - \vec{v}t - \vec{x}_0)
$$
 Eq. (S9)

where \vec{x} (\vec{x}_0) is the instantaneous (initial) displacement of the Ca²⁺ from the center of mass of the Ca^{2+} binding loop.

IV.4 Justification of parameters used in steered molecular dynamics simulations

We set the pulling speed ν and the spring constant of the external force k to effectively estimate Ca^{2+} binding free energy from steered molecular dynamics. We explored different

combinations of spring constant and pulling speed to rationalize the parameter setting. For example, to pull the Ca^{2+} from site III of holoCaM, we used $k = 100$, 1000, 10000, 100000 kJ/mol/nm² and $v = 0.001$, 0.01, 0.02 nm/ps. The simulation time was 4 ns, 0.4 ns and 0.2 ns to ensure that Ca^{2+} was pulled the same distance. In Fig. S10A, with fairly slow pulling speed, using $k = 1000$, 10000 and 10000 kJ/mol/nm², the Ca²⁺ follows the motion of the dummy bead. For $k = 100 \text{ kJ/mol/nm}^2$, the pulling force is so weak to pull the Ca²⁺ out during the entire simulation. For $k = 1000 \text{ kJ/mol/nm}^2$, there exists a lag in the beginning of the simulation, indicating a favored interaction between CaM and $Ca²⁺$ at site III. However, the tail that deviates from the straight line indicates that the thermal fluctuations in the unbound state of Ca^{2+} dominate its motion. By increasing *v* to 0.01 or 0.02 nm/ps, $k = 1000$, 10000 or 100000 kJ/mol/nm² meet the stiff spring approximation, shown in Figs. S9B and S9C. Comparing to $v =$ 0.001 nm/ps, the displacement curves are much smoother.

During the exploration of factors that influence work values, we found that the pulling speed ν is more sensitive. Jarzynski's group [\(29\)](#page-63-12) showed that JE does not depend on the pulling speed. In their study, they showed for the same amount of simulation time, the estimation of the chemical potential of a simple Lennard-Jones fluid did not vary for different switching time (comparable to pulling speed v). However, in practice, the free energy calculation is more efficient with fewer slow pulling trajectories than more fast pulling trajectories given the same amount of simulation time [\(39\)](#page-64-0). For a complex system of pulling Ca^{2+} form a holoCaM-CaMBT compound, the efficiency of convergence of distribution of work values depend on the pulling speed ν . Moreover, it has been shown that slower pulling speed ν reduces the perturbation from the pulling force compared to the level of the thermal fluctuations of the binding site [\(30\)](#page-63-13). Therefore, in this study, we set a relatively slow pulling speed $v = 0.001$ nm/ps with a stiff spring constant k = 10,000 kJ/mol/nm² to guarantee that the Ca^{2+} follows the moving reference bead that accounts for small work. To improve the efficiency of the free energy estimation, furthermore, we also employed cumulative integral (CI) extrapolation method developed by Zuckerman's group [\(31\)](#page-63-14), which is discussed in the next session.

In Fig. S11, the typical force profiles lead to the same conclusion. At $v = 0.001$ nm/ps, *k* $= 10,000$ and 100,000 kJ/mol/nm² demonstrated the same trend in the force profiles except more thermal fluctuations in the latter, indicating a converged behavior of the dissociation of Ca^{2+} in this parameter range. Therefore, $k = 10,000 \text{ kJ/mol/nm}^2$ were used for the pulling simulations.

IV.5 Estimation of binding free energy using Jarzynski's equality and Cumulative Integral extrapolation

After obtaining the work (*w*) distribution from the pulling simulations as explained in IV.3 in the *Supporting Information*, we estimated the binding free energy $\Delta G^{\text{holoCaM-CaMBT}}$ between a one- Ca^{2+} -missing complex of holoCaM-CaMBT (G_U) and a Ca²⁺-loaded complex of holoCaM-CaMBT (G_B) according to JE [\(32\)](#page-63-15),

$$
\Delta G^{\text{holoCaM-CaMBT}} = G_{\text{B}} - G_{\text{U}}
$$

$$
= k_B T ln \left\langle e^{-\frac{W}{k_B T}} \right\rangle
$$

$$
\approx k_B T ln \sum_{n=1}^{M} e^{-w_n / k_B T} / M
$$
 Eq. (S10)

where k_B is Boltzmann constant, T is the temperature, M is number of pulling simulations (M = $100 \sim 150$, w_n is the work in nth pulling calculated as,

$$
w_n(\mathbf{B} \to \mathbf{U}) = \int_0^t \overrightarrow{f_n}(\tau) \cdot \vec{v}_n d\tau
$$
 Eq. (S11)

B and U stands for bound and unbound states of the Ca^{2+} ion, respectively.

However, accurate estimation of the binding free energy ΔG relies on the sufficient sampling of small work [\(32\)](#page-63-15). Zuckerman's group [\(31\)](#page-63-14) showed that using a cumulative integral (CI) to extrapolate the free energy estimation to $n \rightarrow \infty$ can help reduce the required simulation data by 5 to 40 fold [\(30,](#page-63-13) [31\)](#page-63-14), where n is the number of work values. In Table S6, we showed the estimation of change in Ca^{2+} binding free energy ($\Delta\Delta G$) by using JE or CI estimation. We found that JE or CI produced the same signs of the $\Delta\Delta G$, however, the magnitudes were not the same. The discrepancy in the free energy calculation lies in the requirement of sufficient sampling of small work for direct JE calculations, which can be overcome by CI extrapolation [\(28\)](#page-63-11).

V. Analyses

V.1 Calculation of model free order parameter S2

A number of $N \approx 500$ representative coarse-grained structures were selected from the free energy surface of unbound Ng protein using importance sampling method [\(17\)](#page-63-0). The NMR model-free order parameter S_i^2 [\(33\)](#page-63-16) for the backbone N-H bond of residue *i* is calculated on the selected structures by relating to the Root Mean Square Fluctuations (RMSF) of the C_{α} bead through the following empirical relation [\(34\)](#page-63-17),

$$
S_i^2 = 1 - 0.5 * \ln(1 + \frac{RMSF_i}{23.6\text{\AA}} * 10.0)
$$
 Eq. (S12)

$$
RMSF_i = \sqrt{\frac{1}{N} \sum_{j=1}^{N} (r_i^j - \frac{1}{N} \sum_{k=1}^{N} r_i^k)^2}
$$
 Eq. (S13)

where r_i^j is the position of the C_α bead of residue *i* at frame *j* and *N* is the total number of frames. RMSF*ⁱ* was computed after structural alignment.

V.2 Definition of helicity

The fraction of helical structure, or helicity, of Ng (G25-A42) was calculated using the following formula [\(35\)](#page-64-1),

$$
f_H = \frac{1}{N-3} \sum_{i=1}^{N-3} < \Theta(\Delta \phi - |\phi_i - \phi_0|) > \tag{S14}
$$

where $\Theta(x)$ is the Heaviside step function taking value 1 if $x \ge 0$ and value 0 otherwise. *N* is the total number of residues, *i* is the residue index, ϕ_i is the dihedral angle about the residues $i-i+3$ from the simulation, $\phi_0 = 49.50^{\circ}$ is the dihedral angle in a perfect alpha helix and $\Delta \phi = 12.07^{\circ}$. <•••> denotes ensemble average.

V.3 Definition of Z

We defined an order parameter Z to calculate the total number of intermolecular contacts between apoCaM and Ng13-49. For each residue *i* from apoCaM and residue *j* from Ng13-49, the sidechain-sidechain and backbone-backbone contacts $(Z_{ij}^{ss}$ and Z_{ij}^{bb} , respectively) are determined as

$$
Z_{ij}^{ss} = \Theta\left(c - \frac{d_{ij}^{ss}}{\rho_i + \rho_j}\right) \qquad \qquad \text{Eq. (S15)}
$$

$$
Z_{ij}^{bb} = \Theta(c - \frac{d_{ij}^{bb}}{\rho_{HB}})
$$
 Eq. (S16)

 $d_{ij}^{ss}(d_{ij}^{bb})$ is the distance between side-chain (backbone) beads of residue *i* of apoCaM and residue *j* of Ng₁₃₋₄₉ in simulation, ρ_i (ρ_j) is van der Waals radius of residue *i(j)*, $\rho_{HB} = 4.66 \text{ Å}$ is the typical length of a hydrogen bond and cutoff $c = 1.3$. The total number of contacts Z is the summation of backbone-backbone and sidechain-sidechain contacts over all residues $Z =$ $\sum_{ij} (Z_{ij}^{ss} + Z_{ij}^{bb}).$

V.4 Sampling protein configuration for structure analyses

The structures from US simulations are biased and not appropriate for structural analyses directly. We therefore used the Boltzmann reweighting method to sample an ensemble of unbiased structures. The probability of selecting a configuration x is

$$
w(x) = \begin{cases} 1 & \text{if } p(x) \ge \rho \\ 0 & \text{else} \end{cases} \quad \text{where } p(x) = e^{-G(d_{COM}(x))/k_B T} \quad \text{Eq. (S17)}
$$

 $p(x)$ is the probability of the configuration *x* in reweighted ensemble; $G(d_{COM})$ is the reweighted/unbiased free energy obtained from the WHAM analyses along reaction coordinate d_{COM} (G_{min} is scaled to 0); ρ is a random number generated between 0 and 1; k_B is Boltzmann constant and *T* is the temperature. The ensemble generated after Boltzmann reweighting thus obeys the canonical distribution and is employed for subsequent analyses. Each structure from the biased ensemble was challenged by acceptance probability $w(x)$ and a total of 23,722 structures were sampled. This sample achieves a distribution of P^{sample} (d_{COM}). In order to assess the quality of the sampling, we computed the surprisal value compared with the original unbiased distribution $P^{\text{ori}}(d_{\text{COM}})$ defined by the following formula [\(36\)](#page-64-2),

$$
surprisal = \sum_{r} -P^{ori}(r) \ln \left[P^{sample}(r) / P^{ori}(r) \right]
$$
 Eq. (S18)

where the summation is over all the order parameter r (in this case d_{COM}). A surprisal value of 0.14 ensures that the sampled ensemble can well represent the original distribution.

V.5 Self-organized clustering algorithm

We applied the self-organized neural-net clustering algorithm [\(37-39\)](#page-64-3) to determine the structures of the apoCaM-Ng13-49 complexes from the umbrella sampling simulations. In this clustering method a vector with *M* elements represents each conformation *j*, $x_j = [x_{1j}, x_{2j}, ..., x_{Mj}]$, where $j = 1, 2, ... N$ and N is the number of conformations selected for clustering analysis. The element x_{ij} ($i = 1, 2, ..., M$) stands for the Euclidean distance between side-chains of the

polypeptide chain in the conformation *j*. To partition the various conformations into distinct clusters, the clusters are described by the cluster center and the size of the cluster is determined by a radius R_c . A given conformation is assigned to cluster k if the distance between the vector x_j and the center of the k^{th} cluster,

$$
C_k = \frac{1}{M_k} \sum_{j=1}^{M_k} x_j,
$$
 Eq. (S19)

where $C_k = [C_{1k}, C_{2k}, ..., C_{Mk}]$ and M_k is the total number of conformations in C_k .

Conformation *j* belongs to C_k if the Euclidean distance between conformation *j* and the cluster center *k*, $d_{jk} = |x_j - C_k| < R_c$, where R_c is a preassigned value. In the current study we used the native contact pairs from the unbound structure of apoCaM and Ng₁₃₋₄₉ as the M elements and a following cutoff distance $R_c = 25 \sigma (\sigma = 3.8 \text{ Å})$ to categorizes the structures into seven distinct clusters.

V.6 Correlation between intra- and inter- molecular interactions in apoCaM-Ng13-49

binding

In order to better understand the relation between interactions within the two proteins and intermolecular contacts, we built a correlation map between contacts (Fig. S4). The correlation between two contact pairs m and n is computed as in the following equation,

$$
corr_{mn} = \frac{q_m q_n > - \langle q_m \rangle}{\sqrt{\langle q_m^2 \rangle - \langle q_m \rangle^2} \sqrt{\langle q_n^2 \rangle - \langle q_n \rangle^2}}
$$
 Eq. (S20)

where $q_{m(n)}$ is the contact state of the contact pair m (n): 1 if m (n) is a contact, 0 if not. The list of contacts are provided in Table S8. The correlation between those contacts falls into several categories.

- a) Correlation: as shown in diagonal, the contact formation within apoCaM, between apoCaM and Ng_{13-49} , and within Ng_{13-49} are correlated.
- b) Anti-correlation: contact formation between nCaM and Ng₁₃₋₄₉ anti-correlates with contact formation within nCaM; contact formation between cCaM and Ng13-49 anticorrelates contact formation within cCaM. The anti-correlation tells that the contact formation within nCaM or that within cCaM competes with its interaction with Ng13-49. We infer that binding of the target is responsible for the repacking of the nCaM or cCaM by direct competitive interaction with the Ca^{2+} binding loops.
- c) No correlation: there is no apparent correlation among contact formation within nCaM and contact formation within cCaM, indicating the two domains of CaM are relative independent.
- d) Mixed correlation: contact formation between nCaM and cCaM has a mixed correlation with contact formation between $nCaM$ and Ng_{13-49} as well as contact formation between cCaM and Ng_{13-49} .

Supporting Figures

Fig. S1. Overview of the effects of CaM-dependent Ca2+ signaling and effects of CaM binding targets (CaMBTs) on changes in synaptic plasticity. Many of the effects of intracellular Ca^{2+} on synaptic plasticity are mediated through CaM -regulated proteins. Increase in intracellular Ca^{2+} , generated through the activity of NMDA (N-methyl-Daspartate) receptors or voltage-sensitive Ca^{2+} channels, results in the release of CaM that is bound to Ng. CaM mediates the Ca^{2+} stimulation of CaMKII which is required for changes in synaptic plasticity. The structures of Ca^{2+} -free CaM (apoCaM, PDBID: 1CFD) and Ca^{2+} -CaM-CaMKII peptide (PDB ID: 1CDM) are provided. CaM is colored as follows, red \rightarrow nCaM (residue 1 to 76), gray \rightarrow central linker (residue 77 to 82), blue \rightarrow cCaM (residue 83 to 148) and the CaMKII peptide is colored in green.

Fig. S2 Reweighted potential of mean force of apoCaM and Ng13-49. The PMF was reweighted from umbrella sampling simulations using WHAM at varying scaling factors of the inter-molecular nonbonded interaction λ (excluding electrostatic interactions). KB statistical dihedral angle potential was employed. d_{COM} is the distance between center of mass of apoCaM and center of mass of Ng₁₃₋₄₉. σ equals 3.8 Å. T = 1.1 ε / k_B. ε = 0.6 kcal / mol.

Fig. S3 Difference in probability of contact formation between the bound ensemble of the most dominant cluster (cluster 1) and the unbound ensemble. The difference maps for apoCaM intramolecular (a), Ng_{13-49} intramolecular (b) and apoCaM-Ng₁₃₋₄₉ intermolecular (c) contacts are provided. $d_{COM} = 20.0$ s is used for the unbound state. The schematic representation of the helices of apoCaM and Ng13-49 are provided along the axes. The representative structure, which has minimal root mean square deviation from the averaged structure, is in ribbon representation and colored according to the secondary structures: nCaM is in gray, cCaM is in green, acidic region of Ng_{13-49} is in pink and IQ motif of Ng_{13-49} is in orange. The regions of particular interest are encircled by ellipses, rectangles, circles and triangles. The color bar represents the difference in probability of contact formation ΔP between the bound and the

unbound conformations where positive ΔP indicates an increase of contact probabilities in the bound ensemble.

Fig. S4 Correlation map between contacts formation involving apoCaM and Ng13-49 in bound ensemble (cluster 1). The contact pair list is provided in Table S8. The blue color shows strong anti-correlation between involved contact pairs; the red color shows strong correlation; white means no correlation.

Fig. S5 Illustration of the EF-hand β**-scaffold.** The structure is from C-terminus of the crystal structure of holoCaM (PDB: 1CLL). The calcium ions are shown in yellow beads. The $5th$ (+Z) coordinating ligand) and $12th$ (-Z coordinating ligand) residues in the Ca²⁺ binding loops are shown in ball-and-stick representation.

Fig. S6 Probability of contact formation in unbound, encounter and bound ensemble of apoCaM-Ng13-49. In panel (a), the probability map of contact formation for the unbound state was calculated using the ensemble when apoCaM and Ng₁₃₋₄₉ are well separated at $d_{COM} = 20.0$ σ. σ = 3.8 Å. (b) the probability map of contact formation for the encounter complex at d_{COM} = 10.0 σ and (c) in the bound state ($d_{COM} = 2.8$ σ) are provided. The upper and lower triangles show contacts maps for backbone-backbone conatcts and sidechain-sidechain contacts, respectively. The schematic representation of the helices of apoCaM and Ng13-49 are provided along the axes. The region encircled by a circle marks the interaction between the two β-strands in the EF-hand β-scaffold. The color bar represents the probability of contact formation P.

Fig. S7 Calculated helicity of the Ng protein fragment with residual structure matches with experiment using three dihedral angle potentials. The helicity of the fragment G25-A42 was computed from REMD simulations of the intact Ng protein with the coarse-grained side-chain C_{α} model. Three models of dihedral angle potential were used: (A) Structure-Based (SB) potential; (B) Hybrid potential; (C) Karanicolas-Brooks (KB) statistical potential. The two gray lines mark the upper and lower limit of the helicity of the segment from the experiment [\(3\)](#page-62-2); the yellow shade marks the corresponding reduced temperature 1.1 ε / k_B . Helicity was estimated based on the dihedral angles between four consecutive C_{α} beads (see definition in V.2 in *Supporting Information*).

Fig. S8 Illustration of protein reconstruction from coarse-grained configuration to atomistic configuration for the holoCaM-Ng13-49 complex. The backbone of CaM and Ng13-49 is shown in ribbon representation. CaM is colored from red (N-Domain) to blue (C-Domain); Ng₁₃₋₄₉ is shown in magenta. (left) The Sidechain-C_{α} model (SCM), where the backbone atoms are represented by the C_{α} beads (blue balls) and sidechain atoms are represented by the sidechain beads (red balls). (right) All-atomistic (AA) configuration of the holoCaM-Ng13-49 complex with Ca^{2+} (yellow balls) added to the Ca^{2+} binding loops. The other atoms are colored according to the atom names (oxygen atoms are in red, carbon in cyan, hydrogen in white, nitrogen in blue, etc.).

Fig. S9 Schematic illustration of pulling Ca2+ from one of the calcium-binding loop of CaM. The rest of the system including CaM and the solvent molecules is not shown in this illustration for better visualization.

Fig. S10 The displacement of the Ca^{2+} from the Ca^{2+} binding loop during the pulling **simulations of** \bar{Ca}^{2+} **from the** Ca^{2+} **binding site III** of holoCaM. The unit of *k* is kJ/mol/nm² and the unit of *v* is nm/ps. (A) (B) (C) show the case pulling speed $v = 0.001$, $v =$ 0.01 and $v = 0.02$ nm/ps.

Fig. S11 The profile of external force in pulling simulations of Ca^{2+} from the Ca^{2+} binding **site III of holoCaM.** The pulling speed $v = 0.001$ nm / ps. The unit of spring constant *k* is kJ/mol/nm² and the unit of pulling speed ν is nm/ps.

Fig. S12 Work done by external pulling force of pulling Ca2+ at different pulling angles. At each pulling angle zone Ω in [0, 30], [30, 60], [60, 90], [90, 120], [120, 150], [150, 180], the $Ca²⁺$ were pulled from site IV of holoCaM for 100 times. The average work, deviation (upper bar) as well as standard error (lower bar) are provided. The pulling speed $v = 10$ nm / ns and the spring constant of the pulling force $k = 1000 \text{ kJ} / \text{mol} / \text{nm}^2$.

Fig. S13 The running Jarzynski equality (JE) estimate and subsampled block-averaged (BA) estimate are plotted as a function of the number of trajectories used in the estimate. The binding free energies of Ca^{2+} at site III and IV of holoCaM using the running JE method and the BA method are shown in panel (a) and the binding free energies of Ca^{2+} at site III and IV of holoCaM-Ng₁₃₋₄₉ (using the conformation shown in Fig. 5) are shown in panel (b).

Figure S14 Distribution of work for dissociation of Ca2+. (a) and (b) represent the holoCaM and holoCaM-Ng13-49, respectively. (c-f) represent the four conformations of holoCaM-Ng13-49.

Fig. S15 Sequence alignment of neurogranin (Ng). 45 sequences of Ng (see Table S14 for the details of species associated with each sequence) were selected from UniProtKB (www.uniprot.org) by searching the keyword "neurogranin". The sequences were aligned using the "Align" tool in UniProt and visualized in Jalview (40) . The ClustalX color scheme in Jalview was used to highlight the residues and shading intensity of the color is based on the conservation over all the sequences in the alignment. The "acidic region" and the "IQ motif" of Ng are indicated by the grey and black line, respectively. The "DDPG (D/E/A)" motif is indicated inside the black area. The sequence of Ng from mouse (UniProt id P60761) that was used in this study is shown inside the red dotted area.

Fig. S16 Distance map between the Ca²⁺ binding loops (III and IV) of CaM and the "DDPG" motif of Ng. The distance here refers to the closest distance between the corresponding residues. The Ca^{2+} ions (shown as yellow spheres) coordinate (indicated by grey dotted lines) with the first, third, fifth, seventh, ninth (through water molecule (w), shown as a blue sphere) and twelfth residue of the loops. The distance to distance map is based on the holoCaM-Ng13-49 complex structure in Fig. 5a in the main text.

Fig. S17 Illustration of the changes in the binding free energy ΔG and ΔΔG from pH = 7.4 to pH = 6.8. $\Delta G = G_B - G_U$. $\Delta \Delta \tilde{G} = \Delta G^{holoCaM-CaMBT}$ - $\Delta G^{holoCaM}$. B and U stand for bound and unbound states of the Ca²⁺, respectively. The arrows show increase of ΔG for Ca²⁺ in all the three systems.

Supporting Tables

Table S1 Dissociation constant K_d of apoCaM and Ng_{13-49} with several values of strength of **non-electrostatic intermolecular hydrogen bonding and van der Waals interactions λ.** λ is shown in Eq 4. Experimentally measured $K_d = 680$ nM [\(4\)](#page-62-3).

Table S2 The population of major clusters from computer simulations and the correlation coefficient of their computational "apparent chemical shifts" with the data from NMR experiments [\(4\)](#page-62-3).

Table S3 EF-hand angles in several forms of CaM. Definition of EF hand angles is provided in IV.1 from the *Supporting Information*.

Table S4 The composition of the three dihedral potentials for modeling Ng or Ng13-49. SB stands structure-based and KB stands for Karanicolas-Brooks statistical potential [\(12\)](#page-62-11).

Table S6 Difference in binding free energy of Ca2+ (ΔΔG) calculated from non-equilibrium molecular simulations and from the experiments at $pH = 7.4$ **.** Direct use of Jarzynski's equality and a cumulants integral extrapolation was used in the calculation of binding free energy of Ca^{2+} from the simulations.

System		holoCaM-Ng ₁₃₋₄₉ (average)	holoCaM-CaMKII
$\Delta\Delta G$ (kcal/mol) Jarzynski's equality	Site III	8.5 ± 2.7	-0.1
	Site IV	23.9 ± 1.1	-3.4
$\Delta\Delta G$ (kcal/mol) CI extrapolation	Site III	9.2 ± 2.2	-0.5
	Site IV	22.4 ± 0.9	-7.0
$\Delta\Delta G$ (kcal/mol) experiment	Site III/IV	$+2.5$	-3.3

Table S7 Difference in binding free energy of Ca2+ (ΔΔG) for Ca2+ at site III calculated from non-equilibrium molecular simulations at $pH = 6.8$ **.** Direct use of Jarzynski's equality and a cumulants integral extrapolation was used in the calculation of binding free energy of Ca^{2+} from the simulations.

Table S8 Contact pair list used for correlation analysis. The non-specific contact pairs are determined from the difference contact map analysis. A contact pair is selected if the magnitude of the change in the probability of the contact pair in cluster 1 (Fig. S3) is greater than 0.2 from unbound state to the bound state. The contact pairs are listed in a sequence of categories: within nCaM, between nCaM and cCaM, within cCaM, between nCaM and Ng13-49, between cCaM and Ng_{13-49} and within Ng_{13-49} . The sequence of apoCaM is from NMR structure (PDB: 1CFD) and sequence of Ng13-49 is provided in the *Materials and Methods* in the main text.

Residue Index	Residue	Charge on	Error of the	Charge on side-chain	Error of the
1	Name MET	C_{α} 0.847	C_{α} charge 0.002	0.091	side-chain charge 0.001
\overline{c}	ASP	-0.028	0.002	-0.898	0.001
3	CYS	-0.037	0.002	0.035	0.001
$\overline{\mathcal{L}}$	CYS	-0.037	0.002		0.001
				0.033	
5	THR	-0.098	0.002	0.083	0.001
6	GLU	-0.052	0.002	-0.872	0.002
$\overline{7}$	SER	-0.098	0.002	0.066	0.001
8	ALA	-0.022	0.002	0.049	0.001
9	CYS	-0.042	0.002	0.025	0.001
10	SER	-0.077	0.002	0.062	0.001
11	LYS	-0.061	0.003	1.000	0.001
12	PRO	-0.307	0.003	0.316	0.002
13	ASP	-0.010	0.002	-0.926	0.002
14	ASP	-0.037	0.002	-0.927	0.002
15	ASP	-0.040	0.002	-0.930	0.002
16	ILE	-0.070	0.002	0.078	0.001
17	LEU	-0.043	0.002	0.061	0.001
18	ASP	-0.053	0.002	-0.925	0.002
19	ILE	-0.073	0.003	0.080	0.001
20	PRO	-0.301	0.003	0.305	0.002
21	LEU	-0.066	0.002	0.063	0.001
22	ASP	-0.034	0.002	-0.930	0.002
23	ASP	-0.035	0.003	-0.929	0.002
24	PRO	-0.335	0.003	0.320	0.002
25	GLY	-0.182	0.002	0.188	0.001
26	ALA	-0.041	0.002	0.063	0.001
27	ASN	-0.047	0.003	0.011	0.001
28	ALA	-0.043	0.002	0.067	0.001
29	ALA	-0.064	0.002	0.064	0.001
30	ALA	-0.068	0.002	0.068	0.001
31	ALA	-0.057	0.002	0.069	0.001
32	LYS	-0.079	0.002	1.020	0.002
33	ILE	-0.089	0.002	0.103	0.001
34	GLN	-0.077	0.002	0.069	0.001
35	ALA	-0.046	0.002	0.064	0.001
36	SER	-0.091	0.002	0.079	0.001
37	PHE	-0.052	0.002	0.062	0.001
38	ARG	-0.079	0.002	1.027	0.002
39	GLY	-0.181	0.002	0.192	0.001
40	HIS	-0.049	0.002	0.049	0.001
41	MET	-0.069	0.002	0.069	0.001
42	ALA	-0.062	0.002	0.068	0.001
43	ARG	-0.062	0.002	1.020	0.002
44	LYS	-0.074	0.002	1.031	0.002

Table S9 Charge distribution on a coarse-grained side-chain C^α model of full length Ng.

Table S10 Charge distribution on a coarse-grained side-chain C^α model of apoCaM at pH = 7.2, I = 0.15 M.

Residue	Residue	Charge on	Error of the	Charge on	Error of the
Index	Name	C_α	C_{α} charge	side-chain	side-chain charge
	ALA	1.564	0.004	0.293	0.002
2	ASP	-0.038	0.003	-0.889	0.002
3	GLN	-0.050	0.003	0.051	0.001
$\overline{4}$	LEU	-0.068	0.002	0.081	0.001
5	THR	-0.093	0.002	0.073	0.001
6	GLU	-0.051	0.002	-0.915	0.002

Table S11 Charge distribution on a coarse-grained side chain C_α **model of Ng₁₃₋₄₉ at pH = 7.2, I = 0.15 M.**

Residue	Residue	Charge on	Error of the	Charge on	Error of the
Index	Name	Ca	Ca charge	side-chain	side-chain charge
13	ASP	1.574	0.004	-1.659	0.004
14	ASP	-0.020	0.003	-0.921	0.002
15	ASP	-0.034	0.003	-0.911	0.002
16	ILE	-0.071	0.003	0.081	0.001
17	LEU	-0.060	0.002	0.061	0.001
18	ASP	-0.037	0.003	-0.916	0.002
19	$\rm ILE$	-0.539	0.004	0.540	0.002
20	PRO	-0.265	0.003	0.258	0.003
21	LEU	-0.054	0.003	0.064	0.001
22	ASP	-0.042	0.003	-0.931	0.002
23	ASP	-0.818	0.004	-0.133	0.003
24	PRO	-0.282	0.003	0.261	0.002
25	GLY	-0.181	0.003	0.187	0.001
26	ALA	-0.032	0.002	0.054	0.001
27	ASN	-0.046	0.003	0.005	0.001
28	ALA	-0.040	0.003	0.064	0.001
29	ALA	-0.059	0.003	0.060	0.001
30	ALA	-0.066	0.003	0.062	0.001
31	ALA	-0.055	0.003	0.064	0.001
32	LYS	-0.079	0.003	0.996	0.001
33	ILE	-0.089	0.003	0.100	0.001
34	GLN	-0.072	0.002	0.059	0.001
35	ALA	-0.042	0.003	0.062	0.001
36	SER	-0.092	0.003	0.076	0.001
37	PHE	-0.048	0.003	0.057	0.001
38	ARG	-0.077	0.002	1.002	0.002
39	GLY	-0.175	0.002	0.187	0.001

Table S12 Charge distribution on a coarse-grained side chain C^α model of apoCaM at pH = 6.3, I = 0.1 M.

Residue	Residue	Charge on	Error on the	Charge on the	Error on the
Index	Name	C_{α}	C_{α} charge	side-chain	side-chain charge
1	ALA	1.565	0.004	0.292	0.001
$\overline{2}$	ASP	-0.041	0.003	-0.885	0.002
$\overline{3}$	GLN	-0.052	0.002	0.053	0.001
$\overline{4}$	LEU	-0.066	0.002	0.080	0.001
5	THR	-0.095	0.002	0.075	0.001
6	GLU	-0.051	0.002	-0.914	0.002
$\overline{7}$	GLU	-0.072	0.003	-0.871	0.002
8	GLN	-0.040	0.003	0.053	0.001
9	ILE	-0.098	0.003	0.088	0.001
10	ALA	-0.061	0.003	0.058	0.001
11	GLU	-0.061	0.003	-0.868	0.002
12	PHE	-0.047	0.003	0.047	0.001
13	LYS	-0.090	0.003	0.998	0.001
14	GLU	-0.055	0.003	-0.893	0.002
15	ALA	-0.035	0.002	0.058	0.001
16	PHE	-0.055	0.003	0.053	0.001
17	SER	-0.118	0.003	0.070	0.001
18	LEU	-0.060	0.002	0.076	0.001
19	PHE	-0.054	0.003	0.047	0.001
20	ASP	-0.049	0.003	-0.885	0.002
21	LYS	-0.078	0.002	0.986	0.001
22	ASP	-0.049	0.003	-0.914	0.002
23	GLY	-0.196	0.002	0.191	0.001
24	ASP	-0.064	0.002	-0.945	0.001
25	GLY	-0.150	0.002	0.184	0.001
26	THR	-0.062	0.003	0.065	0.001
27	ILE	-0.066	0.003	0.077	0.001
28	THR	-0.107	0.003	0.079	0.001

Table S14. Detail of the neurogranin (Ng) sequences obtained from UniprotKB (www.uniprot.org). 45 sequences of Ng with unique Uniprot ID were used for the sequence alignment (see Fig. S15).

Table S14(a)

Table S14(b)

References

- 1. Jurado, L. A., P. S. Chockalingam, and H. W. Jarrett. 1999. Apocalmodulin. Physiological reviews 79:661-682.
- 2. Bahler, M., and A. Rhoads. 2002. Calmodulin signaling via the IQ motif. FEBS Lett 513:107-113.
- 3. Ran, X., H. H. Miao, F. S. Sheu, and D. Yang. 2003. Structural and dynamic characterization of a neuron-specific protein kinase C substrate, neurogranin. Biochemistry 42:5143-5150.
- 4. Hoffman, L., A. Chandrasekar, X. Wang, J. A. Putkey, and M. N. Waxham. 2014. Neurogranin alters the structure and calcium binding properties of calmodulin. J Biol Chem 289:14644-14655.
- 5. Sugita, Y., and Y. Okamoto. 1999. Replica-exchange molecular dynamics method for protein folding. Chem Phys Lett 314:141-151.
- 6. Sanbonmatsu, K. Y., and A. E. Garcia. 2002. Structure of Met-enkephalin in explicit aqueous solution using replica exchange molecular dynamics. Proteins 46:225-234.
- 7. Cheung, M. S., J. M. Finke, B. Callahan, and J. N. Onuchic. 2003. Exploring the interplay between topology and secondary structural formation in the protein folding problem. J. Phys. Chem. B 107:11193-11200.
- 8. Wang, Q., P. Zhang, L. Hoffman, S. Tripathi, D. Homouz, Y. Liu, M. N. Waxham, and M. S. Cheung. 2013. Protein recognition and selection through conformational and mutually induced fit. Proceedings of the National Academy of Sciences of the United States of America 110:20545-20550.
- 9. Yang, Y., E. Faraggi, H. Zhao, and Y. Zhou. 2011. Improving protein fold recognition and template-based modeling by employing probabilistic-based matching between predicted one-dimensional structural properties of query and corresponding native properties of templates. Bioinformatics 27:2076-2082.
- 10. Debye, P., and E. Hückel. 1923. The theory of electrolytes. I. Lowering of freezing point and related phenomena. Physikalische Zeitschrift 24:185-206.
- 11. Betancourt, M. R., and D. Thirumalai. 1999. Pair potentials for protein folding: Choice of reference states and sensitivity of predicted native states to variations in the interaction schemes. Protein Sci. 8:361-369.
- 12. Karanicolas, J., and C. L. Brooks, 3rd. 2002. The origins of asymmetry in the folding transition states of protein L and protein G. Protein Sci 11:2351-2361.
- 13. Andreasen, T. J., C. W. Luetje, W. Heideman, and D. R. Storm. 1983. Purification of a novel calmodulin binding protein from bovine cerebral cortex membranes. Biochemistry 22:4615-4618.
- 14. Wang, Q., K. C. Liang, A. Czader, M. N. Waxham, and M. S. Cheung. 2011. The Effect of Macromolecular Crowding, Ionic Strength and Calcium Binding on Calmodulin Dynamics. PLoS Comput. Biol. 7:16.
- 15. Dima, R. I., and D. Thirumalai. 2004. Asymmetry in the shapes of folded and denatured states of proteins. J. Phys. Chem. B 108:6564-6570.
- 16. Camacho, C. J., and D. Thirumalai. 1993. Kinetics and Thermodynamics of Folding in Model Proteins. Proceedings of the National Academy of Sciences of the United States of America 90:6369-6372.
- 17. Samiotakis, A., D. Homouz, and M. S. Cheung. 2010. Multiscale investigation of chemical interference in proteins. J Chem Phys 132:175101.
- 18. Anandakrishnan, R., B. Aguilar, and A. V. Onufriev. 2012. H++ 3.0: automating pK prediction and the preparation of biomolecular structures for atomistic molecular modeling and simulations. Nucleic Acids Res 40:W537-541.
- 19. Stewart, J. J. 2007. Optimization of parameters for semiempirical methods V: modification of NDDO approximations and application to 70 elements. J Mol Model 13:1173-1213.
- 20. Kumar, S., D. Bouzida, R. H. Swendsen, P. A. Kollman, and J. M. Rosenberg. 1992. The Weighted Histogram Analysis Method for Free-Energy Calculations on Biomolecules .1. The Method. J Comput Chem 13:1011-1021.
- 21. Grossfield, A. 2013. WHAM: an implementation of the weighted histogram analysis method.<http://membrane.urmc.rochester.edu/content/wham> Version 2.0.4.
- 22. Roux, B. 1995. The Calculation of the Potential of Mean Force Using Computer-Simulations. Comput Phys Commun 91:275-282.
- 23. Pall, S., M. J. Abraham, C. Kutzner, B. Hess, and E. Lindahl. 2015. Tackling Exascale Software Challenges in Molecular Dynamics Simulations with GROMACS. Lect Notes Comput Sc 8759:3-27.
- 24. Lindorff-Larsen, K., S. Piana, K. Palmo, P. Maragakis, J. L. Klepeis, R. O. Dror, and D. E. Shaw. 2010. Improved side-chain torsion potentials for the Amber ff99SB protein force field. Proteins-Structure Function and Bioinformatics 78:1950-1958.
- 25. Jorgensen, W. L., J. Chandrasekhar, J. D. Madura, R. W. Impey, and M. L. Klein. 1983. Comparison of Simple Potential Functions for Simulating Liquid Water. Journal of Chemical Physics 79:926-935.
- 26. Darden, T., D. York, and L. Pedersen. 1993. Particle Mesh Ewald - an N.Log(N) Method for Ewald Sums in Large Systems. Journal of Chemical Physics 98:10089-10092.
- 27. Parrinello, M., and A. Rahman. 1981. Polymorphic Transitions in Single-Crystals - a New Molecular-Dynamics Method. J Appl Phys 52:7182-7190.
- 28. Hess, B., H. Bekker, H. J. C. Berendsen, and J. G. E. M. Fraaije. 1997. LINCS: A linear constraint solver for molecular simulations. J Comput Chem 18:1463-1472.
- 29. Hendrix, D. A., and C. Jarzynski. 2001. A "fast growth" method of computing free energy differences. Journal of Chemical Physics 114:5974-5981.
- 30. Zhang, D., J. Gullingsrud, and J. A. McCammon. 2006. Potentials of mean force for acetylcholine unbinding from the alpha7 nicotinic acetylcholine receptor ligand-binding domain. J Am Chem Soc 128:3019-3026.
- 31. Ytreberg, F. M., and D. M. Zuckerman. 2004. Efficient use of nonequilibrium measurement to estimate free energy differences for molecular systems. J Comput Chem 25:1749-1759.
- 32. Jarzynski, C. 1997. Nonequilibrium equality for free energy differences. Phys Rev Lett 78:2690-2693.
- 33. Lipari, G., and A. Szabo. 1982. Model-Free Approach to the Interpretation of Nuclear Magnetic-Resonance Relaxation in Macromolecules .1. Theory and Range of Validity. J Am Chem Soc 104:4546-4559.
- 34. Berjanskii, M., and D. S. Wishart. 2006. NMR: prediction of protein flexibility. Nat Protoc 1:683-688.
- 35. Kudlay, A., M. S. Cheung, and D. Thirumalai. 2009. Crowding effects on the structural transitions in a flexible helical homopolymer. Phys Rev Lett 102:118101.
- 36. Weinkam, P., E. V. Pletneva, H. B. Gray, J. R. Winkler, and P. G. Wolynes. 2009. Electrostatic effects on funneled landscapes and structural diversity in denatured protein ensembles. Proc Natl Acad Sci U S A 106:1796-1801.
- 37. Carpenter, G. A., and S. Grossberg. 1987. ART 2: self-organization of stable category recognition codes for analog input patterns. Appl Opt 26:4919-4930.
- 38. Cheung, M. S., and D. Thirumalai. 2007. Effects of crowding and confinement on the structures of the transition state ensemble in proteins. J. Phys. Chem. B 111:8250-8257.
- 39. Guo, Z., and D. Thirumalai. 1997. The nucleation-collapse mechanism in protein folding: evidence for the non-uniqueness of the folding nucleus. Folding and Design 2:377-391.
- 40. Waterhouse, A. M., J. B. Procter, D. M. Martin, M. Clamp, and G. J. Barton. 2009. Jalview Version 2--a multiple sequence alignment editor and analysis workbench. Bioinformatics 25:1189-1191.